

Consider an ideal gas of electrons (fermions)

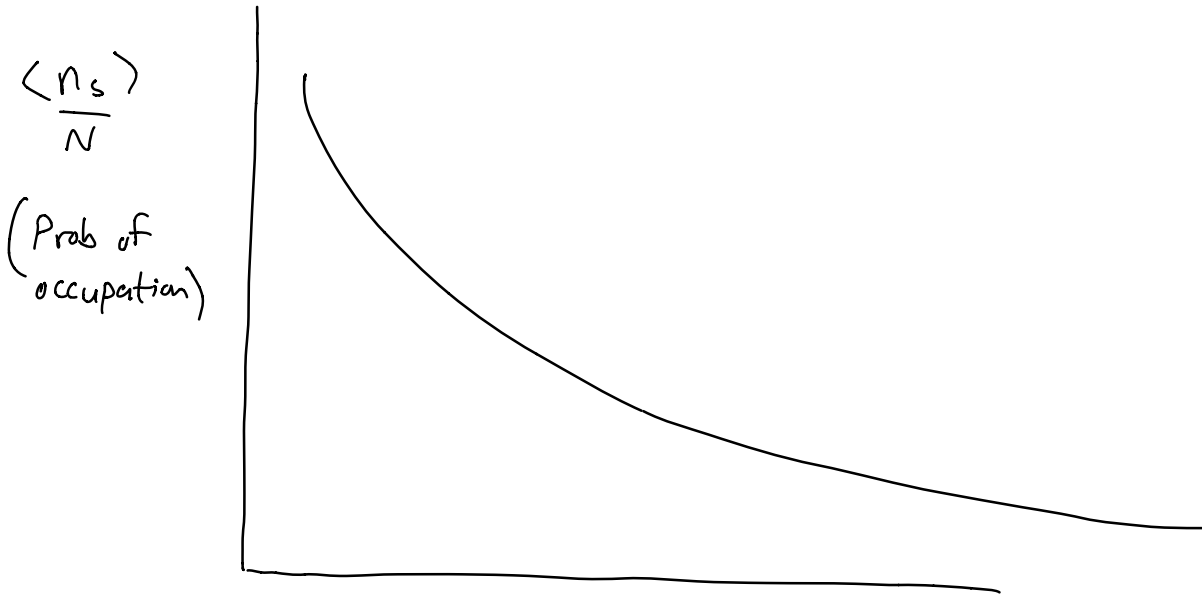
Quantum states given by n_x, n_y, n_z

$$E_s = \frac{\hbar^2 \pi^2}{L^2} (n_x^2 + n_y^2 + n_z^2)$$

s	E_s
$\langle 1, 1, 1 \rangle$	3
$\langle 2, 1, 1 \rangle$	6
$\langle 1, 2, 1 \rangle$	6
$\langle 1, 1, 2 \rangle$	6
$\langle 3, 1, 1 \rangle$	11
\vdots	\vdots

If I add an electron to the gas, which state will it occupy?

At "normal" temperatures



- Particles don't necessarily fill the lowest available energy state

- Why? Entropy

$$\Delta F = \Delta U - T\Delta S$$

Particles will spontaneously jump to higher energies as long as $\Delta F < 0$ ($\Delta U < T\Delta S$)

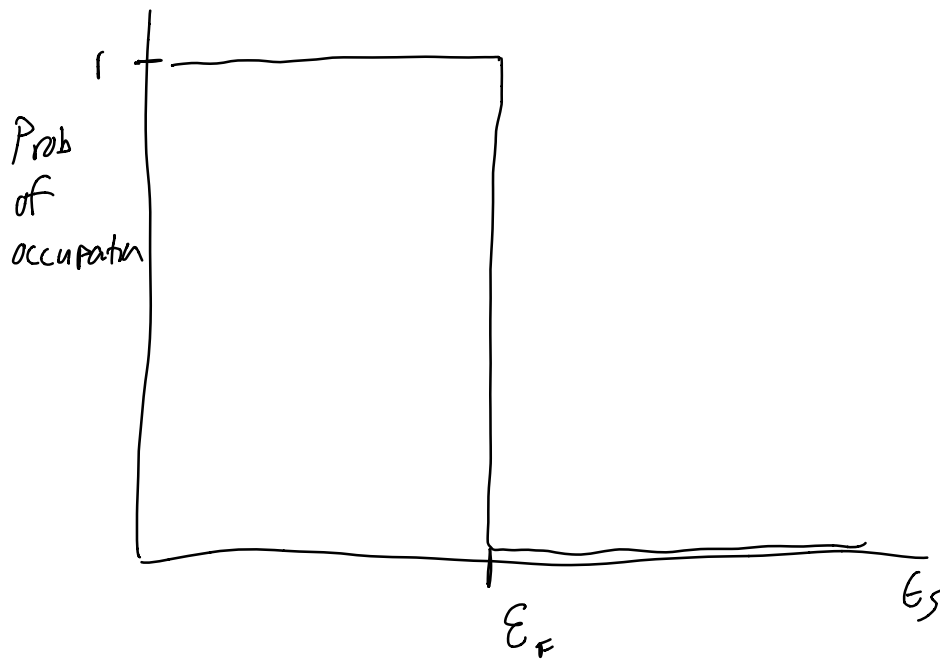
The gas wants to minimize energy while
also maximizing entropy

This means there are some empty
states between the lowest
+ highest energy electron

S	E_s	
$\langle 1, 1, 1 \rangle$	3	Bring in first electron, it goes to ground state
$\langle 2, 1, 1 \rangle$	6	Second electron skips these & goes here instead
$\langle 1, 2, 1 \rangle$	6	
$\langle 1, 1, 2 \rangle$	6	
$\langle 3, 1, 1 \rangle$	11	
\vdots	\vdots	

How energetic is the highest energy electron? Can't say for sure.

Compare this with the $T=0$ scenario



States with $E_s < E_F$ are certain to be occupied

States w/ $E_s > E_F$ are certainly empty

(no unoccupied states between ground
+ E_F)

Particles will spontaneously jump
to higher energies as long
as $\Delta F < 0$ ($\Delta U < T \Delta S$)

$$T \rightarrow 0$$

$$\Delta F < 0 \Rightarrow \Delta U < 0$$

electrons will only move to states
with lower energy; these are
all occupied

Since all states below E_F are filled,
 counting states \longleftrightarrow counting particles

S	E_s	N_s
$\langle 1, 1, 1 \rangle$	3	2 ($\uparrow\downarrow$)
$\langle 2, 1, 1 \rangle$	6	2
$\langle 1, 2, 1 \rangle$	6	2
$\langle 1, 1, 2 \rangle$	6	2
$\langle 3, 1, 1 \rangle$	11	2
\vdots	\vdots	
	E_F	2
	\vdots	0
	\vdots	0
	\vdots	0

The possible combinations of n_x, n_y, n_z
which result in energy E_s

lie on the sphere with radius

$$\left(\frac{2mL^2}{\hbar^2 \pi^2} E_s \right)^{1/2}$$

$$E_s = \frac{\hbar^2 \pi^2}{L^2} (n_x^2 + n_y^2 + n_z^2)$$

↓

$$n_x^2 + n_y^2 + n_z^2 = \frac{2mL^2}{\hbar^2 \pi^2} E_s$$

of states with energy E_s

$$= 2 \cdot \sum_{n_x} \sum_{n_y} \sum_{n_z} (1) \quad ; \quad n_x^2 + n_y^2 + n_z^2 = \frac{2mL^2}{\hbar^2 \pi^2} E_s$$

of states with energy below E_s :

$$= 2 \cdot \sum_{n_x} \sum_{n_y} \sum_{n_z} (1) \quad ; \quad n_x^2 + n_y^2 + n_z^2 \leq \frac{2mL^2}{\hbar^2 \pi^2} E_s$$

$$= 2 \cdot \frac{1}{\Delta n^3} \int_{n_x^2 + n_y^2 + n_z^2 \leq \frac{2mL^2}{\hbar^2 \pi^2} E_s} dn_x dn_y dn_z$$

$$\frac{1}{\Delta n^3} = \frac{1}{1^3} = 1$$

$$dn_x dn_y dn_z \rightarrow n^2 \sin \theta d\theta d\phi$$

$$= 2 \cdot \int_0^{\frac{\pi}{2}} \sin \theta d\theta \int_0^{\frac{\pi}{2}} d\phi \int_0^{\sqrt{\frac{2mL^2}{\hbar^2 \pi^2} E_s}} n^2 dn$$

$$= 2 \cdot \frac{\pi}{2} \cdot \frac{1}{3} \left(\frac{2mL^2}{\hbar^2 \pi^2} E_s \right)^{3/2}$$

$$n(E < E_s) = 2 \cdot \frac{\pi}{6} \cdot \left(\frac{2mL^2}{\hbar^2 \pi^2} E_s \right)^{3/2}$$

$$\text{at } E_s = E_F, \quad n(E < E_F) = N$$

$$N = \frac{\pi}{3} \cdot \left(\frac{2mL^2}{\hbar^2 \pi^2} E_F \right)^{3/2}$$

$$\boxed{E_F = \frac{\hbar^2}{2m} \left(\frac{3\pi^2 N}{V} \right)^{2/3}} \quad ; \quad V = L^3$$

Energy of highest energy electron

$$\textcircled{a} \quad T = 0$$

What is u ?

$$U = 2 \cdot \sum_{n_x} \sum_{n_y} \sum_{n_z} E(n_x, n_y, n_z) \quad ; \quad n_x^2 + n_y^2 + n_z^2 \leq \frac{2mL^2}{\hbar^2 \pi^2} E_F$$

$$= 2 \int \frac{\hbar^2 \pi^2}{L^2} (n_x^2 + n_y^2 + n_z^2) dn_x dn_y dn_z$$

↓

$$= 2 \int_0^{\frac{\pi}{2}} \sin \theta d\theta \int_0^{\frac{\pi}{2}} d\phi \int_0^{\left(\frac{2mL^2}{\hbar^2 \pi^2}\right)^{\frac{1}{2}}} \frac{\hbar^2 \pi^2}{L^2} n^2 \cdot n^2 dn$$

$$= 2 \cdot 1 \cdot \frac{\pi}{2} \cdot \frac{\hbar^2 \pi^2}{L^2} \left(\frac{1}{5}\right) \left(\frac{2mL^2}{\hbar^2 \pi^2}\right)^{\frac{5}{2}}$$

$$\boxed{U = \frac{3}{5} N E_F}$$

Electrons in a metal @ room temp:

$$(C_u)$$

$$\mathcal{E}_F \approx 7 \text{ eV}$$

$$kT \approx 3 \times 10^{-2} \text{ eV}$$

Since the energy of the gas $\left(\frac{2}{5} \mathcal{E}_F\right)$
 $U \gg kT$ (the thermal energy)

we can approximate that $T \approx 0$

- So our gas gets a large amount of internal energy, even at $T=0$

- The gas also exerts pressure

$$dU = TdS - PdV + \mu dN \Rightarrow P = -\frac{\partial U}{\partial V}$$

$$U = \frac{2}{5} \mathcal{E}_F = \left(\frac{2}{5}\right) \frac{\hbar^2}{2m} \left(\frac{3\pi^2 N}{V}\right)^{2/3}$$

$$P = \frac{2}{3} \frac{U}{V}$$

Example: Neutron star

A gas of N neutrons at $T \approx 0$

$$U_{\text{kinetic}} = \frac{3}{5} N \epsilon_F$$

$$U_K = \frac{3}{5} \cdot N \cdot \frac{\hbar^2}{2m_n} \left(\frac{3\pi^2 N}{V} \right)^{2/3}$$

$$M = N m_n \Rightarrow N = \frac{M}{m_n}, \quad V = \frac{4}{3} \pi R^3$$

$$U_K = \frac{3}{10} \frac{\hbar^2 \pi^2}{m_n} \left(\frac{M}{m_n} \right)^{5/3} \left(\frac{9}{4\pi^2 R^3} \right)^{2/3}$$

$$U_{\text{grav}} = -\frac{3}{5} \frac{GM^2}{R}$$

$$U_{\text{tot}} = U_K + U_{\text{grav}} =$$

$$\frac{3}{10} \frac{\hbar^2 \pi^2}{m_n} \left(\frac{M}{m_n} \right)^{5/3} \left(\frac{9}{4\pi^2 R^3} \right)^{2/3} - \frac{3}{5} \frac{GM^2}{R}$$

$$\approx \frac{a}{R^2} - \frac{b}{R}$$



Equilibrium: $\frac{dU_{\text{tot}}}{dR} = 0$ Solve for R

$$M \approx 1 M_{\odot}$$

$$m_n = 1.67 \times 10^{-27} \text{ kg}$$

$$R \approx 12 \text{ km}$$

$$\rho = \frac{M_{\odot}}{\frac{4}{3}\pi R^3} \approx 3 \times 10^{17} \frac{\text{kg}}{\text{m}^3}$$

$$\frac{6 \times 10^{24} \text{ kg}}{R^3} = 3 \times 10^{17}$$

$$1 \text{ spoonfull} \approx \text{Mt. Everest}$$

$$\epsilon_F \approx 6 \times 10^7 \text{ eV}$$

$$T_F = \frac{\epsilon_F}{k} \approx 7 \times 10^{11} \text{ K}$$

$$T_{\text{sun}} \approx 6000 \text{ K}$$