

半导体中载流子的统计分布

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Contents

1	状态密度	3
1.1	三维情况下的自由电子气	3
1.2	状态密度	3
2	费米能级与载流子的统计分布	4
2.1	费米分布函数	4
2.2	导带电子和价带空穴浓度	4
3	本征半导体的载流子分布	5
3.1	本征载流子浓度	5
3.2	本征载流子的费米能级	5
4	杂质半导体的载流子分布	5
4.1	非补偿情形	5
4.2	补偿情形	6

5	简并半导体	7
5.1	简并的出现	7
5.2	简并半导体的载流子浓度	7
5.3	简并条件	7

1 状态密度

1.1 三维情况下的自由电子气

$$\varepsilon_{\vec{k}} = \frac{\hbar^2}{2m} k^2 = \frac{\hbar^2}{2m} (k_x^2 + k_y^2 + k_z^2), k_x, k_y, k_z = 0; \pm \frac{2\pi}{L}; \dots; \pm \frac{2n\pi}{L}$$

每个量子态占据体积为 $\frac{(2\pi)^3}{V}$ \therefore 再考虑自旋, k空间能量状态密度为 $\frac{2V}{(2\pi)^3}$

1.2 状态密度

- 为了描述能带电子状态的分布, 引入态密度 $g(E)$ 表示单位能量间隔内的状态数

$$g(E) = \frac{dZ}{dE} = \frac{dZ}{d\Omega^*} \frac{d\Omega^*}{dk} \frac{dk}{dE}$$

- 状态密度汇总

– 一维 $L^* = 2k \rightarrow 2 \left(\frac{L}{2\pi} dL^* \right) \rightarrow dZ = \frac{2L}{h} \sqrt{\frac{2m_n^*}{E-E_0}} dE$

$$\rightarrow g(E) = \frac{2\pi}{h} \sqrt{\frac{2m_n^*}{E-E_0}}$$

– 二维 $S^* = \pi k^2 \rightarrow dZ = 2 \left(\frac{L}{2\pi} \right)^2 dS^*$

$$\rightarrow g(E) = \frac{4S\pi m_n^*}{h^2}$$

– 三维 $\Omega^* = \frac{4}{3}\pi k^3 \rightarrow dZ = \frac{2V}{(2\pi)^3} \frac{4\pi}{h^3} (2m_n^*)^{3/2} (E-E_0)^{1/2} dE$

$$\rightarrow g(E) = \frac{4\pi V}{h^3} (2m_n^*)^{3/2} \sqrt{E-E_0}$$

- 状态密度有效质量

– 电子状态密度有效质量 导带极值在 $\vec{k} = \vec{k}_c$, 等能面为椭球面

$$2m_{dn}^{3/2} = M(8m_x^* m_y^* m_z^*)^{1/2} \Rightarrow m_{dn} = (M^2 m_l^* m_t^{*2})^{1/3}$$

– 空穴状态密度有效质量 重空穴与轻空穴在价带顶简并

$$(m_{dV}^*)^{3/2} = (m_{lh}^{3/2} + m_{hh}^{3/2}) \Rightarrow (m_{dV}^*) = (m_{lh}^{3/2} + m_{hh}^{3/2})^{2/3}$$

2 费米能级与载流子的统计分布

2.1 费米分布函数

- 费米分布函数 $f(E)$ $f(E) = \frac{1}{1 + \exp\left(\frac{E - E_f}{kT}\right)}$
- 玻尔兹曼分布函数 $f(E)$ ($E - E_f \gg kT$) $f(E) = \frac{1}{\exp\left(\frac{E - E_f}{kT}\right)}$
- 费米能级的物理意义 标志了电子填充水平
- $f_e(E) + f_h(E) = \exp\left(-\frac{E - E_f}{kT}\right) + \exp\left(-\frac{E_f - E}{kT}\right) = 1$

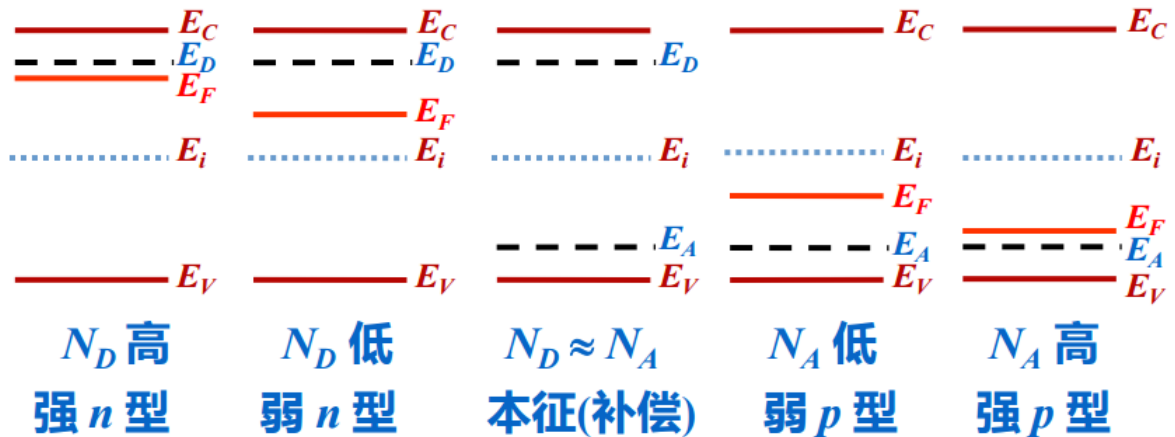


Figure 1: 不同半导体的费米能级

2.2 导带电子和价带空穴浓度

$$\text{导带有效状态密度 } N_c = \frac{2(2\pi m_{dn} kT)^{3/2}}{h^3}$$

$$\text{导带平衡电子浓度 } n = N_c \exp\left(-\frac{E_C - E_f}{k_B T}\right)$$

$$\text{价带有效状态密度 } N_v = \frac{2(2\pi m_{dp} kT)^{3/2}}{h^3}$$

$$\text{价带平衡电子浓度 } p = N_v \exp\left(-\frac{E_f - E_v}{k_B T}\right)$$

$$\therefore np = N_c N_v \exp\left(-\frac{E_c - E_v}{k_B T}\right) = N_c N_v \exp\left(-\frac{E_g}{k_B T}\right), \text{浓度乘积与 } E_f \text{ 无关}$$

3 本征半导体的载流子分布

热激发产生的载流子

3.1 本征载流子浓度

$$\text{电中性条件: } n \cdot p = n_i^2$$

$$n = p = n_i = \sqrt{N_c N_v} \exp\left(-\frac{E_g}{2k_B T}\right)$$

3.2 本征载流子的费米能级

$$E_f = \frac{E_c + E_v}{2} + \frac{k_B T}{2} \ln\left(\frac{N_v}{N_c}\right)$$

$$\frac{E_c + E_v}{2} \gg \frac{k_B T}{2} \ln\left(\frac{N_v}{N_c}\right), \therefore \text{本征费米能级 } E_i \text{ 基本上在禁带中线处}$$

4 杂质半导体的载流子分布

电子占据施主能级的几率

空穴占据受主能级的几率

$$f_D(E) = \frac{1}{1 + \frac{1}{2} \exp\left(\frac{E_D - E_f}{k_B T}\right)}$$

$$f_A(E) = \frac{1}{1 + \frac{1}{2} \exp\left(\frac{E_f - E_A}{k_B T}\right)}$$

4.1 非补偿情形

$$n = p + n_D^+$$

即导带电子浓度= 价带空穴浓度+ 电离施主浓度

$$N_c \exp\left(-\frac{E_c - E_f}{k_B T}\right) = N_v \exp\left(-\frac{E_f - E_v}{k_B T}\right) + \frac{N_D}{1 + 2 \exp\left(-\frac{E_D - E_f}{k_B T}\right)}$$

非补偿情形	电中性条件	电子浓度	费米能级
低温弱电离区	$n_0 = n_D^+$	$n_0 = \left(\frac{N_D N_C}{2}\right)^{1/2} \exp\left(-\frac{E_C - E_D}{2k_B T}\right)$	$E_F = \frac{E_C + E_D}{2} + \frac{k_B T}{2} \ln \frac{N_D}{2N_C}$
强电离区	$n_0 = n_D^+$	$n_0 = N_D$	$E_F = E_C + k_B T \ln \frac{N_D}{N_C}$
过渡区	$n_0 = p_0 + N_D$	$n_0 = N_D + \frac{n_i^2}{N_D}$	$E_F = E_i + k_B T \sinh^{-1}\left(\frac{N_D}{2n_i}\right)$
本征区	$n_0 = p_0 = n_i$	$n_0 = n_i$	$E_F = \frac{E_C + E_V}{2} + \frac{k_B T}{2} \ln \frac{N_V}{N_C}$

4.2 补偿情形

若半导体中同时含有施主、受主杂质，且施主杂质多于受主杂质，低温下，施主杂质将首先填充受主杂质，称为补偿

$N_D > N_A$	电中性条件	电子浓度	费米能级
低温弱电离区	$n_0 = n_D^+ - p_A^-$	$N_A \gg n_0 :$ $n = \frac{N_C(N_D - N_A)}{2N_A} \exp\left(-\frac{E_C - E_D}{k_B T}\right)$ $N_A \ll n_0 :$ $n_0 = \left(\frac{N_D N_C}{2}\right)^{1/2} \exp\left(-\frac{E_C - E_D}{2k_B T}\right)$	$E_F = E_D + k_B T \ln \frac{N_D - N_A}{2N_A}$ $E_F = \frac{E_C + E_D}{2} + \frac{k_B T}{2} \ln \frac{N_D}{2N_C}$
强电离区	$n_0 = N_D - N_A$	$n_0 = N_D - N_A$	$E_F = E_C + k_B T \ln \frac{N_D - N_A}{N_C}$
过渡区	$n_0 = p_0 + N_D - N_A$	$n_0 = \frac{N_D - N_A}{2} +$ $\frac{1}{2}[(N_D - N_A)^2 + 4n_i^2]^{1/2}$	$E_F = E_i + k_B T \sinh^{-1}\left(\frac{N_D - N_A}{2n_i}\right)$
本征区	$n_0 = p_0 = n_i$	$n_0 = n_i$	$E_F = E_i$

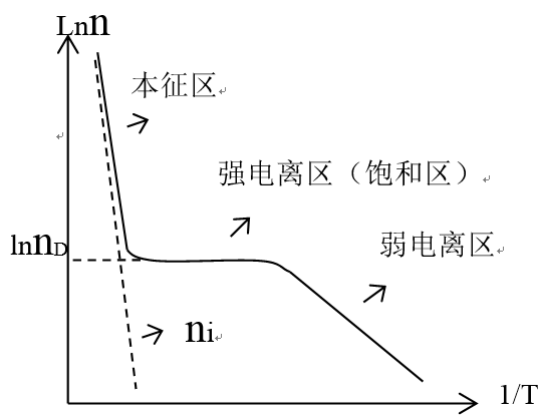


Figure 2: 浓度随温度变化曲线

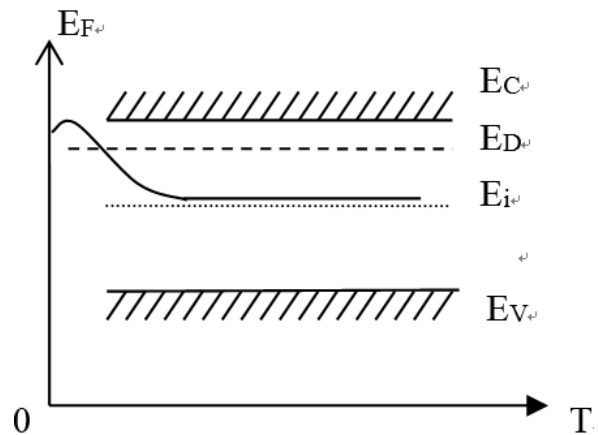


Figure 3: 费米能级随温度变化曲线

5 简并半导体

5.1 简并的出现

E_f 在导带底或价带顶附近，或 E_f 进入导带或价带称为简并情形
此时玻尔兹曼近似不再成立。

5.2 简并半导体的载流子浓度

$$\text{费米积分 } F_{1/2}(\xi) = \int_0^\infty \frac{x^{1/2}}{1 + \exp(x - \xi)} dx$$
$$n_0 = N_C \frac{2}{\sqrt{\pi}} F_{1/2}\left(\frac{E_f - E_c}{k_B T}\right) \quad p_0 = N_V \frac{2}{\sqrt{\pi}} F_{1/2}\left(\frac{E_v - E_f}{k_B T}\right)$$

5.3 简并条件

$$\left\{ \begin{array}{ll} \text{非简并} & E_c - E_f > 2kT \\ \text{弱简并} & 0 < E_c - E_f \leq 2kT \\ \text{强简并} & E_c - E_f \leq 0 \end{array} \right.$$