Perfect Fluidity

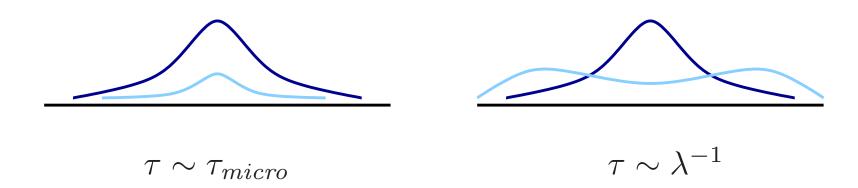
in Cold Atomic Gases?

Thomas Schaefer

North Carolina State University

Hydrodynamics

Long-wavelength, low-frequency dynamics of conserved or spontaneoulsy broken symmetry variables



Historically: Water $(\rho, \epsilon, \vec{\pi})$

Example: Simple Fluid

Continuity equation

$$\frac{\partial \rho}{\partial t} + \vec{\nabla}(\rho \vec{v}) = 0$$

Euler (Navier-Stokes) equation

$$\frac{\partial}{\partial t}(\rho v_i) + \frac{\partial}{\partial x_j} \Pi_{ij} = 0$$

Energy momentum tensor

$$\Pi_{ij} = P\delta_{ij} + \rho v_i v_j + \eta \left(\partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right) + \dots$$

reactive

dissipative

Questions

Existence of solutions (1M\$)

Structure of solutions (Turbulence, Shocks, ...)

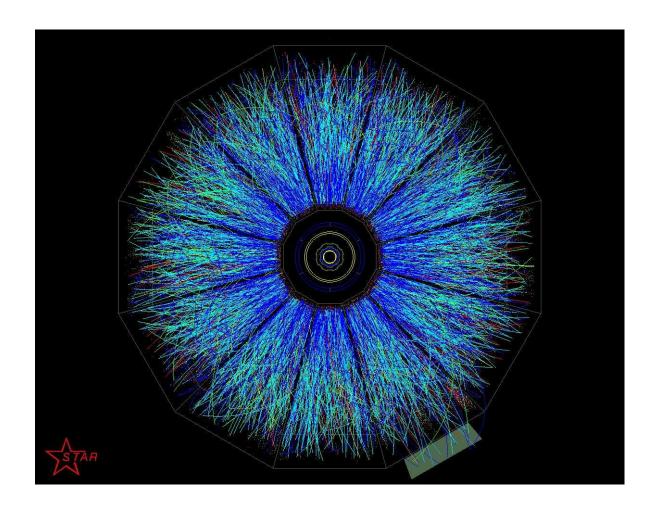
How to find constitutive relations from micro physics

Determine transport coefficients from micro physics

BNL and RHIC



Heavy Ion Collision

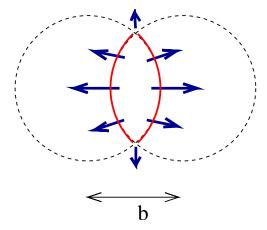


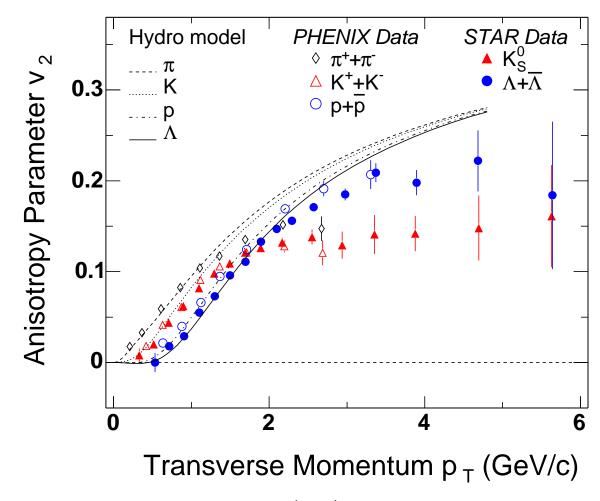
Star TPC

Elliptic Flow

Hydrodynamic expansion converts coordinate space anisotropy

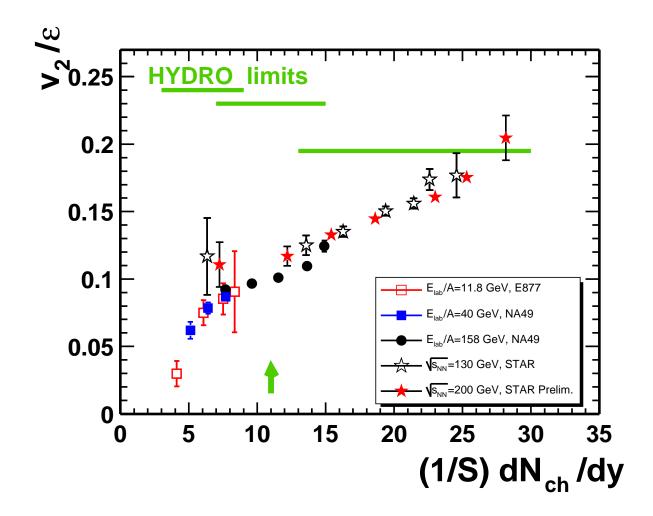
to momentum space anisotropy





source: U. Heinz (2005)

Elliptic Flow II



Requires "perfect" fluidity ($\eta/s < 0.1$?)

(s)QGP saturates (conjectured) universal bound $\eta/s = 1/(4\pi)$?

Viscosity Bound: Rough Argument

Kinetic theory estimate of shear viscosity

$$\eta \sim \frac{1}{3}n\bar{v}ml = \frac{2}{3}n\left(\frac{1}{2}m\bar{v}^2\right)\frac{l}{\bar{v}} = \frac{2}{3}n\epsilon\tau_{mft}$$

Uncertainty relation implies

$$\frac{\eta}{s} \sim \frac{\epsilon \tau_{mft}}{k_B n} \sim \frac{E_{av} \tau_{mft}}{k_B} \leq \frac{\hbar}{k_B}$$

where we have used $s \sim k_B n$.

Danielewicz & Gyulassy (1984)

Validity of kinetic theory as $E\tau \sim \hbar$? Why η/s ? Why not η/n ?

Holographic Duals at Finite Temperature

Thermal (superconformal) field theory

 \Leftrightarrow

 AdS_5 black hole

CFT temperature

 \Leftrightarrow

 \Leftrightarrow

CFT entropy

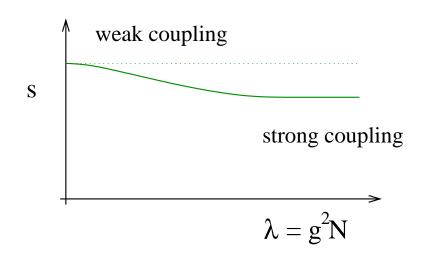
Hawking temperature of black hole Hawking-Bekenstein entropy

 \sim area of event horizon

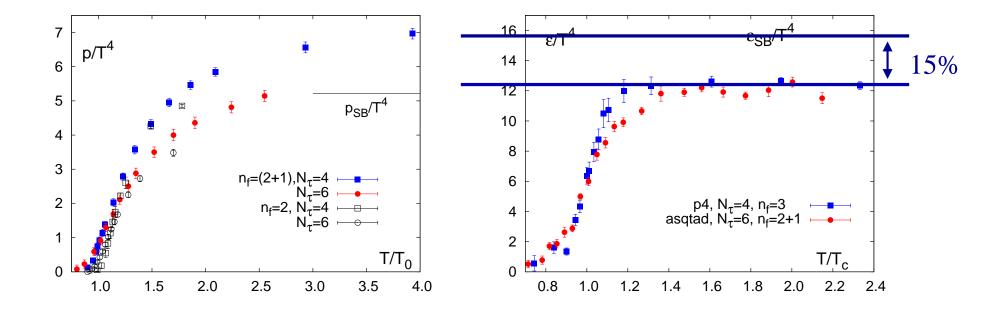
Strong coupling limit

$$s = \frac{\pi^2}{2} N_c^2 T^3 = \frac{3}{4} s_0$$

Gubser and Klebanov



Quark Gluon Plasma Equation of State (Lattice)



Compilation by F. Karsch (SciDAC)

Holographic Duals: Transport Properties

Thermal (superconformal) field theory

 \Leftrightarrow

 AdS_5 black hole

CFT entropy

 \Leftrightarrow

Hawking-Bekenstein entropy

shear viscosity

Graviton absorption cross section

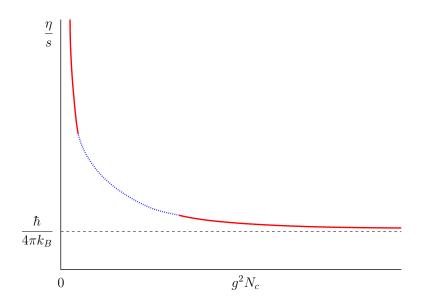
 \sim area of event horizon

 \sim area of event horizon

Strong coupling limit

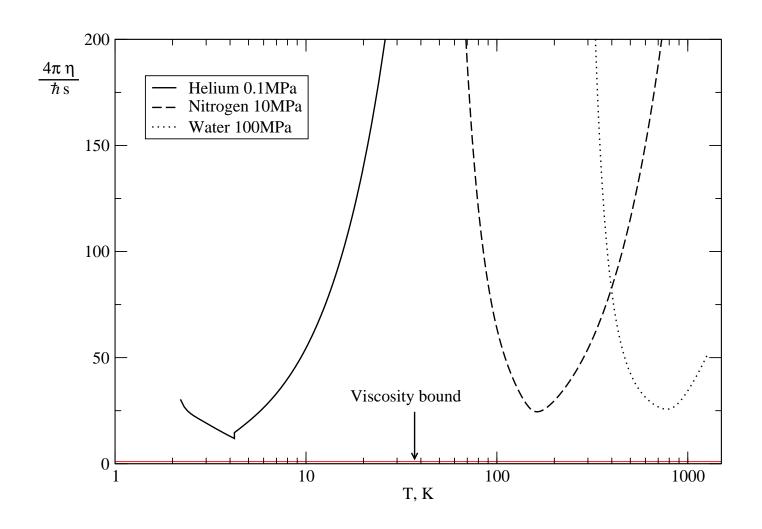
$$\frac{\eta}{s} = \frac{\hbar}{4\pi k_B}$$

Son and Starinets



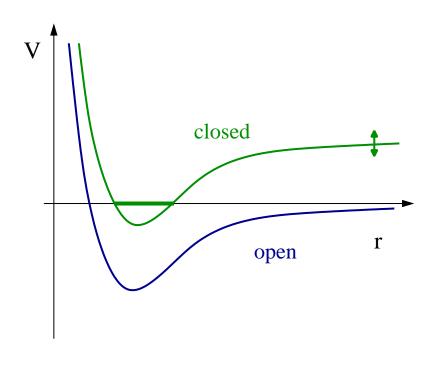
Strong coupling limit universal? Provides lower bound for all theories?

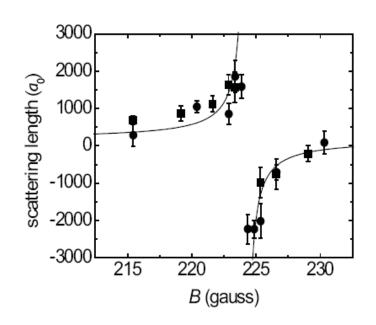
Viscosity Bound: Common Fluids



Designer Fluids

Atomic gas with two spin states: "↑" and "↓"





Feshbach resonance

$$a(B) = a_0 \left(1 + \frac{\Delta}{B - B_0} \right)$$

"Unitarity" limit
$$a \to \infty$$

$$\sigma = \frac{4\pi}{k^2}$$

Why are these systems interesting?

System is intrinsically quantum mechanical

cross section saturates unitarity bound

System is scale invariant at unitarity. Universal thermodynamics

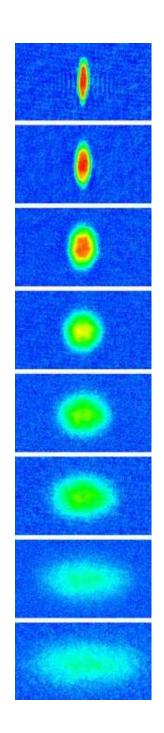
$$\frac{E}{A} = \xi \left(\frac{E}{A}\right)_0 = \xi \frac{3}{5} \left(\frac{k_F^2}{2M}\right)$$

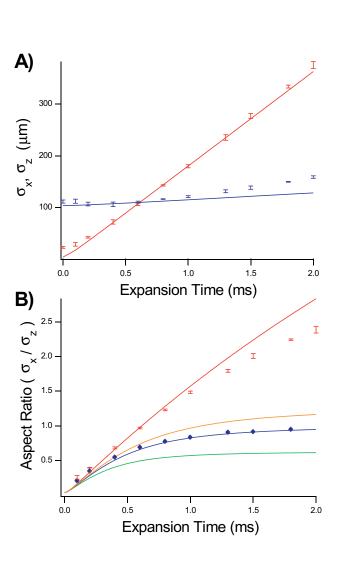
System is strongly coupled but dilute

$$(k_F a) \to \infty \qquad (k_F r) \to 0$$

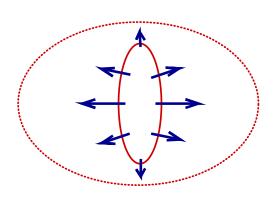
Strong elliptic flow observed experimentally

Elliptic Flow



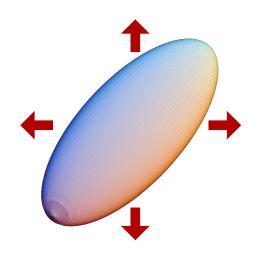


Hydrodynamic expansion converts coordinate space anisotropy to momentum space anisotropy

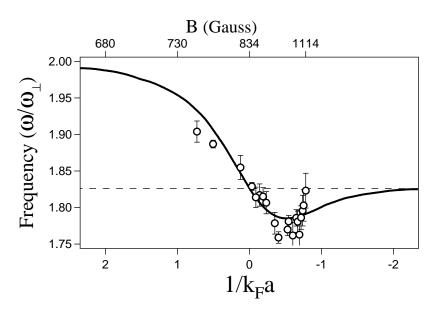


Collective Modes

Radial breathing mode



Kinast et al. (2005)



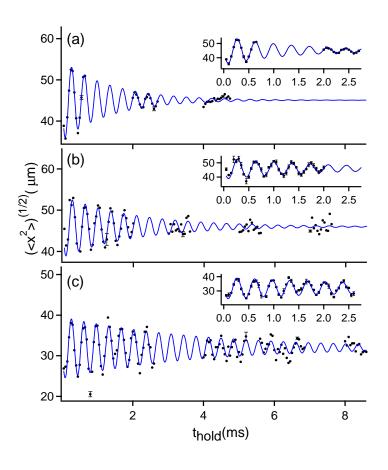
Ideal fluid hydrodynamics, equation of state $P \sim n^{5/3}$

$$\frac{\partial n}{\partial t} + \vec{\nabla} \cdot (n\vec{v}) = 0$$

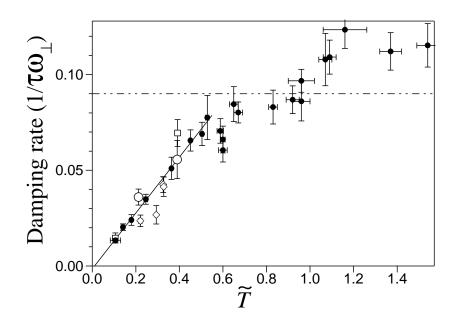
$$\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\frac{1}{mn} \vec{\nabla} P - \frac{1}{m} \vec{\nabla} V$$

$$\omega = \sqrt{\frac{10}{3}} \omega_{\perp}$$

Damping of Collective Excitations



$$T/T_F = (0.5, 0.33, 0.17)$$



 $\tau\omega$: decay time \times trap frequency

Kinast et al. (2005)

Viscous Hydrodynamics

Energy dissipation (η, ζ, κ) : shear, bulk viscosity, heat conductivity)

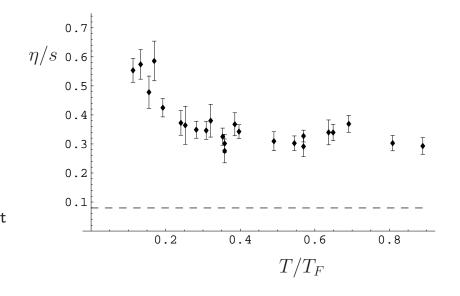
$$\dot{E} = -\frac{\eta}{2} \int d^3x \left(\partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right)^2$$
$$-\zeta \int d^3x \left(\partial_i v_i \right)^2 - \frac{\kappa}{2} \int d^3x \left(\partial_i T \right)^2$$

Shear viscosity to entropy ratio

(assuming
$$\zeta = \kappa = 0$$
)

$$\frac{\eta}{s} = \frac{3}{4} \xi^{\frac{1}{2}} (3N)^{\frac{1}{3}} \frac{\Gamma}{\omega_{\perp}} \frac{\bar{\omega}}{\omega_{\perp}} \frac{N}{S}$$

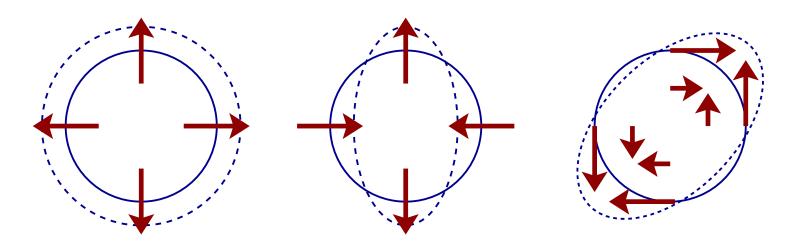
see also Bruun, Smith, Gelman et



al.

Damping dominated by shear viscosity?

Study dependence on flow pattern



Study particle number scaling

viscous hydro: $\Gamma \sim N^{-1/3}$

Boltzmann: $\Gamma \sim N^{1/3}$

Role of thermal conductivity?

suppressed for scaling flows: $\delta T \sim T(\delta n/n) \sim const$

Final Thoughts

Cold atomic gases provide interesting, strongly coupled, model system in which to study sources of dissipation.

$$\eta/s \sim 1/3$$

Smaller than any other known liquid (except for QGP?). Since other sources of dissipation exist, this is really an upper bound.

Conjectured bound has a smooth non-relativistic limit. Note that the $a \to \infty$ limit can also be realized in QCD (by tuning μ, μ_e and m_q).

But: In non-relativistic systems $s \gg n$ possible

Purely field theoretic proofs?

 $\mathcal{N}=4$ SUSY YM is special because there is no phase transition. In real systems there is a phase transition as the coupling becomes large, and the new phase (confined in QCD, superfluid in the atomic system) has weakly coupled low energy excitations, and a large viscosity.

No quasi-particles in sQGP?