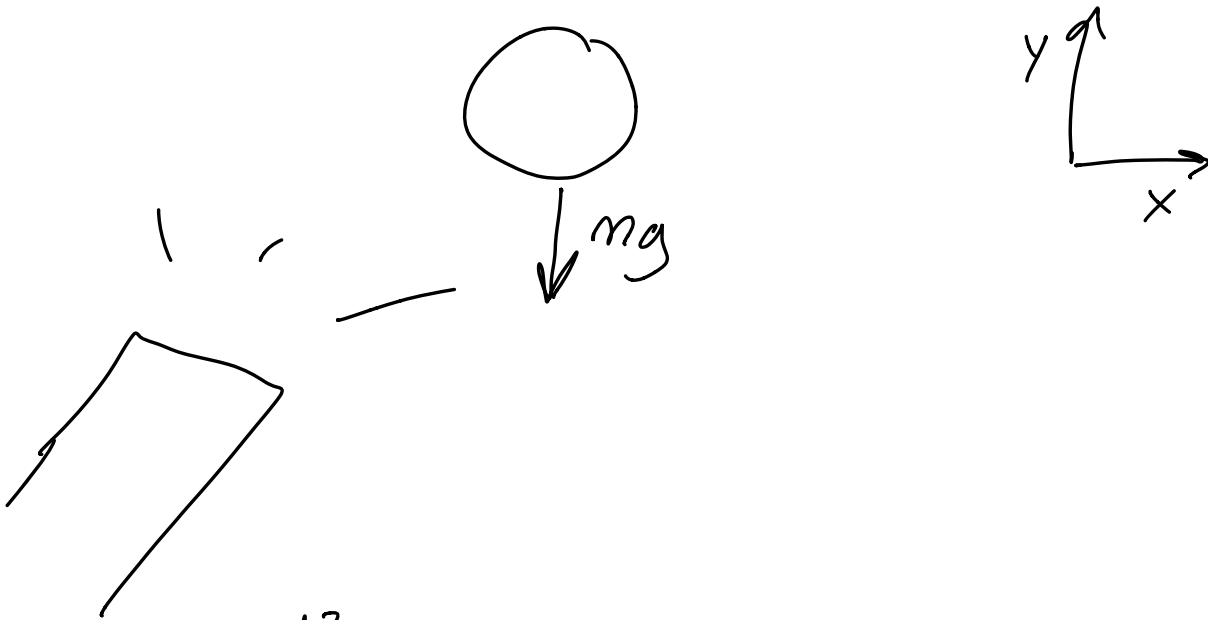


- Review last lecture
- We learned:
 - How to use Euler to solve second-order equations
 - How to apply Euler in more than one dimension



$$\frac{d^2}{dt^2} x = 0$$

$$\frac{d^2}{dt^2} y = -g$$

$$\frac{d^2}{dt^2} (r) \rightarrow \begin{cases} \frac{dv_r}{dt} = F/m \\ \frac{dr}{dt} = v_r \end{cases}$$

$$\frac{dv_x}{dt} = 0$$

$$\frac{dx}{dt} = v_x$$

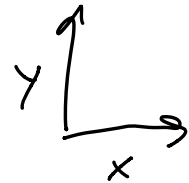
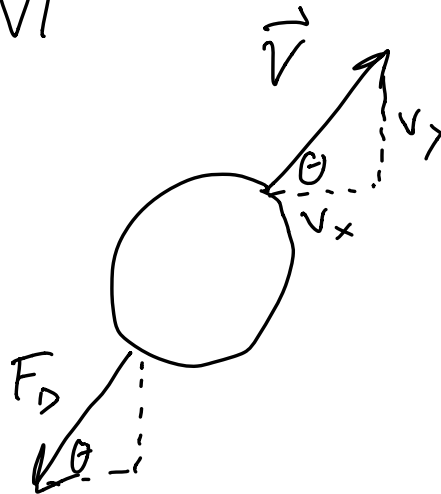
$$\frac{dv_y}{dt} = -g$$

$$\frac{dy}{dt} = v_y$$

- So we've split one second order diffEQ into two first order ones
 - We've done this in two dimensions
- Then we considered drag force

$$F_D = B_2 |\vec{v}|^2$$

$$\vec{F}_D = -\vec{v}$$



$$F_{D,x} = B_2 v^2 \frac{v_x}{v} = B_2 v v_x$$

$$F_{D,y} = B_2 v v_y, \quad v = |\vec{v}| = \sqrt{v_x^2 + v_y^2}$$

$$\frac{dv_x}{dt} = \frac{F}{m} = -\frac{\beta_2 v v_x}{m}$$

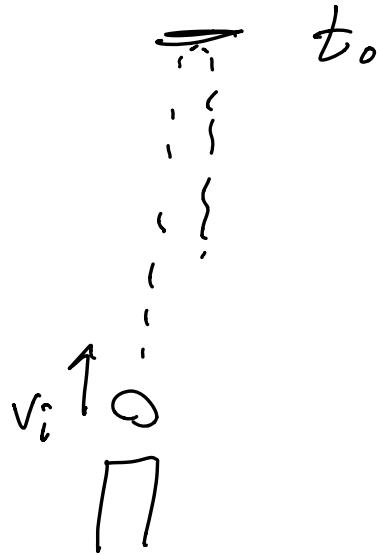
$$\frac{dv_y}{dt} = -g - \frac{\beta_2 v v_y}{m}$$

Units

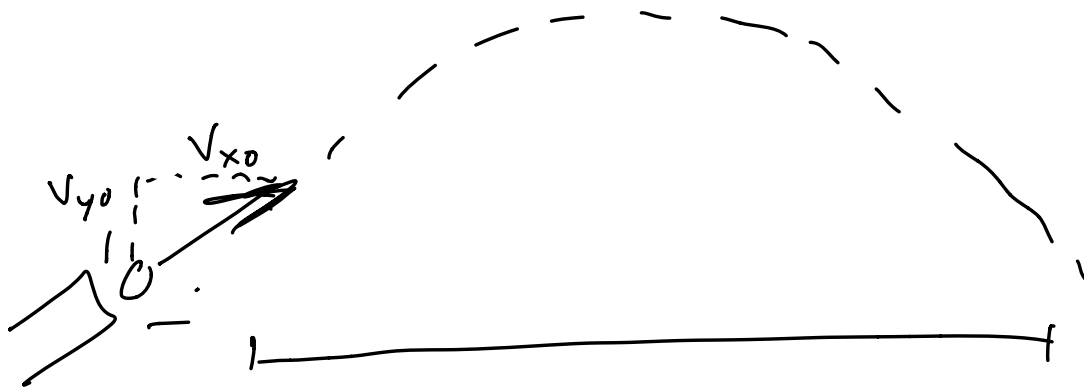
$$\bar{v} = \frac{v}{v_0}, \quad \bar{t} = \frac{t}{t_0}, \quad \bar{x} = \frac{x}{r_0}, \quad \bar{y} = \frac{y}{r_0}$$

$$V_0 = V_{\text{initial}} \Rightarrow \bar{v}(0) = 1$$

$$t_0 = \frac{v_i}{g}$$



$$r_0 = v_0 t_0 = v_i t_0 = \frac{v_i^2}{g}$$

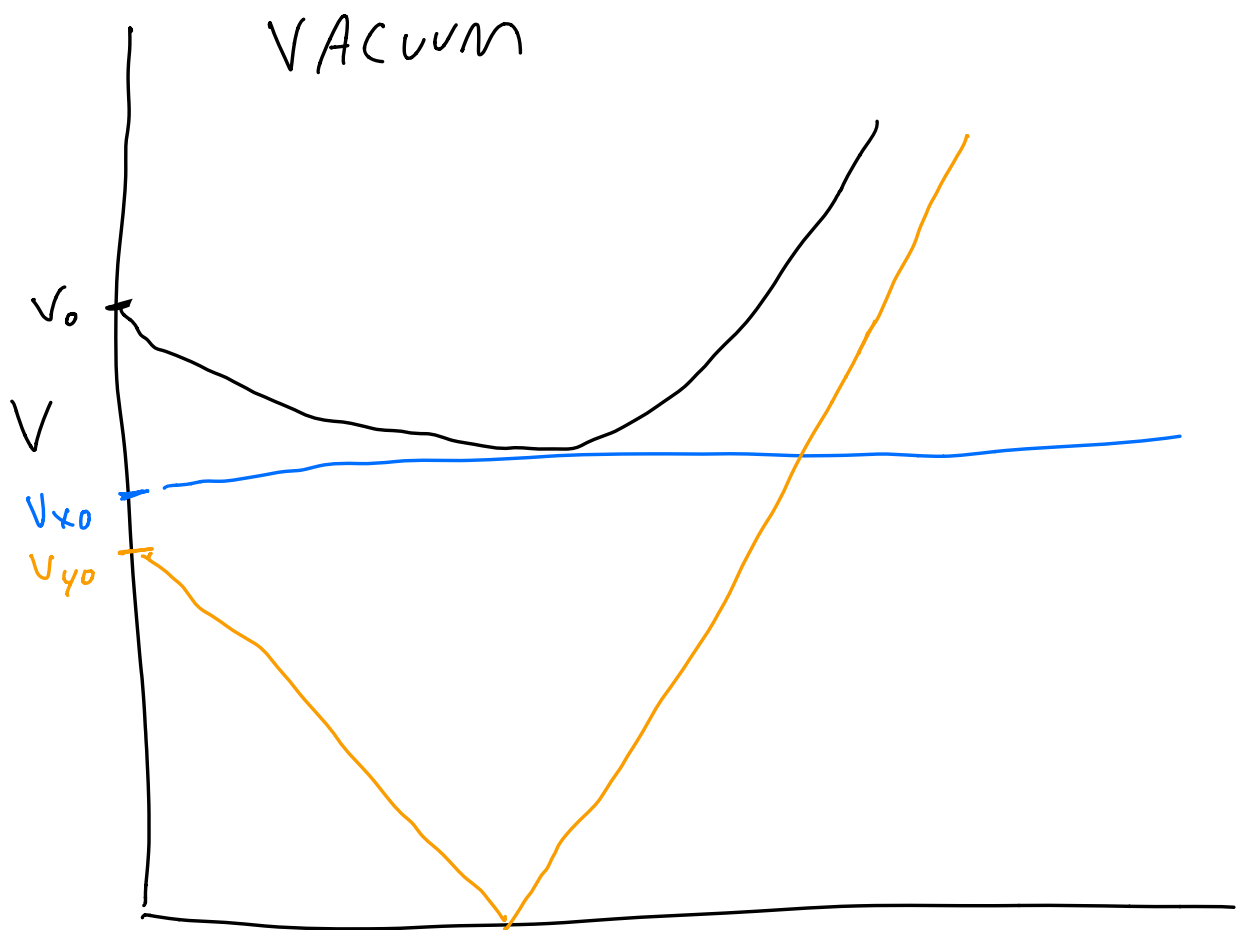


$$\Delta t = \frac{V_{y0}}{g} + \frac{V_{y0}}{g} = 2 \frac{V_{y0}}{g}$$

$$\Delta x = V_{x0} \Delta t = \frac{2 V_{x0} V_{y0}}{g}$$

$$\Delta x \sim \frac{V^2}{g}$$

- Go to jupyter and explain importance of units
 - First do $\theta=90$ (time of flight = $2 \times \text{time unit}$, length unit = $2 \times y_{\text{max}}$)
 - Then $\theta=60$ (time of flight $\sim 1.5 \times \text{time unit}$, length unit $\sim x_{\text{range}} \sim y_{\text{range}}$)
 - Vary initial velocity and show that units adjust accordingly
 - Turn friction on and show units are still reasonable
 - When we normalize, all dynamic variables are of order 1
- We haven't looked at velocities yet? How do we expect v to behave?
- First: In a vacuum



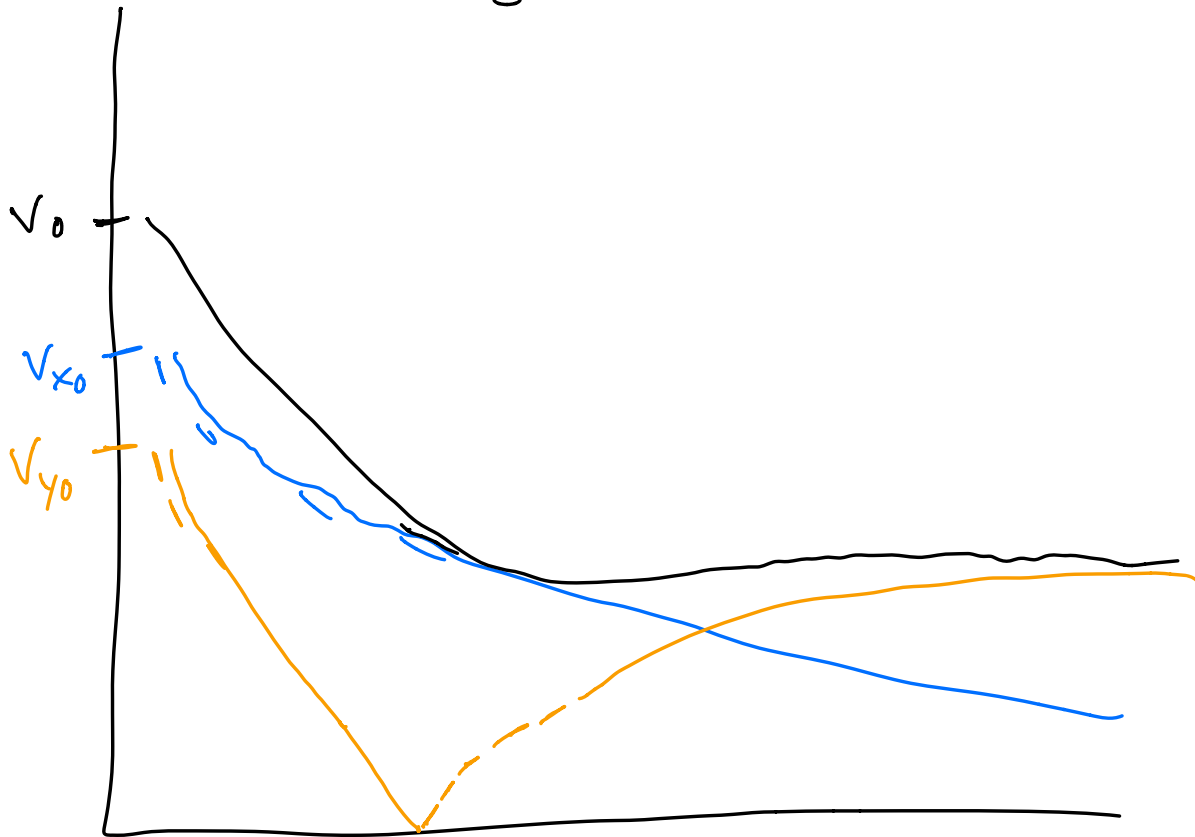
$|\vec{V}|, |V_x|, |V_y|$ t

$V_y \propto t$

$$V = \sqrt{V_x^2 + V_y^2}$$

$$= \sqrt{K + at^2}$$

$$g \neq 0$$



- V_x : No acceleration term, always decreasing but at an ever decreasing rate (because smaller velocity \rightarrow smaller drag force)
 - Eventually $\rightarrow 0$
- V_y : Initially decelerated by BOTH gravity and drag
 - Gravity is constant, drag force decreases
 - Eventually: $V_y = 0$
 - Only for an instant (gravity still working!)
 - Switches direction and speeds up
 - As it does so, drag force gets stronger
 - Reaches terminal velocity eventually
- V :
 - Harder to analyze but know three things:
 - Initially decreasing
 - $V = V_x$ when $V_y = 0$
 - $V \rightarrow V_y \rightarrow V_{\text{term}}$

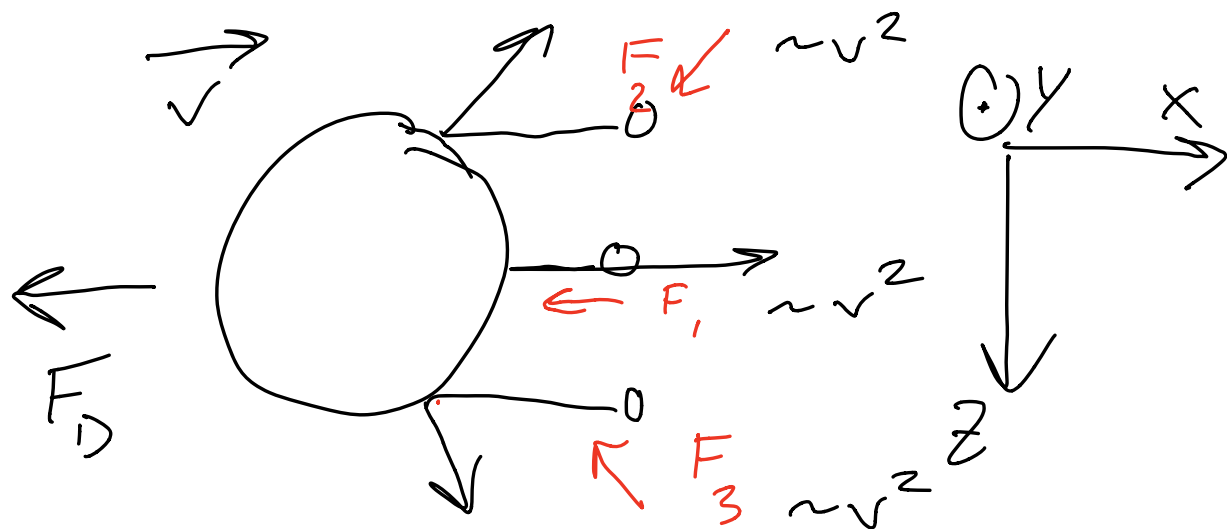
$$V_{\text{term}}: \frac{dv_y}{dt} = 0$$

$$-g + \frac{B_z}{m} v^2 = 0$$

$$v^2 = \frac{mg}{B_z}$$

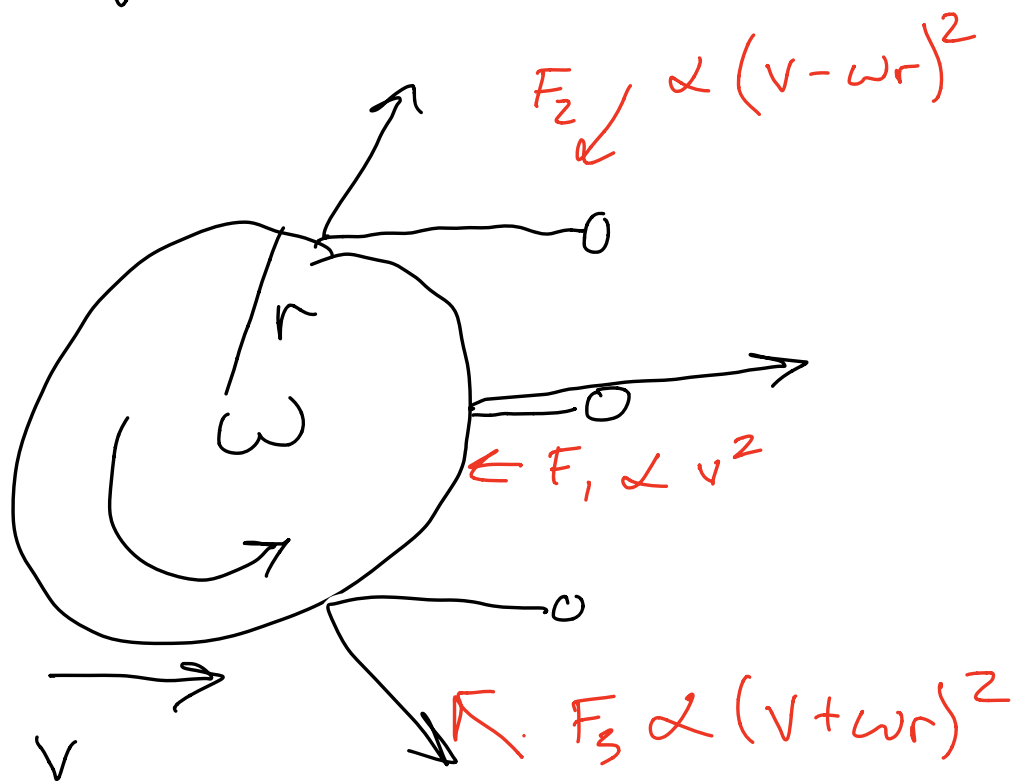
$$V_{\text{term}} = \sqrt{\frac{mg}{B_z}}$$

Magnus



$$\vec{F}_0 = \vec{F}_1 + \vec{F}_2 + \vec{F}_3$$

$$\vec{F}_0 = -\vec{v}$$



$$\vec{F} = \vec{F}_1 + \vec{F}_2 + \vec{F}_3$$

$$= \vec{F}_D + ?$$

Non-cancelling \hat{z} terms

$$\vec{F}_m \propto (v - \omega r)^2 \hat{z} - (v + \omega r)^2 \hat{z}$$

$$\propto (v^2 + \omega^2 r^2 - 2v\omega r - v^2 - \omega^2 r^2 - 2v\omega r)$$

$$\propto -4v\omega r \hat{z}$$

$$\vec{F}_m = -\sum_0 v \omega \hat{z}$$

$$\vec{F} = -\vec{z}$$

$$\vec{v} = \vec{x}$$

$$\vec{\omega} = \vec{y}$$

$$\vec{F} = \vec{\omega} \times \vec{v}$$

$$\vec{F}_m = -\gamma_0 \vec{\omega} \times \vec{v}$$

$$\vec{\omega} \times \vec{v} =$$

$$(\omega_y v_z - \omega_z v_y) \hat{x} +$$

$$(\omega_z v_x - \omega_x v_z) \hat{y} +$$

$$(\omega_x v_y - \omega_y v_x) \hat{z}$$

$$\frac{dv_x}{dt} = -\frac{B_z}{m} v v_x + \frac{S_0}{m} (\omega_y v_z - \omega_z v_y)$$

$$\frac{dv_y}{dt} = -g - \frac{B_z}{m} v v_y + \frac{S_0}{m} (\omega_z v_x - \omega_x v_z)$$

$$\frac{dv_z}{dt} = -\frac{B_z}{m} v v_z + \frac{S_0}{m} (\omega_x v_y - \omega_y v_x)$$

$$\frac{dx}{dt} = v_x, \quad \frac{dy}{dt} = v_y, \quad \frac{dz}{dt} = v_z$$

Curve ball

