Consider a gas of many particles at constent energy E

How many mrcostates?

- Before we proceed, some new tools to help if our analysis

For the QM oscillator, the state of each oscillator is:

- finite (counteble, integer factor of two)

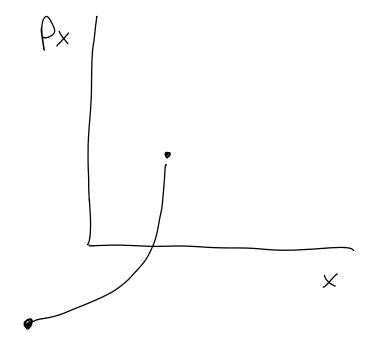
- Specified by a single number $E_{osc} = g \hbar \omega$

For a gas of particles, the state of each particle is specified by 6 numbers (x,y,z,px,py,pz)

In general they are each continuous

We can specify the state of the system by plotting it in a higher dim. phase space

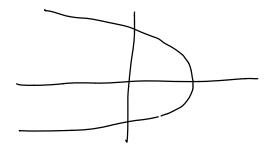
Consider a single particle in 1 dimension State specified by (x, Px) P = Pxi - mgt $X = Xi + \frac{Pxi}{m}t - \frac{1}{2}gt^{2}$



Possible states of system lie on the curve $\frac{1}{2m}P_{x}^{2} + mg_{x} = E$

$$\frac{1}{2m}Px^2 + mgx = E$$

$$P_{x} = \sqrt{2m(E - mg_{x})}$$



Ex: Single particle

mass on the end of a spring

$$H = \frac{P^2}{Zm} + \frac{1}{Z}Kx^2$$

Possible values of (x,px)

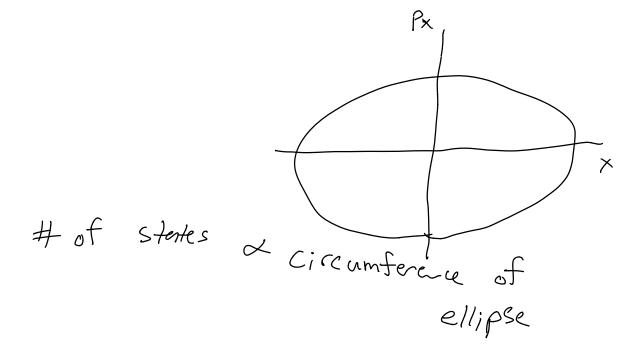
lie on the corne

 $\frac{P^2}{Zm} + \frac{1}{Z}Kx^2 = E$

$$\left(\frac{X}{a}\right)^{2} + \left(\frac{Y}{b}\right)^{2} = 1$$

$$a = \sqrt{2E}$$

$$b = \sqrt{2nE}$$



$$H(x,y,z,p_{x},p_{y},p_{z}) = \frac{1}{2m} \left(p_{x}^{2} + p_{y}^{2} + p_{z}^{2} \right) + \frac{1}{2} \kappa \left(x^{2} + y^{2} + z^{2} \right)$$

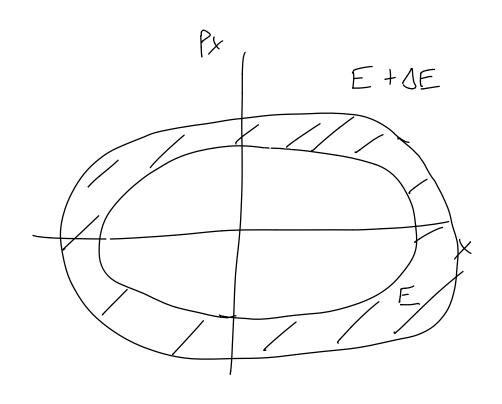
$$\frac{1}{2m} \left(p_{x}^{2} + p_{y}^{2} + p_{z}^{2} \right) + \frac{1}{2} \kappa \left(x^{2} + y^{2} + z^{2} \right) = E$$

Our Hamiltonian Egn forms a "hyper susface" in a 6D space # of microstates & area of hypersurface well talk about how to find the area in a second - Consider the general case N-many particles, in 3 dimensions 6N numbers to specify the state X,, Y,, Z,, Px, Py, P1Z X2, Y2, Z2) P2x, P2y, P2Z

$$H = \left(\sum_{i=1}^{N} \frac{1}{2m} \left(P_{x_i}^2 + P_{y_i}^2 + P_{z_i}^2 \right) \right) + U(x_{x_i}, y_{x_i}, z_{x_i}, x_{z_i}, z_{z_i})$$

To find the area

Consider the Small volume in between E + E + 11E



Volume in phase space >> w

 $\Delta W = \int dx_1 dy_1 dz_1 dx_2 dy_2 dz_2 \dots dp_x_1 dp_y_1 dp_z_1 dp_{x_2} dp_{y_2} dp_$

 $\Delta W = \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E + \Delta E} \int_{E \leq H(r_i, p_i) \leq E} \int_{E \leq H(r$

Total volume bounded by E:

 $W = \int d^3N r d^3N$ $H(r_i, p_i) \leq E$

 $\Delta w = w(E + \Delta E) - w(E) = \left(\frac{\partial w}{\partial E}\right) \Delta E$

$$\Delta w = \sigma(E) \Delta E$$

$$\left(dV = 4\pi c^2 dc \right)$$

So
$$\sigma(E) = \frac{\partial \omega}{\partial E}$$

$$V = \frac{4}{3}\pi r^{3}$$

Ex: The ideal gas
$$Potential = 0 \qquad Fl = \sum_{i=1}^{N} \frac{1}{2m} (Px_i^2 + Py_i^2 + Pz_i^2)$$

$$\omega = \int_{A} d^{3N} d^{3N} \rho$$

H doesn't depend on

$$\omega = \int d^{3N} r \int d^{3N} r$$
H(E)

$$W = V \int_{0}^{1} d^{3N} P$$

$$+1 < E$$

$$P_{1}^{2} + P_{2}^{2} + P_{3}^{2} + P_{4}^{2} + ... + P_{N-1}^{2} + P_{N}^{2} = 2m E$$

In d dimensions

$$V_{d}(R) = \frac{\pi^{d/2}}{\frac{d}{z} \Gamma(\frac{d}{z})} R^{d}$$

$$\Gamma(z) = \int_0^\infty x^{z-1} e^{-x} dx$$

$$\Gamma(s) = 4\Gamma(4)$$

$$\Gamma(4) = 3\Gamma(3)$$
 $\Gamma(5) = 4.3.2.1$
 $= 4!$
 $\Gamma(3) = 2\Gamma(2)$
 $\Gamma(2) = \Gamma(1)$ $\Gamma(z) = (z-1)!$

$$\Gamma(3) = 2\Gamma(2) = 4$$

$$\Gamma(z) = \Gamma(1) \qquad \Gamma(z) = (z-1)!$$

$$V_{J}(R) = \frac{\pi^{d/2}}{\frac{d}{Z} \Gamma(\frac{d}{Z})} R^{d}$$

$$if d = 3$$

$$if d = 3$$

$$V_3(R) = \frac{\pi^{3/2}}{\frac{3}{2} \Gamma(\frac{3}{2})} R^3$$

$$V_3(R) = \frac{11^{3/2}}{\frac{3}{2}} R^3$$

$$V_3(a) = \frac{4}{3} \pi R^3$$

$$V_{z}(R) = \frac{\pi}{11} R^{2}$$

$$V_{J}(R) = \frac{\pi^{d/2}}{\frac{d}{d} \Gamma(\frac{d}{d})} R^{d}$$

$$d = 3N$$

$$R = (2mE)^{V_2}$$

$$V_{d}(\sqrt{zmE}) = \frac{\pi^{\frac{3N}{2}}}{2} (2mE)^{3N/2}$$

$$W(E) = V^{N} \frac{3\frac{N}{2}}{11} (2mE)^{3N/2}$$

$$\sigma(E) = \frac{\partial w}{\partial E} = \frac{\sqrt{N}}{3N} \left(\frac{3N}{2}\right) \left(2mE\right)^{\frac{3N}{2}-1} \left(2m\right)$$

$$(2mE)^{\frac{3N}{2}-1}(2m)=(2m)^{\frac{3N}{2}}E^{\frac{3N}{2}-1}$$

$$\sigma(E) = \frac{\sqrt{N_{\text{T}}^{3N}} \frac{3N}{2} (2m)^{\frac{3N}{2}} E^{\frac{3N}{2}-1}}{\Gamma(\frac{3N}{2})}$$

$$E^{\frac{3N}{2}-1} = E^{\frac{3N}{2}} = C \approx E^{\frac{3N}{2}}$$

$$\sigma(E) = \frac{\sqrt{N \pi^{\frac{3N}{2}} (2mE)^{\frac{3N}{2}}}}{\Gamma(\frac{3N}{2})}$$

$$\Omega \propto O(E)$$

$$[O] = [L]^{3N} [P]^{3N}$$

We are counting phase space volume

SL is a number (dim-less)

$$V = \int dx dy dz \qquad H = \frac{1}{4} \int dx dy dz$$

$$D = \frac{1}{4x^3} \int dx dy dz \qquad H = \frac{P}{\Delta P}$$

Split phase space into cells of "volume" (DXDP)

$$\Omega = \frac{1}{(2 \times 2 P)^{3N}} \frac{\sqrt{N \pi^{\frac{3N}{2}} (2mE)^{\frac{3N}{2}}}}{\Gamma(\frac{3N}{2})}$$

Since only the relative amount of se matters, the proportionality constant doesn't really matter

The quantity DXDP has an absolute minimum

Let's pick $\triangle x + \triangle p$ so that $\triangle x \triangle p = x$

$$\Omega = \frac{1}{\left(\frac{1}{h}\right)^{3N}} \frac{\sqrt{N \pi^{\frac{3N}{2}} \left(2mE\right)^{\frac{3N}{2}}}}{\Gamma\left(\frac{3N}{2}\right)}$$

if the particles are indistinguishable, then we've over-counted

- we can swap 2 particles w/o changing the State of the system

- number of duplicates

$$\Omega = \frac{1}{\sqrt{1 + \frac{3N}{N!} \Gamma(\frac{3N}{2})}}$$

-1 ~ $V^N U^{3N/Z}$