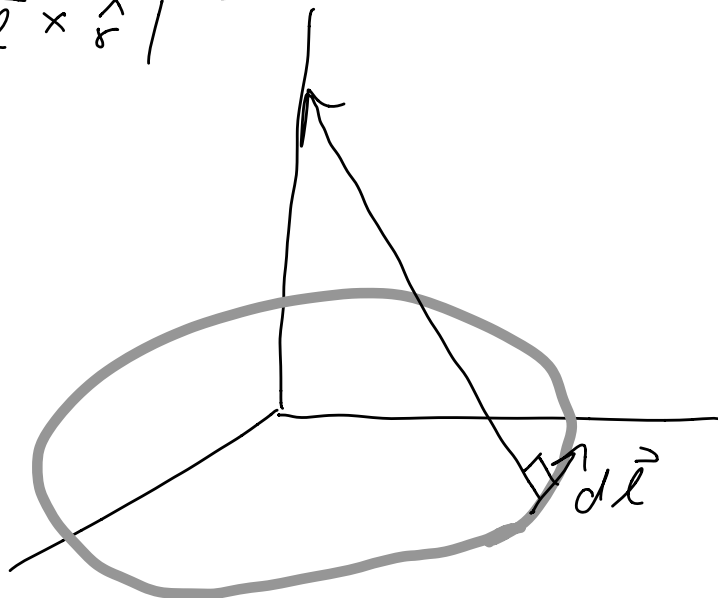


$$dB_z = |d\vec{B}| \frac{R}{\sqrt{R^2 + z^2}}$$

$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{d\vec{\ell} \times \hat{r}}{r^2}$$

$$|d\vec{\ell} \times \hat{r}| = |d\vec{\ell}| |\hat{r}| \sin(\delta)$$

$$|d\vec{\ell} \times \hat{r}| = |d\vec{\ell}| \quad \delta = \frac{\pi}{2}$$



$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{|d\vec{\ell}|}{r^2} \hat{z}$$

$$|d\vec{r}| = R d\phi \quad \left(\Theta = \frac{S}{R}\right)$$

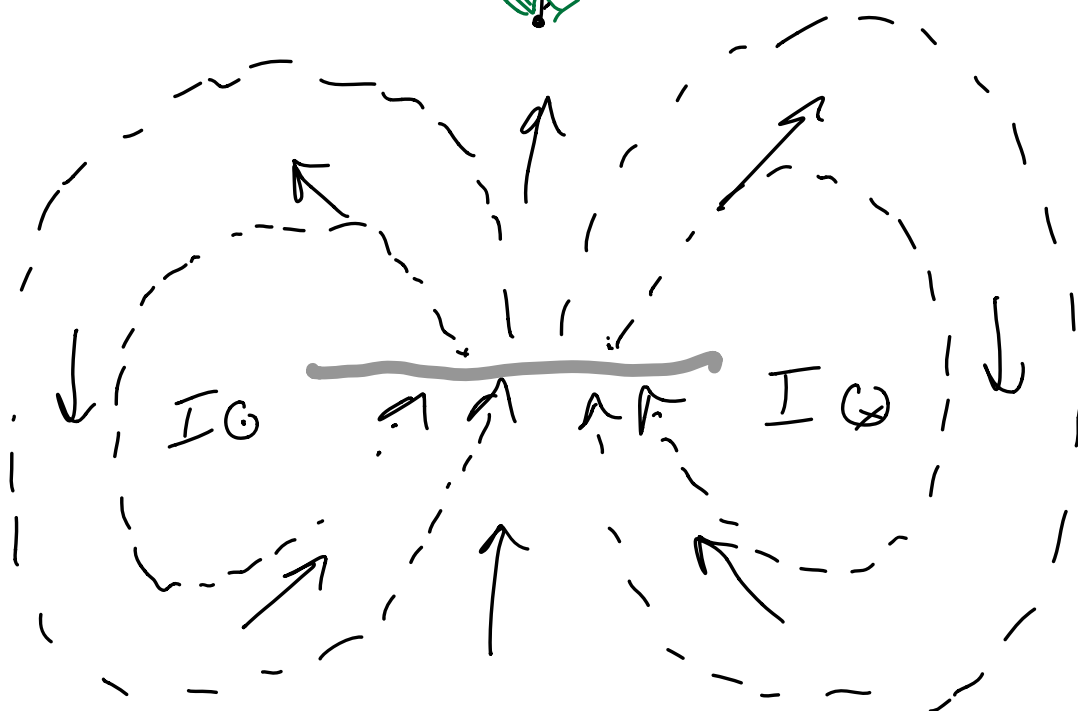
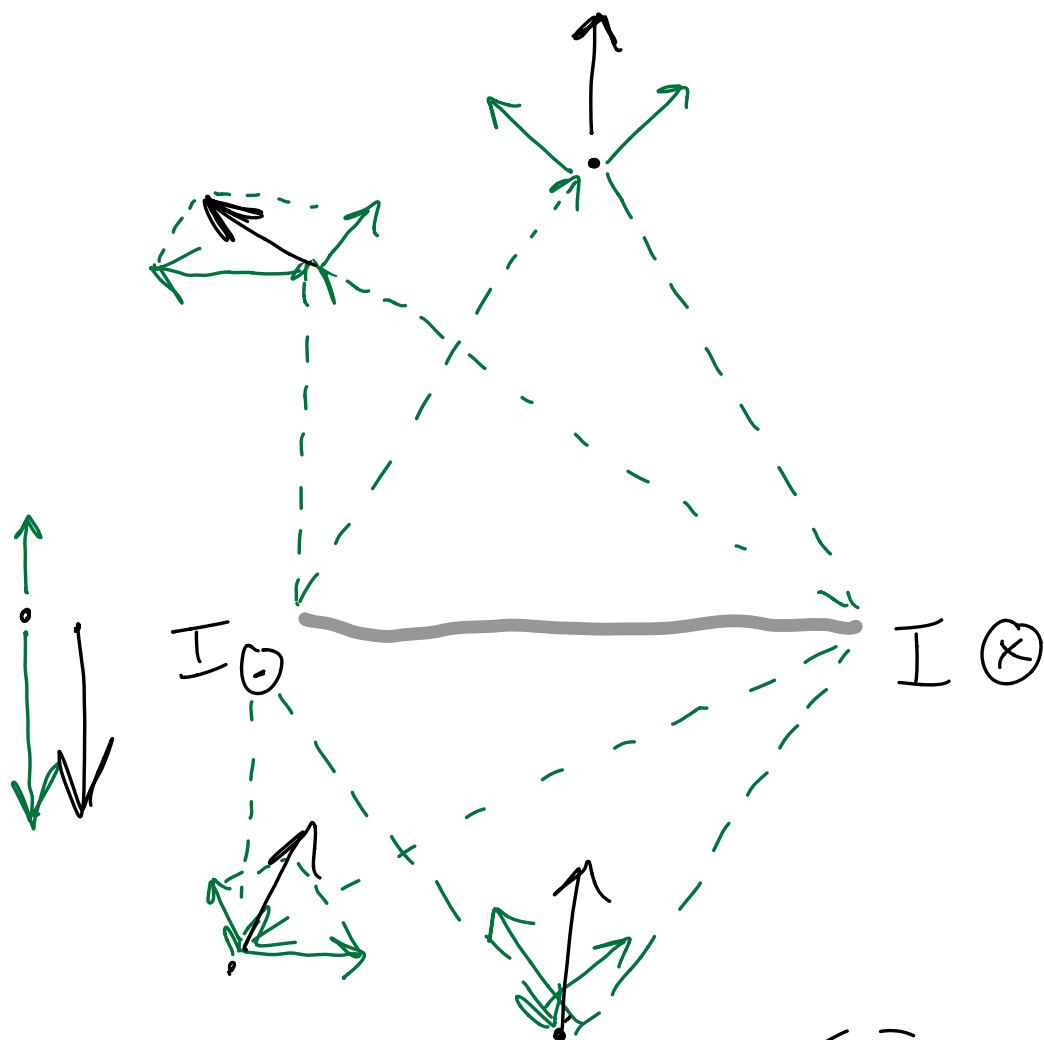
$$|d\vec{B}| = \frac{\mu_0}{4\pi} \frac{I R d\phi}{R^2 + z^2}$$

$$dB_z = |d\vec{B}| \frac{R}{\sqrt{R^2 + z^2}}$$

$$dB_z = \frac{\mu_0}{4\pi} \frac{I R^2 d\phi}{(R^2 + z^2)^{3/2}}$$

$$\vec{B} = \frac{\mu_0 I R^2}{4\pi (R^2 + z^2)^{3/2}} \int_0^{2\pi} d\phi \hat{z}$$

$$\vec{B} = \frac{\mu_0}{4\pi} \frac{2\pi R^2 I}{(R^2 + z^2)^{3/2}} \hat{z}$$



$$\vec{B} = \frac{\mu_0}{4\pi} \frac{2\pi R^2 I}{(R^2 + z^2)^{3/2}} \hat{z}$$

if $z \gg R$

$$(R^2 + z^2)^{3/2} = z^3 \left(\left(\frac{R}{z} \right)^2 + 1 \right)^{3/2}$$

$$\approx 0 \quad \nearrow$$

$$\approx z^3$$

$$|\vec{B}| \approx \frac{\mu_0}{4\pi} \frac{2\pi R^2 I}{z^3}$$

$$\pi R^2 = A$$

$$|\vec{B}| = \frac{\mu_0}{4\pi} \frac{2 A I}{r^3}$$

Look Familiar?

$$|\vec{E}| = \frac{1}{4\pi\epsilon_0} \frac{2 q s}{r^3}$$

Magnetic dipole

$$\frac{1}{4\pi\epsilon_0} \longrightarrow \frac{\mu_0}{4\pi}$$

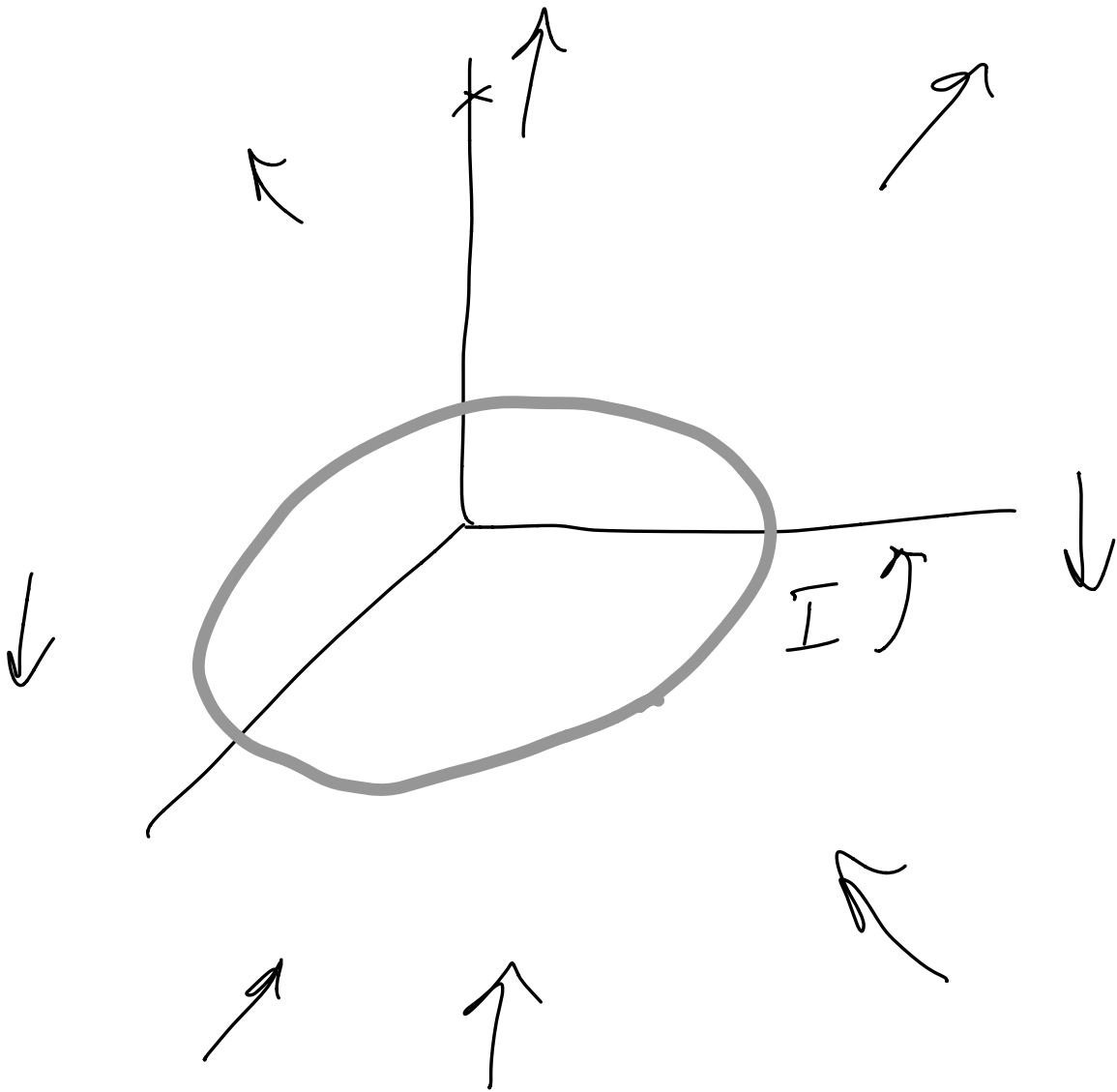
$$q s = p \longrightarrow I A = \mu$$

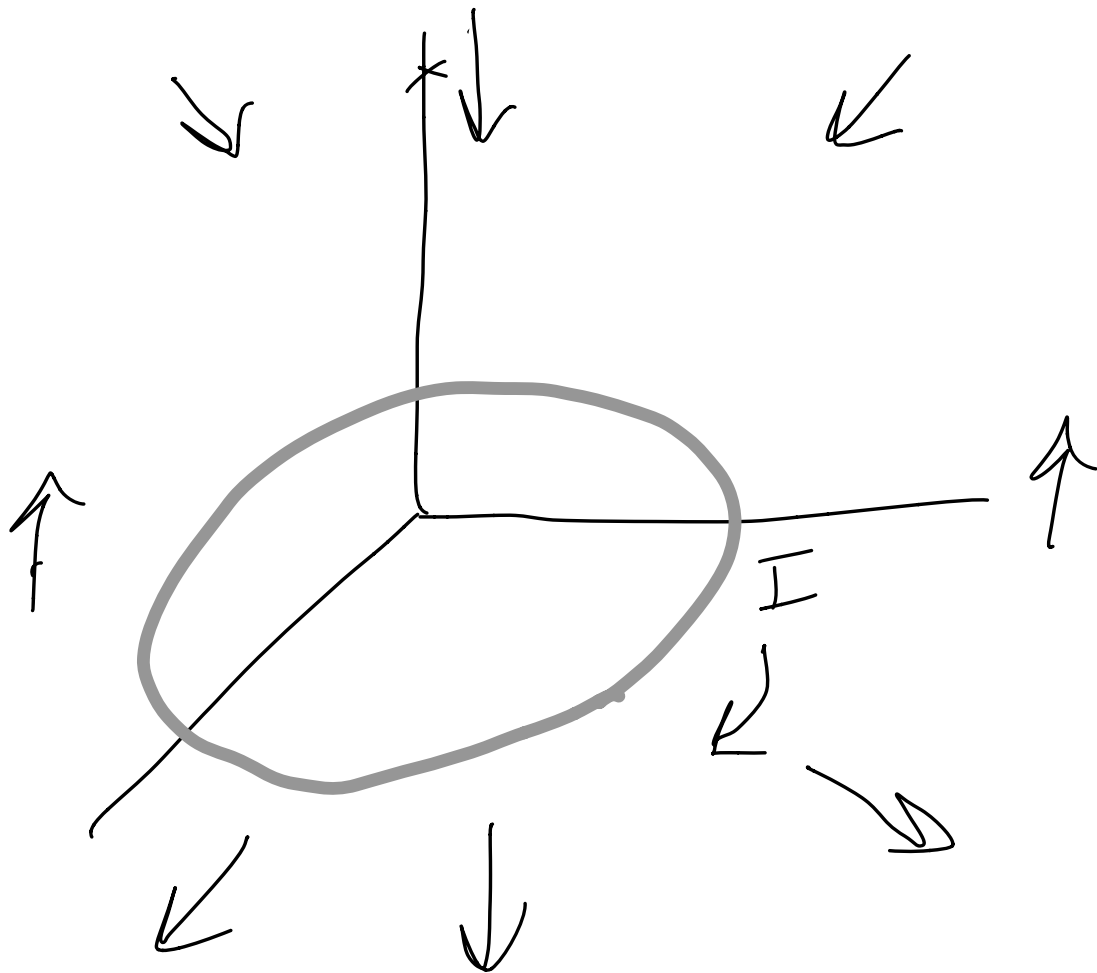
Dipole, on axis:

$$|\vec{B}| = \frac{\mu_0}{4\pi} \frac{2 \mu}{r^3}$$

Direction of dipole Field?

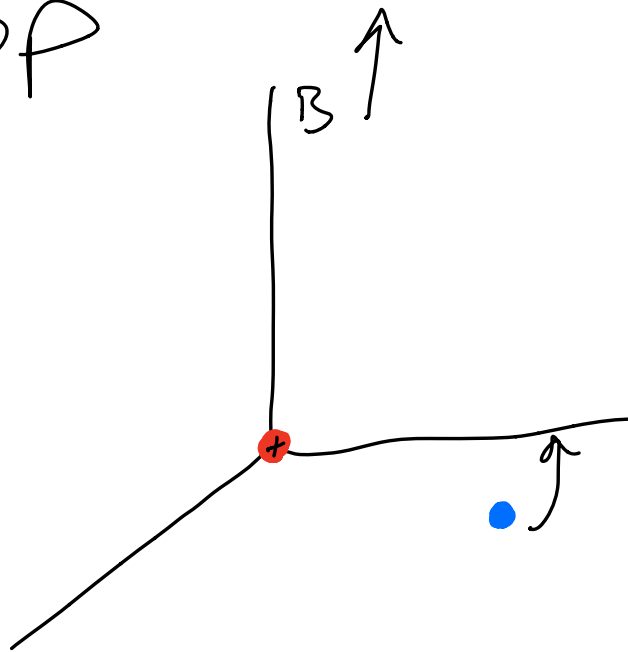
- Depends on current

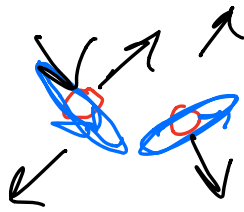




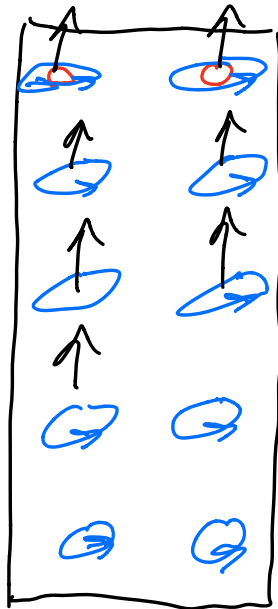
Magnetic Materials

An atom is a
tiny current
loop

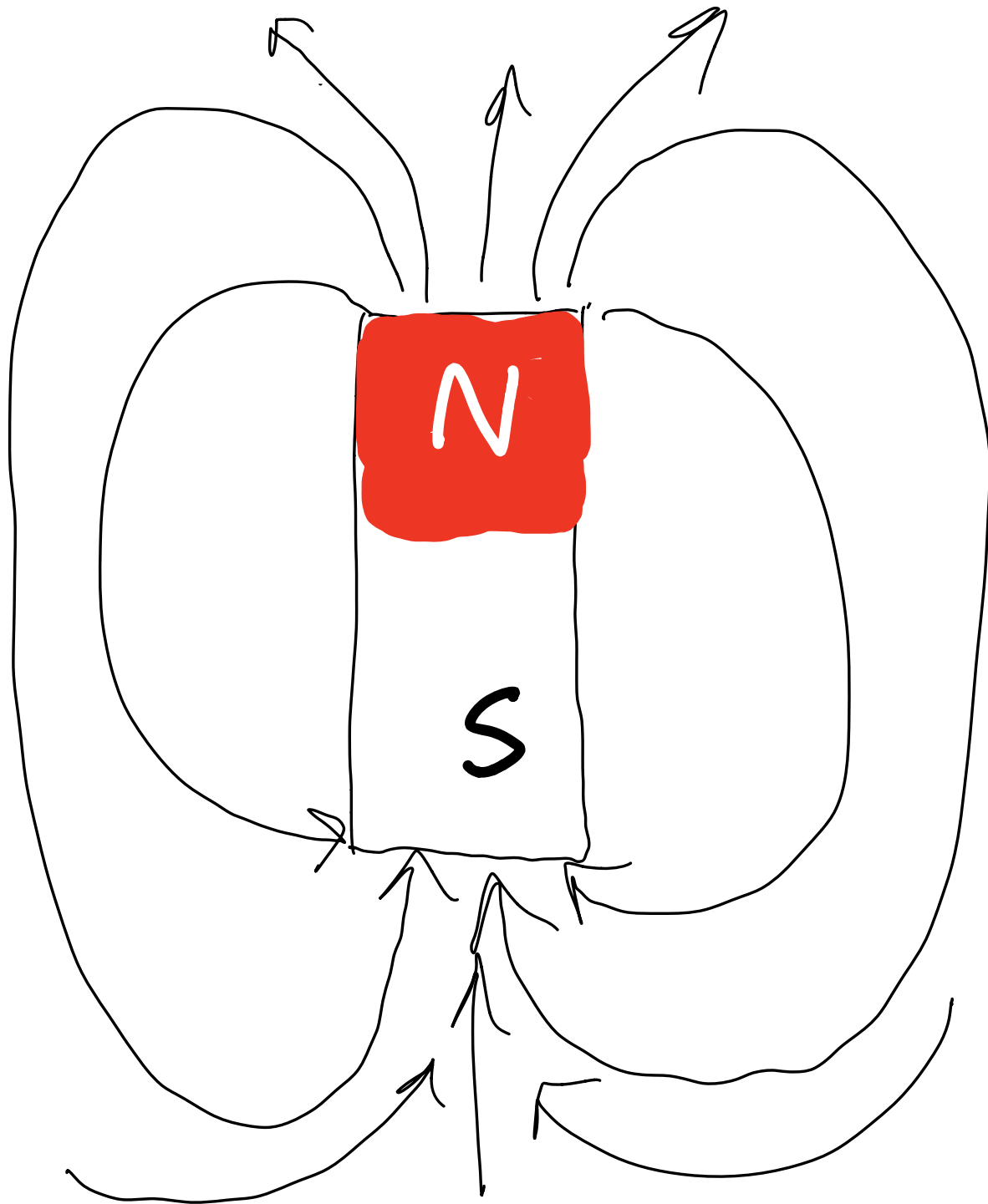




I_n most materials
 \vec{B} cancels



Net \vec{B}



Field lines: Toward S
Away N

