

**A measurement of $Wb\bar{b}$ production and a search for monophoton signals of dark matter
using the CMS detector at the CERN LHC**

by

Thomas Mastrianni Perry

A dissertation submitted in partial fulfillment of
the requirements for the degree of

Doctor of Philosophy

(Physics)

at the

UNIVERSITY OF WISCONSIN-MADISON

2016

Date of final oral examination: 3 August 2016

The dissertation is approved by the following members of the Final Oral Committee:

Wesley Smith (Advisor), Bjorn Wiik Professor, Physics

Sridhara Dasu, Professor, Physics

Matt Herndon, Professor, Physics

Yang Bai, Professor, Physics

David C. Schwartz, Professor, Chemistry

CONTENTS

	Page
List of Figures	iv
List of Tables	v
1 Introduction	1
1.1 The Standard Model	1
1.1.1 <i>Standard Model Particles</i>	1
1.1.2 <i>Electroweak symmetry breaking</i>	3
1.1.3 <i>Cross sections and decay rates</i>	5
1.1.4 <i>QCD and Proton Structure</i>	9
1.2 Previous $Wb\bar{b}$ Measurements	10
1.3 Dark Matter	11
1.3.1 <i>Experimental motivations</i>	11
1.3.2 <i>Simplified theoretical models</i>	12
1.3.3 <i>Previous DM Searches</i>	13
2 Phenomenology of Processes	14
2.1 The Standard Model process $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$	14
2.1.1 $pp \rightarrow W$	14
2.1.2 $W \rightarrow \ell\nu$	15
2.1.3 $g \rightarrow b\bar{b}$	16
2.2 The Standard Model process $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$	17
2.2.1 $pp \rightarrow Z/\gamma$	17
2.2.2 $Z \rightarrow \nu\bar{\nu}$	18
2.2.3 $pp \rightarrow Z\gamma$	18
2.3 Dark matter models of $pp \rightarrow \gamma + \text{invisible}$	19
3 The LHC and CMS	21

3.1	The Large Hadron Collider	21
3.1.1	<i>LHC pre-acceleration</i>	22
3.1.2	<i>LHC acceleration</i>	22
3.2	The Compact Muon Solenoid Detector	26
3.2.1	<i>Geometry</i>	27
3.2.2	<i>Magnet</i>	28
3.2.3	<i>Tracking System</i>	29
3.2.4	<i>Electronic Calorimeter</i>	30
3.2.5	<i>Hadronic Calorimeter</i>	32
3.2.6	<i>Muon System</i>	34
3.2.7	<i>Trigger and Data Acquisition</i>	35
4	Event Simulation and Reconstruction	37
4.1	Simulation of Events	37
4.1.1	<i>Monte Carlo Event Generation</i>	37
4.1.2	<i>Monte Carlo Generators</i>	39
4.1.2.1	<i>Matrix Element Calculators</i>	39
4.1.2.2	41
4.1.3	<i>Detector Simulation</i>	41
4.2	Reconstruction of Events	42
4.2.1	<i>Track and Primary Vertex Reconstruction</i>	42
4.2.2	<i>Electron ID and Reconstruction</i>	43
4.2.3	<i>Muon ID and Reconstruction</i>	44
4.2.4	<i>Lepton Isolation</i>	45
4.2.5	<i>Photon ID and Reconstruction</i>	46
4.2.6	<i>Missing Transverse Energy</i>	46
4.2.7	<i>Jet ID and Secondary Vertices</i>	47
5	$Wb\bar{b}$ Cross Section Measurement	49
5.1	Event Selection	49
5.2	Simulated samples	50

5.3	Background Estimation	52
5.3.1	<i>QCD</i>	52
5.3.2	<i>W+jets: light and charm component</i>	55
5.3.3	<i>Single top backgrounds</i>	55
5.3.4	<i>Z$\ell\bar{\ell}$ backgrounds</i>	59
5.4	Analysis Strategy	59
5.5	Systematic uncertainties	63
5.6	Signal Extraction	64
5.7	Cross Section and Comparisons	66
6	Monophoton Analysis	75
6.1	Event Selection	75
6.2	Estimation of Background Contributions	78
6.2.1	<i>Reweighting</i>	79
6.2.2	<i>Vγ Estimates</i>	79
6.2.3	<i>Electron Mis-ID</i>	80
6.2.4	<i>Non-collision Backgrounds</i>	81
6.2.5	<i>Minor SM Processes</i>	81
6.3	Results	82
6.3.1	<i>pp $\rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ Cross Section Measurement</i>	83
6.3.2	<i>Limits on Dark Matter</i>	85

LIST OF FIGURES

1.1	Illustrations of SM couplings	5
1.2	One-loop corrections to vertices and propagator	6
2.1	Feynman diagrams for $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$	14
2.2	Feynman diagrams for $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$	17
2.3	Feynman diagrams for dark matter monophoton	20
3.1	The LCH accelerator complex at CERN	23
3.2	The LHC dipole	24
3.3	The CMS Detector	27
3.4	The CMS Tracking System	30
3.5	The CMS Electromagnetic Calorimeter	32
3.6	The CMS Hadronic Calorimeter	34
3.7	The CMS Muon System	36
4.1	Example Parton Distribution Function from CTEQ6M	38
5.1	QCD shape for $Wb\bar{b}$ analysis	54
5.2	Wjj control region for the $Wb\bar{b}$ measurement	56
5.3	Diagram for single top t -channel control region	57
5.4	Single-top control region for the $Wb\bar{b}$ measurement	58
5.5	$Zb\bar{b}$ control region for the $Wb\bar{b}$ analysis	60
5.6	Zjj control region for the $Wb\bar{b}$ measurement	61
5.7	Three-step fitting procedure for $Wb\bar{b}$ measurement	72
5.8	More variables in the $Wb\bar{b}$ signal region	73
5.9	Cross section comparison for $Wb\bar{b}$ and generators	74
6.1	Signal region distributions in the $pp \rightarrow \gamma + \text{invisible}$ analysis	82
6.2	Exclusion plots in $m_\chi - m_M$ plane	87
6.3	Exclusion limits on DM-nucleon cross section	87

6.4	Expected lower limit for EFT cutoff parameter Λ	88
-----	---	----

LIST OF TABLES

1.1	Fermion couplings	2
1.2	Fundamental particle decay channels and rates	8
5.1	$Wb\bar{b}$ Cutflow	50
5.2	Systematic uncertainties in $Wb\bar{b}$	70
5.3	Measured breakdown of samples in $Wb\bar{b}$	71
5.4	Measured cross sections in the muon, electron, and combined lepton channels.	71
6.1	Monophoton Cutflow	78
6.2	Estimated yields for $pp \rightarrow \gamma + \text{invisible}$	83
6.3	Systematic uncertainties in $pp \rightarrow \gamma + \text{invisible}$	84
6.4	Upper limits on DM cross section for vector $m_M = 10$ TeV	85
6.5	Upper limits on DM cross section for axial-vector $m_M = 10$ TeV	85

1 INTRODUCTION

One of the goals of modern society is to experimentally investigate particles and interactions, and to achieve this goal, the governments from nearly 100 different countries, states and territories are presently funding over 13,000 scientists, engineers and technicians to build, operate, maintain and analyze data from the Large Hadron Collider (LHC) at the European Center for Nuclear Research (CERN). This thesis presents analyses of data taken with the Compact Muon Solenoid (CMS) detector using proton-proton collisions provided by the LHC during its operation in 2012 and 2015. Using the 2012 data, a measurement of the production cross section for the standard model (SM) process $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$ is performed, where ℓ is an electron or a muon. Using the 2015 data, the monophoton signature is examined in the context of the SM process $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ and as a search for evidence of dark matter (DM).

1.1 The Standard Model

1.1.1 Standard Model Particles

The SM Lagrangian uses the Yang-Mills construction on the group $SU(3) \times SU(2) \times U(1)$. Strong interactions are described by the symmetry $SU(3)$ (color) which requires eight gauge bosons, known as gluons. Electroweak interactions are described by $SU(2) \times U(1)$ are moderated by the massive W^\pm and Z bosons and the massless photon (γ).

Table 1.1: Fermions can exist in doublet and singlet configurations which are listed below along with the gauge fields the configuration interacts with, denoted by the symmetry the field respects.

Fermion type	Interactions		
	$SU(3)$	$SU(2)$	$U(1)$
All quarks (doublet)	Yes	Yes	Yes
All leptons (doublet)	No	Yes	Yes
<i>u</i> -type quarks (singlet)	Yes	No	Yes
<i>d</i> -type quarks (singlet)	Yes	No	Yes
Charged leptons (singlet)	No	No	Yes

The matter and antimatter components are spin- $1/2$ fermions, which come in two families, known as leptons and quarks. Quarks and leptons come in three families, known as generations for the quarks and flavors for the leptons. The leptons are colorless and therefore do not couple to gluons, but the quarks do and additionally come in three color varieties for each generation. The singlet configurations contain only electrically charged fermions.

All quarks and leptons can exist in $SU(2)$ doublet configurations, and the different I_3 values further break quarks and leptons into types. For quarks, there are *u*-type (u, c, t) = (up, charm, top) which have charge $+2/3$, and there are *d*-type (d, s, b) = (down, strange, bottom) which have charge $-1/3$. Leptons are either charged (e, μ, τ) = (electron, muon, tauon) or neutral (ν_e, ν_μ, ν_τ) = (electron-, mu-, tau-neutrino), and all charged fermions are arranged such that mass increases with successive generations within a given type. In units where $\hbar = c = 1$, the top ($m_t = 173$ GeV) and bottom ($m_b = 4$ GeV) are the heaviest quarks of their respective types, with the bottom weighing three orders of magnitude greater than the lightest quark, up ($m_u = 2$ MeV).

The mass separation for the charged leptons ($m_\tau = 1.7$ GeV, $m_e = 0.5$ MeV) also spans multiple orders of magnitude, but while the observations of neutrino oscillations indicate that neutrinos have mass, only upper limits on the values they may have have been set on the order of MeV.

1.1.2 Electroweak symmetry breaking

Fermions and gauge bosons acquire mass through the Brout-Englert-Higgs mechanism in which a scalar field is introduced to the SM Lagrangian in the doublet configuration. This generates mass for the Higgs boson (m_H) as well as for the W^\pm and Z bosons, while leaving the photon massless,

$$M_Z = \frac{v}{2} \sqrt{g_1^2 + g_2^2}, \quad M_W = \cos \theta_W M_Z \quad , \quad (1.1)$$

where the Weinberg angle, $\theta_W = g_1/g_2$, is the ratio between the $U(1)$ and $SU(2)$ coupling strengths, and $v \propto m_H$ is vacuum expectation value of the Higgs field. The \pm in W^\pm aligns with the electric charge of the W boson, and the Z and γ are both electrically neutral and orthogonal.

The vector boson couplings to the fermions are then expressed as three types of currents. Electromagnetic interactions involve couplings between the photon and charged particles, and any particle which interacts with the photon can also interact with the Z boson. The neutral current interactions have additional couplings to the lepton and quark doublets and charged current interactions are between the doublets and the W bosons. Couplings to each of the bosons are scaled by a coupling factor

which decides the relative strengths of the interactions. The coupling factors are related to the strengths of the unbroken $U(1)$ and $SU(2)$ couplings as

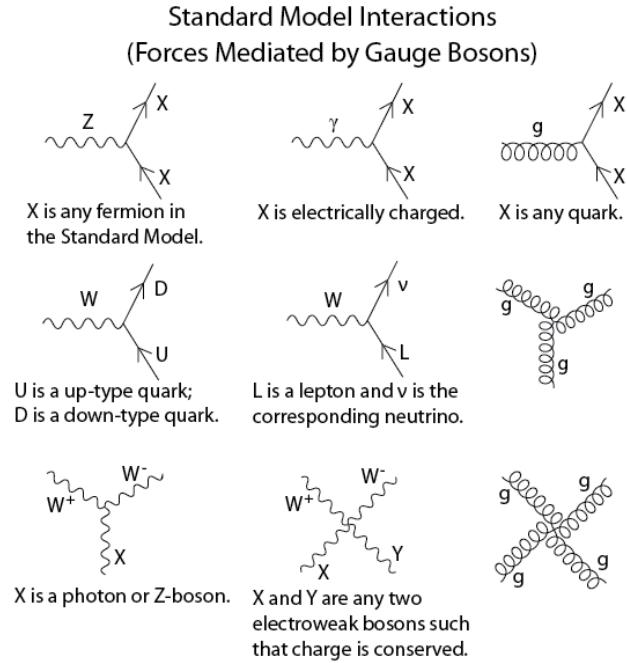
$$e = \frac{g_1 g_2}{\sqrt{g_1^2 + g_2^2}}, \quad g_Z = \frac{e}{\cos \theta_W \sin \theta_W}, \quad g_W = \frac{e}{\sqrt{2} \sin \theta_W} \quad (1.2)$$

In addition to generating mass terms for the bosons, the Higgs field gives rise to fermion mass via Yukawa couplings between the Higgs and fermion doublets. It is because the mass of the Higgs boson is tied to the masses of the gauge bosons and charged fermions that the discovery of the Higgs boson in 2012 was of such importance. The result was announced just a few months after I moved to Geneva to do research at CERN, and it provided the first measurement of a parameter in the SM that had been previously unknown, $m_H = 125$ GeV.

To describe the scattering of fermions and gauge bosons in the SM, diagrams with the appropriate initial and final state particles are constructed from the vertices illustrated in Figure 1.1. External lines correspond to real particles which are observed in the final state and are on-shell. Internal lines correspond to virtual particles that are not observed in the final state and therefore may or may not be on-shell. A diagram is tree-level or leading order (LO) if it uses the smallest number of vertices possible to depict an interaction having the correct initial and final states.

Renormalization is necessary to account for corrections to the tree-level propagators and vertices which arise from contributions of virtual particles connecting to form closed internal loops. The lowest order diagrams with one internal loop, are illustrated in Figure 1.2. Because these diagrams contain the next fewest number of

Figure 1.1: The couplings for the SM gauge interactions.

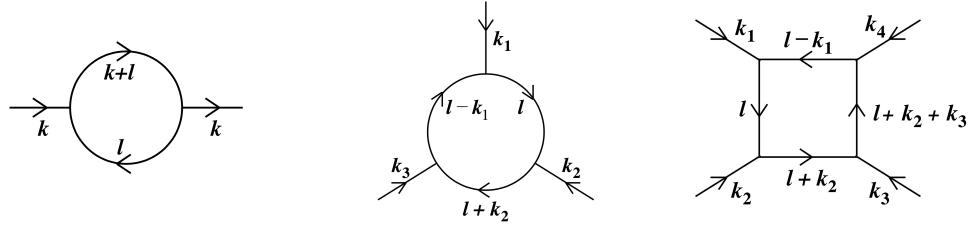


possible vertices above LO, they are termed next-to-LO (NLO). Renormalization is accomplished by introducing an energy scale, and assuming that the couplings are small compared to this scale. The coupling constants are therefore functions of energy and are quoted at a particular renormalization scale, μ_R . Renormalization works Next-to-NLO

1.1.3 Cross sections and decay rates

Quantum mechanics dictates that only predictions of probability are possible, and the final probability of observing a particular interaction is dependent on many variables, including the energies, types and angular momenta of the incoming and outgoing

Figure 1.2: Renormalization takes place via a modification of the coupling parameters g_1, g_2, g_3 to account for the contributions that loops of virtual particles make on the propagators and vertices. Below are diagrams keeping track of the flow of momentum for one-loop corrections to the propagator and three and four point vertices.



particles as well as the masses of the propagators and the orientation and efficiency of the detector.

A quantity typically measured is therefore the interaction cross section, σ , and for the scattering of two incoming particles going to n' particles, $2 \rightarrow n'$, in the CM frame, the differential is

$$d\sigma = \frac{1}{4|\mathbf{k}_1|_{\text{CM}}} |\mathcal{T}|^2 d\text{LIPS}_{n'}(k_1 + k_2) \quad (1.3)$$

where the scattering matrix element, \mathcal{T} , is defined using Equation ??, as

$$\langle f | i \rangle = (2\pi)^4 \delta^4 \left(\sum k_{\text{in}} - \sum k_{\text{out}} \right) i \mathcal{T} \quad (1.4)$$

and the Lorentz-invariant measure of the phase space for the n' outgoing particles is

$$d\text{LIPS}_{n'}(k) = (2\pi)^4 \delta^4 \left(k - \sum_{j=1}^{n'} k'_j \right) \prod_{j=1}^{n'} dq'_j \quad (1.5)$$

with the Lorentz-invariant differential dq .

The cross section is used to calculate the rate at which a process occurs, but is not the only relevant factor in determining the overall production rate. The production rate of a given final state is also dependent on the incoming rate of possible interactions and is known as luminosity, \mathcal{L} . Luminosity has the units of inverse area per unit time and the total number of events produced is therefore proportional to $\int \mathcal{L} dt$. In any real detector, final state particles are collected only within a finite solid angle and the number of particles scattered into a given solid angle, Ω , is given by

$$\frac{dN}{d\Omega} = \mathcal{L} \frac{d\sigma}{d\Omega} . \quad (1.6)$$

It is also possible for particles to decay as $1 \rightarrow n'$. Massive particles decay to lighter ones in both the fermion and boson sectors, with all massive bosons able to spontaneously decay via the diagrams in Figure 1.1. Of the charged fermions, only the first generation is stable for each type, and neutrinos are not known to spontaneously decay, but oscillate between flavors while propagating in free space. Like the differential cross section, the differential decay rate is a function of the scattering amplitude and has integration measure $d\text{LIPS}$,

$$d\Gamma = \frac{1}{2E} |\mathcal{T}|^2 d\text{LIPS}_{n'}(k) . \quad (1.7)$$

The differential decay rate is inversely proportional to the energy of the particle, $E = \sqrt{m^2 + p^2}$. This means that comparatively heavy particles will decay faster than comparatively light ones and that energetic particles will appear to live longer for a

Table 1.2: Below are listed the decay channels and rates for each of the unstable fundamental particles. At CMS, with the detection apparatus located a finite distance away from the interaction vertex, particles such as the W , Z and Higgs bosons, as well as the t and tau, decay before reaching the first layer of the detector.

Particle	Primary decay modes(s)	Total rest-frame $d\Gamma$	Typical decay location
W	$W \rightarrow \ell\nu$	X	Before reaching CMS
Z	$Z \rightarrow f\bar{f}$ (for $2M_f < M_Z$)	X	Before reaching CMS
τ	$\tau \rightarrow W\nu_\tau$	X	Before reaching CMS
μ	$\mu \rightarrow W\nu_\mu$	X	After leaving CMS
t	$t \rightarrow W^+b$	X	Before reaching CMS
b	$b \rightarrow W^-c$	X	Inside CMS
c	$c \rightarrow W^+s$	X	Inside CMS
s	$s \rightarrow W^-u$	X	Inside CMS

stationary observer due to relativistic time dilation effects. The total decay rate of a given particle is found by summing the decay rates from each of the contributing processes, and the primary decay channels and rates for the fundamental particles are given in Table 1.2.

At CMS, the heaviest quark and the heaviest lepton both decay before reaching the detector volume. This makes b quarks the heaviest fundamental particles which can be seen to decay inside the detector, and therefore an object of interest. Additionally, their heavy mass means that they couple strongly with the Higgs boson which still has many properties that are under investigation. The W and Z bosons are both so massive that they decay before reaching the innermost layers of the detector and are often identified by their decay products pointing back to a common vertex.

1.1.4 QCD and Proton Structure

Feynman diagrams are used to describe interactions between fundamental particles, but at the LHC, collisions take place between protons, which are composite.

One feature of the $SU(3)$ symmetry of the strong force is that gluons carry one unit of color and one unit of anticolor while the quarks carry one unit of color charge. This is what allows gluons to interact with each other as well as with quarks. That quark confinement is necessitated by the $SU(3)$ structure has not been conclusively determined, but observationally, a free gluon or quark has never been observed. Instead, quarks appear as bound in colorless (singlet) combinations called hadrons which are further classified as mesons ($q\bar{q}$) or as baryons (qqq or $\bar{q}q\bar{q}$), and are held together by gluons. Evidently, the binding energy of the quarks has a form such that after a distance of roughly 10^{-15} meters, the energy stored in the gluon field is greater than the energy needed to create a quark-antiquark pair, bringing the pair into existence. This process of energetic quarks creating particles as they separate is called hadronization and is an important effect at the LHC.

Protons are a type of baryon and at low energy, may combine with a single electron to form a neutral hydrogen atom. At higher energies, the internal structure of the proton becomes more evident, and it contains three valence quarks, uud , which are constantly exchanging gluons. When probed at high enough energy, or equivalently, at short enough length scales, these gluons can also each split into a $q\bar{q}$ pair which typically reannihilate with each other. With gluons inside the proton splitting into quarks and coupling with other gluons, this forms a ‘sea’ of quarks and gluons, and as protons are accelerated to energies of GeV or TeV as is the case at the LHC, the

fraction of the momentum of the proton attributed to the gluons becomes higher than that attributed to the valence quarks.

A proton-proton collider was therefore a sensible choice for the LHC. The physics goals of the project are to measure quantities associated with a wide range SM processes and to continue the search for evidence of new physics. Quarks interact with all of the SM gauge bosons as well as with the Higgs boson and the proton contains the lightest quarks of each type in addition to the gluons and sea. Colliding proton beams thus allow for the interactions between many different initial particle configurations to be explored, and with the exception of the neutrinos which interact only via the weak exchange of the Z boson and escape the detectors, all other fundamental SM particles have been directly observed at CERN.

1.2 Previous $Wb\bar{b}$ Measurements

The production of W bosons in association with b jets has been studied at a center-of-mass energy of 7 TeV using data samples with up to 5 fb^{-1} of integrated luminosity, by the ATLAS and CMS experiments ??, as well as by the Collider Detector at Fermilab (CDF) and D0 Collaborations using $p\bar{p}$ collisions provided by the Tevatron ?? at Fermilab at $\sqrt{s} = 1.96 \text{ TeV}$. Neither of the experiments at the Tevatron directly examined a final state with more than one b -tagged jet, but both published results on the process $p\bar{p} \rightarrow Wbj \rightarrow \ell\nu bj$ where j signifies any jets. The CDF Collaboration observed a cross section that was a factor of two higher than the best NLO prediction available at the time, and the D0 Collaboration observed a cross section that was

20% – 40% higher than various NLO predictions. The ATLAS Collaboration also examined a final state having only one b -tagged jet in the process $pp \rightarrow Wb \rightarrow \ell\nu b$. They measured cross sections for final states with exactly one jet and exactly two jets, and their observed cross section in the one-jet channel was 70% higher than the NLO prediction. The agreement improved in the two-jet channel, where the observed cross section was 30% higher than the NLO calculation predicts.

The first measurement of a final state having two identified b jets and the products of a W decay was made by the CMS Collaboration at $\sqrt{s} = 7$ TeV in the process $pp \rightarrow Wb\bar{b} \rightarrow \mu\nu b\bar{b}$ where exclusively two jets were allowed, both of which b -tagged. In this measurement, the cross section was found to be within 4% of the NNLO prediction. The analysis presented in this thesis extends the $pp \rightarrow Wb\bar{b} \rightarrow \mu\nu b\bar{b}$ analysis to $\sqrt{s} = 8$ TeV and additionally incorporates W decays $W \rightarrow e\nu$, with the goal of providing a measurement of this process at higher energy than ever before and comparisons with the latest simulations and predictions.

1.3 Dark Matter

1.3.1 Experimental motivations

Albert Einstein’s theory of general relativity, GR, has many experimental predictions which run counter to human intuition. GR predicts that massive objects warp a four-dimensional spacetime and thus feel mutual attraction. This has as a consequence, the prediction that even massless objects such as photons will experience a net deviation in their path near a massive object as a result of gravity and this effect was famously

verified by Arthur Eddington through the observation of stars around the sun during a full solar eclipse. More recently, the direct detection of gravitational waves by the LIGO Collaboration also aligns with the GR predictions of distorted spacetime around colliding black holes. The time distortion effects due to the varying strengths of Earth's gravitational field on the surface and at the GPS satellites, provide precise tests of the quantitative predictions of GR. However, though these tests and others provide evidence that GR is an accurate theory of gravity, some basic predictions related to gravitational interactions do not agree with observations, motivating the concept for DM.

The first observational evidence for DM came from an analysis of the speeds of galaxies in the Coma cluster by Fritz Zwicky. The magnitude of the angular velocities of the galaxies was too great to be explained by the visible matter alone and DM is now believed to outweigh visible matter in a ratio of 5 : 1 throughout the universe and 10 : 1 throughout the Milky Way galaxy.

1.3.2 Simplified theoretical models

The defining features of DM are that it is massive and appears to interact on large scales only via the gravitational force. On the galactic and supergalactic scales, DM is distributed along similar structures as is visible matter, and it surrounds visible matter in extended halos.

Because DM has not yet been observed, the models of DM being considered in this thesis are simplified and based on minimal assumptions, the first being that DM is even capable of interacting with hadrons and is thus possible to produce at the

LHC. While the visible sector of particles is diverse, the models used in this analysis consist of a single DM particle, χ , which is assumed to be a fermion and may be different from $\bar{\chi}$.

One way DM could couple to the SM is via the addition of a $U(1)$ symmetry that gives rise to a vector gauge mediator, M . If some quarks are also charged under $U(1)$, then DM may be produced in the s channel as $f\bar{f} \rightarrow M \rightarrow \chi\bar{\chi}$. If M conserves parity in $f\bar{f} \rightarrow M$, it is said to have a vector coupling, and if it violates parity, it is termed axial-vector. In these models, M is assumed not to couple to leptons, but an effective field theory (EFT) model is also considered in this analysis which estimates a direct interaction between DM and photons. This coupling is mediated by a vertex $\gamma\gamma\chi\bar{\chi}$. and allows for DM production via the channel $pp \rightarrow \gamma \rightarrow \gamma\chi\bar{\chi}$.

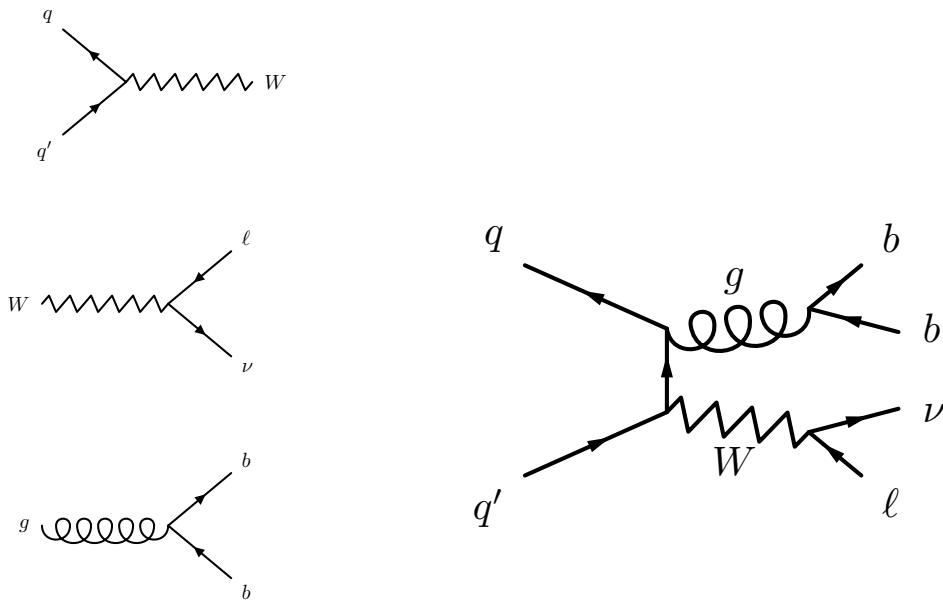
1.3.3 Previous DM Searches

A search for particle DM ? used 19.6 fb^{-1} of pp data collected in 2012 by the CMS Collaboration at $\sqrt{s} = 8 \text{ TeV}$, but no evidence was found and a new upper limit for the DM cross section was set. A similar search was performed by the ATLAS experiment ? using 3.2 fb^{-1} of pp collision data at $\sqrt{s} = 13 \text{ TeV}$ collected in 2015, but again no evidence for new physics was found. This imposed limits on the possible masses DM particles and the mediator mass as $m_\chi > 150 \text{ GeV}$ or $m_M > 710 \text{ GeV}$.

2 PHENOMENOLOGY OF PROCESSES

2.1 The Standard Model process $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$

Figure 2.1: The Feynman diagram for the process $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$ is illustrated below, and is composed from the individual vertices illustrated on the left, each of which is described in Section 2.1.



2.1.1 $pp \rightarrow W$

The W boson couples to all charged fermions and can be created during the collision of a quark-antiquark pair with a relative charge difference of e . In the proton are quarks and the most prevalent valence quark is the u . Therefore in a pp collision, the channel by which most W bosons are produced is via a the annihilation of a valence

u quark from one proton with a \bar{d} from the sea of the other, $u\bar{d} \rightarrow W^+$. Quarks of higher generation can also be found inside the sea as the result of gluons splitting into $q\bar{q}$ pairs, but all interactions are modified by a coefficient in the CKM matrix and higher generation mixing is thus suppressed. In this thesis, all modes of $pp \rightarrow W^\pm$ production are considered.

2.1.2 $W \rightarrow \ell\nu$

Just as the W boson can be created by the collision $q\bar{q}' \rightarrow W$, it can also decay as $W \rightarrow q\bar{q}'$. This is known as hadronic W decay and can be a useful analysis channel for experimentalists, especially for decay products with energies approaching the TeV scale. Leptonic W decay, $W \rightarrow \ell\nu$, is also an important channel for experimentalists and is the one considered in this analysis. Because leptons constitute a negligible fraction of the sea, the detection of leptons at high energy after a pp collision is often a good indicator of the decay of a massive gauge boson, $W \rightarrow \ell\nu$ or $Z \rightarrow \ell\bar{\ell}$.

The W boson is much heavier than any of the leptons and therefore decays with roughly equal probability to any of $e\nu_e, \mu\nu_\mu, \tau\nu_\tau$. From Table 1.2, tauons created at CMS subsequently decay before reaching the detector, so for this analysis, the decay channel of the W investigated is $W \rightarrow \ell\nu$ where $\ell \in e, \mu$.

To reconstruct muons from the decay products, the transverse mass, m_T variable is often used. It is defined by $m_T^2 = m^2 + p_x^2 + p_y^2$ where p_i is the component of the momentum along the i axis, and in the case of a massive particle decaying to two

massless particles, can be rewritten as

$$m_T^2 = 2p_{T,1}p_{T,2}(1 - \cos\phi) \quad (2.1)$$

where ϕ is the angle between the particles and $p_{T,j}$ is the component of the momentum of the particles in the transverse plane.

In the decay of a W boson, a neutrino is produced, but can not be detected. The CMS detector is designed to capture the energy from all of the other particles produced in a collision, so neutrinos are accounted for as E_T^{miss} . The variable E_T^{miss} is the transverse component of the negative vector sum of all of the energy identified as having come from a particular interaction vertex, known as the primary vertex. Therefore the transverse mass of the W boson is

$$(m_T^W)^2 = 2p_T^\ell E_T^{\text{miss}}(1 - \cos\phi) \quad (2.2)$$

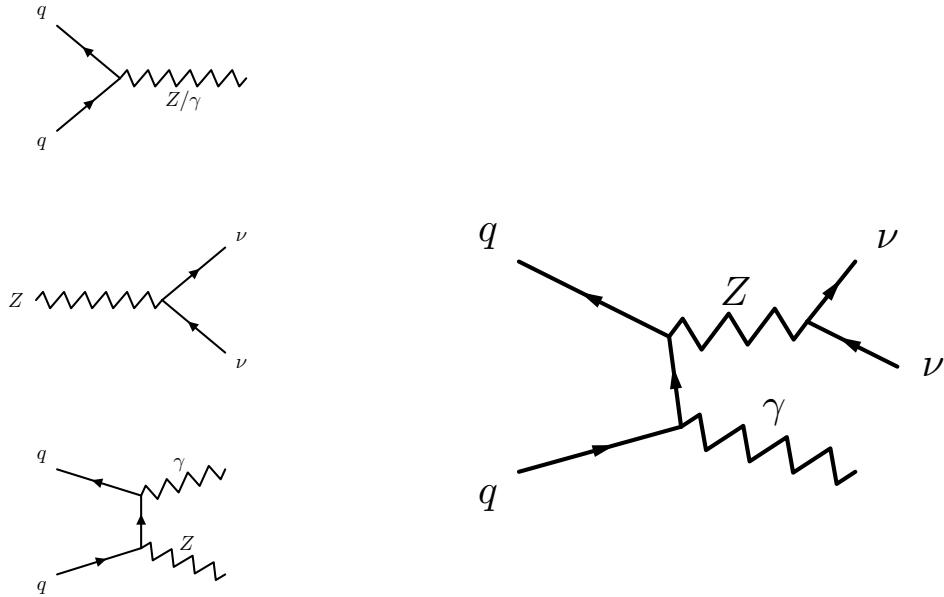
where ϕ is the angle between the lepton and E_T^{miss} .

2.1.3 $g \rightarrow b\bar{b}$

Because quarks couple strongly to gluons and $q\bar{q}' \rightarrow W$ has been shown to be an important production channel in pp collisions, it is possible for one of the initial state quarks to radiate a gluon. This is called initial state radiation, ISR, and if the gluon is produced with enough energy, it is capable of splitting to a quark-antiquark pair. In particular, a $g \rightarrow b\bar{b}$ vertex can be added to either of the incoming quarks to form $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$.

2.2 The Standard Model process $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$

Figure 2.2: The Feynman diagram for the process $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ is illustrated below, and is composed from the diagrams illustrated on the left, each of which is described in Section 2.2.



2.2.1 $pp \rightarrow Z/\gamma$

Similar to the W boson, the Z boson and the photon can also each be produced via the collision of quarks in the process $q\bar{q} \rightarrow Z/\gamma$. Unlike interactions with the W boson, interactions with Z/γ conserve parity invariance and do not transport charge. Any interaction which can happen as mediated by a photon can also happen with the exchange of a Z boson, but for collisions at $\sqrt{s} < M_Z = 90$ GeV, the Z can not be made on-shell. In this low energy regime γ exchange dominates, but in 2015, the

LHC ran at $\sqrt{s} = 13$ TeV and the relative mass difference between the Z and the γ played a negligible role in their relative rates of production.

2.2.2 $Z \rightarrow \nu\bar{\nu}$

The only particle which the Z boson can couple to but the photon can not is the neutrino. At $\sqrt{s} = 13$ TeV, the mass differences between the five lightest flavor of quark and the six leptons are negligible and the Z boson can decay into any kinematically allowed pairs, $Z \rightarrow f\bar{f}$. Including the three color possibilities for each quark, these are $3 \times 5 + 6 = 21$ final states, each of which has roughly the same branching fraction. Therefore only approximately $2/7$ of Z decays happen in the leptonic channel $Z \rightarrow \ell\bar{\ell}$. and of these decays, approximately $2/3$ happen as $Z \rightarrow \nu\bar{\nu}$.

2.2.3 $pp \rightarrow Z\gamma$

With the photon coupling only to electrically charged particles, the only place where a vertex containing a photon could be attached to either of the upper two left diagrams in Figure 2.2 is on one of the quarks. Photons are massless and therefore stable and so are a final state observable. Like the gluon in $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$, the photon is an example of ISR. In the CM frame of the colliding $q\bar{q}$, which is approximately the lab rest frame for colliding beams of equal energy as is the case at the LHC, conservation of momentum dictates that the Z boson and photon should have equal and opposite momenta.

Unlike any of the other fermions, neutrinos are electrically neutral and therefore only interact via the weak force. So while the cross sections for most fermion-fermion

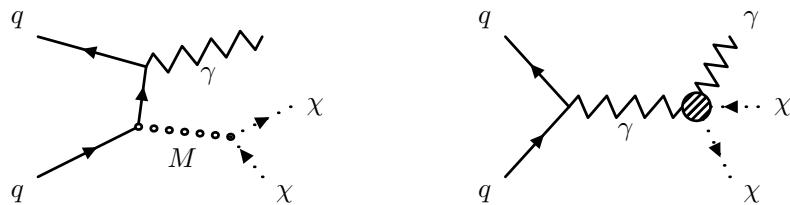
interactions involve contributions from the comparatively stronger electromagnetic and strong forces, the neutrino cross section contains contributions from only the W and Z bosons at tree-level and is much smaller than that of the charged fermions. This makes the detection of neutrinos very difficult in general, and impossible to do with present technology given the extreme backgrounds present in a collider setting.

In the case where the ISR photon is recoiling against a Z boson which decays to neutrinos, no direct detection of the Z boson or of its decay products is possible, leaving only the photon visible in the final state. This is called the monophoton signature, where a photon is observed recoiling against apparently nothing, and while the monophoton signature is predicted to be observed as a result of SM process as in $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$, if the observed monophoton cross section is measured to be higher than predicted, it could also be an indicator of physics beyond the SM (BSM). Specifically, the monophoton signature used in searches for dark matter.

2.3 $pp \rightarrow \gamma + \text{invisible}$

The existence of particle dark matter is well motivated, and the simplified model theories of DM used in this thesis allow for interactions which can result in the monophoton signature. One of the classes of models considered is a $U(1)$ gauge theory in which $\chi\bar{\chi}$ is produced via a vector or axial-vector mediator M which couples to quarks. The tree-level process in this model which leaves a monophoton signature is illustrated on the left of Figure 2.3, and the relevant parameters governing the cross section of this interaction are the masses of the two particles, m_χ and m_M ,

Figure 2.3: Feynman diagrams for the DM process $pp \rightarrow \gamma + \text{invisible}$ using simplified models are illustrated below. On the left is the $U(1)$ gauge model in which DM production is mediated by M which can be either vector or axial-vector. On the right is the diagram for DM production using an EFT model of the $\gamma\gamma\chi\bar{\chi}$ coupling.



and the strengths of the couplings between M and quarks, g_{Mq} , and between M and DM, $g_{M\chi}$. An EFT describing the four-point interaction vertex $\gamma\gamma\chi\bar{\chi}$ is also considered and illustrated on the right side of Figure 2.3. In this theory, the coupling is a function of two parameters, k_1 and k_2 , and is moderated by a mass scale, Λ . The other parameter in the EFT is m_χ , and by measuring the cross section for $pp \rightarrow \gamma + \text{invisible}$ in comparison with the SM prediction, estimations or limits can be set on the parameters used in either of these two models.

3 THE LHC AND CMS

The measurements presented in this thesis are performed on data of proton-proton collisions collected by the Compact Muon Solenoid (CMS) detector and provided by the Large Hadron Collider (LHC) at the European Center for Nuclear Research (CERN). The LHC was designed to probe physics at the scale of TeV and is capable of operating at multiple energy scales. As measured in the CM frame of protons colliding inside CMS, the LHC operated at $\sqrt{s} = 8$ TeV in 2012 and $\sqrt{s} = 13$ TeV in 2015. The measurement of the $pp \rightarrow Wb\bar{b}$ cross section is performed using 19.8 fb^{-1} of integrated luminosity collected at $\sqrt{s} = 8$ TeV and the measurement of the $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ cross section and the extensions to set limits on DM models uses 2.3 fb^{-1} of data collected at $\sqrt{s} = 13$ TeV.

3.1 The Large Hadron Collider

The LHC is a single-ring, double-bore particle accelerator and collider located on the border of France and Switzerland outside Geneva. It was built using the existing 26.7 km of tunnels from the Large Electron Positron collider and hosts four primary experiments, located at four interaction points where beams of hadrons are made to cross. Of the four experiments, two (CMS and ATLAS) are built for studying SM processes and searching for new physics in general, one (ALICE) is designed to investigate quark-gluon plasma resulting from the high energy collisions of heavy ions such as lead, and one (LHCb) was built for the study of b-mesons and CP violation.

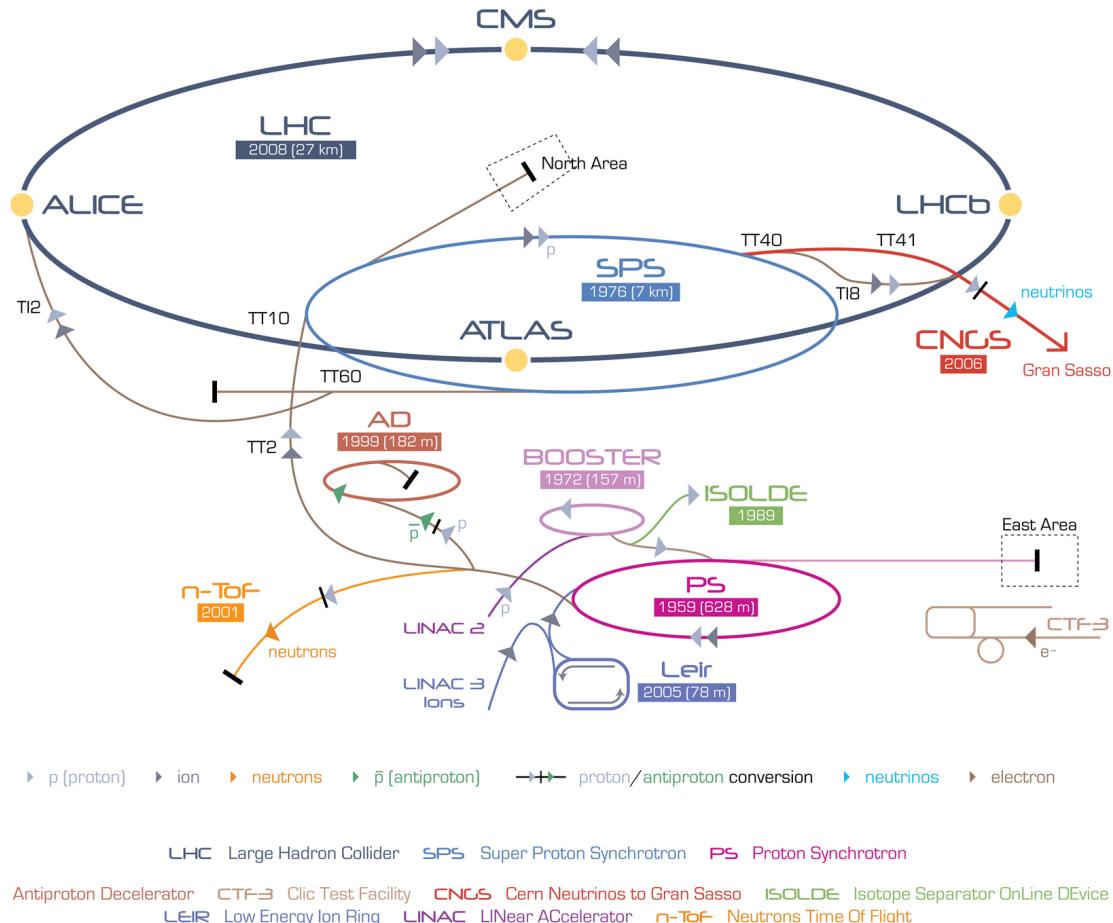
3.1.1 LHC pre-acceleration

To accelerate protons to their collision energy, a multi-stage procedure is used and the major components of the accelerator infrastructure are illustrated in Figure 3.1. First, protons are separated from electrons in neutral hydrogen gas before entering the linear accelerator (LINAC2) which brings them up to an energy of 50 MeV using a series of oscillating electric potentials. In this process, rather than having a continuous stream of accelerating protons, the protons are grouped into bunches, and the beam retains this structure of distinct groups of protons separated by gaps throughout the acceleration procedure. After the LINAC2, protons enter the Proton Synchrotron Booster (BOOSTER) where they are accelerated to 1.4 GeV and prepared for injection into the Proton Synchrotron (PS). Inside the PS, bunches are accelerated to 26 GeV before being injected into the Super Proton Synchrotron (SPS) where they are further accelerated to 450 GeV. After the SPS, bunches of protons are sent into the LHC.

3.1.2 LHC acceleration

The work of accelerating and containing the protons which form the beam of the LHC is done by superconducting magnets. They are cooled to a temperature of 1.9 K using liquid helium and are housed in the LHC dipole apparatus diagrammed in Figure 3.2. The dipole contains two beam pipes which are each surrounded by superconducting coils of Niobium Titanium (NbTi) which carry oscillating currents when in operation. These constitute RF cavities operating at 400 MHz and having the ability to circulate proton bunches in opposing directions between the two beam pipes with a spacing of 25 ns between bunches. The magnets are capable of reaching a strength of over 8 T,

Figure 3.1: Before protons are released into the LHC for final acceleration and collision, they pass through the LINAC2, BOOSTER, PS, and SPS, undergoing acceleration at each stage.



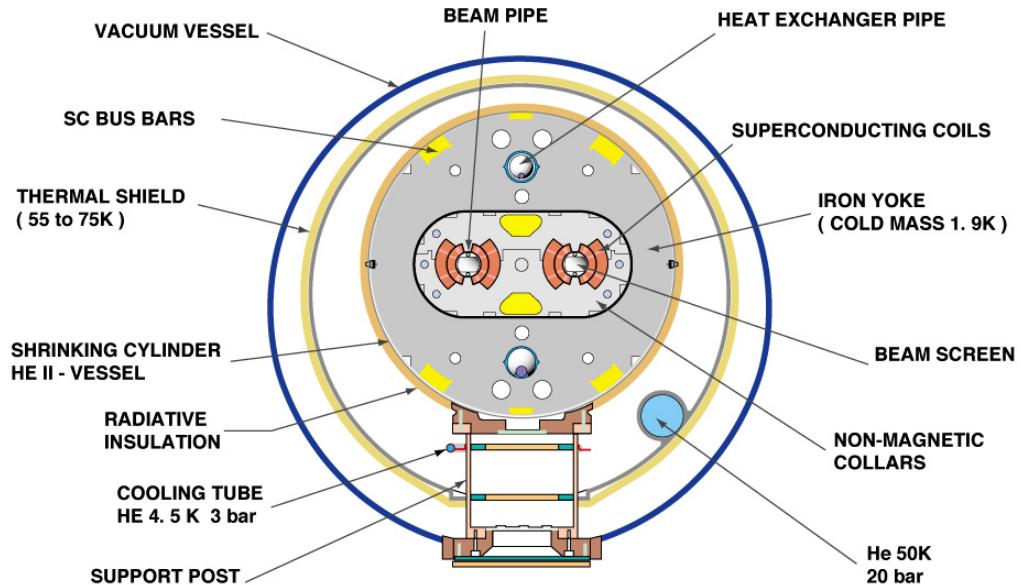
a constraint imposed by the desired energy scale of the accelerator and the radius of the existing LEP tunnels in which the LHC was built.

The Lorentz-invariant magnetic force on a particle of charge q , moving at velocity \mathbf{v} in a magnetic field \mathbf{B} is

$$\mathbf{F}_B = q\mathbf{v} \times \mathbf{B} \quad (3.1)$$

Figure 3.2: Below is a cross section of the LHC dipole apparatus. It contains two beam pipes, each surrounded by superconducting magnetic coils which are held in place by an iron yoke. The system is cooled to a temperature of 1.9 K and is thermally isolated as well as protected from radiation.

CROSS SECTION OF LHC DIPOLE



CERN AC _HE107A_ V02/02/98

and the relativistic version of Newton's second law states

$$\mathbf{F} = \frac{d\mathbf{p}}{dt} = \frac{d}{dt} (\gamma m_0 \mathbf{v}) \quad (3.2)$$

where γ is the relativistic correction factor and m_0 is the rest mass of the particle. These can be solved to find that the magnetic field required to hold a proton in planar

circular motion as in the LHC is

$$\mathbf{B} = \frac{p}{qR} \quad (3.3)$$

where $R = 4.3$ km. At individual beam energies of 4 and 6.5 TeV in 2012 and 2015 respectively, the rest mass of the proton constitutes a negligible fraction of the proton momentum and the minimum magnetic field required by the magnets is [todo - unit conversion].

The rate at which a particular collision process occurs at the LHC is proportional to the cross section of that interaction and the luminosity of the colliding beams as given in Equation 1.6. Assuming a Gaussian beam distribution, the machine parameters determine \mathcal{L} as

$$\mathcal{L} = \frac{N_b^2 n_b f_{\text{rev}} \gamma_r}{4\pi \epsilon_n \beta^*} \mathcal{F}(\theta) \quad (3.4)$$

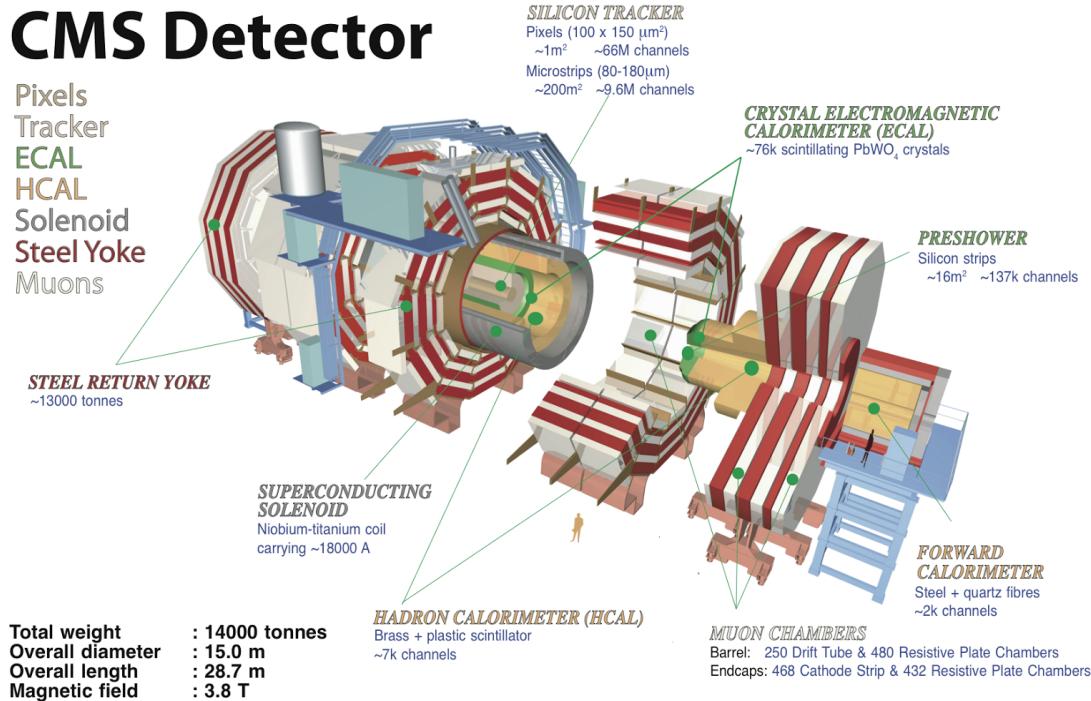
where N_b is the number of particles per bunch, n_b is the number of bunches per beam, f_{rev} is the revolution frequency of the bunches, $\gamma_r = E_p/m_p$ is the relativistic gamma factor for protons at energy E_p , ϵ_n is the normalized emittance which characterizes bunch width, β^* is a measure of the betatron oscillation envelope, and $\mathcal{F}(\theta)$ is a relativistic geometrical correction factor which is a function of the angle at which the beams cross. In addition to pushing the energy frontier, the LHC also has a significantly greater \mathcal{L} than previous hadron colliders. [Reference to Tevatron]

3.2 The Compact Muon Solenoid Detector

The CMS detector was built at Interaction Point 5 on the LHC ring to collect particle collision data exploiting the full physics reach of the LHC. The analysis of these data includes the discovery of the Higgs boson[REF] and high precision measurements of SM processes, as well as searches for physics beyond the standard model. To be able to perform such precision measurements, CMS was designed with four main subdetectors that work in concert and with a superconducting solenoid. The tracking and most of the calorimetric detectors are inside the solenoid while the muon detectors are outside. When running, the solenoid produces a 3.8 T uniform magnetic field in its interior, and has a uniform 2 T field over the bulk of the detector external to the solenoid.

The innermost of the subdetectors is the tracker which uses silicon pixel and strip detectors to record the tracks of charged particles passing through it. The tracks are used in conjunction with the 3.8 T magnetic field to measure the momentum of these particles and this information is used for identifying the pp interaction vertex as well as locating secondary vertices from the decay of heavy flavor quarks such as the b or c. Outside the tracker is the electromagnetic calorimeter (ECAL), which is designed to have good energy resolution in recording the electromagnetic interactions of charged particles such as electrons or photons over a wide range of angles. The hadronic calorimeter (HCAL) is outside the ECAL and is designed to absorb energy which comes in the form of neutral hadrons and provide good resolution in missing transverse energy, E_T^{miss} . Outside the calorimeters is the solenoid and steel return yolk, and the outermost layers of the detector are dedicated to the efficient detection

Figure 3.3: The CMS detector consists primarily of a tracker and electromagnetic and hadronic calorimeters which are mostly located inside a 3.8 T field provided by a superconducting solenoid, as well as a muon detection system located outside the solenoid.



of muons. The overall length of CMS is 21.6 m, with a radius of 7.3 m and a total weight of 12500 tons.

3.2.1 Geometry

The coordinate system used by CMS is one in which the z-axis is aligned with the beam pipe, the y-axis is pointing upward vertically and the x-axis points radially

inward toward the center of the LHC ring. The detector itself is mostly cylindrically symmetric about the beam pipe so cylindrical coordinates are also used. In this system, r is the radial distance as measured from the beam pipe, the azimuthal angle, ϕ , is measured up from the x-axis in the x-y plane, and the polar angle, θ , is measured down from the z-axis. The angle θ is commonly replaced by pseudorapidity,

$$\eta = -\ln(\tan \theta/2) \quad (3.5)$$

since the distribution of particles is roughly constant as a function of η . For the calorimeters, "barrel" refers to the region of $|\eta| < 1.4442$, and "endcap" to the region $3.0 > |\eta| > 1.566$. Instrumentation cables are run through the gap between the barrel and endcap, so this area has detecting components. The HCAL forward region covers $3.0 < |\eta| < 5$ and the tracker extends to $|\eta| < 2.5$.

3.2.2 Magnet

To precisely measure the momentum of a charged particle, it is necessary to measure radius of curvature of that particle as it moves through a magnetic field. The momentum resolution varies as

$$\frac{\delta p}{p} \sim \frac{1}{L^2 B} \quad (3.6)$$

where L is the length of the track of the particle through a magnetic field of strength B . For particles at high energy, this requires a very strong magnetic field which is achieved by the superconducting solenoid in CMS. The solenoid operates at 3.8 T

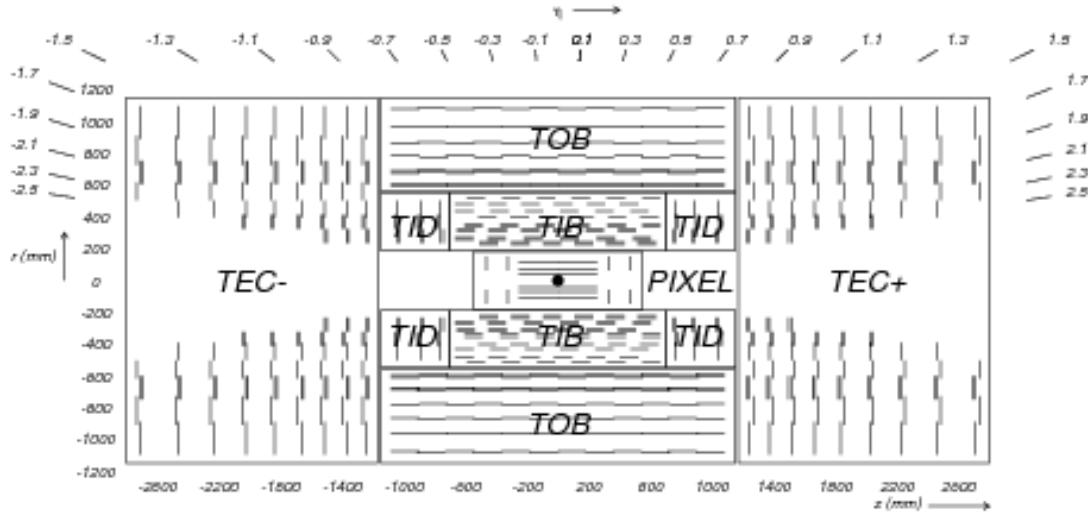
with a bore of 3 m in radius and 12.5 m in length and is constructed from four layers of NbTi superconductor. The steel yoke which provides physical support for the CMS structure and serves as an absorber for the muon system is fully saturated by the fringe magnetic field from the solenoid.

3.2.3 Tracking System

The inner tracking system of CMS is designed to provide precise and efficient measurements of the trajectories of charged particles produced during collisions, as well as a precise reconstruction of secondary vertices. The tracker has a length of 5.8 m and a radius of 1.25 m in a cylindrical structure surrounding the interaction point, as illustrated in Figure 3.4.

At the core of the tracker and closest to the beam line are three concentric cylindrical layers of hybrid pixel detector modules which are complemented by two discs of pixel modules on each end and extend a to a distance of 10 cm from the beam line. [TODO - Explain - silicon] In total, the pixel component of the tracker covers an area of about 1 m² with 66 million pixels. External to the pixel detector are the tracker inner barrel and discs (TIB/TID) which are made from silicon strips and extend out to a distance of 55 cm. There are four layers of strips in the TIB, with 3 discs at each end. The tracker outer barrel (TOB) is composed of 6 layers of micro-strip sensors and extends in z between ± 118 cm and to a radius of 116 cm. At the end of the z range for the TOB are the tracker end caps (TEC) which cover the ranges $124 < |z| < 282$ cm and $22.5 \text{ cm} < r < 113.5$ cm. Each TEC is composed of 9 discs, each carrying up to 7 rings of silicon micro-strip detectors. In total, the

Figure 3.4: Below is a schematic of the CMS tracking system where each line represents a detector module. The system is made from silicon pixels and silicon microstrips distributed into four sections, TIB, TID, TOB, TEC.



tracker contains 9.3 million strips which cover an area of 198 m^2 and extends to an acceptance of $|\eta| < 2.5$. For tracks with momentum on the order of 100 GeV, the momentum resolution is around 1-2% up to $|\eta| < 1.6$ and degrades to around 10% with increasing η .

3.2.4 Electronic Calorimeter

The electronic calorimeter (ECAL) is a homogeneous calorimeter made from nearly 76000 crystals of lead tungstate (PbWO_4) mounted in the barrel and endcap sections with a preshower detector located in front of the endcaps, arranged as shown in Figure 3.5 In the barrel, avalanche photodiodes are used as photodetectors, and in

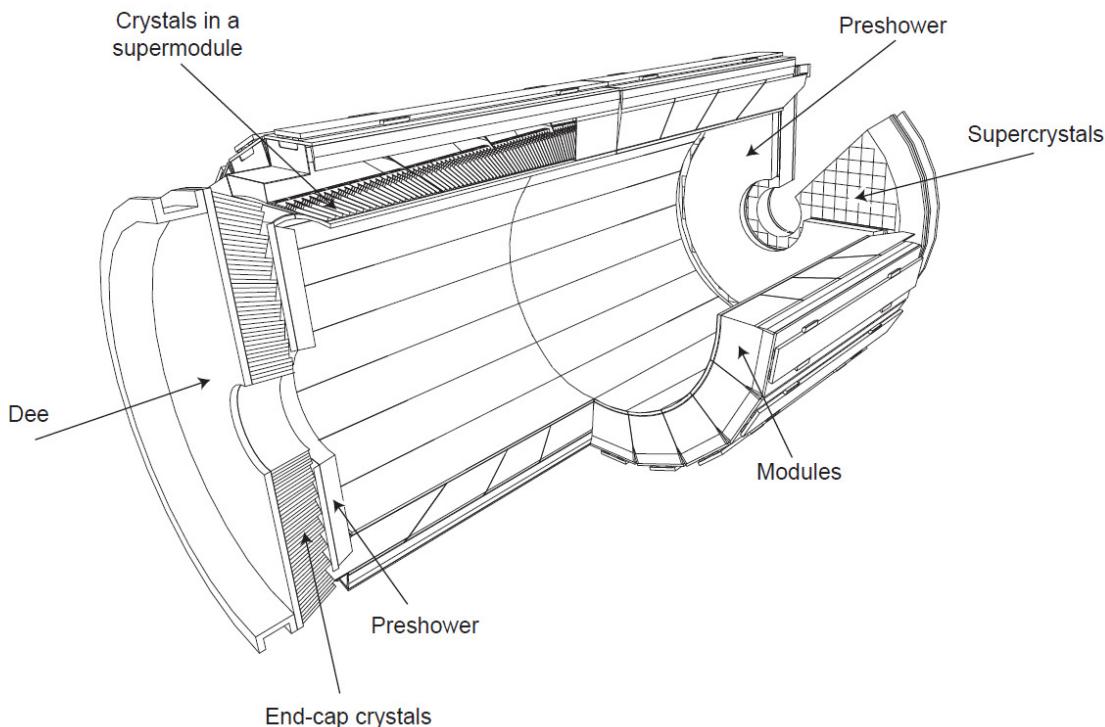
the endcap vacuum phototriodes are used. The material PbWO₄ was chosen for its properties of being dense, optically transparent and radiation hard. The radiation length inside the ECAL is typically less than 1 cm with a Moliere radius of 2.2 cm and about 80% of the light is emitted from a crystal within the first 25 ns. Since the length of a given crystal is on the order of 20 cm, most photons and electrons deposit all of their energy within the ECAL, and do not reach the HCAL.

The use of PbWO₄ crystals allows for excellent position and timing resolution with the energy resolution given by

$$\left(\frac{\delta_E}{E}\right) = \left(\frac{2.8\%}{\sqrt{E}}\right)^2 + \left(\frac{0.12}{E}\right) + (0.30\%)^2. \quad (3.7)$$

In this expression, the first term comes from the statistical error in the measurement which arises from the stochastic nature of electromagnetic shower evolution and the second term represents the error in the measurement which results from noise in the electronics or energy deposits from additional soft interactions. The ECAL provides stable and accurate measurements of energies over a range from 1 GeV to 1 TeV, with the upper limit set by the energy at which electromagnetic showers penetrate through the ECAL into the HCAL. [REF][TODO - how much?] With time, after undergoing a heavy bombardment of high energy radiation, the PbWO₄ crystals physically deteriorate and develop nonuniform light transmission properties. [REF] This is monitored and corrected for using a laser calibration system that probes for changes in crystal transparency.

Figure 3.5: Below is a diagram of the ECAL, which sits between the tracker and HCAL in CMS. It is made from PbWO₄ crystals throughout the volume with avalanche photodiodes in the barrel and vacuum phototriodes in the endcaps.



3.2.5 Hadronic Calorimeter

Situated mostly between the ECAL and the superconducting solenoid is the hadronic calorimeter (HCAL) which plays a crucial role in the measurement of hadron jets and particles such as neutrinos which escape the detector and result in apparent missing transverse energy. The HCAL is designed to contain the energy of neutral particles which pass through the ECAL and is therefore made from dense materials such as

steel and brass interleaved with scintillating material. Because the HCAL is designed to fit between these two components, it takes the shape of a hollow cylinder of inner radius 1.77 m and outer radius 2.95 m and one half of the HCAL is illustrated in Figure 3.6.

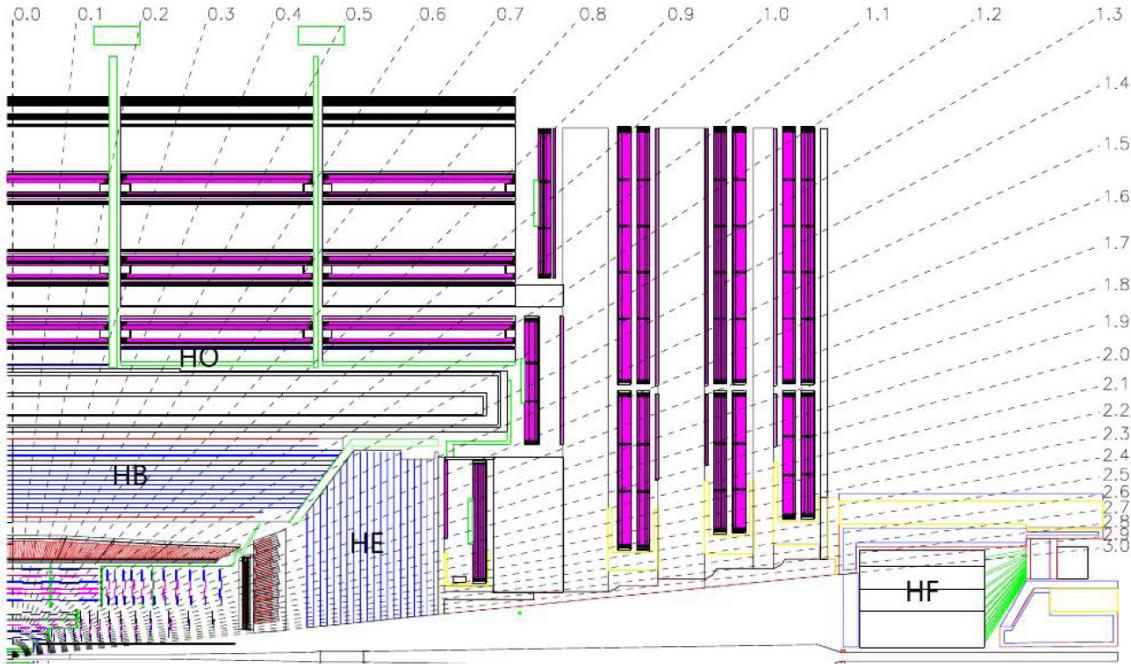
The barrel of the HCAL (HB) extends to $|\eta| < 1.3$ and is constructed from brass absorber plate wedges aligned parallel to the beam axis and mounted in an overlapping configuration, with a smaller amount of steel used in the inner and outermost wedges for structural stability. The endcap of the HCAL (HE) extends this coverage to $|\eta| < 3.0$ and is complemented by the forward hadron calorimeter (HF) which is made from the comparatively radiation-hard steel plates embedded with quartz fibers. Inside the barrel region there is an additional layer of the HCAL, the outer calorimeter (HO), which is located just outside the solenoid and uses it as an absorber for energetic showers which start late in the HB.

In the HB, HO and HE, light from particle showers [TODO - not exactly ..] inside scintillators and collected by quartz fibers and then used as an estimate of the total energy of the shower. In the HF, this estimate is made using the Cherenkov radiation from particles with energy above 190 keV collected by the quartz fibers. For the two cases, the energy resolution takes the same functional form

$$\left(\frac{\delta_E}{E} \right) = \left(\frac{A}{\sqrt{E}} \right)^2 + (B)^2 \quad (3.8)$$

where A is 90% (172%) in the HB/HO/HE (HF) and relates to the stochastic uncertainty of shower evolution and B is 4.5% (9.0%) and comes from uncertainties in calibration.

Figure 3.6: A schematic layout of the HCAL, which complements the ECAL in providing a measurement of the total energy produced in a collision. The HCAL is made from brass and steel plates, embedded with quartz fibers.



3.2.6 Muon System

Muons play a central role in the physics program outlined by CMS and the muon detection system is positioned as the outermost layer of the detector. Unlike the other charged leptons, muons typically pass through the ECAL and HCAL and deposit only a fraction of their energy, so a dedicated muon system is necessary in order to determine the momentum of these particles. The muon system is composed of three different kinds of gaseous detectors, drift tubes (DTs), resistive plate chambers

(RPCs) and cathode strip chambers (CSCs) and their layout is illustrated in Figure 3.7.

The barrel region of the muon system is covered by DTs in the range $|\eta| < 1.2$ and the endcaps are covered by CSCs in the range $0.9 < |\eta| < 2.4$. The RPCs are located in the range $|\eta| < 1.6$ and provide fast, independent and highly segmented transverse momentum measurements of muons.

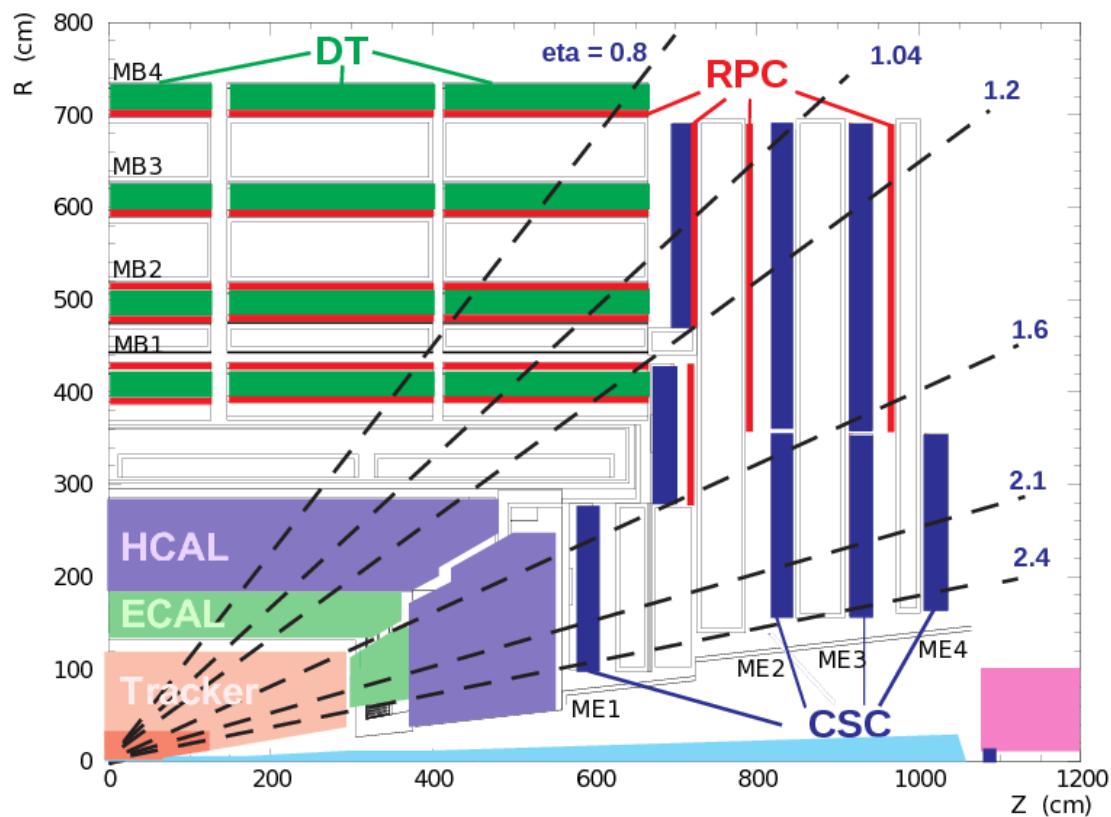
The DT system is composed of 4 stations which form concentric cylinders about the beam line and contain 172000 sensitive wires. As charged particles enter the DTs, they ionize the Ar/CO₂ gas mixture, knocking off electrons which then are attracted to the positively charged wires.

The CSCs are less sensitive to uneven magnetic fields and high particle rates so are therefore used in the endcaps. They are made from crossed arrays of positively charged wires and negatively charged strips in gas and are composed of six layers, giving them precise timing as well as positional information. As an upgrade between the 2012 and 2015 data taking periods, a fourth layer of CSCs was added to the CMS detector, adding to the three which were present in 2012.

The RPCs are built from two sheets held at opposite charges and separated by a gas volume. As muons move through the chamber, electrons are ionized from the gas and attracted to small metallic strips which they reach after a small but well known time delay. The timing resolution of RPCs is on the order of 1 ns.

3.2.7 Trigger and Data Acquisition

Figure 3.7: The CMS muon system uses DTs, RPCs, and CSCs to provide muon detection up to $\eta < 2.4$. Shown below is the geometrical arrangement of the different muon subsystems and how they fit with the rest of the CMS detector.



4 EVENT SIMULATION AND RECONSTRUCTION

4.1 Simulation of Events

Vital to the analysis of the data gathered using the CMS detector are accompanying predictions to be compared against. Good predictions can be used not only for direct comparison against data as in the case of a cross section measurement, but can also be used in the design of future detectors and experiments, or for optimizations of parameters in blinded analyses. Predictions are made by simulating pp collisions and the subsequent decays and interactions that take place inside the CMS detector volume using Monte Carlo (MC) techniques. The first step in producing these simulations is the generation of the collision event itself, and the second is in simulating the interaction of the collision products with the detector.

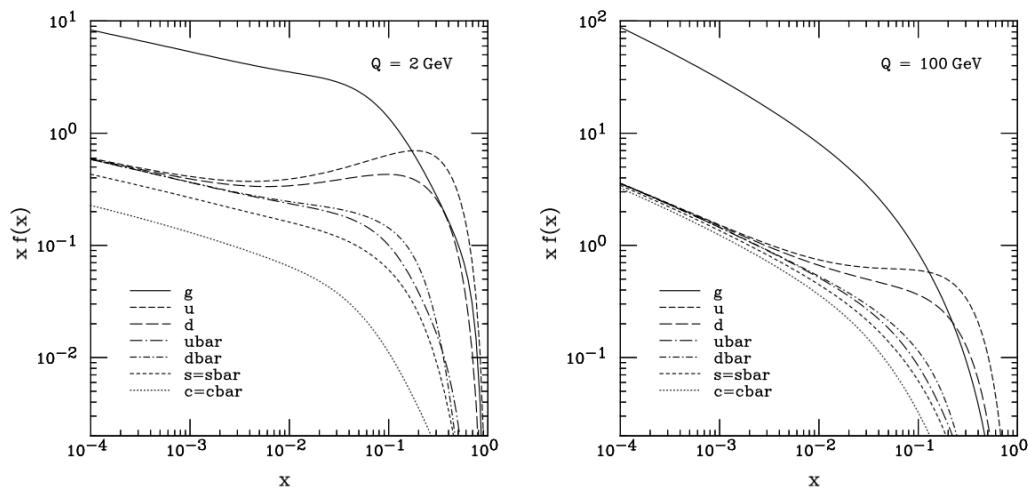
4.1.1 Monte Carlo Event Generation

There are two complimentary methods used in producing a simulation of the collision event, the direct calculation of the scattering amplitude (also known as the matrix element, ME), and the showering of particles as they decay, hadronize and radiate. Particles produced in the in the collision before showering and hadronization are referred to as prompt.

As discussed in Section 1.1.4, protons are composite objects which form as the result of strong interactions between bound quarks. Protons are therefore modeled

using parton distribution functions (PDFs) which describe the probability of finding a given constituent particle, or parton, to have a given fraction of the momentum of the proton. Parton densities vary as a function of the energy used to probe the proton and the quantity of energy absorbed by the proton is known as the momentum transfer, Q . The PDF is thus actually a set of density functions, one for each parton taken into account and for each momentum transfer, and examples of the CTEQ6M parton distribution are shown in Figure 4.1. Weighted with probabilities from the PDF set and summed over, MEs are calculated explicitly from Feynman diagrams which have initial state particles found in the PDF and the desired final state.

Figure 4.1: Below are proton PDFs shown at two different values of momentum transfer, Q . The horizontal axis shows the momentum fraction carried by the parton and the vertical axis shows the parton density.



Radiation is also important to correctly model. At any point in the collision any colored particle can radiate a gluon, and any charged particle can radiate a photon. Gluon radiation from initial state partons is always present in pp collisions at the LHC and results in jets, columnated showers of particles which are the products of quark hadronization and gluon splitting. The effects of radiation and parton showering are simulated using a Markov process in which vertices iteratively are added to partons with probabilities based on the coupling strengths, energies of the participants and the generation of random numbers. For the constituents of the proton which did not participate in the hard interaction, quarks must be created from the vacuum to enforce confinement. This produces low energy, soft, radiation and is known as the underlying event. The underlying event must be simulated along with the hadronization effects for all colored particles.

4.1.2 Monte Carlo Generators

Each of the MC generators used in this thesis are designed to serve one of two purposes. They either calculate MEs for hard scattering processes or they compute corrections such as hadronization or NLO effects.

4.1.2.1 Matrix Element Calculators

For a given $2 \rightarrow n'$ scattering process, the differential cross section is a function of the Lorentz-invariant phase space as in Equation 1.3. To calculate the cross section within a finite phase space, $d\sigma$ is integrated and MADGRAPH numerically does this as LO through the sampling of random numbers. The phase space can be interpreted

as a multidimensional hypercube spanning all degrees of freedom for all final state particles and Equation 1.3 is integrated over using the VEGAS package to calculate a weight, dw , for each point sampled in the phase space. The average of the weights converges towards $\int dw$.

To produce events with the frequency predicted by the theory being modeled, MADEVENT uses the Von Neumann method to unweight events. For each event, a random number, g , is generated between 0 and 1 and compared to the ratio dw/dw_{\max} where dw_{\max} is the largest event weight sampled. If $dw/dw_{\max} > g$ then the event is kept and is otherwise rejected. Accepted events generated with MADGRAPH/MADEVENT in this way have the same frequency and follow the same kinematic distributions as predicted from the input Equation 1.3. MADGRAPH computes events at leading order LO,

A Monte Carlo for FeMtobarn (MCFM) also uses VEGAS to perform phase space integrations, but does so at NLO instead of LO. With a much larger phase space to integrate over, MCFM performs optimizations on the integration and produces weighted events. With weighted events, an estimation of a cross section to NLO accuracy can be made, but is made only at the level of partons interacting, and in particular does not include corrections due to hadronization effects. Estimations made by MCFM thus need to be corrected for hadronization effects before being compared with observed cross sections.

4.1.2.2

Pythia performs hadronization using the Lund string model in which quarks are confined to the ends of strings and gluons are represented as kinks on that string. As quarks separate, the string breaks and creates a $q\bar{q}$ pair, thus building confinement directly into the model. The underlying event is modeled in Pythia as a set of $2 \rightarrow 2$ processes which are correlated with each other via the color connections present in the proton.

and the set of parameters used by Pythia to perform calculations is referred to as the tune.

4.1.3 Detector Simulation

After events have been produced, they are passed to Geant4 for simulation of the passage of particles through the physical mass of CMS. The Geant4 toolkit includes a full model of CMS, including all of the subdetectors as well as the inert material from the support structure and readout electronics. The magnetic field is emulated using data from measurements on the real field and Geant4 uses all of this information to register hits in the simulated detector as a consequence of the interaction between particles produced in the simulated event and the simulated material of CMS. Additionally, hits are added to the simulated detector taking into account the rates of background noise, and the final output is emulated data which is stored in the same way as would be data as taken from the real detector.

4.2 Reconstruction of Events

Real data collected from the detector and simulated data output from Geant4 consist of time-correlated energy deposits in the various subdetectors of CMS. As a result of the coordinated designs of the subdetectors, the final-state particles which arise from pp collisions at the LHC can be individually identified and reconstructed using the combined information from the entirety of CMS. The associated global event description from this particle-flow (PF) reconstruction provides excellent performance for the identification of electrons and muons, as well as for vertex identification and the evaluation of E_T^{miss} . As a result of the PF reconstruction of an event, particles are identified and placed into the mutually exclusive categories: charged and neutral hadrons, photons, electrons and muons.

4.2.1 Track and Primary Vertex Reconstruction

The subdetector closest to the interaction vertex is the tracker, which records precise information about the trajectories of charged particles as they pass through it. Combined with the magnetic field, this allows for the measurement of the momenta of these particles as well as a means of identifying the the location of the primary interaction.

Tracks are identified via an iterative process. The first tracks to be reconstructed are those which pass strict seeding criteria, designed to have a moderate efficiency, but negligibly small fake rate. Then the detector hits associated with these tracks are masked and the remaining hits are used to form track seeds with slightly relaxed

criteria. This operation is repeated, with every iteration imposing more complex and time-consuming seeding, filtering and track fitting algorithms.

Because bunches of protons instead of single protons are made to cross in the LHC, multiple collisions can take place during the same bunch crossing. The vertex with the highest scalar sum transverse momentum, p_T , of tracks and passing further quality selections based on the goodness of fit for the tracks and the number of tracks associated with a given vertex is chosen as the primary vertex (PV). The other vertices are referred to as pileup vertices and the associated collision products as pileup.

In $pp \rightarrow W b\bar{b} \rightarrow \ell\nu b\bar{b}$ events, the two b quarks and the lepton from the W decay all leave energy deposits in the tracker, thus making the choice of PV unambiguous. However, in the $pp \rightarrow \gamma + \text{invisible}$ events, the only visible final state object is a photon, and photons do not leave hits in the tracker. This makes the identification of the PV in the monophoton analysis difficult and motivates the using of variables that are less sensitive to correct PV identification.

4.2.2 Electron ID and Reconstruction

Electrons are reconstructed using tracker hits and ECAL deposits. The seed of an electron candidate is selected as an energy deposit in the ECAL with $E_T > 4$ GeV having nearby deposits in the tracker. As electrons move in magnetic fields, they emit bremsstrahlung radiation tangential to their flight path and this radiation both appears in the detector, and alters the course of the electron.

The effects of this radiation are taken into account via the Gaussian Sum Filter

(GSF) track fitting algorithm. This algorithm uses weighted sums of Gaussian functions to describe electron energy loss and thus allows for non-Gaussian corrections to the fitting of tracks. In the CMS detector, bremsstrahlung from electrons results in the emission of photons in an extended strip in the ϕ direction and electron superclusters (SCs) are made by including the energy deposits from these photons in the ECAL as part of the candidate electron object.

Further requirements during the reconstruction of the electron improve the purity of selection. The SC and the GSF track are required to be separated by no more than $|\eta| < 0.02$ and $|\phi| < 0.15$ and the fraction of energy deposited in the HCAL directly behind the SC, and the SC is required to be no more than 15%.

4.2.3 Muon ID and Reconstruction

Muon identification is performed using two reconstruction and filtering methods to produce ‘tracker muons’ and ‘standalone muons’ which are combined to form ‘global muons.’ Tracker muons are identified starting with a track, $p_T > 0.5$ GeV and $p > 2.5$ GeV, which is then extrapolated to the muon system. If the distance between the extrapolated track and the nearest hit in one of the muon chambers is less than 3 cm, a tracker muon is identified. Tracker muons are also identified if the pull between the extrapolated track and the matched station hit is less than four, where pull is defined as the distance between the track and the station hit divided by the uncertainties on both measured quantities. Tracker muons are built from the inside of the detector towards the outside, and standalone muons are built in the other direction. Only hits in the muon stations are used to reconstruct standalone muons,

with the additional constraint that the path reconstructed from the hits points back toward the interaction region. Thus, the tracker muon algorithm is well-suited for the identification of low- p_T muons by having low thresholds and requiring only one track and one station hit, while the standalone muon algorithm is aimed at high- p_T muons which have the energy to penetrate multiple layers of muon stations to form tracks which can be traced back to the interaction. Global muons are required to pass the criteria for both standalone muons and tracker muons, and, starting with the standalone muons, the global muon trajectory is refit using information from both the muon stations and the tracker, yielding an improved energy resolution than either one.

4.2.4 Lepton Isolation

After having been identified and reconstructed, leptons are also assessed in terms of their isolation. Isolation is an important variable for differentiating prompt leptons from leptons resulting from secondary decays or hadronization effects and is defined as

$$I = \frac{\sum p_T^{\text{charged}} + \max(0, \sum p_T^\gamma + \sum E_T^{\text{neutral}} - 0.5 \cdot p_T^{\text{PU}})}{p_T^\ell}, \quad (4.1)$$

with the sum running over the PF candidates (hadrons, electrons, photons) in a cone of size $\Delta R < 0.4$ (0.3) around the muon (electron) direction. The last term in the numerator, p_T^{PU} , is a correction for pileup effects and is based on the scalar sum of transverse momenta of charged particles not associated with the primary vertex that are within the isolation cone.

4.2.5 Photon ID and Reconstruction

Photons are reconstructed using the same ECAL clustering algorithms as are used for electrons. This allows for the simultaneous reconstruction of photons that have and have not split to $e\bar{e}$ pairs. The size of the SC is determined dynamically and the center is determined to be the barycenter of the distribution, with weights assigned using the logarithm of the fractional energy deposits of the ECAL crystals clustered in the SC. The angular width of the cluster is σ_{inj} and photons tend to have narrow deposits.

In an ideal tracker, photons would not interact at all and objects that leave signatures similar to those of photons could be rejected through the rejection of tracks. However, some photons do convert to $e\bar{e}$ pairs inside the tracker volume which leave tracks, so the rejection of tracks is not a perfect way to distinguish between photons and electrons.

For a photon object to be considered isolated, scalar sums of the transverse momenta of PF charged hadrons, neutral hadrons, and photons within a cone of $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} < 0.3$ around the candidate photon must individually fall below bounds defined for 80% signal efficiency. Only the PF candidates that do not overlap with the electromagnetic shower of the candidate photon are used in the isolation sums.

4.2.6 Missing Transverse Energy

The transverse mass of the W boson is defined in Equation 4.1, and is a natural discriminator against non- W final states such as QCD multijet events, that have a

lepton candidate and E_T^{miss} , but a relatively low value of m_T . In calculating m_T , the E_T^{miss} is corrected for noise in the ECAL and HCAL ?, and corrections to limit the effects of pileup are also included ?.

4.2.7 Jet ID and Secondary Vertices

?? The reconstruction of jets is accomplished using the anti- k_t clustering algorithm on particles identified in the PF. The anti- k_t algorithm is both infrared and collinear safe, meaning that it is stable against soft (low energy) radiation getting clustered into individual jets, and also stable against hard (high energy) jets splitting collinearly and affecting the shape of the jet.

Jets are corrected in simulation and in data to remove energy believed to come from elsewhere than the PV, thus removing the luminosity dependence of the jet. Jets are also corrected to have a response that is independent of η by studying dijet events and calibrating the jets to anti-align. To make the jet response independent of the p_T of the jet, an absolute correction is applied, and in data, one further correction on the relative energy scale is applied. After all of these corrections are applied, simulated jets are observed to have sharper energy resolution than is observed, so jets in MC are smeared in energy.

Bottom quarks have a relatively long lifetime and are the heaviest fundamental particle that has been seen to decay inside the volume of the CMS detector. A b quark produced in a pp collision at CMS therefore has enough time to hadronize into a jet before decaying, and such jets are called b -jets. The identification, or tagging, of b -jets is focused around the vertex associated with the b -hadron which, since it is not

the PV but is still a vertex associated with the event, is called a secondary vertex, SV. The tagging of b -jets is accomplished using a multivariate analysis technique in which information about the number and energy of tracks associated with the SV and their corresponding alignment with the SV, as well as the presence and energies of soft leptons is combined into a single discriminator value.

5 $Wb\bar{b}$ CROSS SECTION MEASUREMENT

The first of the two processes examined in this thesis is the SM production of $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$ at $\sqrt{s} = 8$ TeV, described in Section 2.1. Throughout this section, ℓ or lepton denotes an electron or a muon.

5.1 Event Selection

Two decay channels of the W boson are considered, $W \rightarrow \mu\nu_\mu$ and $W \rightarrow e\nu_e$, and events are selected using single-muon (single-electron) triggers with a loosely isolated muon (electron) with transverse momentum $p_T > 24$ (27) GeV and pseudo-rapidity $|\eta| < 2.1$ (2.5). Individual particles emerging from each collision are reconstructed with the particle-flow (PF) technique described in Section 4.2.

Both the muon and electron candidates are required to have $p_T > 30$ GeV, $|\eta| < 2.1$ and to originate from the PV of the event, as defined in Section 4.2.1. These leptons must also pass tight ID and isolation requirements, with $I < 0.12$ (0.10) for selected muons (electrons) as defined in Equation 4.1. Additionally, to reject some of the background events contributed by processes such as $pp \rightarrow t\bar{t}$ and $pp \rightarrow Zj$, events with more than one lepton $p_T > 10$ GeV and $|\eta| < 2.4$ are rejected. Jets are constructed from PF candidates using the anti- k_t clustering algorithm ? with a distance parameter of 0.5, and jets are b -tagged using the multivariate discriminator at an operating point with a efficiency of 40% and a misidentification probability of 0.1% for light jets and 1% for charm jets described in Section ???. Exactly two b -tagged jets with

Table 5.1: Listed below are the raw number of events in data passing the selection listed in the first column as well as all selections in higher rows. The second column is for the $W \rightarrow \mu\nu$ decay channel, and the third column is for $W \rightarrow e\nu$.

Selection	Events passing selection	
	$W \rightarrow \mu\nu$	$W \rightarrow e\nu$
$p_T^{\ell_1} > 30$ GeV and $ \eta^{\ell_1} < 2.1$	X	X
$p_T^{\ell_2} \not> 10$ GeV or $ \eta^{\ell_2} \not< 2.1$	X	X
$p_T^{j_1,j_2} > 25$ GeV and $ \eta^{j_1,j_2} < 2.4$	X	X
$p_T^{j_3} \not> 25$ GeV or $ \eta^{j_3} \not< 4.7$	X	X
b -tag j_1, j_2	X	X

$p_T > 25$ GeV and $|\eta| < 2.4$ are required to be present in selected events, to remove contamination in the signal region from charm and light flavor jets. To reduce the contribution from $pp \rightarrow t\bar{t}$ events, events with a third jet with $p_T > 25$ GeV and $|\eta| < 4.7$ are rejected. The effects of the various cuts with regard to the total number of events passing selections on the data is illustrated in Table 5.1.

5.2 Simulated samples

After all selection requirements detailed in Section 5.4 are applied, the contributing processes to the overall yield are the associated production of a massive vector boson and jets ($V + \text{jets}$), as well as diboson (WW , WZ , ZZ), $t\bar{t}$, single top quark, $\gamma + \text{jets}$, and QCD multijet production. The corresponding contributions are estimated from simulation, except for the QCD background, which is estimated from data as described in Section 5.3.1.

Simulated samples of $V + \text{jets}$, $\gamma + \text{jets}$ and $t\bar{t} + \text{jets}$ are generated at tree-level with MADGRAPH 5.1 ?? using the CTEQ6L ? Parton Distribution Function (PDF) set. These samples are interfaced with PYTHIA 6.4 ? using the Z2* tune for hadronization. The most recent PYTHIA Z2* tune is derived from the Z1 tune ?, which uses the CTEQ5L parton distribution set, whereas Z2* adopts CTEQ6L ?. The k_t -MLM ?? matching scheme is used. For the signal distributions, the $Wb\bar{b}$ component of an inclusive $W + \text{jets}$ sample is used, with the shapes of the distributions taken from a dedicated high-statistics generated sample of exclusive $Wb\bar{b}$. The shape of the $Wb\bar{b}$ signal distribution is obtained by separating the $W + \text{jets}$ simulated sample into three subsamples labeled as $Wb\bar{b}$, $Wc\bar{c}$, and $Wusdcg$. These three categories are described below. The separation is done at the particle list level. If an event contains a b jet, from the matrix element or parton shower, it falls into the $Wb\bar{b}$ category. A b quark at generator level requires the presence of a b hadron within a cone of radius $R = 0.4$ with respect to the jet axis. The jets are constructed using generator-level information using all stable particles in the event excluding neutrinos. Jets with a distance smaller than $R = 0.5$ with respect to a lepton are removed from the event. If an event contains no b jets, but an even, non-zero, number of charm jets, again from the matrix element or parton shower, it falls into the $Wc\bar{c}$ category. The remaining events fall into the $Wusdcg$ category. The energy of the selected leptons at the generator level is corrected for final-state radiation by summing up the four-momenta of all the photons generated within a cone of radius $R = 0.1$ around the lepton. Generated leptons which do not originate in simulation from the primary interaction vertex are not considered for selection.

Single top quark event samples are generated at next-to-leading order (NLO) with POWHEG 2.0 [????](#) using the CTEQ6M PDF set. Hadronization is performed using PYTHIA 6.4 with the Z2* tune. Diboson samples are generated and hadronized with PYTHIA 6.4 at leading order (LO) using the CTEQ6L PDF set and the Z2* tune.

The cross sections for the $V + \text{jets}$ and $\gamma + \text{jets}$ processes are normalized using the predictions for inclusive W, Z and photon production from FEWZ 3.1 [??](#) evaluated at next-to-next-to-leading order (NNLO). Single top quark and diboson production cross sections are normalized to the NLO cross section predictions from *MCFM* 7.0 [??](#) using the MSTW08 NLO PDF set [?](#). The $t\bar{t}$ cross section used is $241.5 \pm 8.5 \text{ pb}$ and was determined from data collected by the ATLAS and CMS experiments [???](#) at the LHC at $\sqrt{s} = 8 \text{ TeV}$.

Events induced by additional simultaneous pp interactions are simulated using events generated with PYTHIA 6. During the 2012 data taking, the average pileup rate was 21 interactions per bunch crossing and the simulated number of pileup interactions have been reweighted to match this distribution in the data.

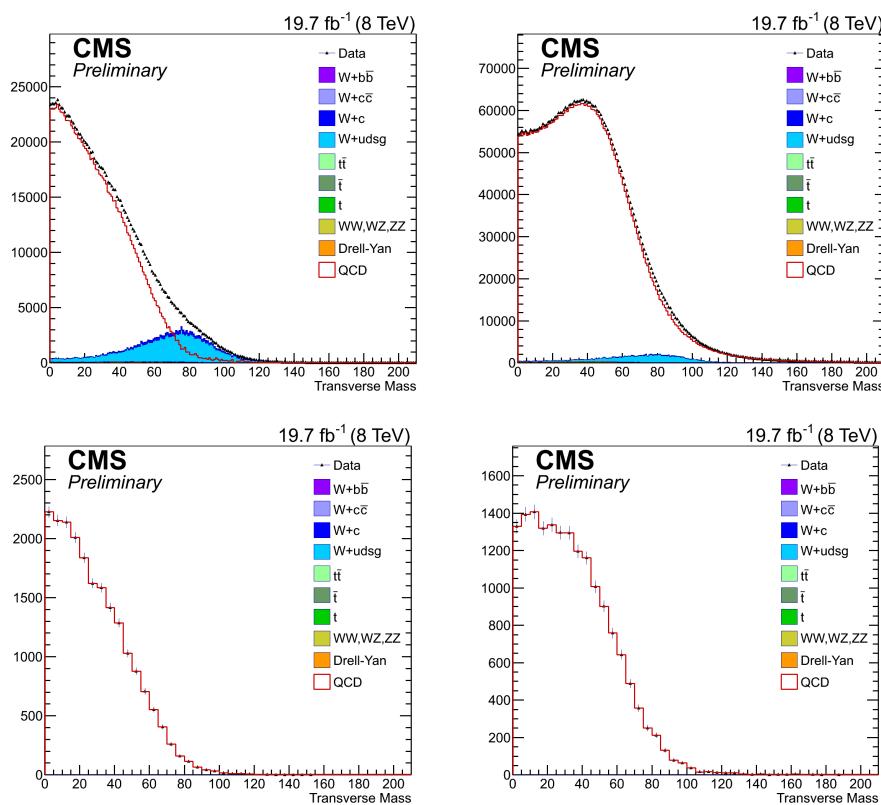
5.3 Background Estimation

5.3.1 QCD

The QCD multijet sample is derived using a data-driven method. The shapes of the distributions for QCD multijet events are taken as the difference between the data sample and the sum of the other simulated backgrounds in a region of phase space

enriched in multijets as illustrated in Figure 5.1. This region is found using the same selection requirements as those in the signal region, but requiring the muon (electron) to be antiisolated: $I > 0.20$ (0.15). In the fiducial regions used in this analysis, no correlation is observed between I and m_T , validating the use of an inverted isolation requirement to obtain the QCD background shape. This is not the case for the ΔR distance between the two b -tagged jets, $\Delta R(b, \bar{b})$, or the lepton p_T . The shape of the QCD distribution for these variables is therefore taken from an $m_T < 30$ GeV sideband and validated against QCD multijet simulation. The normalization of the QCD background in these variables is set to the final normalization resulting from the fit to the m_T variable which was derived using the inverted isolation requirement. To obtain an initial estimate of the number of QCD multijet events passing signal region selections, the shape derived in the antiisolated region, is put in the signal region and scaled by $(d_{20} - m_{20})/q_{20}$ where d_{20} is the yield in data in the range $0 < m_T < 20$, m_{20} is the combined yield from the simulated samples in this range, and q_{20} is the corresponding unnormalized yield of QCD multijet events. This has the effect of normalizing the QCD sample such that the combination of the QCD and the simulated backgrounds has the same total yield as data in the range $0 < m_T < 20$. If $d_{20} < m_{20}$, the QCD contribution is taken to be negligible. The relative uncertainty in the yield of QCD multijet events is estimated to be $\pm 50\%$, taking into account both the fit result and the extrapolation from $0 < m_T < 20$ to the high- m_T range. This relative uncertainty also covers shape mismodelings of the multijet contribution in the final sample.

Figure 5.1: The shape for the QCD is found by inverting the lepton isolation and subtracting MC from the data. Shown above is the data, MC background and extrapolated QCD shape (difference between data and MC backbrouunds) for both the muon and electron channels in the $W + jj$ and $W + b\bar{b}$ phase spaces. The requirement of two well-identified b tags essentially eliminates all MC backgrounds in the $W + b\bar{b}$ region, leaving the QCD shape the same as that of the data.



5.3.2 $W + \text{jets}$: light and charm component

$W + \text{light jets}$ is the dominant background in the Wjj phase space, which is found using identical selections as are used in the signal region with the exception of the b tag requirement. In the $W + jj$ phase space, no requirements on b tags are made. This control region therefore serves as a cross check on the reconstructed objects observed in the signal region before the added complication of b -tagging has been introduced. In Figure 5.2 is shown the p_T of the identified lepton along with the E_T^{miss} and m_T in both decay channels. Agreement between simulation and data is on the order of 10%.

The process $pp \rightarrow Wb$ with a single b quark produced is CKM suppressed by two generations if light quarks are interacting. The process $pp \rightarrow Wc$ is only suppressed by one generation but is still found to contribute negligible rate in the signal region due to the requirement of a second b tag along with a veto on a third jet. A second jet can be added to $ppwc$ via ISR or FSR, and light quarks require less energy to produce than heavy quarks, making them more probable. Therefore, the dominant contribution to $ppwc$ in the signal region comes from the mistag of a c jet for a b jet, along with the misidentification of an ISR/FSR light jet. However, the contribution to events in the signal region from high energy $g \rightarrow c\bar{c}$ is not negligible and moreover these events have kinematics closely related to that of the signal.

5.3.3 Single top backgrounds

The single top t -channel control region is defined by the signal selection requirements without the third and forward jet vetos, and with the leading jet required to be

Figure 5.2: Selecting for a tight ID muon with $p_T > 30$ GeV and exactly two central jets passing loose ID, we recover the distributions shown above. Shown in the upper left (right) is the momentum of the leading lepton in the muon (electron) channel. The missing transverse energy is shown in the center, and transverse mass is given in the bottom two distributions. The shaded band in ratio plots shows statistical uncertainty.

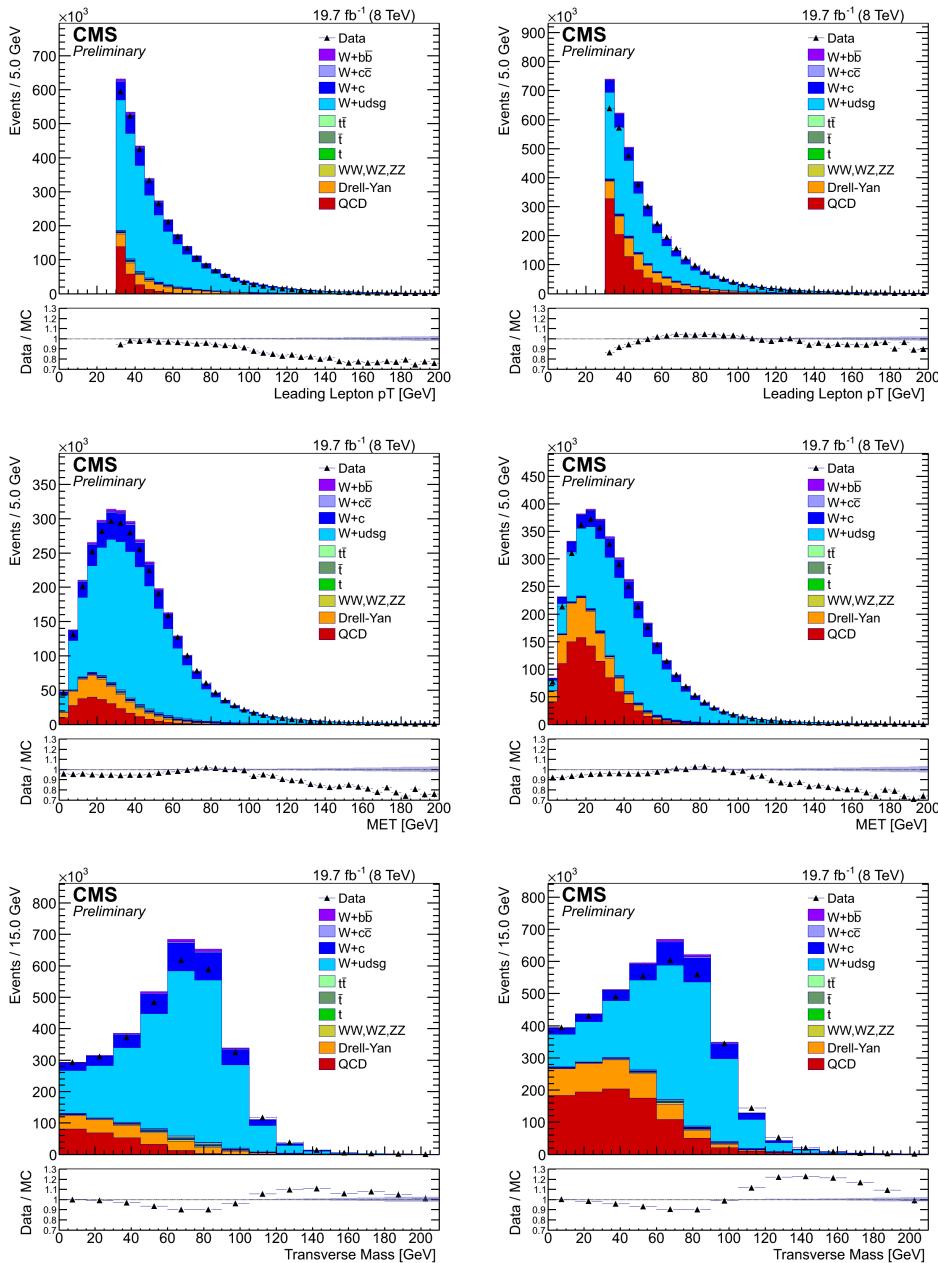
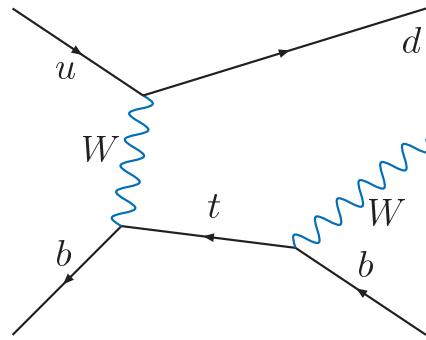
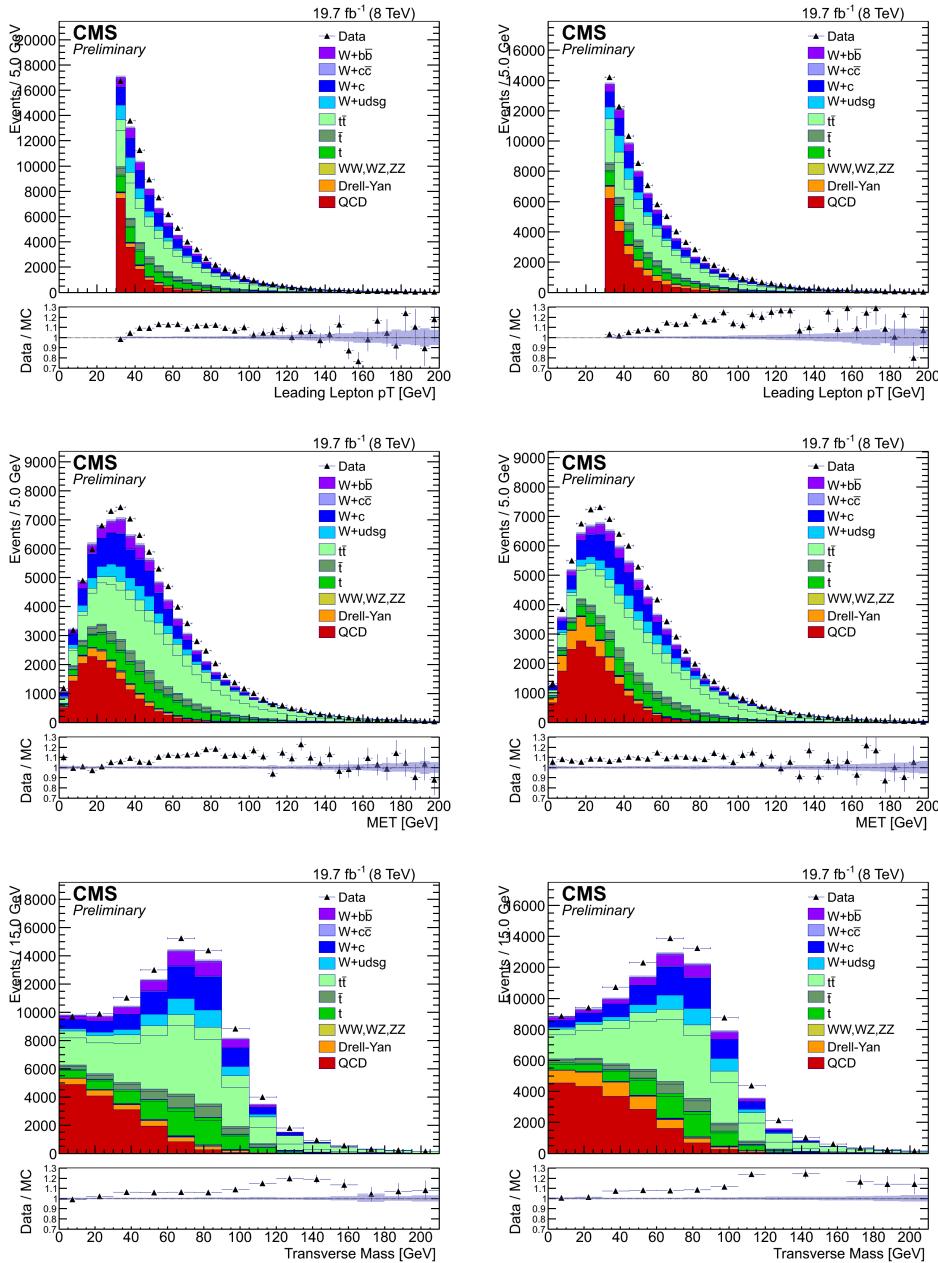


Figure 5.3: Below is the diagram for the process attempting to be isolated in the single top t -channel control region.



central ($|\eta| < 2.4$) and tightly b -tagged, while the subleading jet has no b requirement and must fall within $2.4 < |\eta| < 5.0$. The reasons for these selections are based on examination of Figure 5.3. The b quark is a product of the t decay and is recoiling against a W , making it a high- p_T jet. The other jet comes from an initial state quark that exchanges a W and recoils against the t . Because it is the product of an initial state quark which is most likely to be light, the final state jet is also most likely to be light. The fact that it is recoiling against the much more massive t means that it is likely be forward by conservation of momentum. As can be seen in Figure 5.4, there are many backgrounds contaminating the purity of this phase space, but agreement between data and simulation is on the order of 5-10%.

Figure 5.4: The single top control region is defined by one b -tagged central jet and one forward jet. Shown above are distributions in the single top control region. Left plots are in the muon decay channel and right plots are in the electron decay channel.



5.3.4 $Z\ell\bar{\ell}$ backgrounds

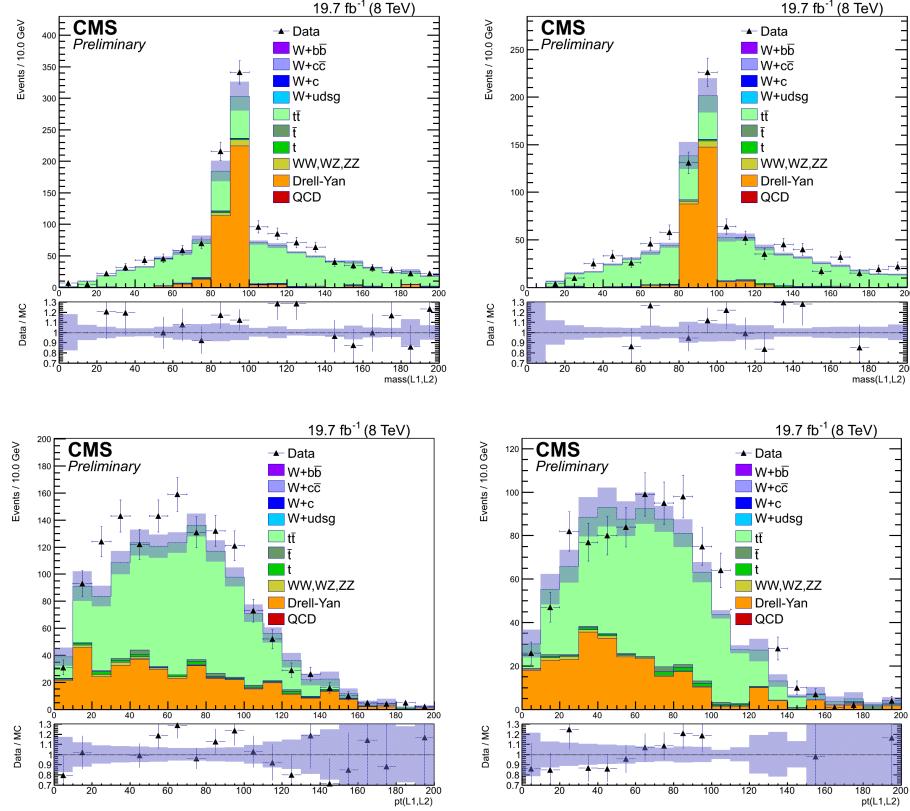
The Drell-Yan background is validated in a control region where the $Wb\bar{b}$ selection requirements are applied, but the lepton veto is inverted, requiring two isolated, same-flavor leptons which are presumed to be the decay products of the Z boson. This is referred to as the $Zb\bar{b}$ region and distributions of the mass and transverse momentum of the dilepton pair is shown in Figure 5.5. Contamination from $t\bar{t}$ is evident. A cleaner Drell-Yan phase space is found by requiring exactly two jets but placing no b tag requirement and is referred to as Zjj . Figure 5.6 shows the same distributions as Figure 5.5 in this phase space.

5.4 Analysis Strategy

The $Wb\bar{b}$ yield is measured using a likelihood fit to the m_T distribution in the signal region. With the exception of QCD processes, the initial shapes and normalizations of all contributions in the fit are taken from simulation. Consequently, it is important to verify that the simulation describes the data. The dominant background in the signal region arises from the $t\bar{t}$ process. The data and simulation are thus compared in two $t\bar{t}$ -dominated control regions: one characterized by a pair of opposite flavor leptons, and the other by the presence of three or more jets. The simulation is reweighted to describe the data in the control regions and then is used to predict the m_T distributions in the signal region.

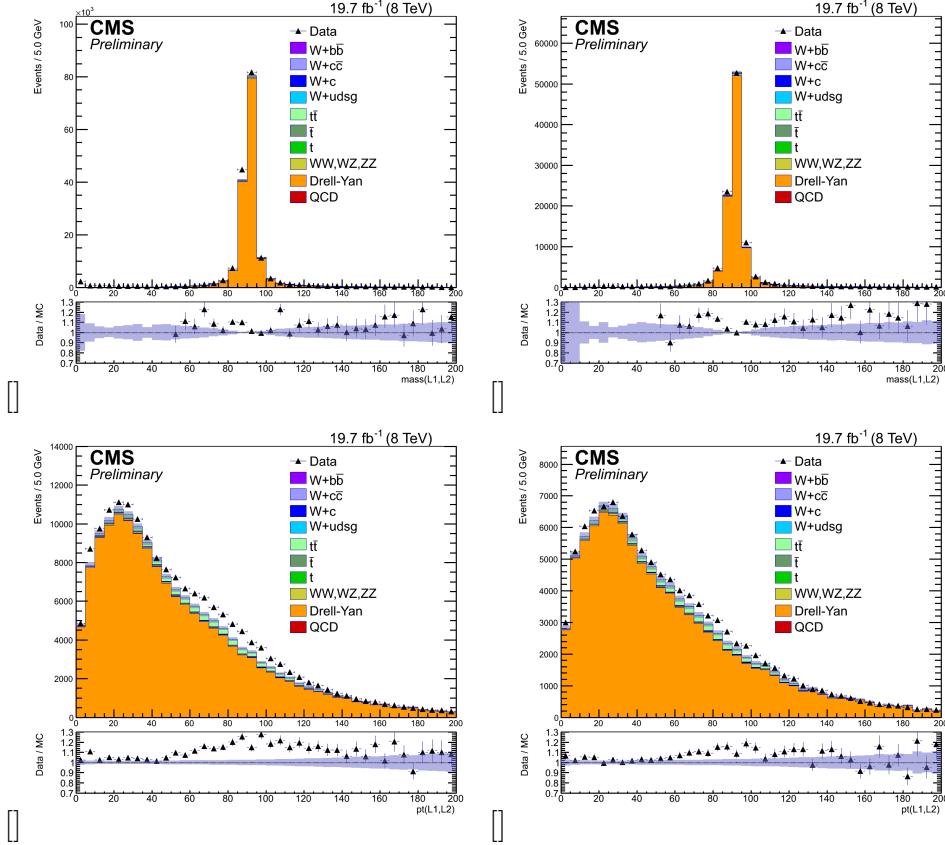
The signal region requires a muon (electron) with $p_T > 30$ GeV, $|\eta| < 2.1$, and satisfying $I < 0.12$ (0.10). Exactly two b -tagged jets with $p_T > 25$ GeV and $|\eta| < 2.4$

Figure 5.5: Above are distributions of the mass and transverse momentum of the dilepton pair in the $Zbb\bar{b}$ phase space. Left plots are in the muon decay channel and right plots are in the electron decay channel.



are also required. Events with additional leptons with $p_T > 10$ GeV and $|\eta| < 2.4$ or a third jet with $p_T > 25$ GeV and $|\eta| < 4.7$ are rejected. The $t\bar{t}$ -multijet control region is obtained using the same selection criteria as for the signal region, but requiring at least three jets in the event with $p_T > 25$ GeV and $|\eta| < 2.4$ instead of vetoing events which have more than two jets. These are to capture events in which one of the W bosons from the t decay decays leptonically and the other decays hadronically,

Figure 5.6: Above are distributions of the mass and transverse momentum of the dilepton pair in the Zjj phase space. Left plots are in the muon decay channel and right plots are in the electron decay channel.



producing jets. The $t\bar{t}$ -multilepton control region uses similar selection criteria as the signal region, but changing the lepton requirement from vetoing events which contain a second lepton, to requiring two isolated leptons of different flavor, both with $p_T > 30$ GeV and $|\eta| < 2.1$. This selection is to capture $t\bar{t}$ events where both W bosons decay leptonically, while avoiding same-flavor dileptons which are characteristic of Z boson decay. In the $t\bar{t}$ -multilepton region, the m_T variable is calculated with respect to the

electron in the electron channel and the muon in the muon channel.

The normalizations and shapes of the simulated backgrounds are allowed to vary in the fit within the uncertainties listed in Table 5.2. The uncorrelated normalization uncertainties are uncertainties on the cross section of the given sample.

Two major parameters in the simulations significantly affect the normalization of the simulated distributions: the b -tagging efficiency and the jet energy scale (JES). Both control and the signal regions show similar sensitivity to the b -tagging efficiency, and its adjustment affects all the regions in a correlated manner. Because the $t\bar{t}$ production has more than two jets in the final state, the rejection of events with a third jet makes it sensitive to JES. The effect on the leading jets is moderate but JES variations lead to significant migration of jets into and out of the veto region. The $t\bar{t}$ -multijet control region, since it has no veto on a third jet, is less sensitive to JES variations than the $t\bar{t}$ -multilepton control region. The same is true for the $Wb\bar{b}$ process, which at LO has only two jets in the final state from the ISR gluon splitting and therefore is not affected by the jet veto requirement.

The fit procedure consists of three consecutive steps where the simulated distributions in two control regions and the signal region are fit to data using the m_T variable, which is chosen because it has a well known shape for $W + \text{jets}$ production that allows for reliable signal extraction. First, the fit is performed in the $t\bar{t}$ -multijet region. In this control region, the uncertainty is dominated by the uncertainty in the b -tagging efficiency so the fit results in a correction of the b -tagging efficiency that is measured separately in the muon and electron channels and then combined. The simulation is corrected by 14% using this result and the corrected simulated samples

are fit to the data in the $t\bar{t}$ -multilepton region. The result of the second step is used to adjust JES by 3.4% and as a result of this procedure, the simulation is expected to better describe the $t\bar{t}$ contribution. The final step is to extract the number of $Wb\bar{b}$ events from the fit to m_T in the signal region.

5.5 Systematic uncertainties

The main sources of the systematic uncertainties are listed in Table 5.2. The size of the variation is shown for each uncertainty source, together with its effect on the measured cross section and is included as a nuisance parameter in the fit. Some of the uncertainties affect only the normalization in the respective contributions. These include the uncertainties on the theoretical cross section of a given sample, which are uncorrelated between samples and are included as log-normal constraints on the rate. For any sample, the effect on the final normalization from uncertainties which only effect normalization is

$$N_f = N_i \prod a_j^{x_j} \quad (5.1)$$

where a is the input bound given in the “Variation” column of Table 5.2 and x is the parameter being fit for the j nuisances. The uncertainty due to the b-tagging efficiency and the uncertainty due to the JES are allowed in principle to have shape dependencies in this analysis, but only affect the normalizations of the samples in the m_T variable in practice.

The uncertainties which affect both the normalization and the shape of the m_T distributions are those listed in the table under “Shape” and are incorporated into the

fit via binned templates. These templates are obtained by varying the source of the given uncertainty and reprocessing the simulated sample. Taking h_0 as the unshifted template and h_j^\pm as the up and down templates obtained through reprocessing after shifting nuisance j , the final fitted shape h is given by

$$h = h_0 + \sum_j [A(x)h_j^- + B(x)h_0 + C(x)h_j^+] \quad (5.2)$$

$$A(x) = x(x+1)/2, B(x) = -x^2, C(x) = x(x-1)/2 \quad (5.3)$$

The 50% uncertainty in the QCD background is taken as a conservative estimate and ultimately has a 2-3% contribution to the uncertainty on the measured cross section. The b-tagging efficiency and JES rescaling uncertainties are taken from their respective fits. The renormalization and factorization scale uncertainties are estimated by simultaneously changing the renormalization and factorization scales, μ_R and μ_F , up and down by a factor of two. The PDF uncertainties are estimated from the change in acceptance found by varying the PDF set following the LHAPDF/PDF4LHC prescription [????](#) considering PDF sets from CTEQ, MSTW, NNPDF, and HERA Collaborations.

5.6 Signal Extraction

The fit in the $t\bar{t}$ -multijet region is used to obtain b-tag rescaling factors separately for the muon and electron channels in order to better describe the b-tagging efficiency in the simulation as described in Section [??](#). The results of the fit are presented the two

plots at the top of Fig. 5.7. The central values of the b-tag rescaling factors, 1.12 ± 0.08 (muon channel) and 1.16 ± 0.08 (electron channel), are combined to 1.14 ± 0.08 . The simulation is reweighted accordingly for the next fit and the uncertainty on this fit sets the one standard deviation bound on the b-tag efficiency rescale factor in subsequent fits.

A fit to the $t\bar{t}$ -multilepton region adjusts the JES, as described in Section ???. As a result, the simulated m_T distributions change normalization. The best fit suggests changing the normalization by approximately 3.4% from its central value which corresponds to 1.3 standard deviations in JES. The middle plots in Fig. 5.7 show the results of the fits in the $t\bar{t}$ -multilepton event enhanced control region for the muon (left) and the electron (right) channels. The JES is therefore shifted by 1.3 standard deviations in the simulation with the uncertainty taken from the fit. Thus the simulation is tuned to describe the $t\bar{t}$ control regions and is used to extract the signal yield in the signal region.

The results of the fit in the $Wb\bar{b}$ signal region are shown in the bottom of Fig. 5.7. All background contributions are allowed to vary in the fit within their uncertainties, while the $Wb\bar{b}$ normalization remains a free parameter of the fit. The correlations across all simulated samples are taken into account as shown in Table 5.2. The composition of the event sample in the signal region is summarized in Table 5.3. Events coming from the production of a Higgs boson in association with a vector boson constitute a negligible fraction of the overall event yield and are not considered.

Distributions for variables other than those being directly fit are also produced by applying the results from the three fits to the simulated samples. Distributions of

$\Delta R(b\bar{b})$ and p_T^ℓ combining both lepton flavors are presented in Fig. 5.8.

5.7 Cross Section and Comparisons

The cross section for the process $\sigma(pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b})$, is derived from the signal strength measurement as obtained from the fit. The cross section is written as

$$\sigma(pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}) = \frac{N_{\text{signal}}^{\text{Data}}}{A \cdot \epsilon \cdot \mathcal{L}} = \frac{N_{\text{signal}}^{\text{Data}}}{(N_{\text{signal}}^{\text{MC}}/N_{\text{generated}}^{\text{MC}}) \cdot \mathcal{L}} = \alpha \sigma_{\text{gen}}$$

where $N_{\text{signal}}^{\text{Data}}$ is the number of observed signal events, $N_{\text{signal}}^{\text{MC}}$ is the number of expected signal events from simulation, $N_{\text{generated}}^{\text{MC}}$ is the number of generated events in the fiducial region, A , ϵ are the acceptance and efficiency correction factors, α is the measured signal strength in the given lepton channel, and σ_{gen} is the simulated fiducial cross section of the signal sample.

In this analysis, the fiducial cross section was calculated in the following manner: MADGRAPH is used to compute the $Wb\bar{b}$ cross section with fiducial cuts applied. Then a k-factor for inclusive W production is applied, obtained from the ratio of the inclusive W cross sections calculated with FEWZ (at NNLO using the five-flavour CTEQ 6M PDF set) and with Madgraph. The product $A \cdot \epsilon$ is 13 (11)% in the muon (electron) channels and results from the combined effect of the efficiency from lepton identification requirements (80%), and b tag efficiency (40% per jet). The uncertainty on this product is 10% as listed in the bottom row of Table 5.2, which was calculated by varying the PDF set using the LHAPDF/PDF4LHC **?????** prescription considering PDF sets from CTEQ, MSTW, NNPDF, and HERA as well as varying the choice

of scales μ_F , μ_R simultaneously up and down by a factor of two. Varying the PDFs and the choice of scale is a way to estimate the dependence of the measured cross section on the choices of these parameters.

The $Wb\bar{b}$ cross section is measured within a fiducial volume, which is defined by requiring leptons with $p_T > 30\text{GeV}$ and $|\eta| < 2.1$ and exactly two b -tagged jets of $p_T > 25\text{GeV}$ and $|\eta| < 2.4$. The measured cross sections are presented in Table 5.4. The combination of the muon and electron measurements is done using a simultaneous fit to both channels, taking into account correlations between different sources of uncertainties.

The measured cross sections are compared to theoretical predictions from MCFM ?? with the MSTW 2008 PDF set, as well as from MADGRAPH 5 interfaced with PYTHIA6 in the four- and five-flavour schemes and MADGRAPH 5 with PYTHIA8 ? in the four-flavour scheme. In the five-flavour scheme, the PDF set CTEQ 6L was used and PYTHIA6 was run using TuneZ2*. The two four-flavour samples were produced using a NNLO PDF set interfaced with PYTHIA (version 6 in one sample, version 8 in the other) in the CUETP8M1 tune.

Comparisons between the results of calculations performed under different assumptions provide important feedback on the functioning and validity of the techniques employed. Differences in predictions arising from the modelling of b quarks as massive or massless are possible, as are variations in predictions arising from the use of different showering packages (PYTHIA6 vs. PYTHIA8) or matrix element generators (MADGRAPH vs. MCFM). In the phase space explored here, these predictions are all very close in their central value and agree with each other well within their respective

uncertainties.

The MCFM cross section calculation is performed at the level of parton jets and thus requires a hadronization correction. The multiplicative hadronization correction factor 0.81 ± 0.07 is calculated using the MADGRAPH + PYTHIA6 sample and agrees well with a similar factor calculated in the 7TeV Z+b analysis calculated as 0.84 ± 0.03 ?. This factor is determined by taking the ratio of the predicted cross sections in a given sample after having undergone a simulation of the hadronization process, and using jets identified at the particle list level. The correction factor is obtained for jets computed excluding neutrinos from the particle list, as particles such jets are closer in kinematics to particle jets at the detector level. The uncertainty reflects both the statistics of the MADGRAPH + PYTHIA6 sample as well as a comparison with the MADGRAPH + PYTHIA8 sample.

The MCFM and four-flavour MADGRAPH predictions do not account for $Wb\bar{b}$ production where the $b\bar{b}$ system comes from multiple parton scattering. CMS simulations of MADGRAPH + PYTHIA events that include double parton interactions (DPIs) reproduce the $W +$ jets data ?, therefore a MADGRAPH + PYTHIA8 sample of a W boson produced in association with a $b\bar{b}$ pair coming from DPS was generated to study the effect on the fiducial cross section. Using this dedicated sample, an additive correction σ_{DPS} is estimated to be 0.06 ± 0.06 pb, where the uncertainty is conservatively assigned to be 100% of the value.

The uncertainty in the theoretical cross sections arising from the choice of PDF is also accounted for, using the LHAPDF/PDF4LHC ???? prescription in which PDF sets from CTEQ, MSTW, NNPDF, and HERA are considered. Uncertainties in the

theoretical cross section due to the choice of scale are also estimated by varying the scales μ_F , μ_R simultaneously up and down by a factor of two.

The resulting cross section predictions in the fiducial phase space at the hadron level and including the estimated hadronization and DPS corrections when needed are compared in Fig. 5.9 with the measured value. Within one standard deviation the predictions agree with the measured cross section. The results also agree within one standard deviation with previously published $Wb\bar{b}$ measurements at 7 TeV, where data are found to be well described by the same predictions.

Table 5.2: Breakdown of the main sources of systematic uncertainty in the $Wb\bar{b}$ signal region. The column labeled “Variation” indicates the bounds on the normalization change of a given sample due to a variation of the uncertainty by one standard deviation. The last column indicates the contribution of the given systematic to the overall uncertainty in the measured cross section. The uncertainty labeled “b-tag eff rescaling” is the uncertainty associated with the rescaling of the b-tagging efficiency. UES refers to the scale of energy deposits not clustered into jets, and MES and EES refer to the muon and electron energy scales. The uncertainty labeled as "Id/Iso/Trg" is the uncertainty associated with the efficiency of the lepton identification, isolation, and trigger. The uncertainties in the integrated luminosity ? and on the acceptance due to PDF uncertainties and scale choices are not included in the fit, and are treated separately.

	Uncertainty	Variation	Effect on the measured cross section
Normalization	$t\bar{t}$	3.5%	3.8%
	Single top	5.4%	2.5%
	$Wusdcg$	13.2%	< 2%
	$Wc\bar{c}$	13.2%	< 2%
	Diboson	8.1%	< 2%
	Drell–Yan	7.9%	< 2%
	$\gamma+$ jets	10.0%	< 2%
	QCD	50%	2-3%
Shape	b-tag eff rescaling	8.4%	9.2%
	JES rescaling	0-6%	3.8%
	MES	0-3%	< 2%
	UES	0-3%	< 2%
	EES	0-3%	< 2%
	Id/Iso/Trg	0-4%	< 2%
Luminosity			2.6%
Scales (μ_R, μ_F)			10%
PDF choice			1%

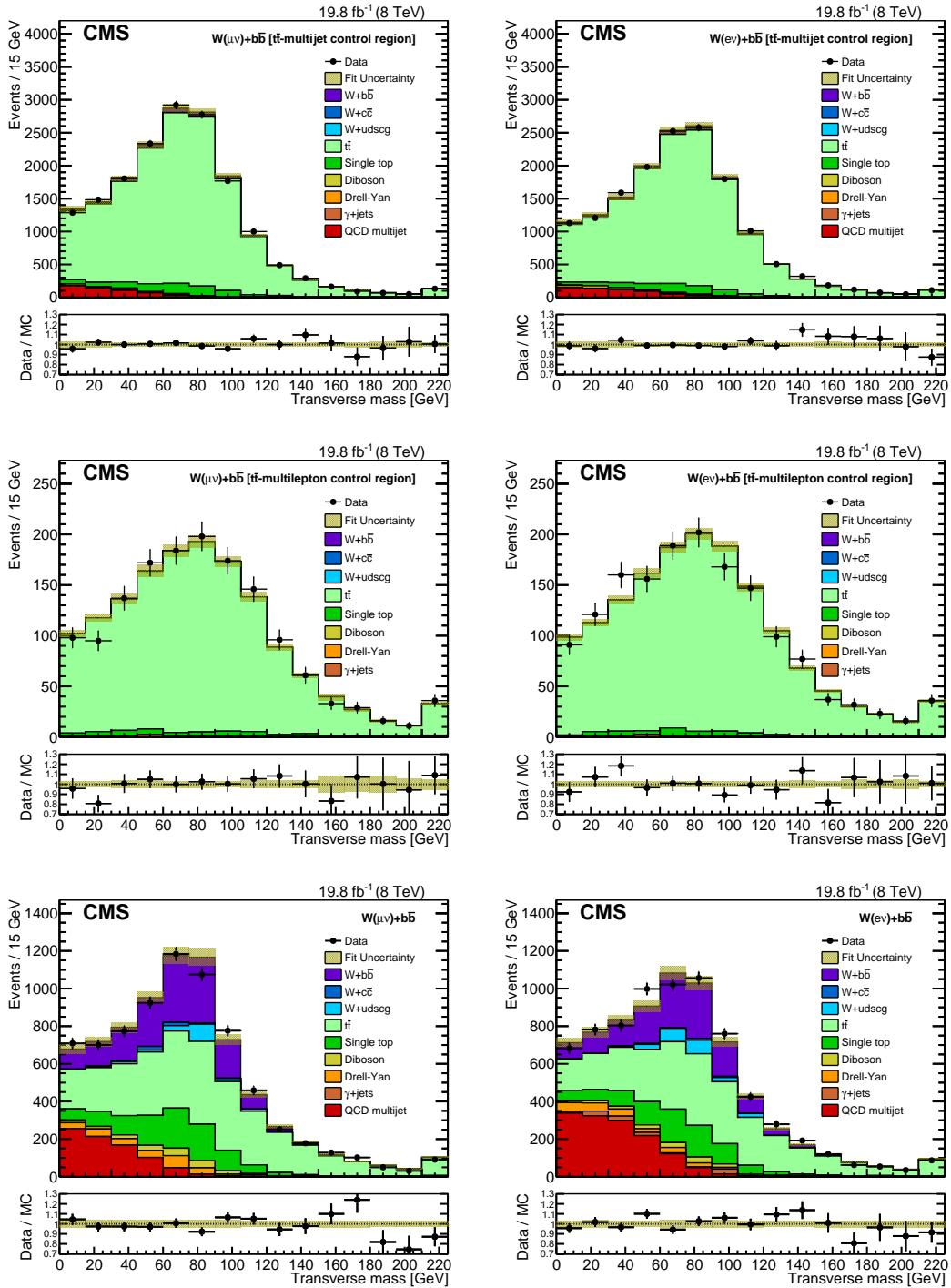
Table 5.3: Initial and final yields obtained in the $Wb\bar{b}$ signal region. The uncertainties in the signal strength represent the total uncertainty of the fit.

	Muon		Electron	
	Initial	Fitted	Initial	Fitted
Data	7432		7357	
$Wb\bar{b}$	1323	1712	1121	1456
$Wc\bar{c}$	60	61	36	37
$Wusdcg$	182	179	220	217
$t\bar{t}$	3049	3296	2640	2864
Single top	958	1008	820	865
Drell-Yan	261	265	220	224
Diboson	175	181	139	144
$\gamma+$ jets	-	-	98	105
QCD	1109	803	1654	1373
Total MC	7116	7505	6948	7284
Signal strength	1.21 ± 0.19		1.37 ± 0.23	
Combined	1.26 ± 0.17			

Table 5.4: Measured cross sections in the muon, electron, and combined lepton channels.

Channel	$\sigma(pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b})$ pb
Combined	$0.64 \pm 0.03(\text{stat}) \pm 0.10(\text{syst}) \pm 0.06(\text{theo}) \pm 0.02(\text{lumi})$
Muon	$0.62 \pm 0.04(\text{stat}) \pm 0.11(\text{syst}) \pm 0.06(\text{theo}) \pm 0.02(\text{lumi})$
Electron	$0.70 \pm 0.05(\text{stat}) \pm 0.15(\text{syst}) \pm 0.07(\text{theo}) \pm 0.02(\text{lumi})$

Figure 5.7: The transverse mass distributions (upper) in the $t\bar{t}$ -multijet phase space after fitting to obtain the b -tag efficiency rescale factors, (middle) in the $t\bar{t}$ -multilepton event enhanced control region after fitting to find the appropriate JES, and (lower) in the $Wb\bar{b}$ signal region after fitting simultaneously muon and electron decay channels. The lepton channels are shown separately with the muon sample on the left and the electron sample on the right. The last bin contains overflow events. The shaded area represents the total uncertainty in the simulation after the fit.



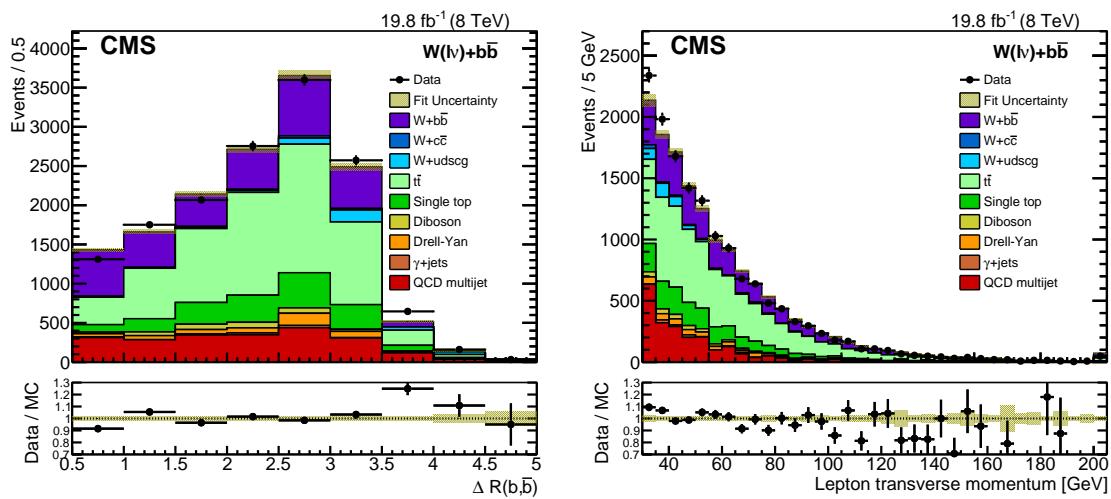
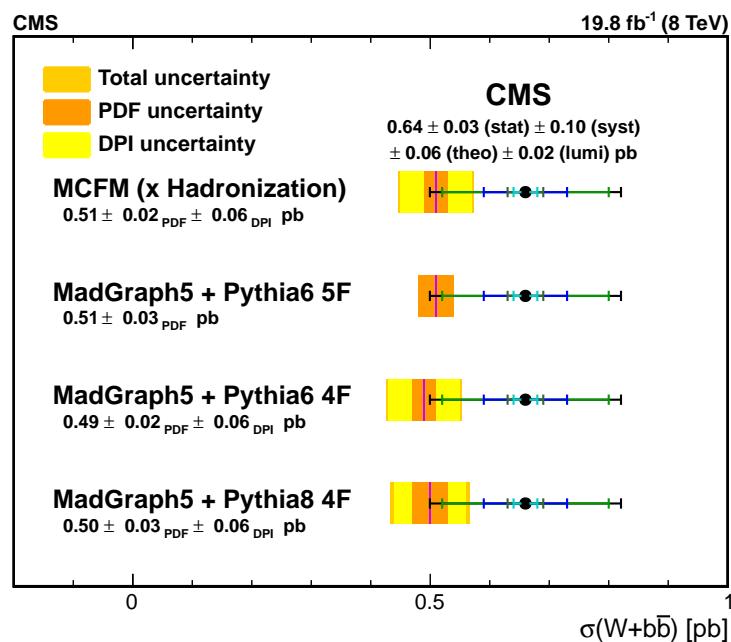


Figure 5.8: Distributions of $\Delta R(b\bar{b})$ and p_T^l after applying the results from the fits to the simulation. The QCD background shape is taken from an $m_T < 30$ GeV sideband and the muon and electron channels have been combined in these distributions. The last bin contains overflow events and the shaded area represents the total uncertainty in the simulation after the fit.

Figure 5.9: Comparison between the measured $pp \rightarrow Wb\bar{b} \rightarrow \ell\nu b\bar{b}$ cross section and various QCD predictions. The orange band indicates the uncertainty in the given sample associated with PDF choice and the yellow band represents the uncertainty associated with DPI. The labels 4F and 5F refer to the four- and five-flavour PDF schemes, and in the case of the MADGRAPH + PYTHIA 6 (5F) sample, the effects of DPI are already included in the generated samples so the DPI correction is not needed. The measured cross section is also shown with associated uncertainties.



6 MONOPHOTON ANALYSIS

The second analysis presented in this thesis is of the monophoton final state described in Section 2.2. The data are analyzed in the context of the SM process $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ in which the E_T^{miss} is interpreted as coming from the invisible decay of the Z boson, $Z \rightarrow \nu\bar{\nu}$. The data are also analyzed as a dark matter search under the interpretation that the E_T^{miss} arises from the annihilation of incoming particles into DM and the γ is initial state radiation recoiling against this process. Under this interpretation, limits are set on the cross section of DM as a function of the mediator mass for vector and axial-vector models. This analysis is performed using protons colliding at $\sqrt{s} = 13$ TeV provided by the LHC and detected by the CMS detector.

6.1 Event Selection

The sample of data analyzed was collected using a trigger that requires at least one photon HLT SC candidate with $p_T^\gamma > 165\text{GeV}$. Because photons interact electromagnetically, they are expected to deposit all of their energy in the ECAL, while jets typically have a neutral component and deposit some energy in the HCAL as well. To increase photon efficiency, the trigger therefore also requires at least 90% of the energy deposited in the calorimeters to be deposited in the ECAL. This trigger is 98% efficient at selecting photons which pass the other analysis selections. Events passing the trigger are further required to have at least one PF photon with $p_T^\gamma > 175\text{GeV}$ in the barrel fiducial region ($|\eta| < 1.44$).

To distinguish photons from electrons, which leave a similar signature of energy

deposits in the ECAL and HCAL, candidate photons are required to not have any associated track seeds in the pixel detector. To distinguish photons from jets, selections based on calorimetric information and isolation are applied. The fraction of energy deposited in the ECAL compared to the total deposit in the calorimeters is tightened relative to the trigger to be 95% and the shower shape variable describing the spread of the energy deposits in the η direction, $\sigma_{in\eta}$, is required to be $\sigma_{in\eta} < 0.0102$. Additionally, the photon is required to pass the isolation requirements described in Section ??.

Because photon objects are not reconstructed from tracks, there is an ambiguity in identifying the collision vertex that the photon originates from in the presence of pileup collisions. Association of a vertex to the photon candidate impacts the photon in two ways. First, the photon momentum direction is defined by the straight line which connects the ECAL cluster position and the identified vertex. Additionally, the isolation sum uses only the PF charged hadrons having tracks associated to the vertex. While the first effect is minor and is not relevant for this analysis, the second will cause photon candidates that are actually not isolated to appear isolated, if the vertex is misassigned. In practice, photon momentum is always computed with respect to the PV, but for the charged hadron isolation sum, all vertices are considered, and the maximum value of the isolation sum is used as a conservative estimate of the true isolation sum.

To reduce the contribution of backgrounds arising from occurrences in the CMS detector which did not originate from collisions, the energy pulse which seeded the photon cluster is required to be within ± 3 ns of the time expected for particles from

a collision, and the cluster must not be so narrow that it is consistent with a cluster formed by a single crystal. To reduce contamination from beam halo, the ECAL crystals not associated with the photon candidate are examined for evidence of the passage of a minimum-ionizing particle (MIP) roughly parallel to the beam axis (beam halo tag). If at least 4.9 GeV of energy is found deposited along this trajectory, the event is rejected. This value was determined by through optimizing the signal to background photon efficiencies, with a 95% identification for prompt photons and a 20% misidentificatin rate for deposits originating in beam halo events.

The candidate events are required to have $E_T^{\text{miss}} > 170$ GeV. The azimuthal opening angle between the candidate photon and E_T^{miss} is required to be greater than 2 radians to ensure that the main source of E_T^{miss} is not photon energy mismeasurement.

Because jet energy mismeasurement can also give rise to E_T^{miss} , events are rejected if the minimum azimuthal opening angle between E_T^{miss} and up to four leading jets ($\min\Delta\phi(E_T^{\text{miss}}, j)$) is less than 0.5 radians. As was the case in the $Wb\bar{b}$ analysis, jets are reconstructed using the PF algorithm, but in this analysis the jet clustering cone size is $\Delta R < 0.4$ radians.

Finally, events are also vetoed if they contain a charged lepton (an electron or a muon) with $p_T > 10$ GeV that is separated from the photon by $\Delta R > 0.5$ radians.

The effects of the various cuts with regard to the total number of events passing selections on the data is illustrated in Table 6.1. After applying all of the selection criteria, 77 candidate events are found in data.

Table 6.1: Listed below are the raw number of events in data passing the selection listed in the first column as well as all selections in higher rows.

Selection	Events passing selection
PF photon, $p_T^\gamma > 175$ GeV, $ \eta^\gamma < 1.44$	X
$ t^{\text{seed}} < 3$ ns	X
$MIP < 4.9$ GeV	X
$E_T^{\text{miss}} > 170$ GeV	X
$\Delta\phi(E_T^{\text{miss}}, \gamma) > 2$ radians	X
$\Delta\phi(E_T^{\text{miss}}, j) > 0.5$ radians	X
Veto charged lepton	X

6.2 Estimation of Background Contributions

The dominant SM processes contributing to this phase space of the candidate events are the associated productions of a Z or W boson with a high-energy photon. If the Z boson decays into a neutrino-antineutrino pair, the final state exhibits a high- E_T photon and large missing transverse energy. Similarly, if the W boson decays into a lepton-neutrino pair and the lepton is outside of the detector acceptance or fails reconstruction, the event appears to be $\gamma + E_T^{\text{miss}}$. Together, these two processes account for approximately 75% of the events as estimated using Monte Carlo (MC) simulations. Hard-scattering events are generated with MADGRAPH5_AMC@NLO version 2.2 ? at leading order (LO) in QCD, with NNPDF3.0 LO ($\alpha_s = 0.130$) as the parton distribution function. Parton shower and hadronization is performed by PYTHIA8.2 ?. Generated particles are processed through the full GEANT-based simulation of the CMS detector ?? and event reconstruction used for data. Minimum-bias simulations are overlaid to model pileup interactions.

6.2.1 Reweighting

To account for differences arising from imperfect modeling of the data in the simulation, a total correction factor $\rho = 0.99 \pm 0.06$ is applied to all MC-based estimates. This is the product of individual correction factors which are each taken as the ratio of the efficiency measured in data and in simulation. The efficiency for photon identification is measured and provided centrally using $Z \rightarrow e\bar{e}$ events as 0.99 ± 0.016 for photon identification measured using $Z \rightarrow e\bar{e}$ events. The photon seed trigger efficiency is also measured using $Z \rightarrow e\bar{e}$ events, and is found to be 1.00 ± 0.0246 using jet triggers as a reference. The efficiencies for worst isolation, beam halo tag and lepton veto were measured using events in data which were triggered as having at least one muon and on MC using a combination of Drell-Yan, $t\bar{t}$ and VV samples requiring $Z\gamma \rightarrow \mu\bar{\mu}\gamma$ events to be identified by the dimuon pair with mass in the range $61 < m_{\mu\mu} < 121$ GeV. The measured relative efficiencies are 1.00 ± 0.05 .

Generated samples are weighted on an event by event basis with a product of two factors. The first factor matches the distribution of the generator-level photon p_T to that calculated at next-to-next-to-leading order (NNLO) in QCD using the DYRES ? calculator and the second factor, taken from Refs. ??, further corrects this distribution to account for electroweak next-to-leading order (NLO) effects.

6.2.2 $V\gamma$ Estimates

After accounting for the event selection efficiency difference between data and MC, respectively 42.1 ± 6.3 and 10.7 ± 1.5 events are estimated from $Z\gamma \rightarrow \nu\bar{\nu}\gamma$ and $W\gamma \rightarrow \ell\nu\gamma$. Four sources of systematic uncertainty on $Z\gamma$ and $W\gamma$ estimates are

considered: PDF and scale uncertainties are found using the the LHAPDF/PDF4LHC recommendations of varying the scales up and down by a factor of 2 and using PDf sets from CTEQ, MSTW and NNPDF to be 5.37% and 8.9% respectively. Electroweak correction uncertainties are estimated conservatively as the quoted uncertainties which are 11% for $Z\gamma$ and 7% for $W\gamma$. Scale factors are estimated to have an uncertainty of 6% which mostly arises from the statistical limitations of the data samples used. The systematic uncertainty due to jet/ E_T^{miss}/γ energy scale and pileup, is estimated at 6.2% by shifting the energy of the respective PF object and observing the relative change in the number of events passing selections. As a crosscheck, the total contribution from $Z\gamma \rightarrow \nu\bar{\nu}\gamma$ is estimated in data using a sample of $Z\gamma \rightarrow \ell\bar{\ell}\gamma$ candidates, where the leptons from the decay of the Z boson are removed and considered as E_T^{miss} ?. This provides an estimate of 64.6 ± 17.6 , where the uncertainty is dominated by the size of the sample.

6.2.3 Elecron Mis-ID

The most important SM background comes from events where electrons are misidentified as photons, mainly in the $W \rightarrow e\nu$ process. Seeding efficiency in the pixel detector for electron tracks is $\epsilon = 0.982 \pm 0.004$ for electrons with $p_T > 100\text{GeV}$. This efficiency is measured in data using the tag-and-probe method ? on $Z \rightarrow e\bar{e}$ events, and is verified with MC simulation. Electrons from W boson decay that are not seeded appear as isolated photons accompanied with large E_T^{miss} from the escaping neutrino. This class of events is modeled by an electron proxy event sample selected in data using criteria that are identical to those described in Sec. 6.1, except the photon

candidate is required to have a pixel seed. The number of electron proxy events is then scaled by $(1 - \epsilon)/\epsilon$ where ϵ is the efficiency to yield an estimated contribution of 7.4 ± 1.2 from electron misidentification events. The dominant uncertainty in the estimate is the statistical uncertainty in the tag-and-probe fit, and is assessed by generating a large ensemble of toy dielectron mass distributions on which the fit procedure is repeated. The standard deviation of the number of $Z \rightarrow e\bar{e}$ events obtained from the fits is then propagated to the uncertainty in the efficiency.

6.2.4 Non-collision Backgrounds

Non-collision backgrounds, from things such as detector noise, cosmic rays, and beam halo, are estimated from the time distribution of the cluster seeds since each process exhibits a distinctive time distribution when the cluster is in the ECAL barrel. Templates for anomalous signals, cosmic ray muons, and beam halo events are obtained by inverting the shower shape and beam halo tag requirements, and are fitted to the timing distribution of the candidate sample. The only nonnegligible residual contribution to the candidate sample is found to arise from the beam halo, with an estimated 5.9 ± 4.7 events coming directly from the template fit.

6.2.5 Minor SM Processes

The SM processes $W \rightarrow \ell\nu\gamma$, $Z \rightarrow \ell\bar{\ell}\gamma$, $W(l\nu)$ and $\gamma + jets$ are generated with MADGRAPH5_aMC@NLO at LO [2] with up to 2 jets and then processed with PYTHIA 6.426 generator [3] for showering and hadronization, with the NNPDF3.0 LO($\alpha_s = 0.130$) parton distribution function. The total background expectation from

these processes is 3.05 ± 0.67 events, where the uncertainty includes the statistical and systematic uncertainty due to scale factor and jet/ E_T^{miss}/γ energy scale.

6.3 Results

After applying the full selection criteria, 77 events in 2.32 fb^{-1} of data remain. Table 6.2 shows the estimated number of events and uncertainty from each background for the full 2015 run. The p_T spectrum and PF E_T^{miss} of the full combination of selected candidate events and estimated backgrounds can be seen in Figure 6.1 along with distributions of p_T/E_T^{miss} and the number of jets.

Figure 6.1: The photon p_T and E_T^{miss} distribution for the candidate sample, compared with estimated contributions from SM backgrounds are shown. Here QCD γ refers to $\gamma+\text{jet}$ background and the background uncertainty includes statistical and systematic error.

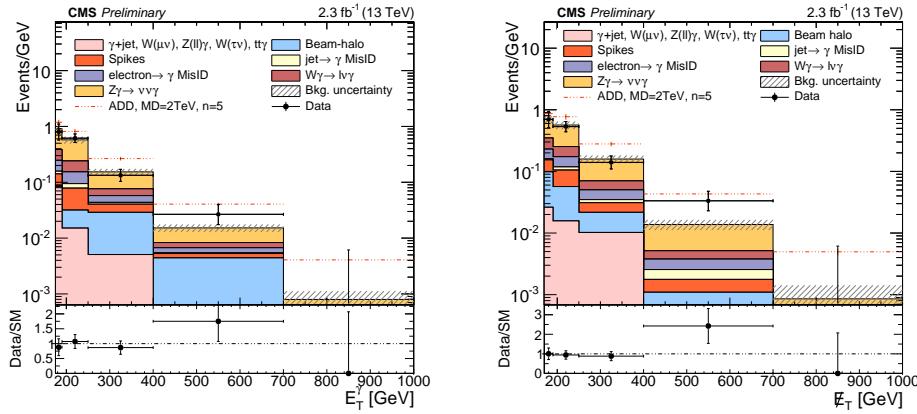


Table 6.2: Summary of estimated backgrounds and observed total number of candidates for 2.32 fb^{-1} of 2015 data. The category Others includes $W \rightarrow \mu\nu$, $Z \rightarrow \ell\bar{\ell}\gamma$ and $t\bar{t}\gamma$

Process	Estimate
$Z\gamma \rightarrow \nu\bar{\nu}\gamma$	42.10 ± 6.31
$W\gamma \rightarrow \ell\nu\gamma$	10.69 ± 1.49
$W \rightarrow e\nu$	7.80 ± 1.78
$jet \rightarrow \gamma$ fakes	3.36 ± 1.13
Beam halo	5.9 ± 4.7
Others	3.05 ± 0.67
Total Expectation	72.9 ± 8.30
Data	77

6.3.1 $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ Cross Section Measurement

The $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ cross section for $p_T^\gamma > 175$ GeV in the range $|\eta| < 1.4$ is calculated using the formula

$$\sigma(pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma) = \frac{N_{data} - N_{BG}}{A \times \epsilon \times L} \quad (6.1)$$

where N_{data} is the observed number of events, N_{BG} is the number of estimated background events, A is the geometrical and kinematic acceptance of the selection criteria, ϵ is the selection efficiency within the acceptance, and L is the integrated luminosity. The product of $A \times \epsilon_{MC}$ is estimated from LO MADGRAPH simulation and a correction factor, ρ , described in Section 6.2.1 is applied to account for the difference between the efficiency in the data and Monte Carlo:

$$A \times \epsilon = A \times \epsilon_{MC} \times \rho. \quad (6.2)$$

The product of $A \times \epsilon_{MC}$ is estimated to be 0.314 ± 0.002 (stat) ± 0.048 (syst) and rho is 0.99 ± 0.06 .

The photon energy scale, jet and E_T^{miss} energy scale and resolution, and pileup related contributions are considered as sources of systematic uncertainty in the acceptance calculation. The uncertainty on the photon energy scale is about 1.5% and the uncertainty from variations in the E_T^{miss} energy scales is 5%. Contributions from the jet energy scale are accounted for in the uncertainty on the E_T^{miss} . The uncertainty on the integrated luminosity is 2.7% ?. A summary of the systematic uncertainties are shown in Table 6.3.

Table 6.3: Summary of systematic unceratinties for signal and different background sources.

Sources	$Z\gamma \rightarrow \nu\bar{\nu}\gamma$ [%]	$W\gamma \rightarrow \ell\nu\gamma$ [%]	j faking γ [%]	e faking γ	γj	Other bkg [%]
Luminosity	2.7	2.7	-	-	2.7	2.7
PDF and Scale	5.37	8.9	-	-	-	-
EWK corrections	11	7	-	-	-	-
j faking γ	-	-	30	-	-	-
e faking γ	-	-	-	20	-	-
$j, E_T^{\text{miss}}, \gamma$ energy scale	6	6	-	-	6	6
Scale Factors	6	6	-	-	6	6

The measured cross section for $pp \rightarrow Z\gamma \rightarrow \nu\bar{\nu}\gamma$ for photon $p_T > 175$ GeV within rapidity range $|\eta_\gamma| < 1.4$ is

$$64.06 \pm 12.14(\text{stat}) \pm 12.88(\text{syst}) \pm 1.72(\text{lumi}) \text{ fb.} \quad (6.3)$$

The NNLO therotical cross section is 65.55 ± 0.02 fb where the uncertainty includes only the scale variaitons. The measured cross section agrees well with the NNLO theoretical cross section and this agreement with the SM prediction constrains possible

DM models.

6.3.2 Limits on Dark Matter

Interpreting these results as setting limits on the cross section of a DM particle as a function of DM mass, Tables 6.4 and 6.5 show 90% confidence level (CL) upper limits on the production cross sections provided for the Vector and the Axial-vector model for a mediator mass of 10 TeV.

Table 6.4: Observed (expected) 90%CL upper limits on the DM production cross section σ for vector mediator mass 10 TeV.

Mass DM [GeV]	σ [fb]
1	3.821(3.242)
10	3.820(3.244)
50	3.827(3.249)
150	3.826(3.254)
500	3.588(3.052)
1000	3.370(2.862)

Table 6.5: Observed (expected) 90%CL upper limits on the DM production cross section σ for axial vector mediator mass of 10 TeV

Mass DM [GeV]	σ [pb]
1	3.782(3.211)
10	3.785(3.213)
50	3.793(3.213)
150	3.754(3.192)
500	3.488(2.961)
1000	3.30(2.814)

Figure 6.2 shows the upper limits on the ratio of the cross section with respect to the theoretical predictions ($\mu = \sigma^{95\%}/\sigma_{Th}$) for the vector and axial-vector mediator

scenarios on the m_χ - m_M plane. The solid red and black curves are the expected and observed exclusion contours. The uncertainty on the expected upper limit includes the experimental uncertainties. For the simplified DM model considered, a mediator mass of up to 600 GeV is excluded for $m_\chi < 10$ GeV.

Exclusion contours in Figure 6.2 are also translated into the m_χ - $\sigma_{SI/SD}$ plane as shown in Figure 6.3 where $\sigma_{SI/SD}$ are the spin-independent/dependent DM-nucleon scattering cross sections ?. For these contours, we use 90% CL to do a direct comparison with the limits from the direct detection experiments ???.

For DM EFT model with a contact interaction of type $\gamma\gamma\chi\bar{\chi}$, upper limits are placed on the production cross section, which are then translated into the lower limits on the suppression scale Λ . The 95% CL observed and expected lower limits on Λ as a function of dark matter mass m_χ are shown in Figure 6.4. Values of Λ up to 542GeV are excluded at 95% CL.

Figure 6.2: 95% CL upper limits on $\mu = \sigma/\sigma_{Th}$ in the m_χ - M_M plane for vector and axial-vector mediator, assuming couplings $g_q = 0.25$ and $g_\chi = 1$. The solid red and black curves are the expected and observed exclusion contours. The dotted black contours around the observed limit and the dotted red contours around the expected limit represent the one standard deviation theoretical uncertainties in the cross section and the combination of the statistical and experimental systematic uncertainties, respectively.

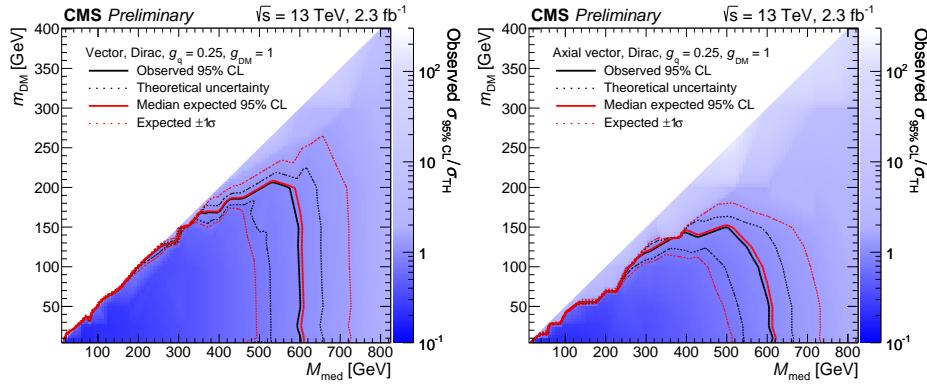


Figure 6.3: The 90% CL exclusion limits on the χ -nucleon scattering cross section in a simplified model of dark matter production involving a vector and axial-vector operator as a function of the dark matter mass m_χ .

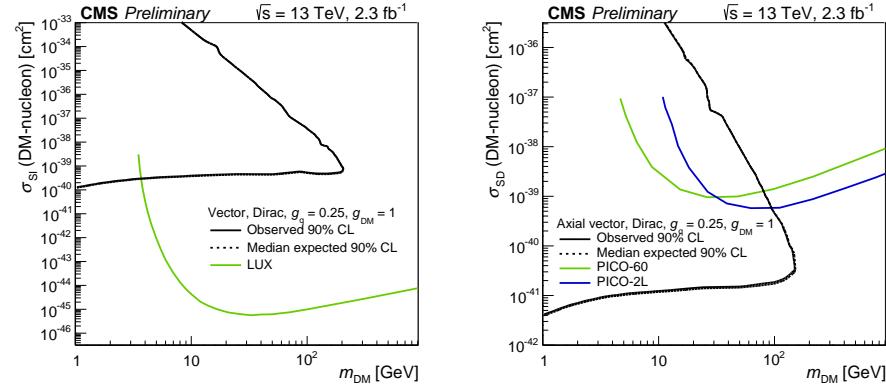


Figure 6.4: (a) The 95% CL observed and expected lower limits on Λ for a dimension-7 operator EFT model with a contact interaction of type $\gamma\gamma\chi\bar{\chi}$ as a function of dark matter mass m_χ .

