

Multi-meson systems in lattice QCD

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- Multi-hadron systems in QCD
- $n < 13$ pions and kaons
 - Two and three body interactions
 - Pion and kaon condensates
 - In-medium screening of the $Q\bar{Q}$ potential

Many body* QCD

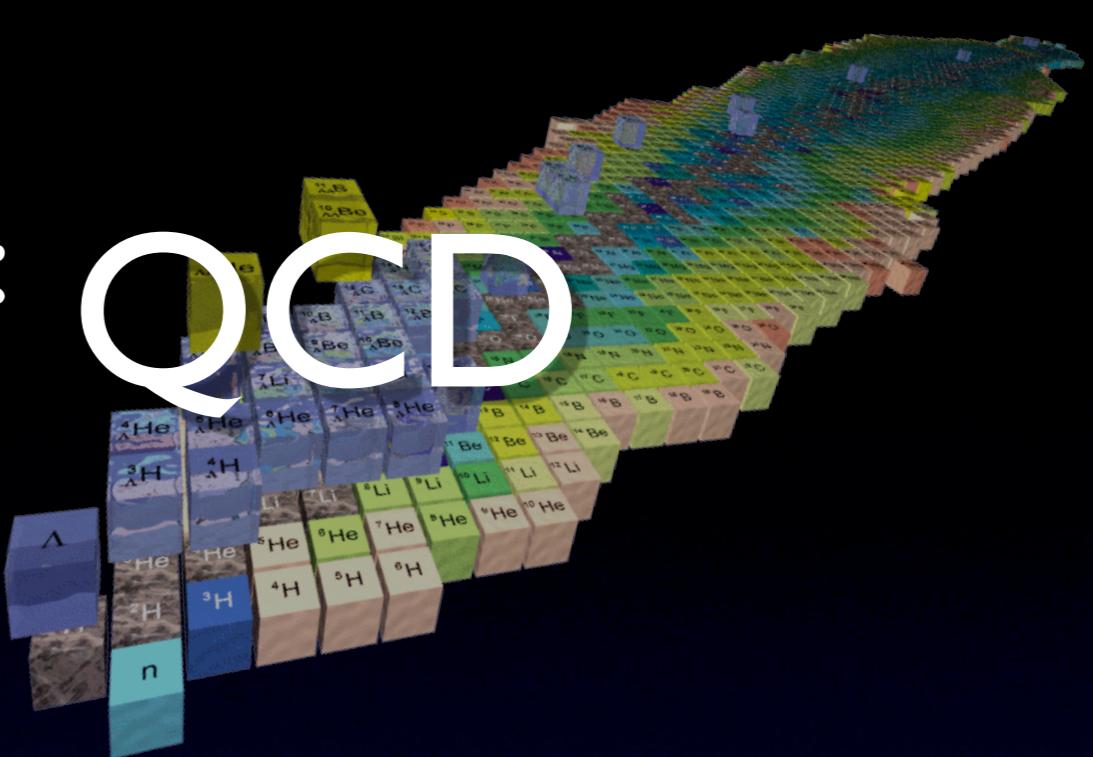
* Many body: more than two hadrons

Many body* QCD

- Nuclear physics from QCD

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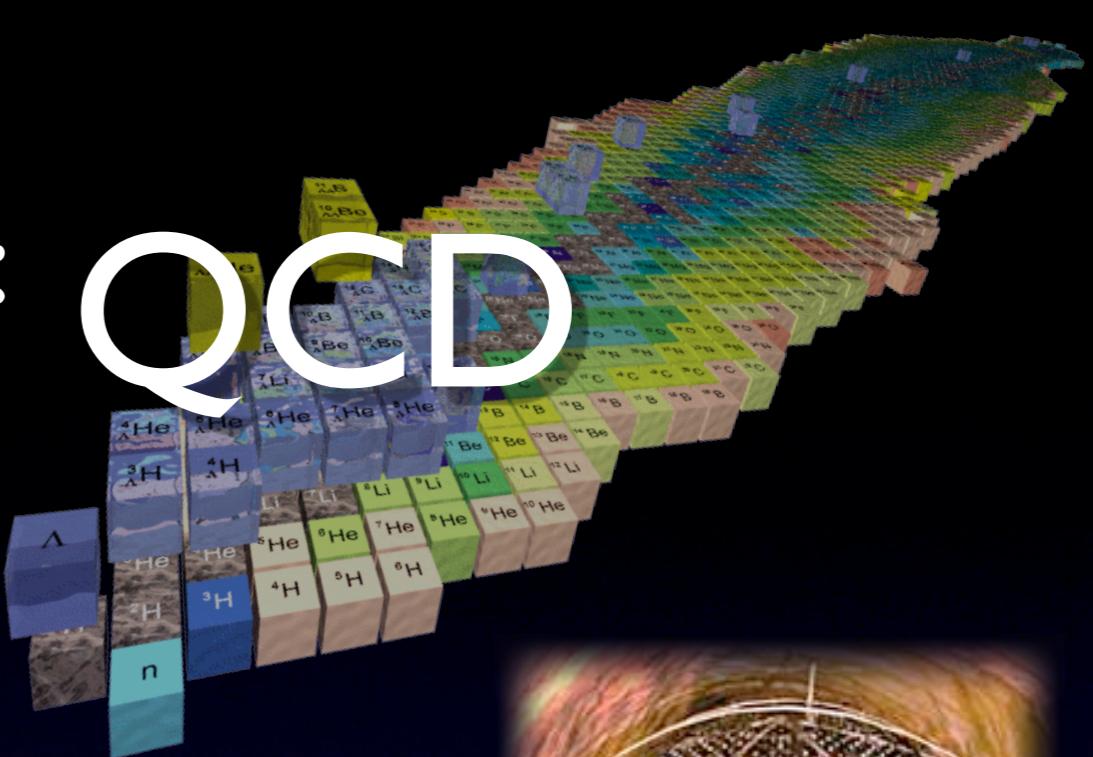
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- Nuclear physics from QCD
- Spectra

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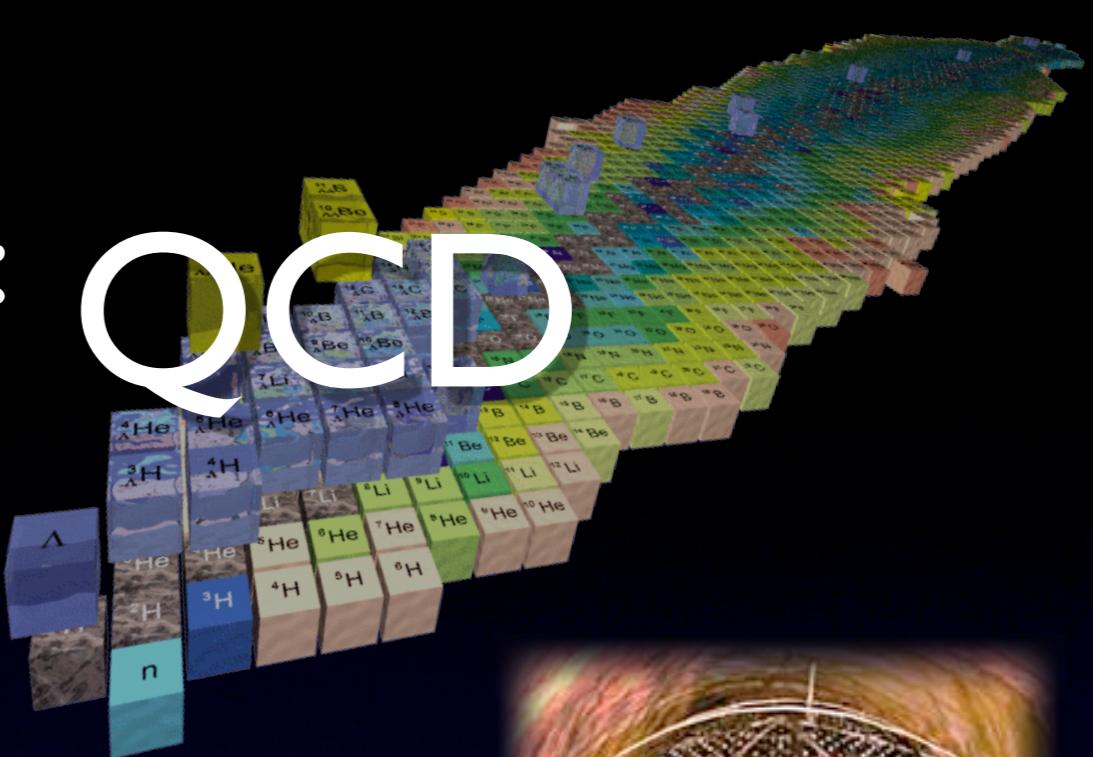


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 - Interactions and decays



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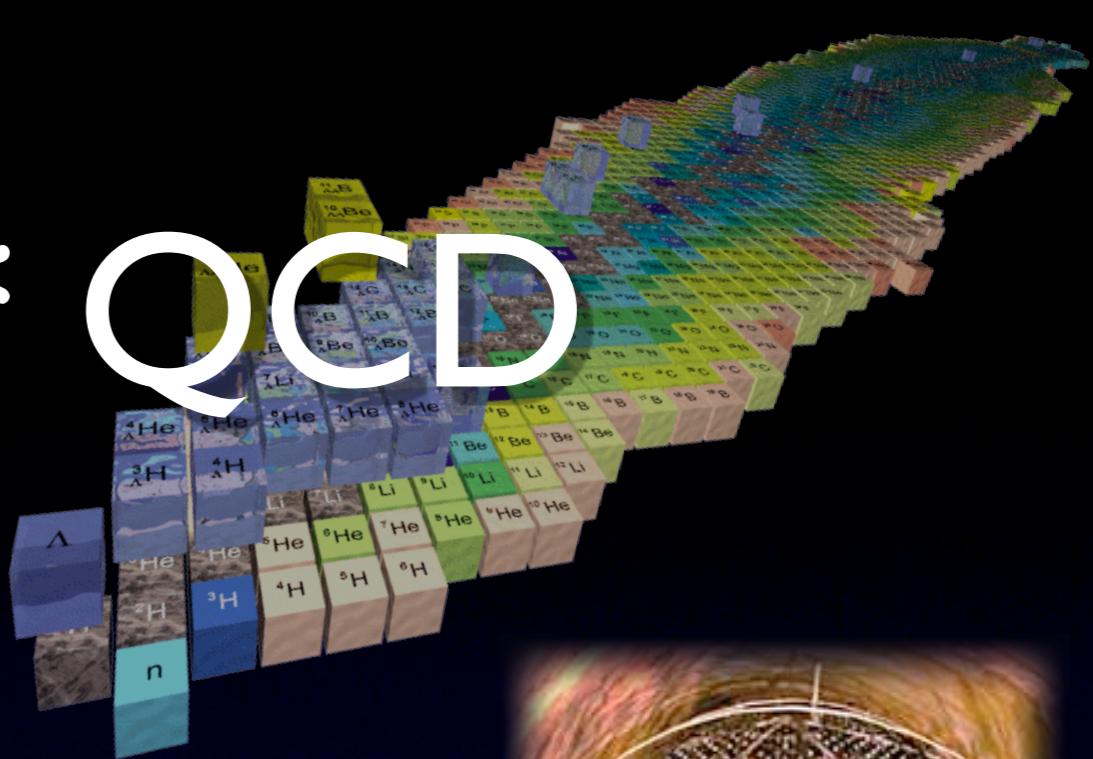


- Nuclear physics from QCD
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 - m_q dependence

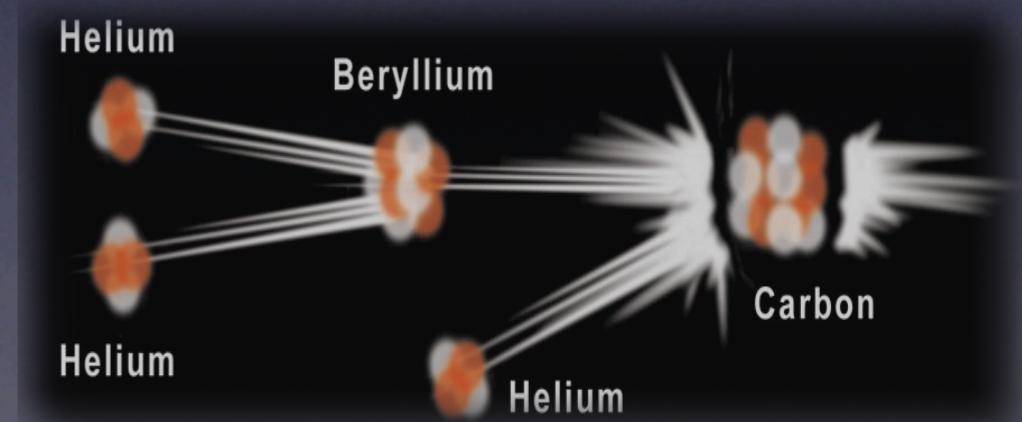


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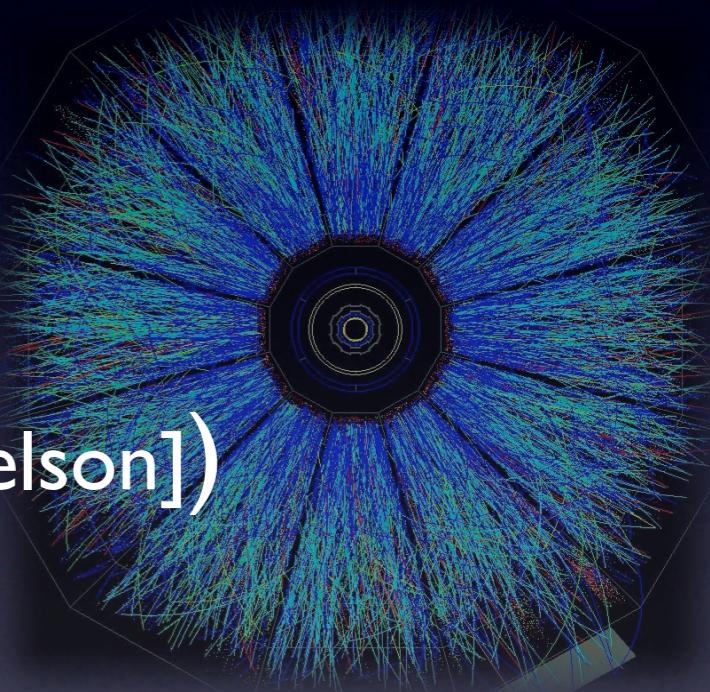
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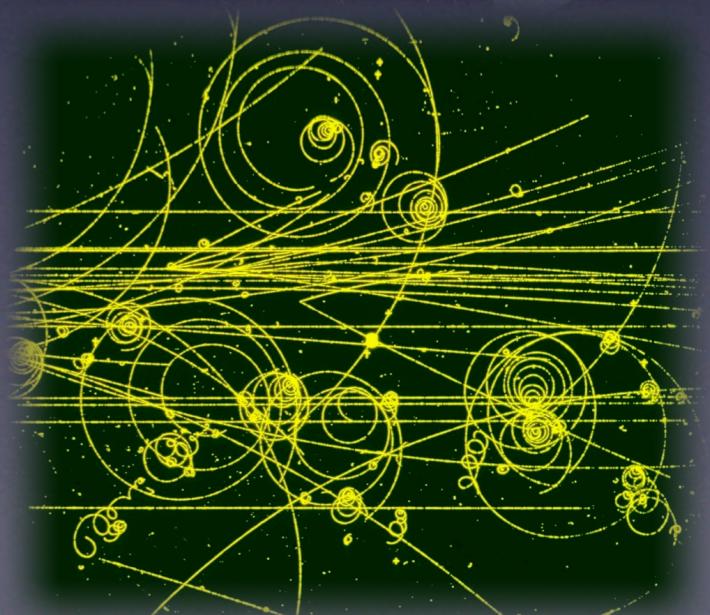
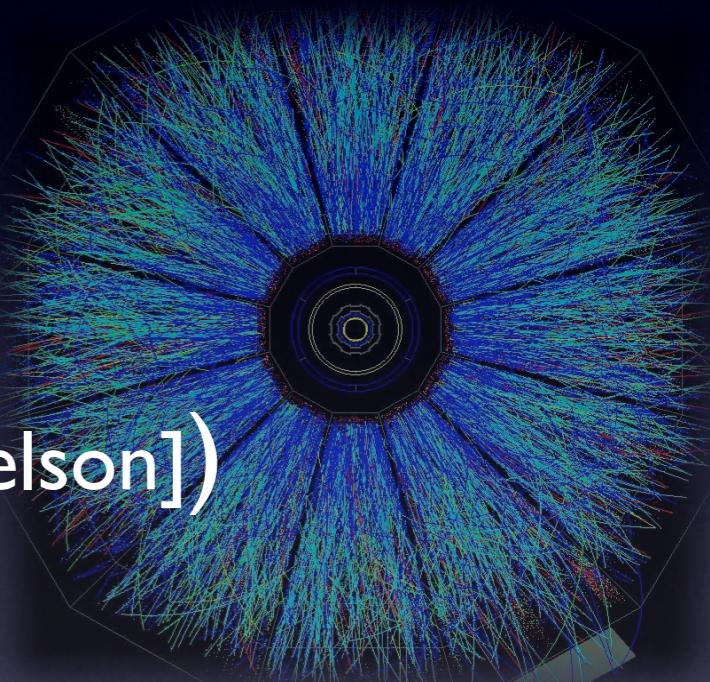
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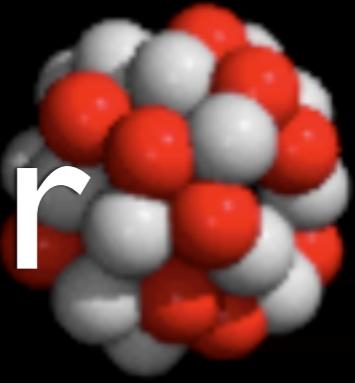


Many body QCD

- π and K condensation
- EOS in n stars (kaons [Kaplan&Nelson])
- Three π (and K) interferometry (HBT) at SPS and RHIC
- Many body decays:



Many body correlator



- How can we compute the mass of ^{208}Pb in QCD?

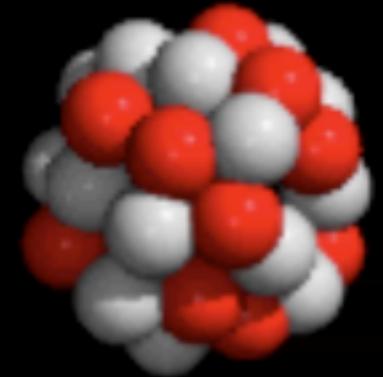
$$\langle 0 | T q_1(t) \dots q_{624}(t) \bar{q}_1(0) \dots \bar{q}_{624}(0) | 0 \rangle$$

- Large time behaviour gives GS energy (up to EW effects)

$$\stackrel{t \rightarrow \infty}{\rightarrow} \# \exp(-M_{Pb} t)$$

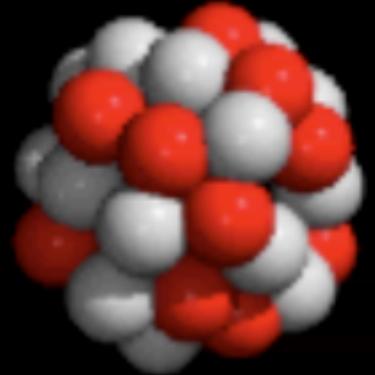
- But...

Many difficulties



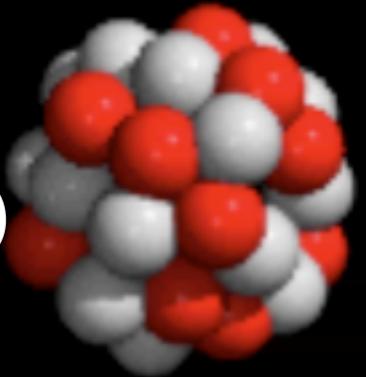
- Interpolators and contractions
- Statistical noise
- Signals for very massive states
- Small energy splittings
- Numerical precision

Many difficulties



- Interpolators and contractions
- Statistical noise
- Signals for very massive states
- Small energy splittings
- Numerical precision

Signal-to-noise ratio



- QCD functional integrals done by importance sampling: propagators
- Variance in correlator determined by

$$\sigma^2 \langle C \rangle = \langle C C^\dagger \rangle - |\langle C \rangle|^2$$

- For nucleon:

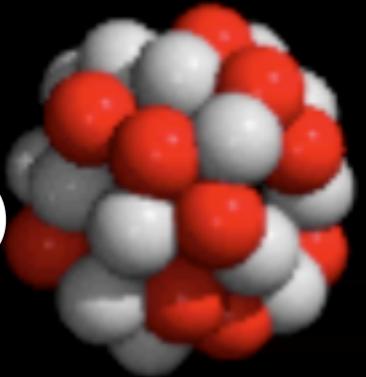
$$s/n \sim \exp[-(M_N - 3/2m_\pi)t]$$

- For nucleus A:

$$s/n \sim \exp[-A(M_N - 3/2m_\pi)t]$$



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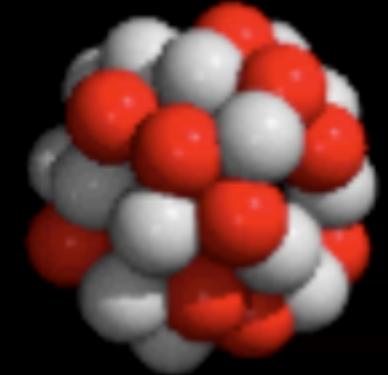
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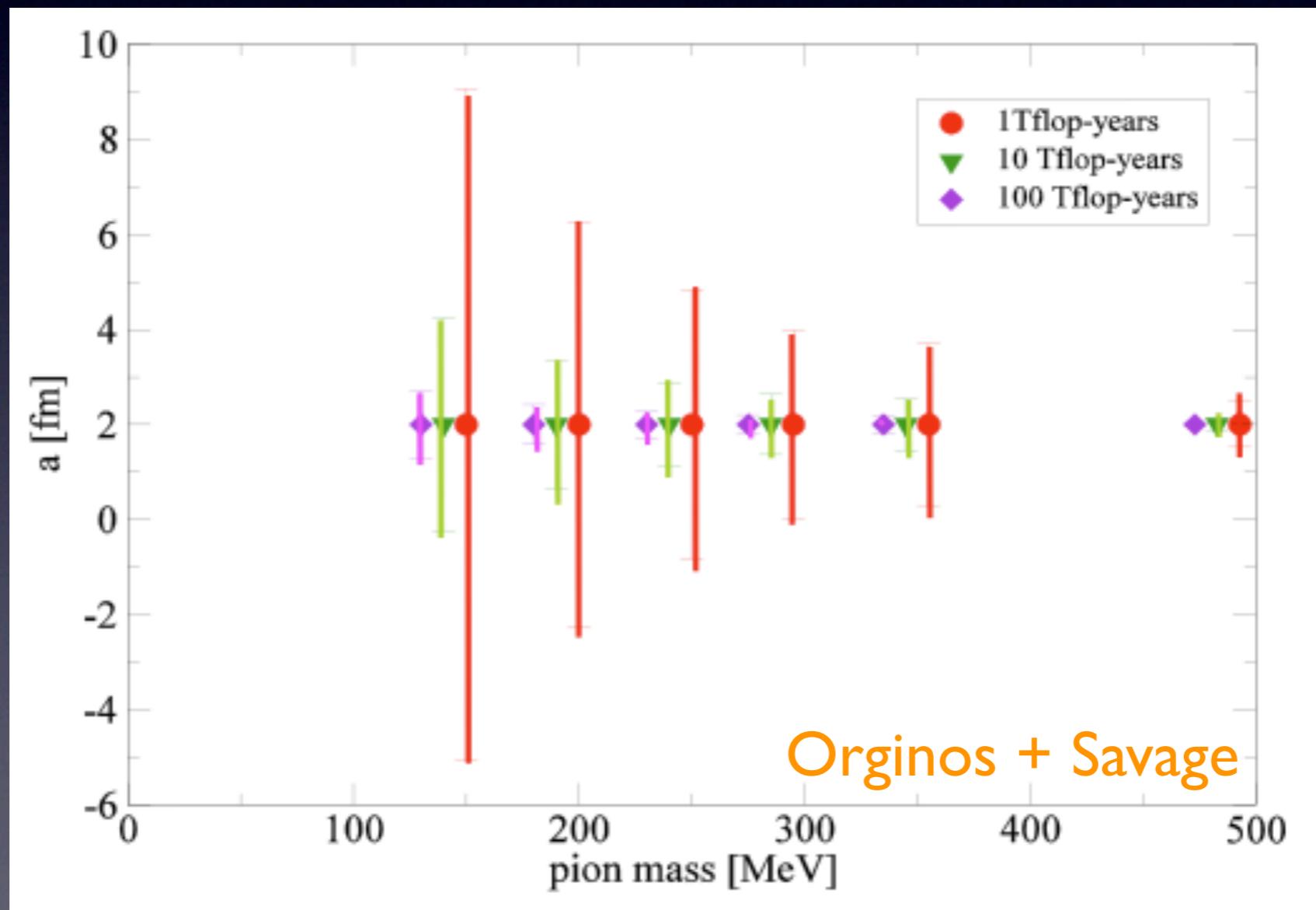
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Signal-to-noise



- Exponentially bad problem
 - ex: NN scattering length



Phys Rev D76:074507, 2007

Phys. Rev. Lett. 100:082004, 2008

Phys Rev D77:057502, 2008

Phys Rev D78:014507, 2008

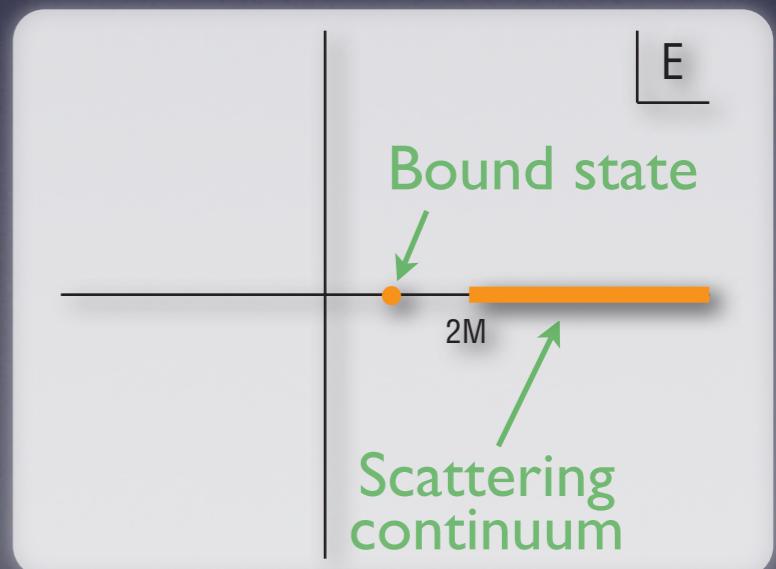
arXiv:0807.1856

+ ..

Multi-meson systems: Poor man's nuclear physics

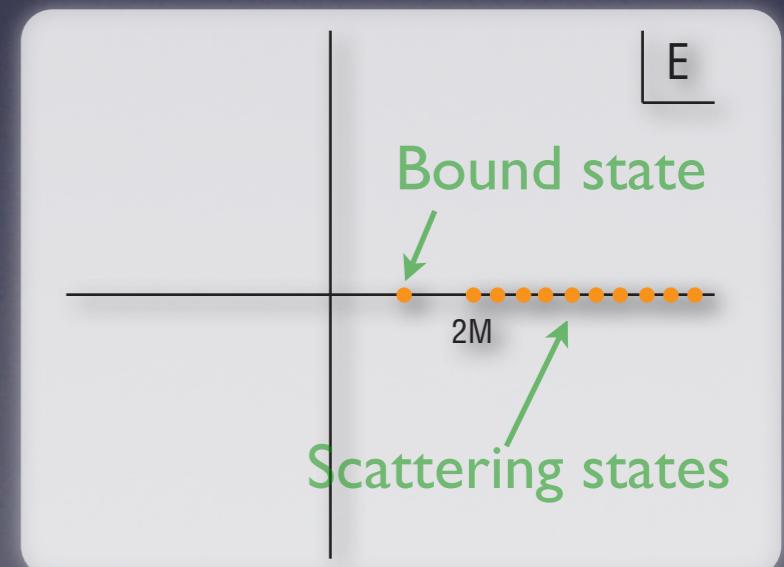
Hadron scattering

- Maiani-Testa: extracting multi-hadron *S-matrix elements from Euclidean lattice calculations of corresponding Green functions is impossible*
- Lüscher: volume dependence of two-particle energy levels \Rightarrow scattering phase-shift up to inelastic threshold



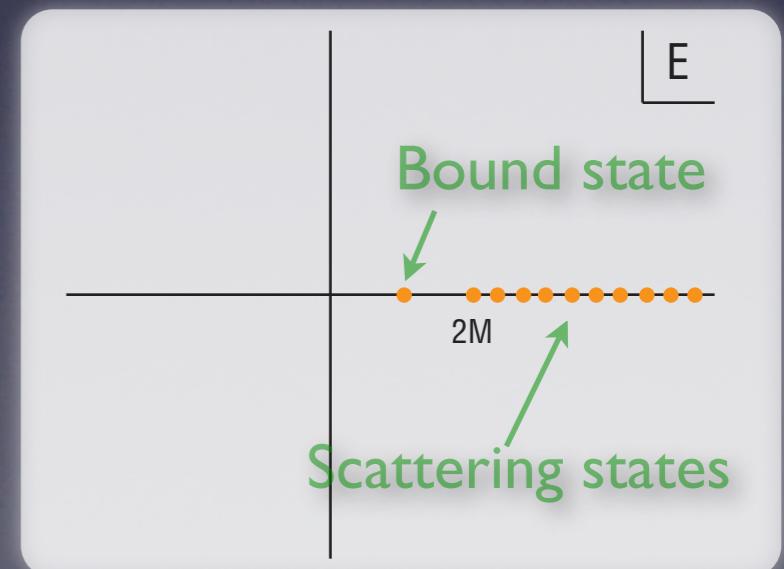
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Hadron scattering

- Maiani-Testa: extracting multi-hadron S-matrix elements from Euclidean lattice calculations of corresponding Green functions is impossible
- Lüscher: volume dependence of two-particle energy levels \Rightarrow scattering phase-shift up to inelastic threshold
- Exact relation provided $r \ll L$
- Used for $\pi\pi$, KK , NN , ΛN





Multi-boson energies

[Beane,WD & Savage; WD & Savage]

- Large volume expansion of GS energy of n meson system to l/L^7
 - 2 & 3 body interactions (N body: $L^{-3(N-l)}$)
 - Relativistic up to particle production
 - $n=2$: reproduces expansion of Lüscher
- Can include higher PW, higher body, excited states, fermions

Multi-boson energies

[WD+Savage arXiv:0801.0763]

$$\begin{aligned} \Delta E_n = & \frac{4\pi \bar{a}}{M L^3} {}^n C_2 \left\{ 1 - \left(\frac{\bar{a}}{\pi L} \right) \mathcal{I} + \left(\frac{\bar{a}}{\pi L} \right)^2 [\mathcal{I}^2 + (2n-5) \mathcal{J}] \right. \\ & - \left(\frac{\bar{a}}{\pi L} \right)^3 [\mathcal{I}^3 + (2n-7) \mathcal{I} \mathcal{J} + (5n^2 - 41n + 63) \mathcal{K}] \\ & + \left(\frac{\bar{a}}{\pi L} \right)^4 [\mathcal{I}^4 - 6\mathcal{I}^2 \mathcal{J} + (4+n-n^2) \mathcal{J}^2 + 4(27-15n+n^2) \mathcal{I} \mathcal{K} \right. \\ & \left. \left. + (14n^3 - 227n^2 + 919n - 1043) \mathcal{L} + 16(n-2)(T_0 + nT_1) \right] \right\} \\ & + {}^n C_3 \frac{1}{L^6} \hat{\eta}_3^L + {}^n C_3 \frac{6\pi \bar{a}^3}{M^3 L^7} (n+3) \mathcal{I} + \mathcal{O}(L^{-8}) \end{aligned}$$

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$\mathcal{I}, \mathcal{J}, \dots$
geometric
constants

Many mesons in LQCD

- Consider $n \pi^+$ correlator ($m_u = m_d$)

$$C_n(t) = \left\langle 0 \left| \left[\sum_{\mathbf{x}} \bar{d} \gamma_5 u(\mathbf{x}, t) \bar{u} \gamma_5 d(0, 0) \right]^n \right| 0 \right\rangle$$
$$\rightarrow A e^{-E_n t}$$

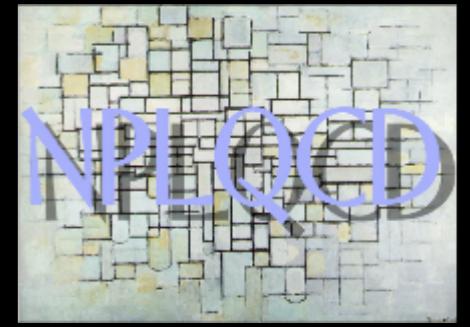
- $n!^2$ Wick contractions: $(13!)^2 \sim 10^{19}$

$$C_3(t) = \text{tr} [\Pi]^3 - 3 \text{ tr} [\Pi] \text{ tr} [\Pi^2] + 2 \text{ tr} [\Pi^3]$$

$$\Pi = \sum_{\mathbf{x}} \gamma_5 S(\mathbf{x}, t; 0) \gamma_5 S^\dagger(\mathbf{x}, t; 0)$$



- π^+ contractions: only a single quark propagator

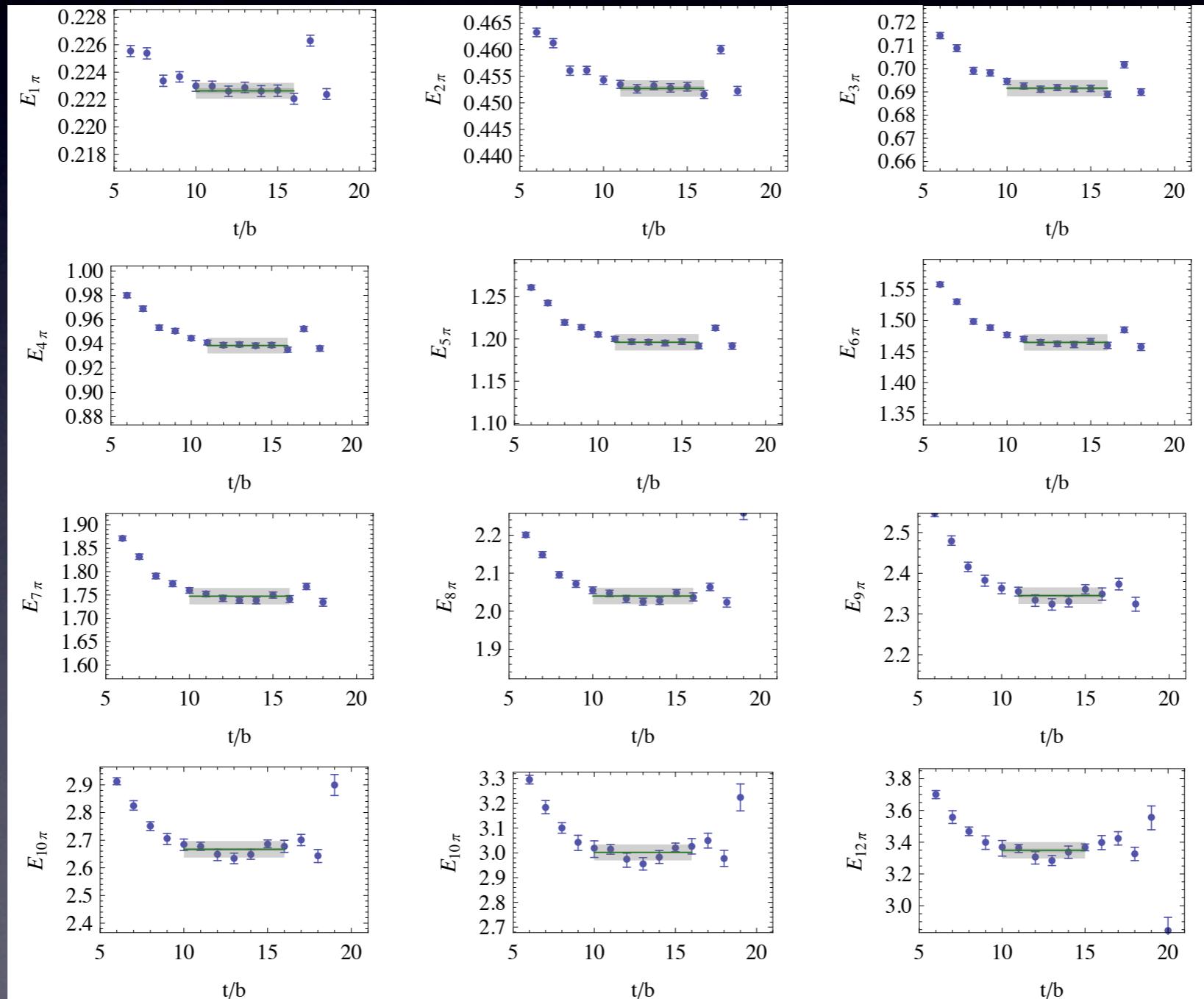


Lattice details

- Calculations use MILC gauge configurations
 - $L=2.5$ fm, $a=0.12$ fm, *rooted staggered*
 - also $L=3.5$ fm and $a=0.09$ fm
- NPLQCD: domain-wall quark propagators
 - $m_\pi \sim 291, 318, 352, 358, 491, 591$ MeV
 - 24 propagators / lattice in best case
- $l_z=n=1,\dots,12$ pion and ($S=n$) kaon systems

n-meson energies

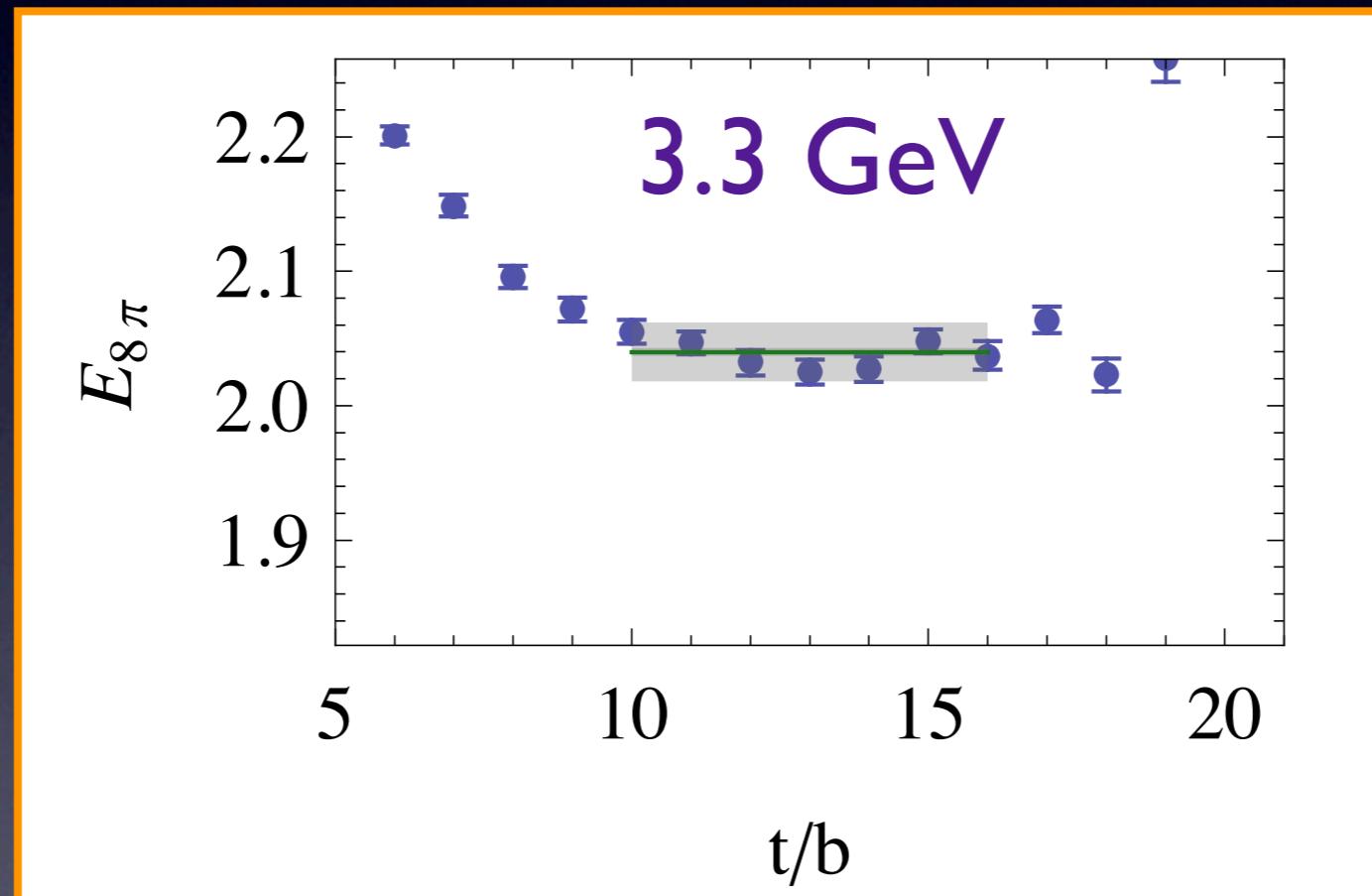
- Effective energy plots: $\log[C_n(t)/C_n(t+l)]$



$m_\pi = 352 \text{ MeV}$

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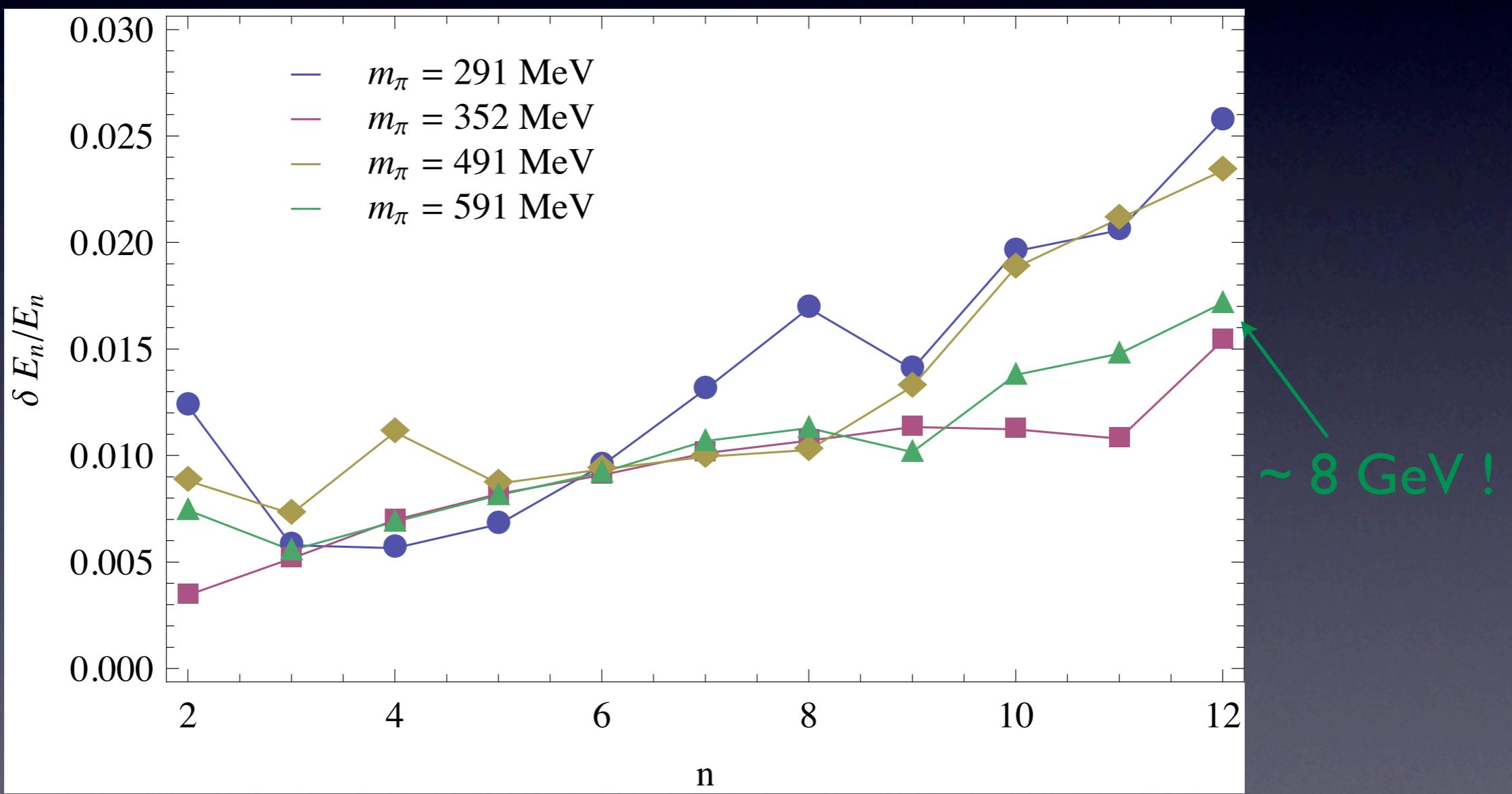
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n-meson energies

- Clean signals for $n=1, \dots, 12$





Multi-boson energies

- Multiple extractions of \bar{a} and $\hat{\eta}_3^L$

$$\begin{aligned}\Delta E_n = & \frac{4\pi \bar{a}}{M L^3} {}^n C_2 \left\{ 1 - \left(\frac{\bar{a}}{\pi L} \right) \mathcal{I} + \left(\frac{\bar{a}}{\pi L} \right)^2 [\mathcal{I}^2 + (2n-5)\mathcal{J}] \right. \\ & - \left(\frac{\bar{a}}{\pi L} \right)^3 [\mathcal{I}^3 + (2n-7)\mathcal{I}\mathcal{J} + (5n^2 - 41n + 63)\mathcal{K}] \\ & + \left(\frac{\bar{a}}{\pi L} \right)^4 [\mathcal{I}^4 - 6\mathcal{I}^2\mathcal{J} + (4+n-n^2)\mathcal{J}^2 + 4(27-15n+n^2)\mathcal{I}\mathcal{K} \\ & \quad \left. + (14n^3 - 227n^2 + 919n - 1043)\mathcal{L} + 16(n-2)(T_0 + nT_1) \right] \Big\} \\ & + {}^n C_3 \frac{1}{L^6} \hat{\eta}_3^L + {}^n C_3 \frac{6\pi \bar{a}^3}{M^3 L^7} (n+3) \mathcal{I} + \mathcal{O}(L^{-8})\end{aligned}$$

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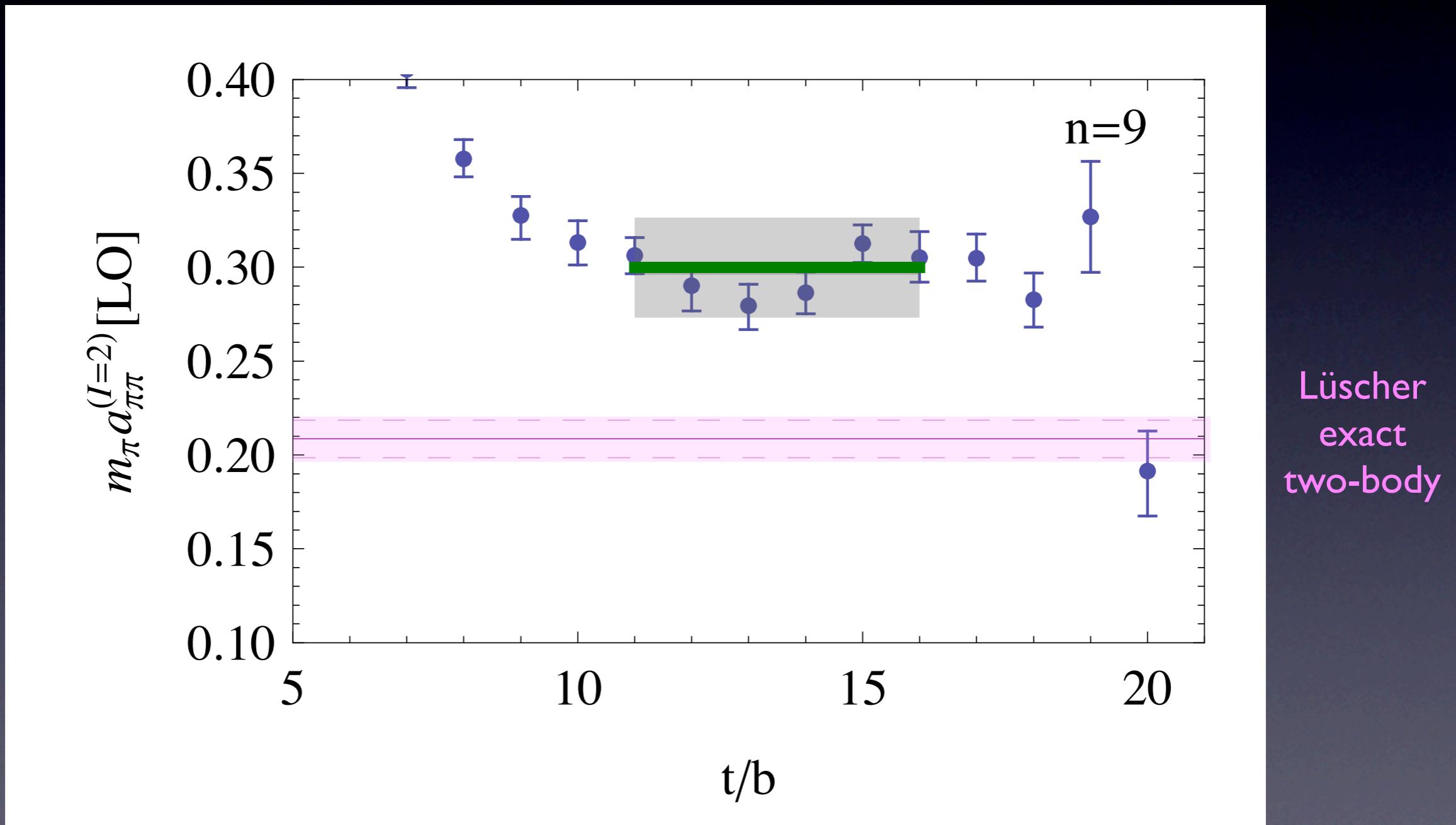
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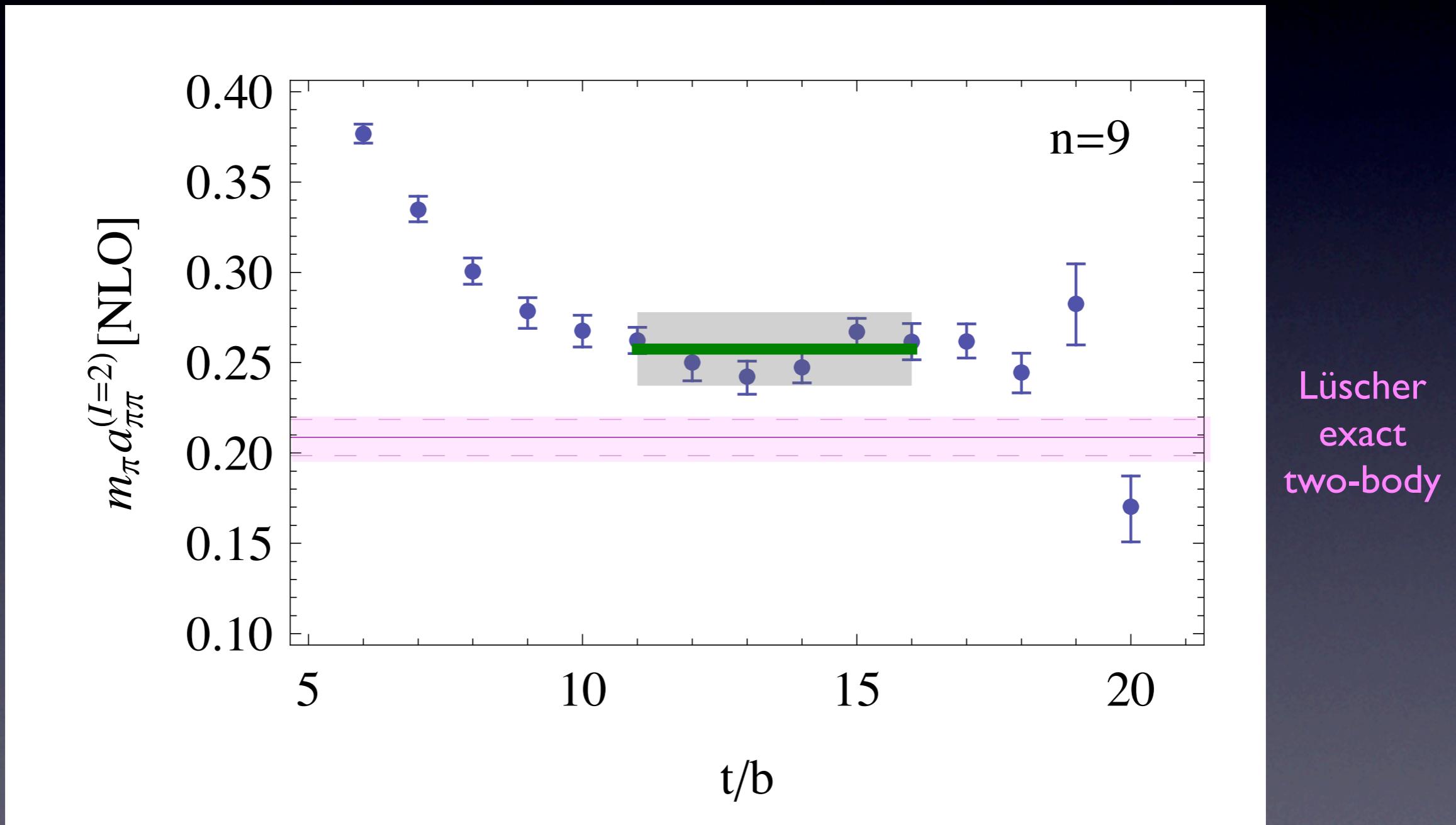
Pion scattering

- Extractions of $m_\pi a$ from four orders in L



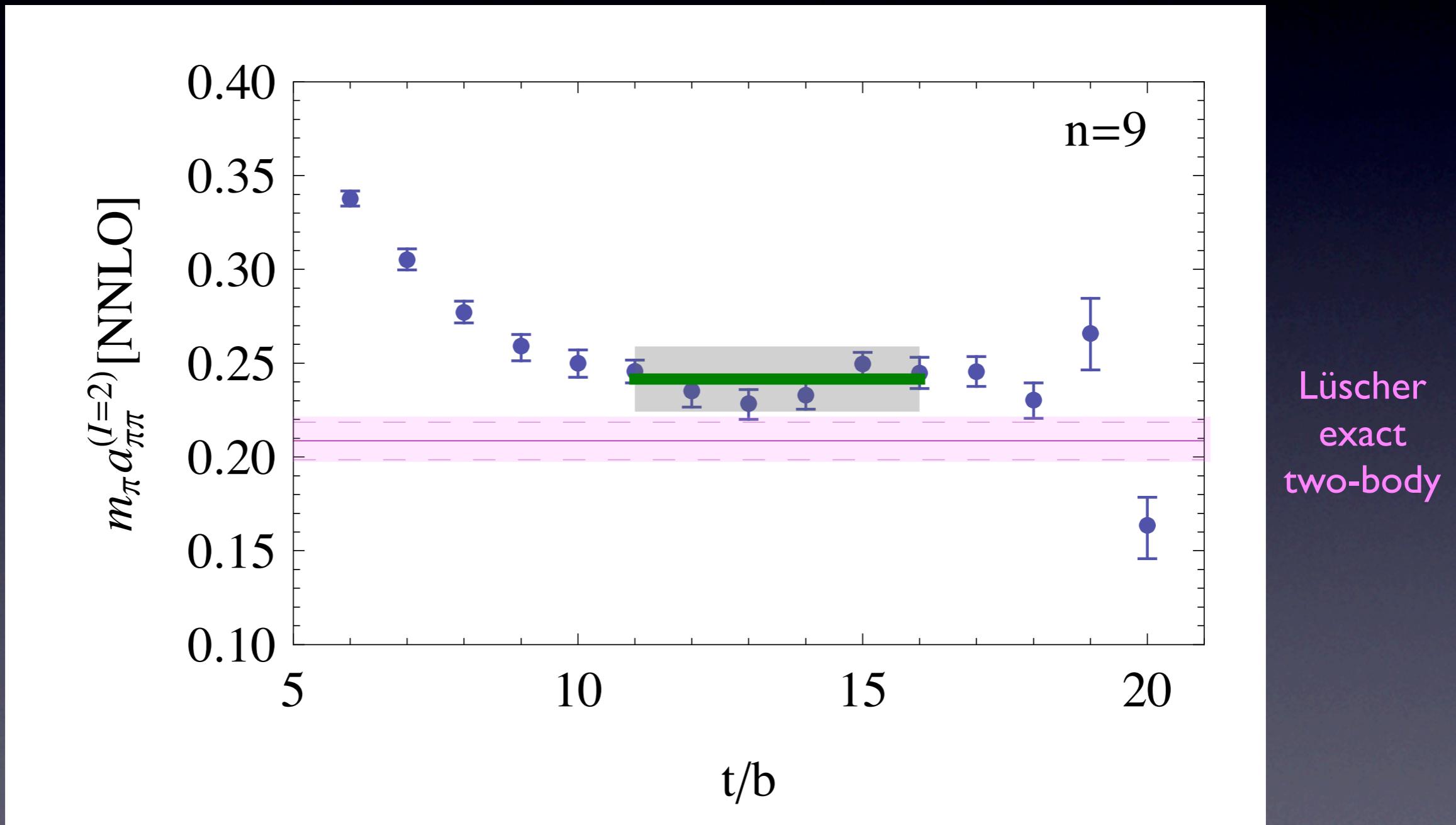
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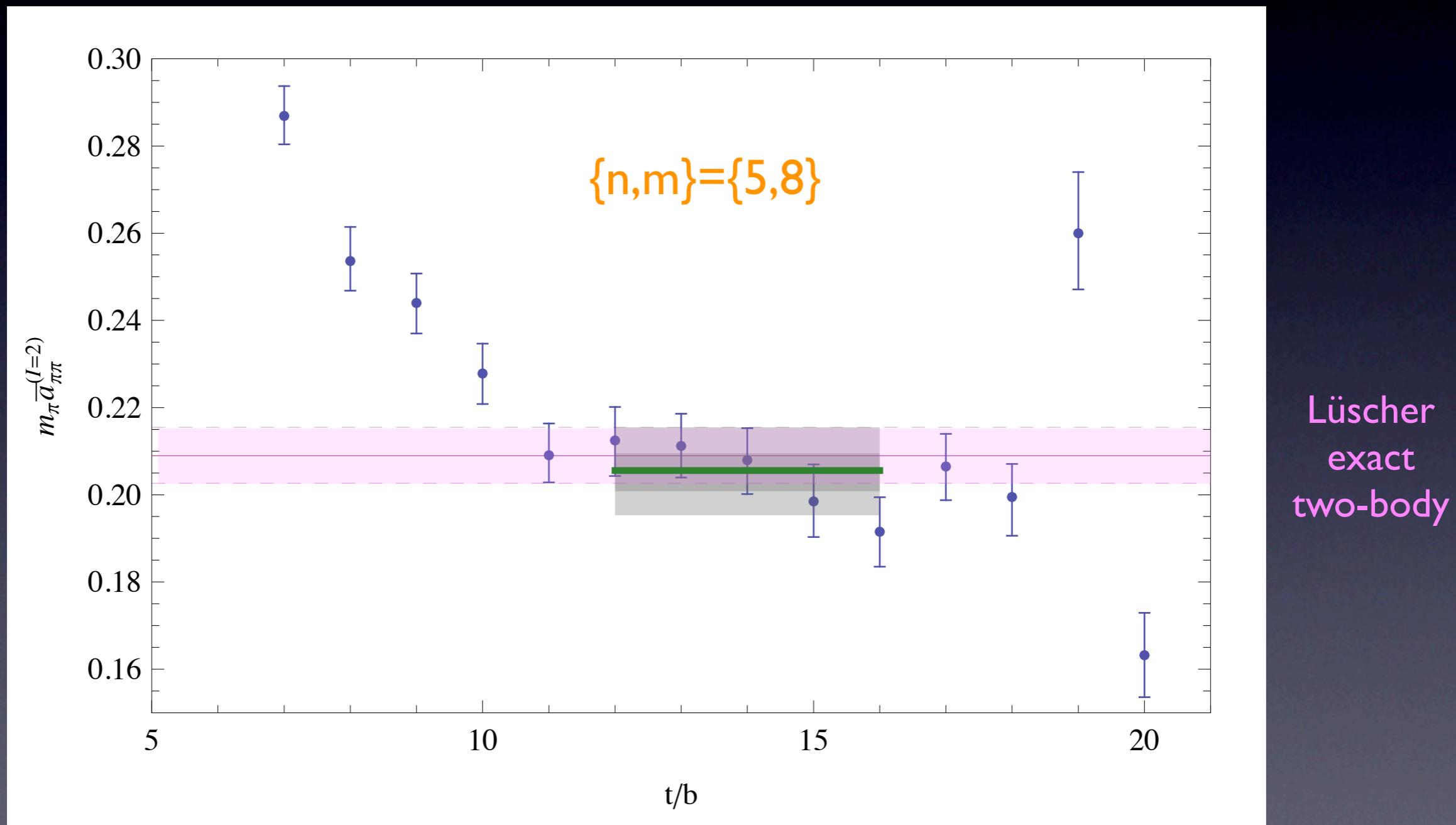
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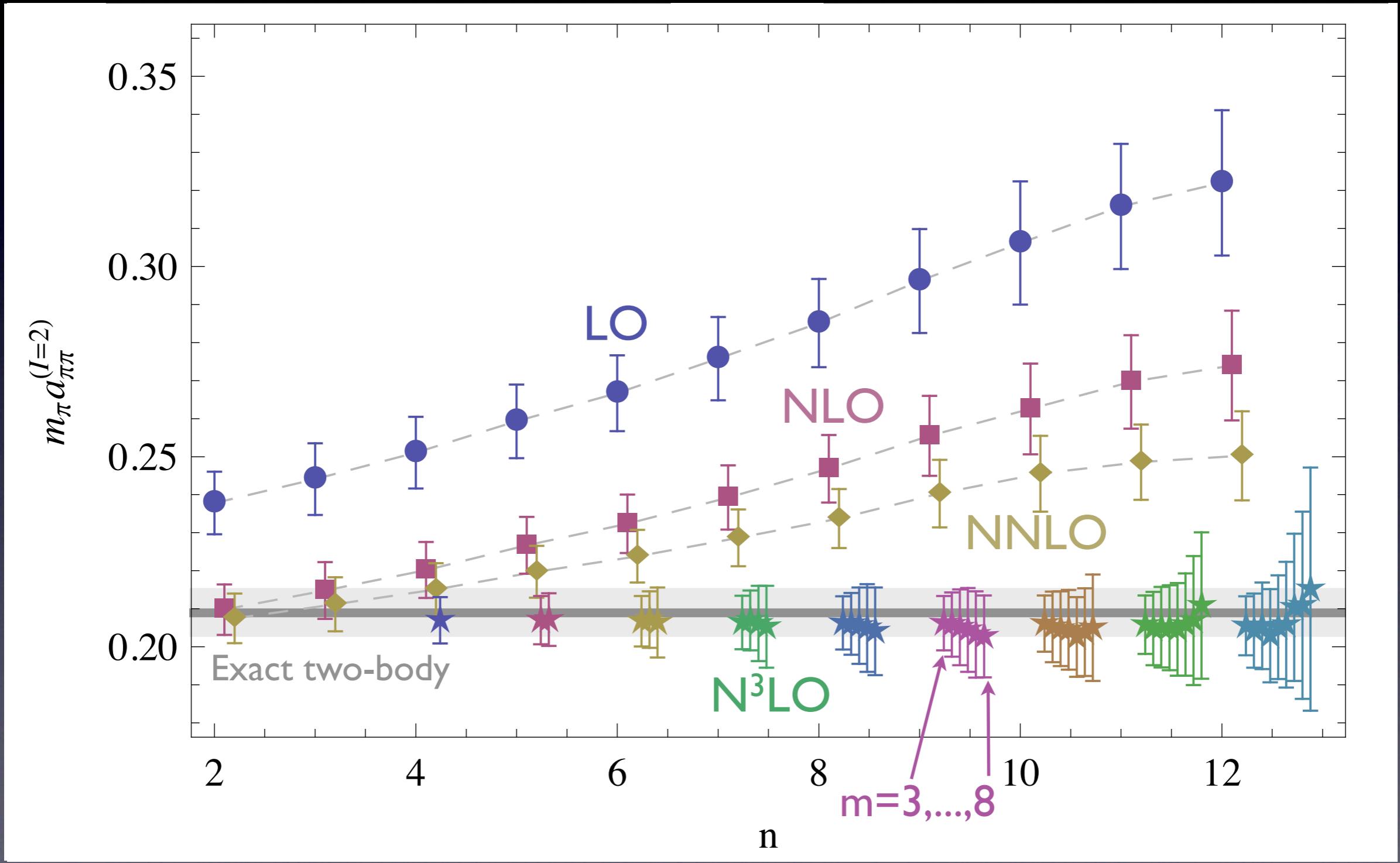


N³LO

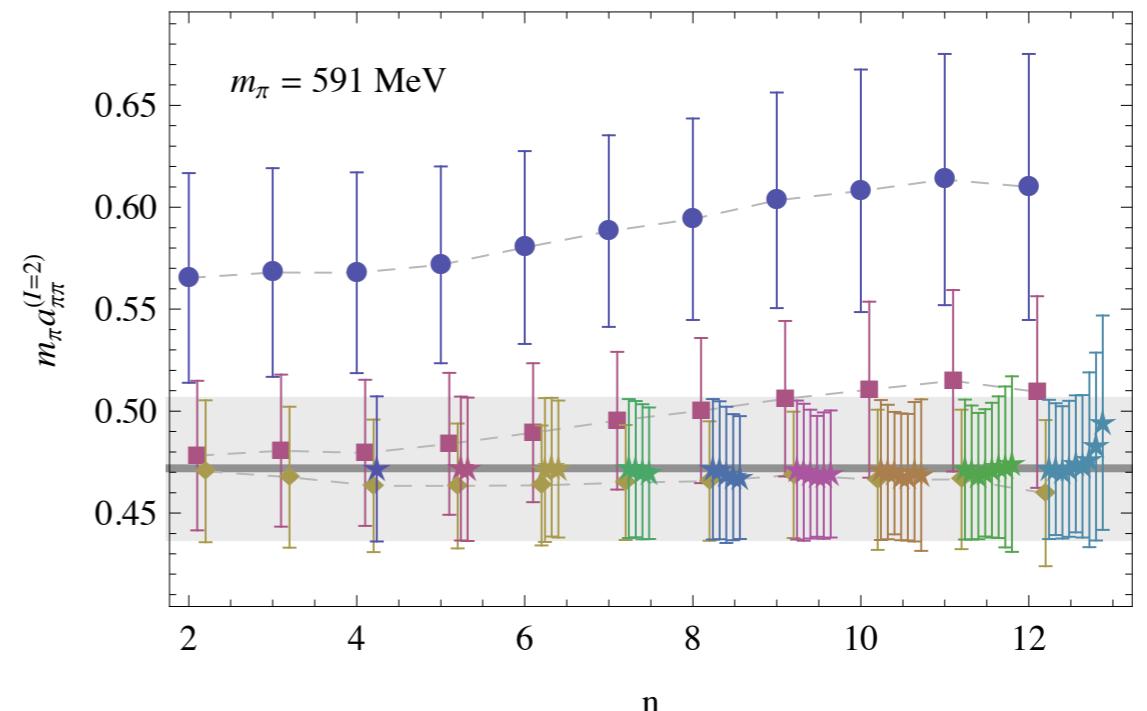
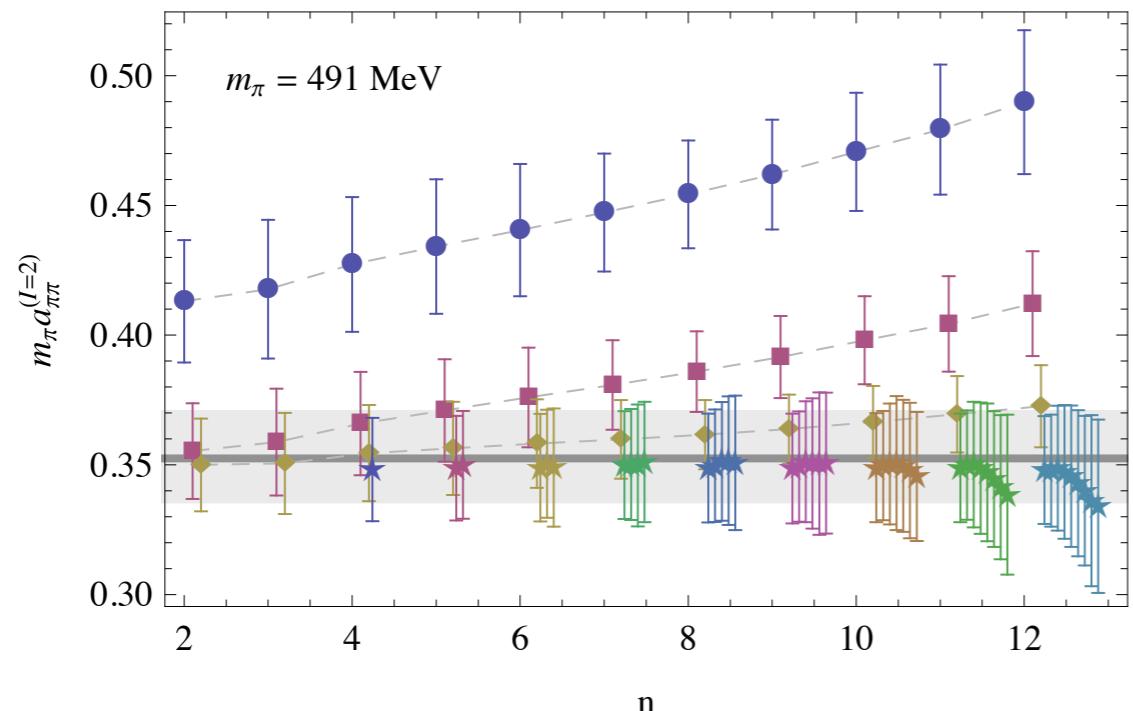
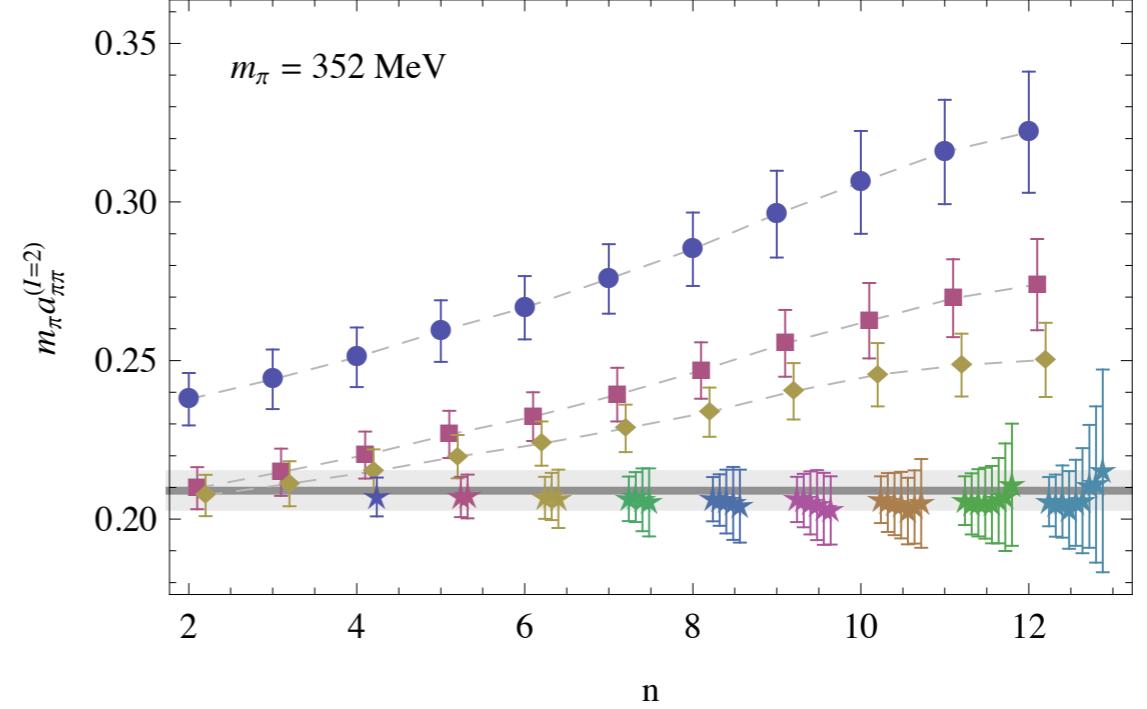
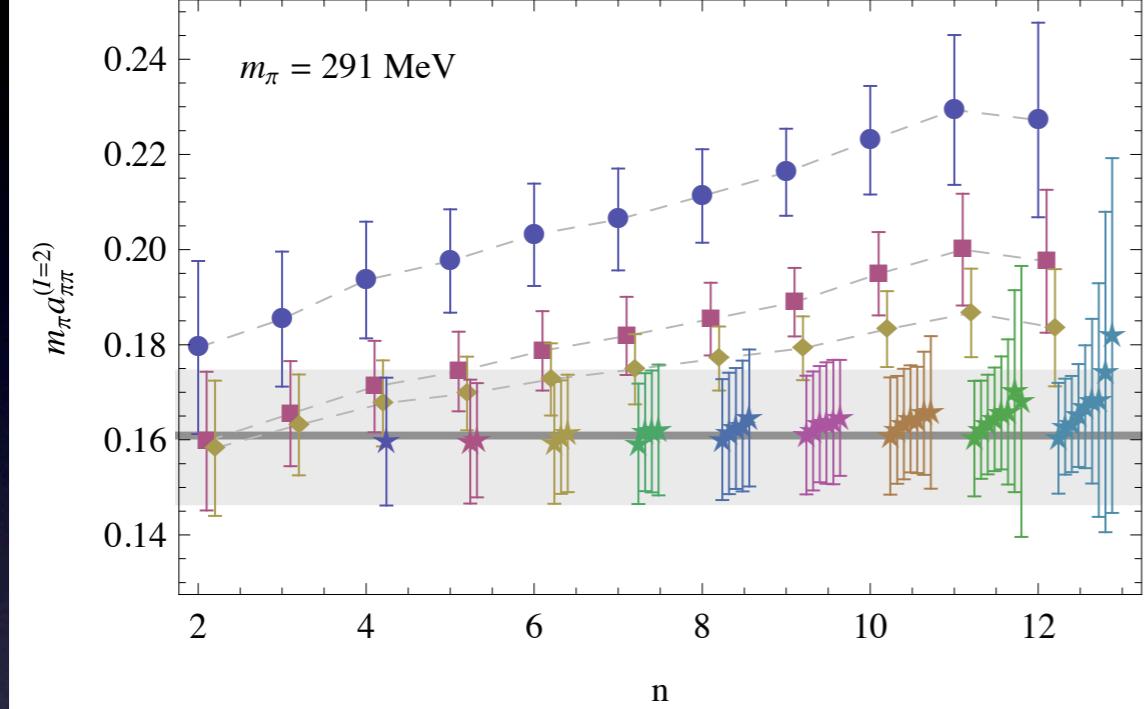
- Two energies to cancel 3-body: 45 combinations



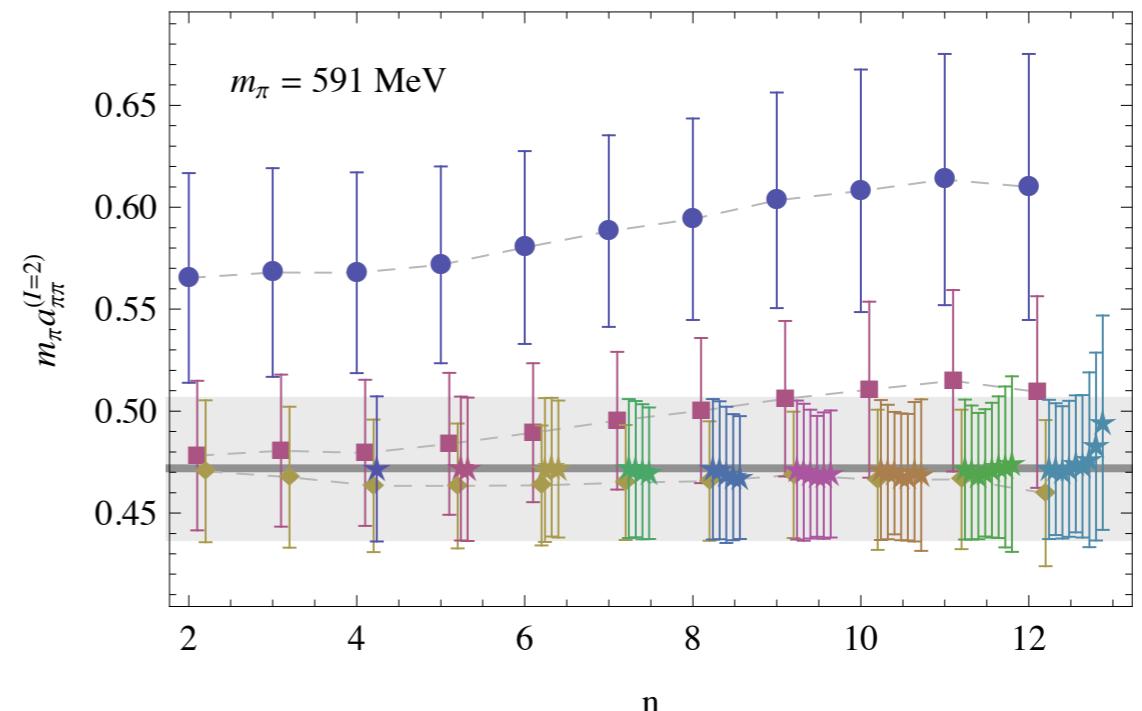
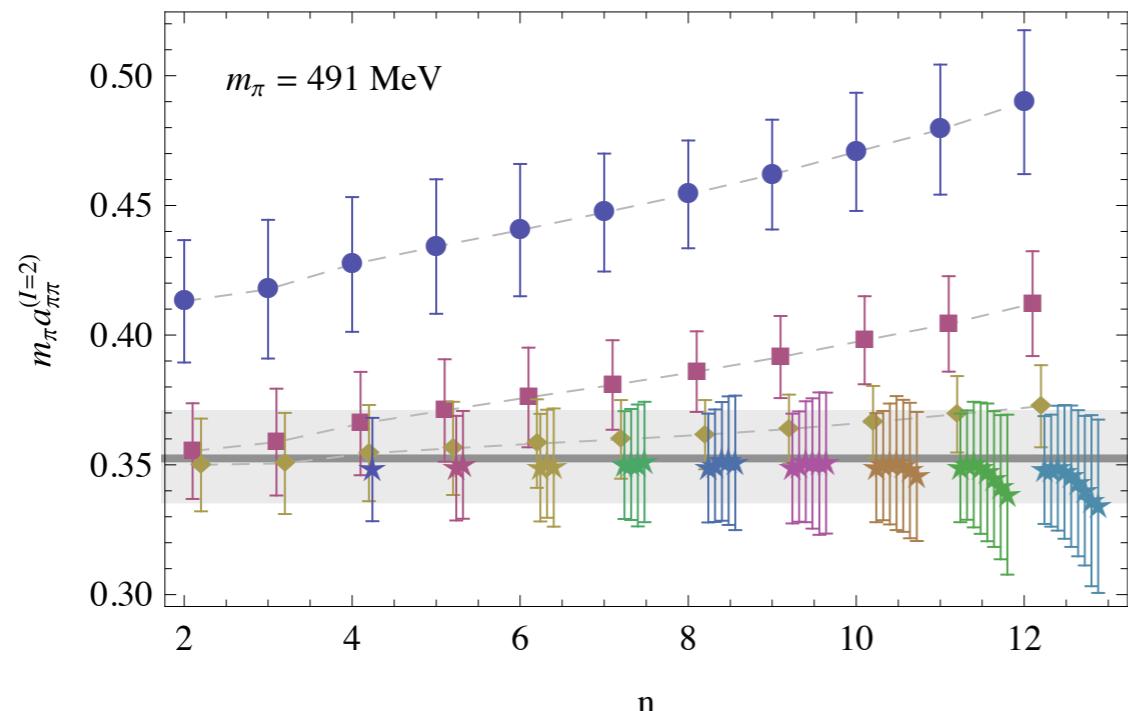
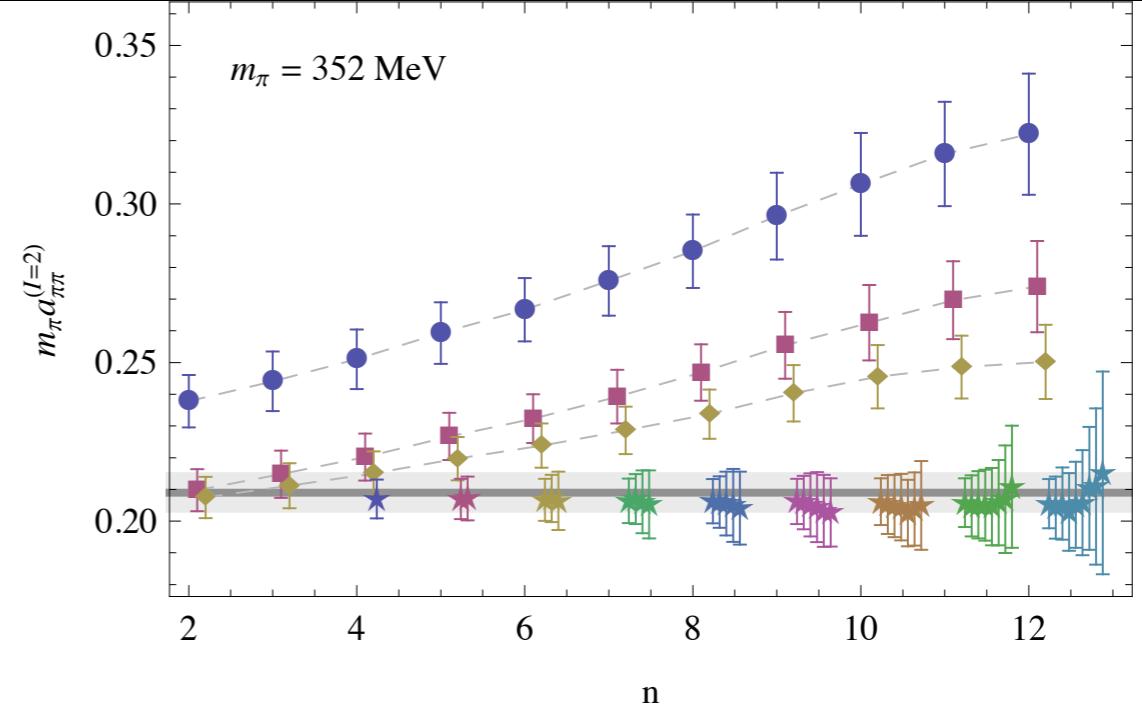
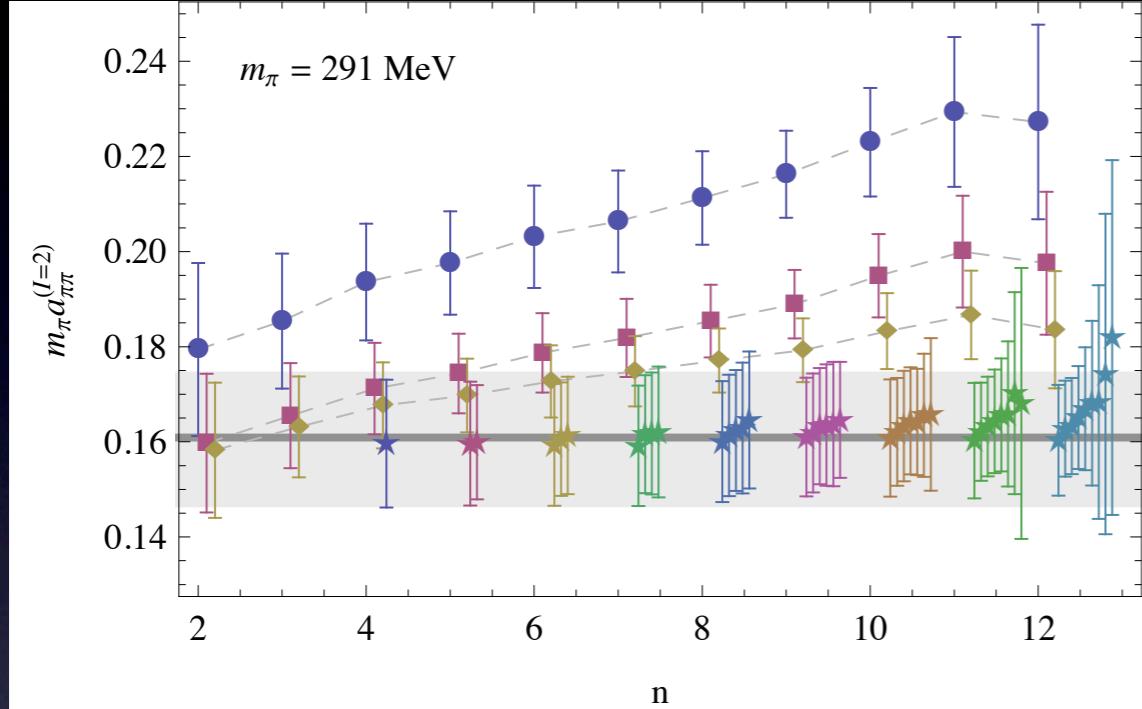
Pion scattering



Pion scattering



Pion scattering



Expansion shows no sign of breakdown?

Scattering lengths

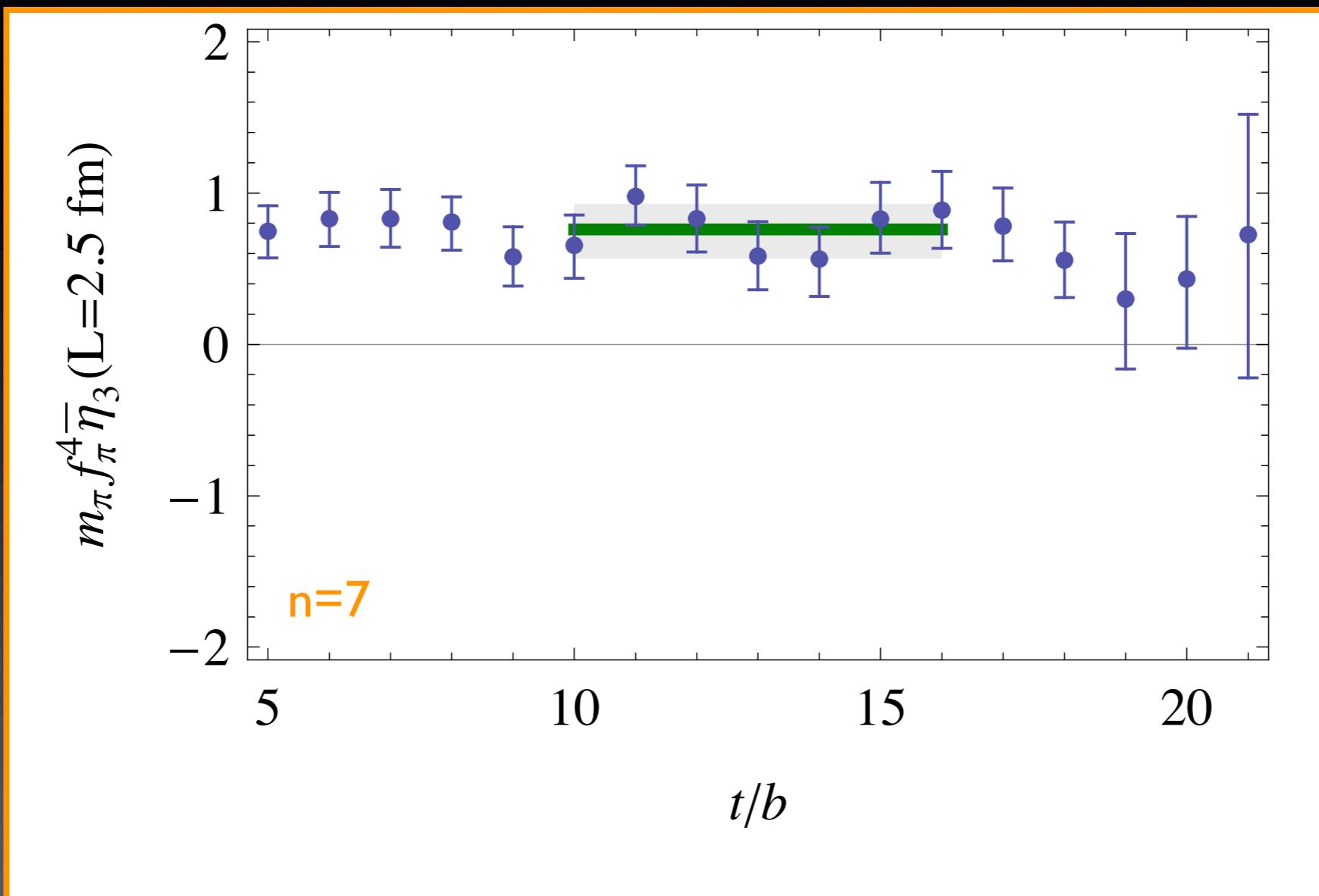
- Scattering lengths equally well extracted for two mesons or ten mesons
- Well described by analytic prediction
- Shows presence of contribution that scales as $\binom{n}{3}$
- varies by two-orders of magnitude

Three meson interactions

- At $1/L^6$, point-like three-boson interaction must occur [Braaten, Nieto '95]
 - RGI 3BI: $\bar{\eta}_3^{(L)}$ physically meaningful
 - Depends logarithmically on L
- Naive dimensional-analysis $m_\pi f_\pi^4 \bar{\eta}_3^{(L)} \sim 1$
- Combinations of energy shifts isolates the RGI interaction



$\pi^+\pi^+\pi^+$ interaction



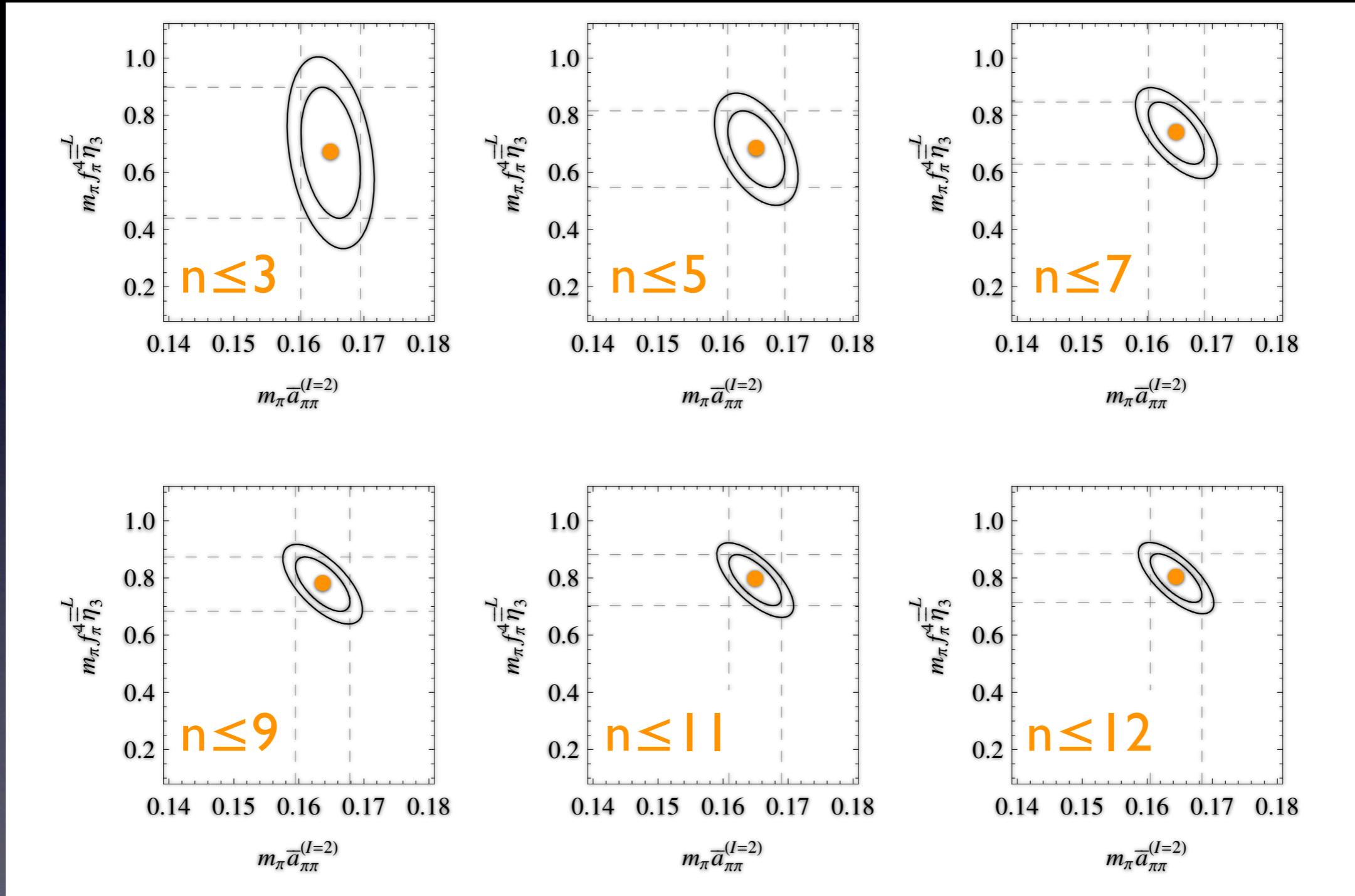
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n correlations

- Fit effective energies to extract two parameters: a , η_3 from l/L expansion
- Use 12 eff. energies in n - t -correlated analysis
 - Large correlation matrix: correlated χ^2
 - Reduces uncertainties as n pion correlators “explore more of the lattice”

Three-body

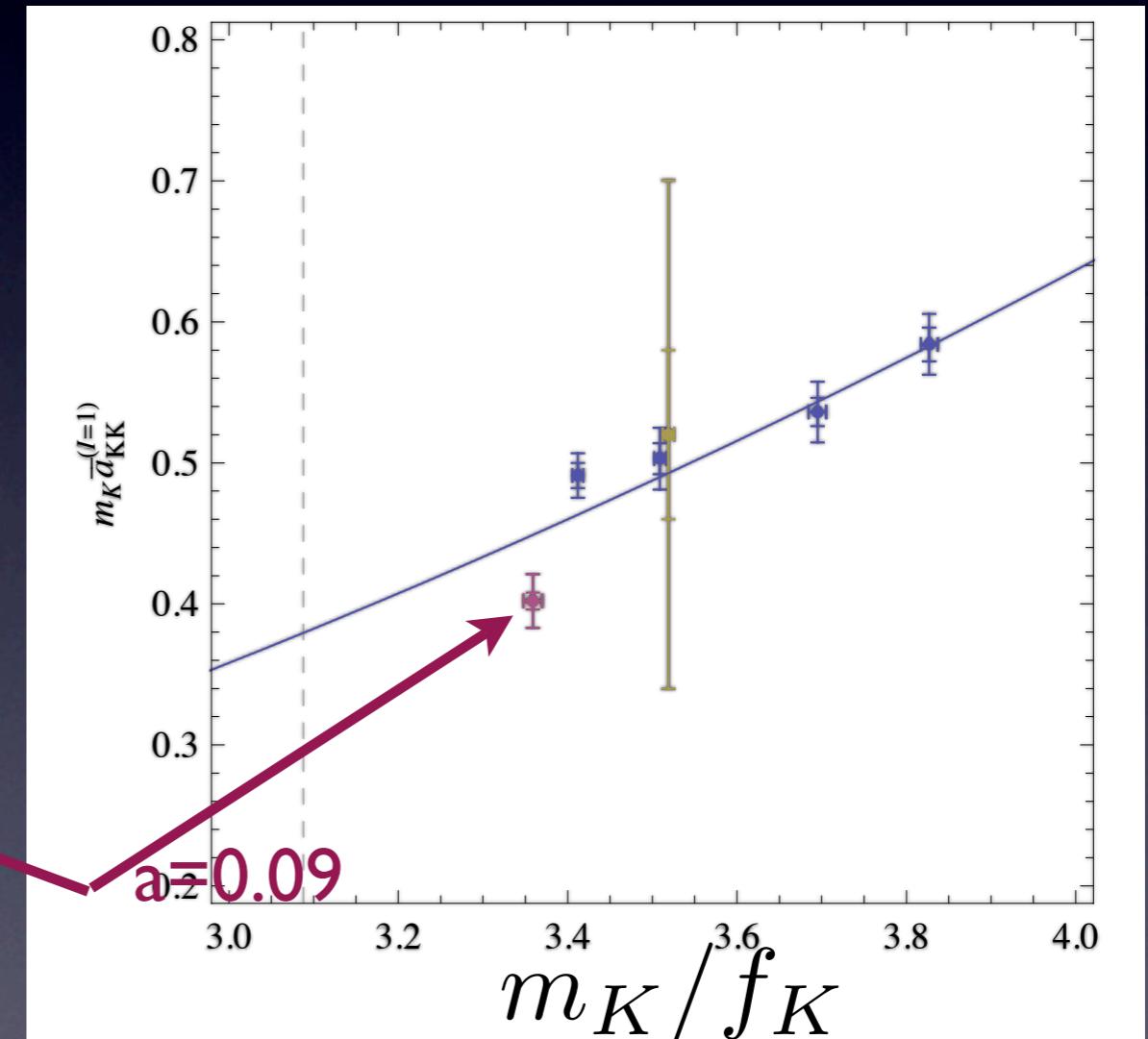
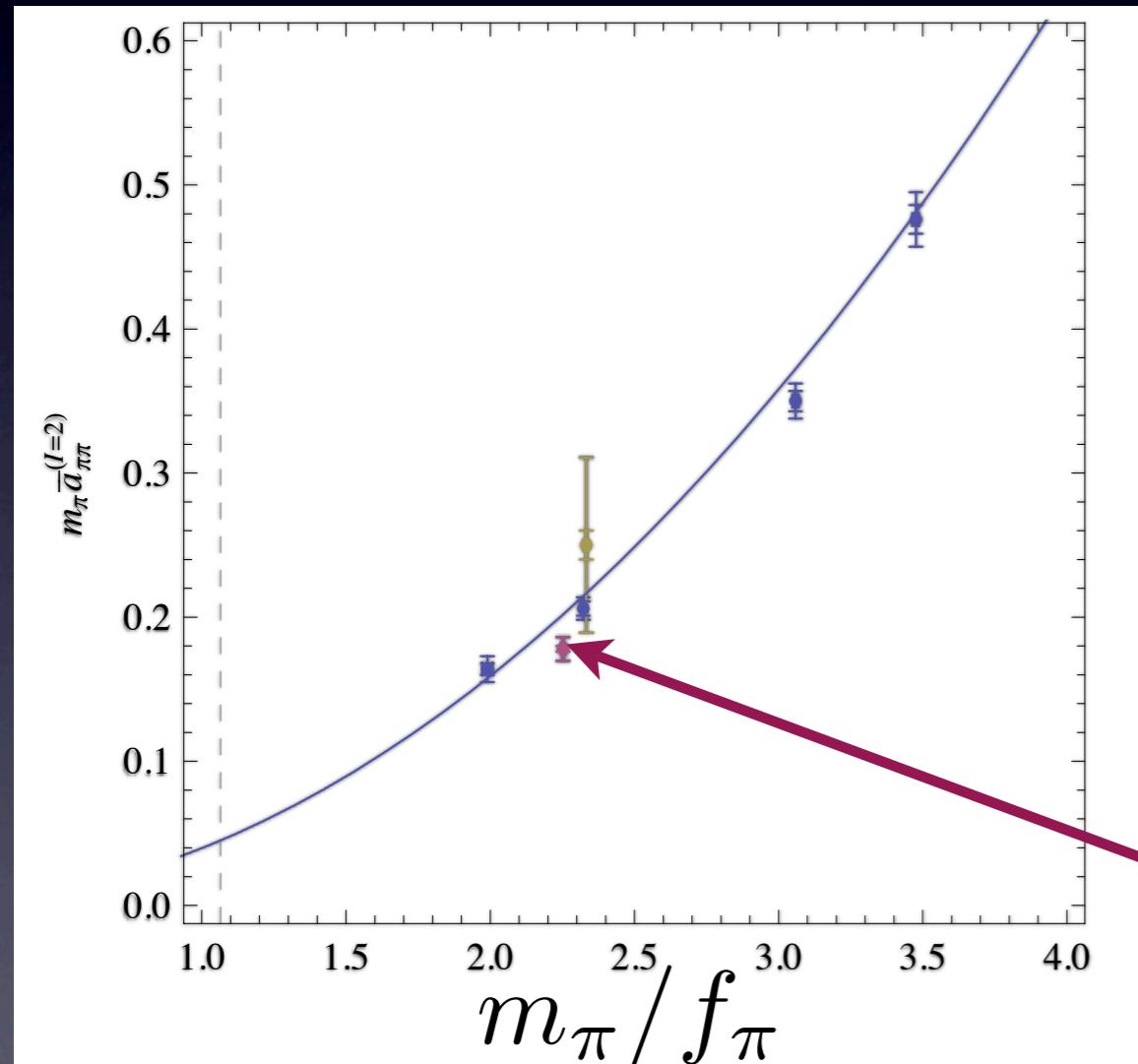
n correlations



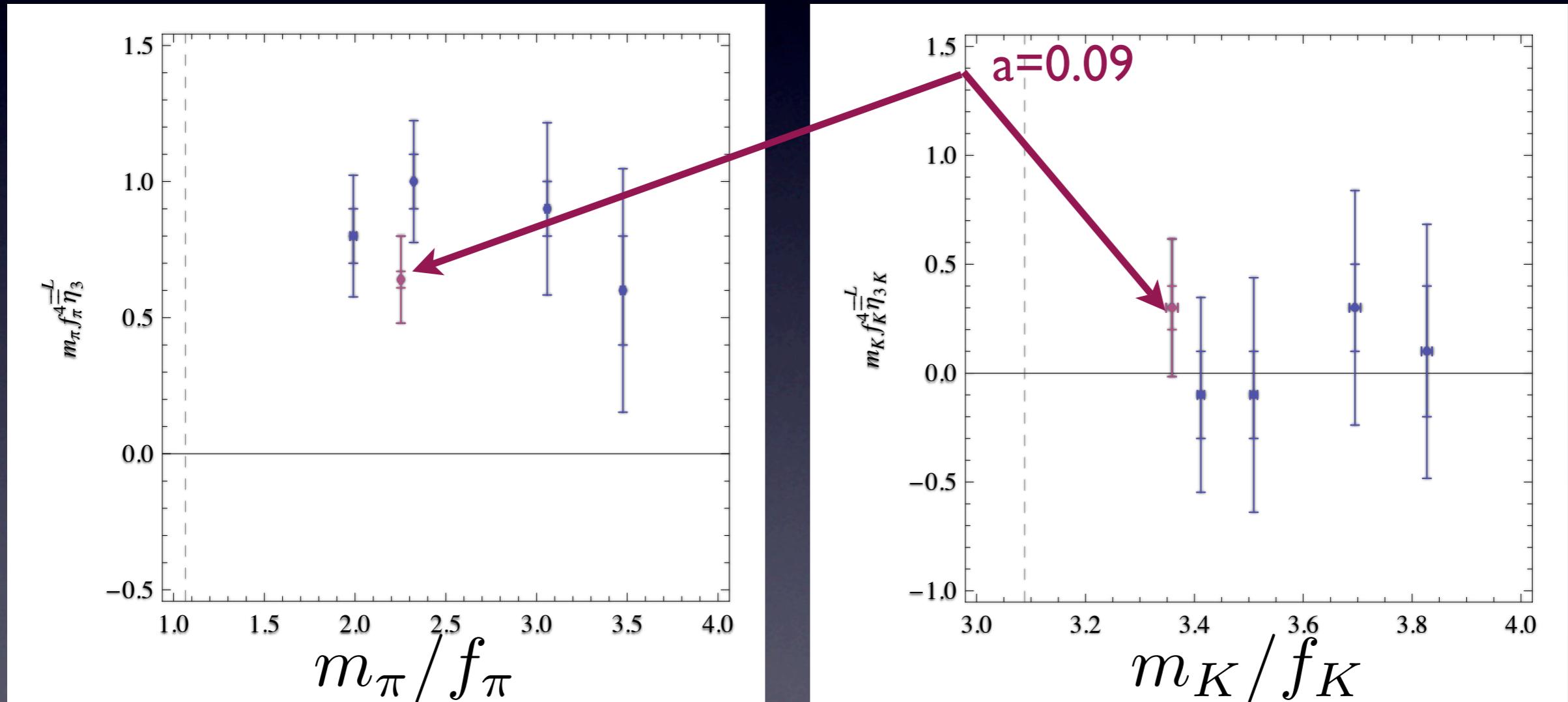
Two body

$2\pi^+$ and $2K^-$ interaction

curves: Weinberg



$3\pi^+$ and $3K^-$ interaction



Naïve dimension analysis: 1

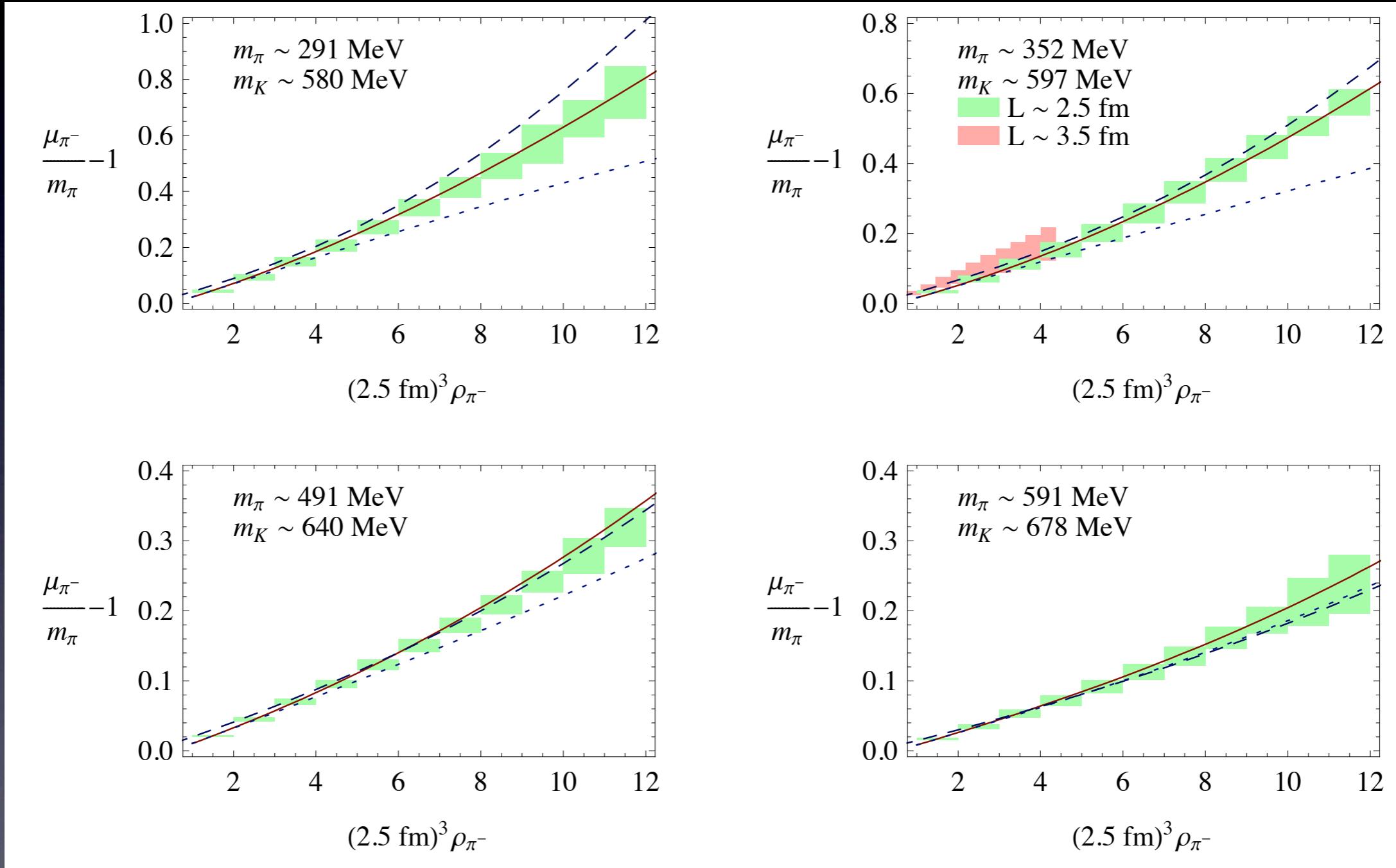
Equation of State

- I/L expansion: analytic form of EOS
- Chemical potential

$$\mu = \left. \frac{d E}{d n} \right|_{V \text{const}}$$

- $\mu(\rho)$ numerically using finite difference
- Compare with LO χ PT [Son & Stephanov]

Isospin Chemical Potential

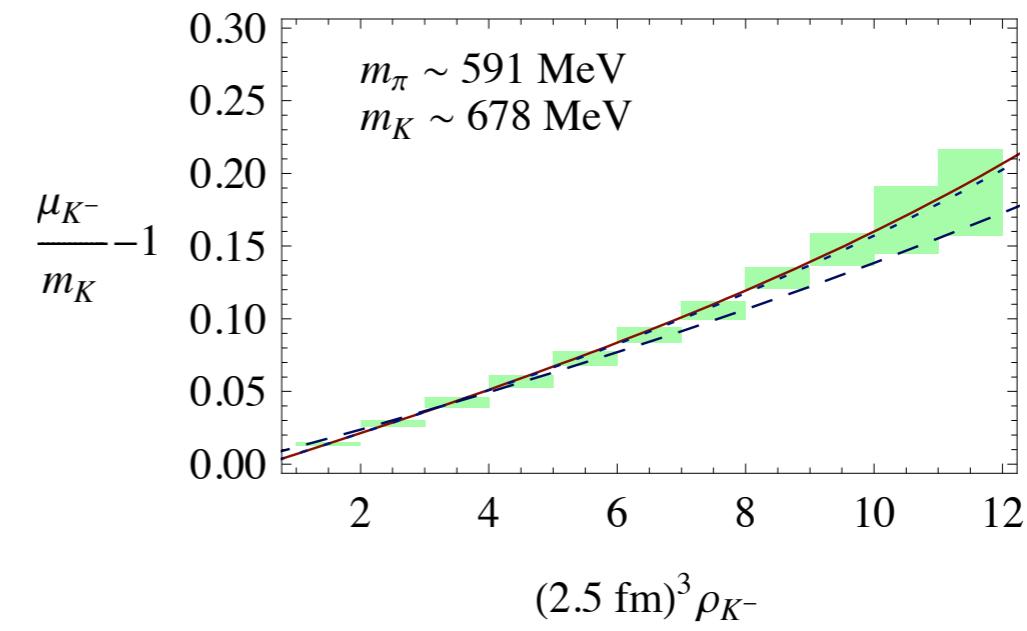
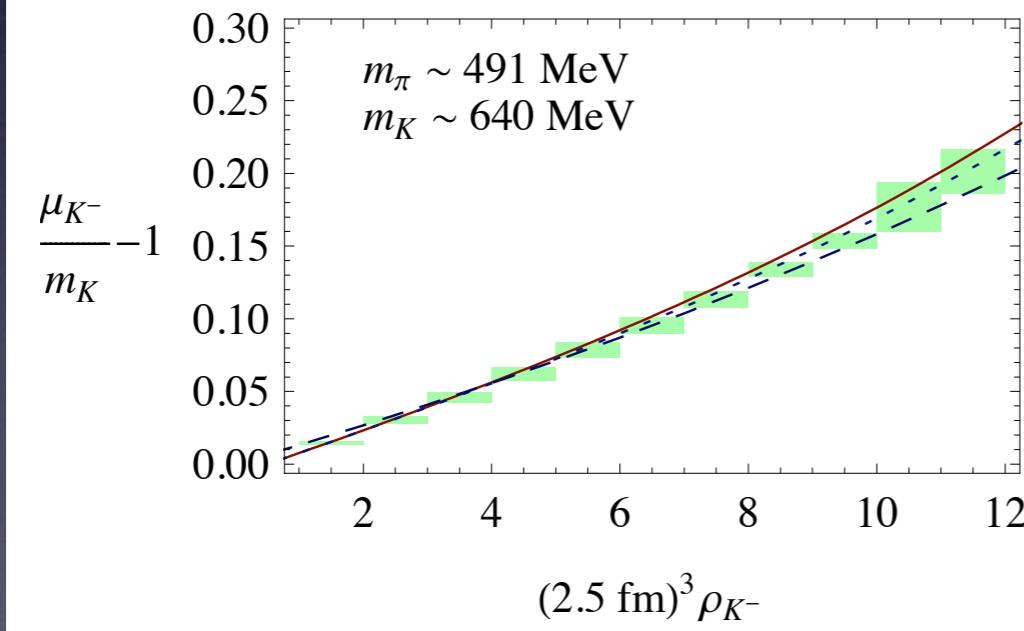
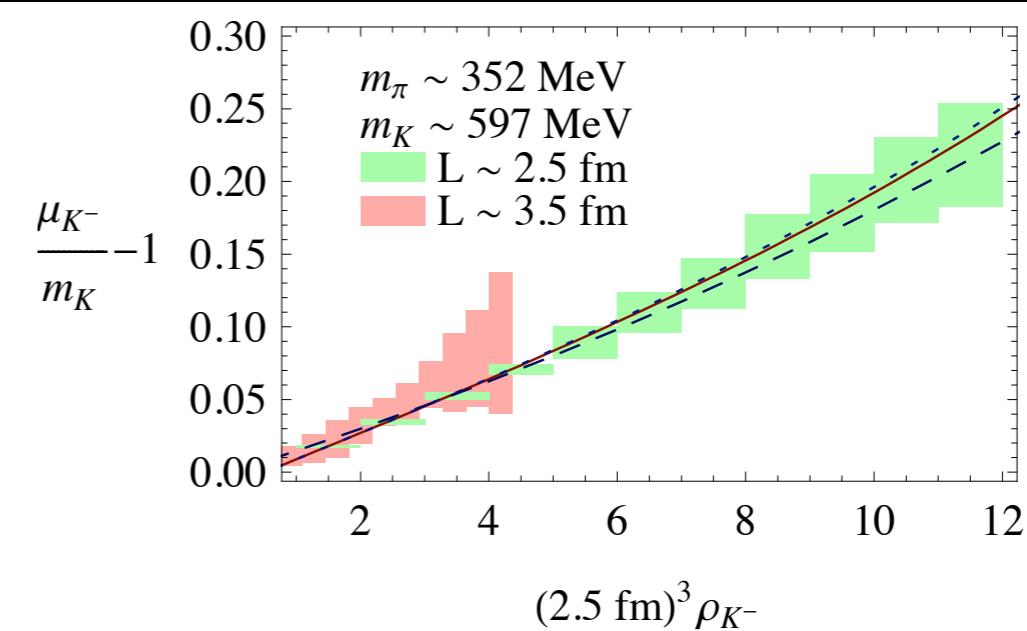
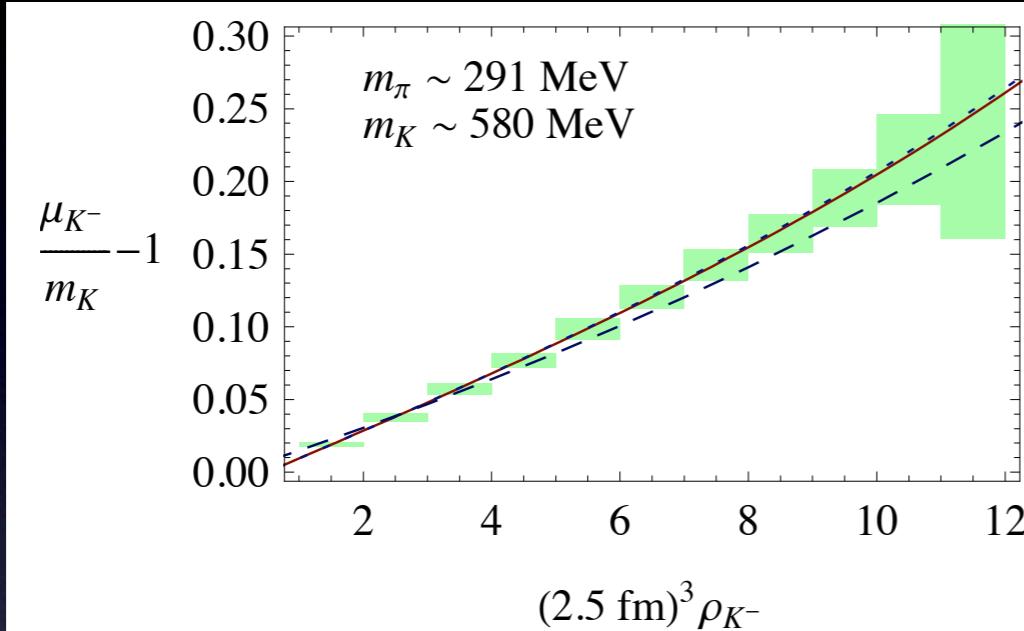


— 2+3 body fit

..... No 3 body

- - - LO χ PT

Kaon Chemical Potential



— 2+3 body fit

..... No 3 body

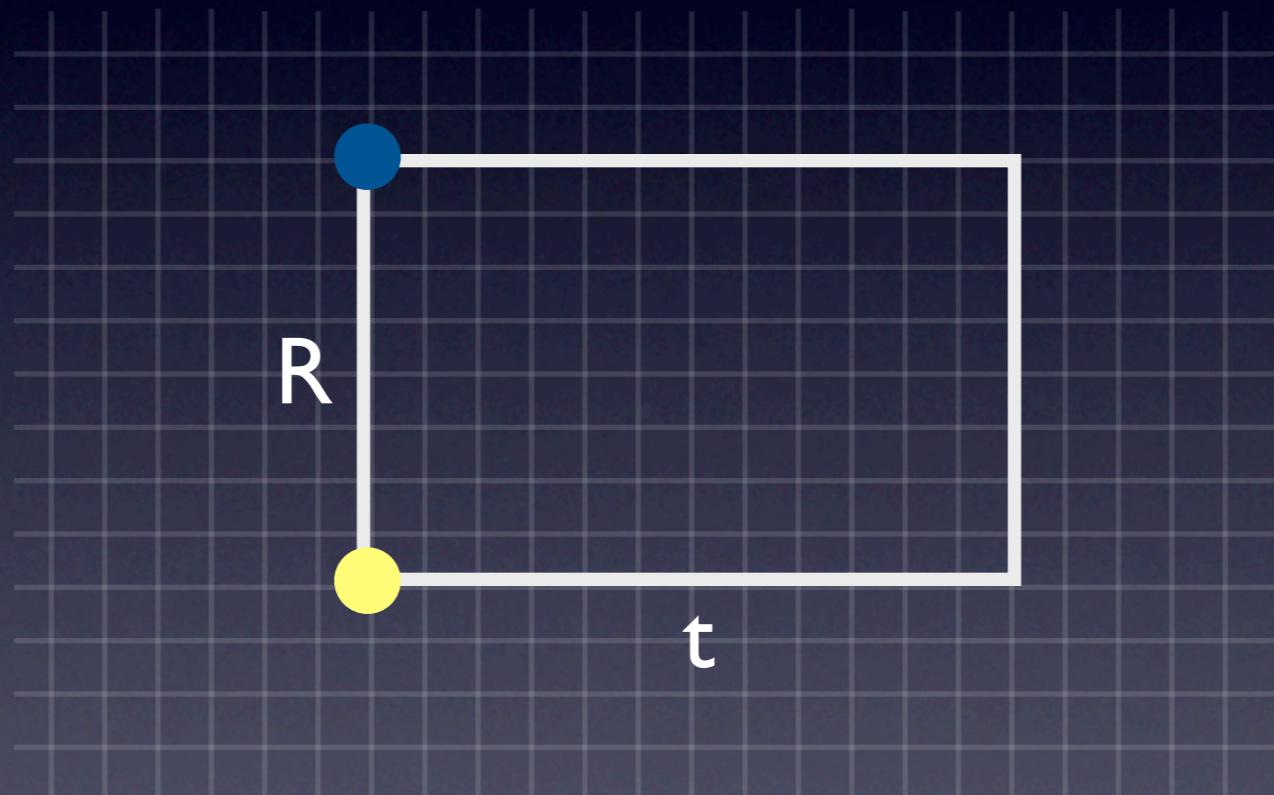
- - - LO χ PT

Chemical potentials

- π^- and K^- chemical potentials
 - Important contribution from $\pi \pi \pi$
 - Results consistent with LO χ PT
 - supports Kaplan/Nelson analysis for n-stars
 - Other μ_l results [de Forcrand et al, Kogut&Sinclair]

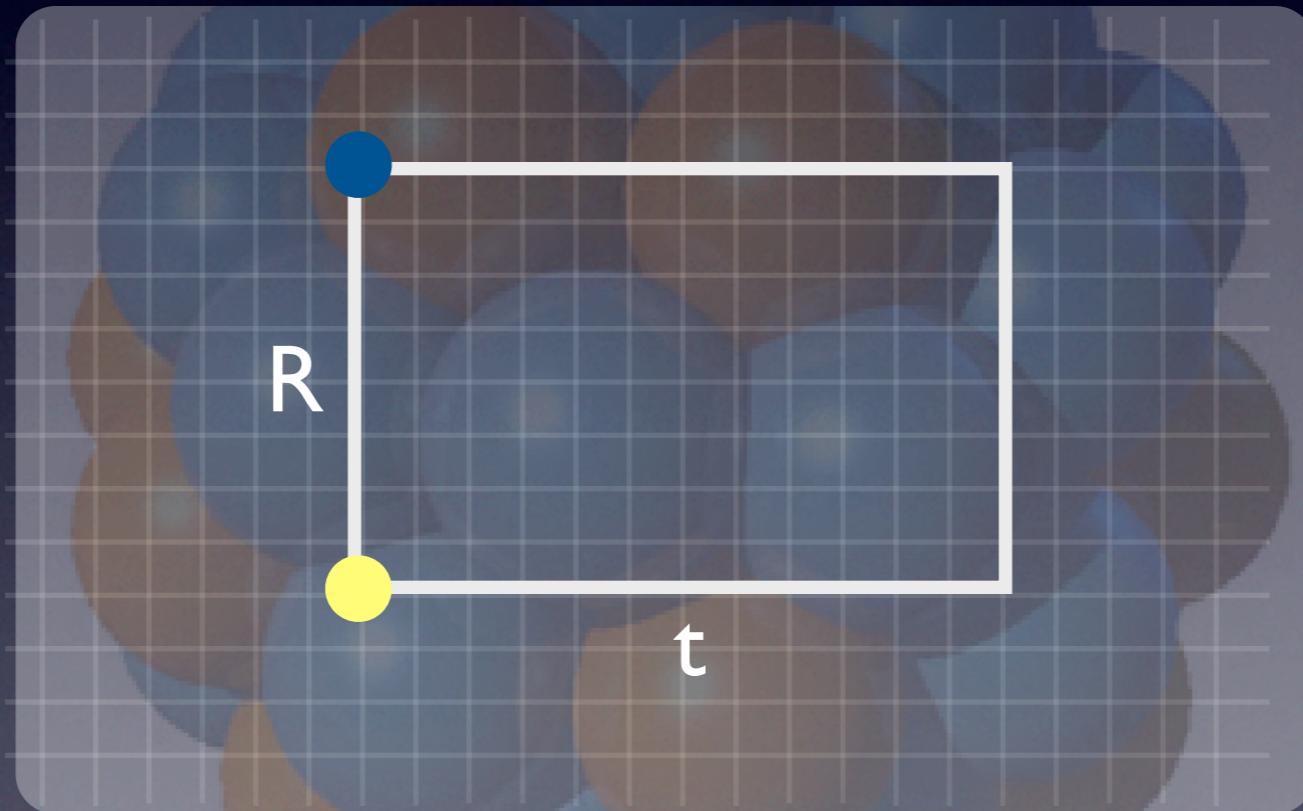
In medium $Q\bar{Q}$ potential

- Static quark potential



In medium $Q\bar{Q}$ potential

- Static quark potential



- Modified by condensate? Hadronic screening?

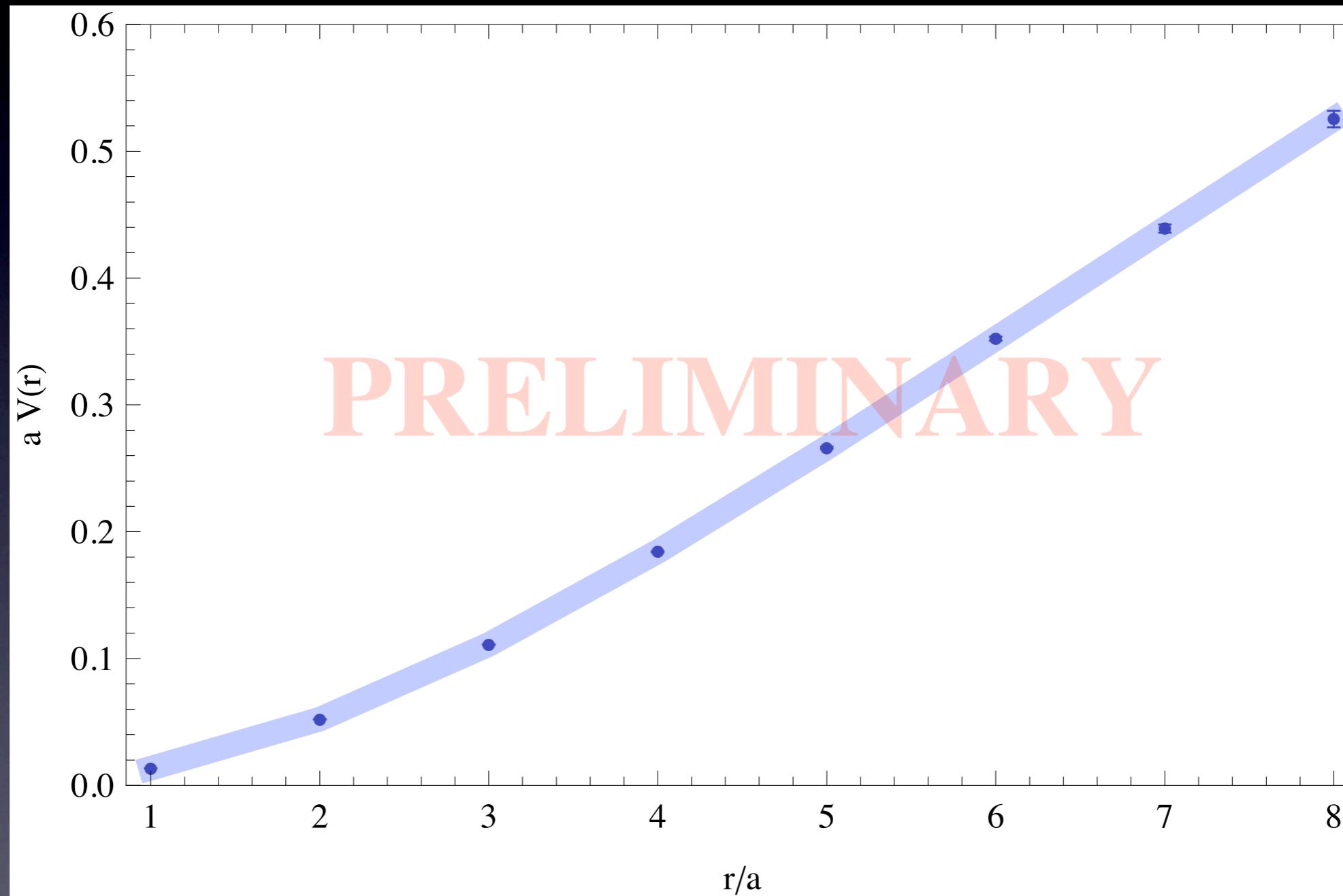
In medium $Q\bar{Q}$ potential

$$C_V(r, t) = \langle 0 | W[R, t_0, t] | 0 \rangle \\ \longrightarrow Z \exp [-V(r)(t - t_0)]$$

$$C_V(r, t; n) = \left\langle 0 \left| W[R, t_0, t] \left[\sum_{\vec{x}} \bar{d} \gamma_5 u(\vec{x}, t) \bar{u} \gamma_5 d(\vec{0}, 0) \right]^n \right| 0 \right\rangle$$

$$R(r, t; n) = \frac{C_V(r, t; n)}{C_V(r, t) C_n(t)} \\ \longrightarrow \# \exp [-\delta V(r, n)(t - t_0)]$$

Vacuum Static Potential



- Gauge field smeared for optimal signal at large r

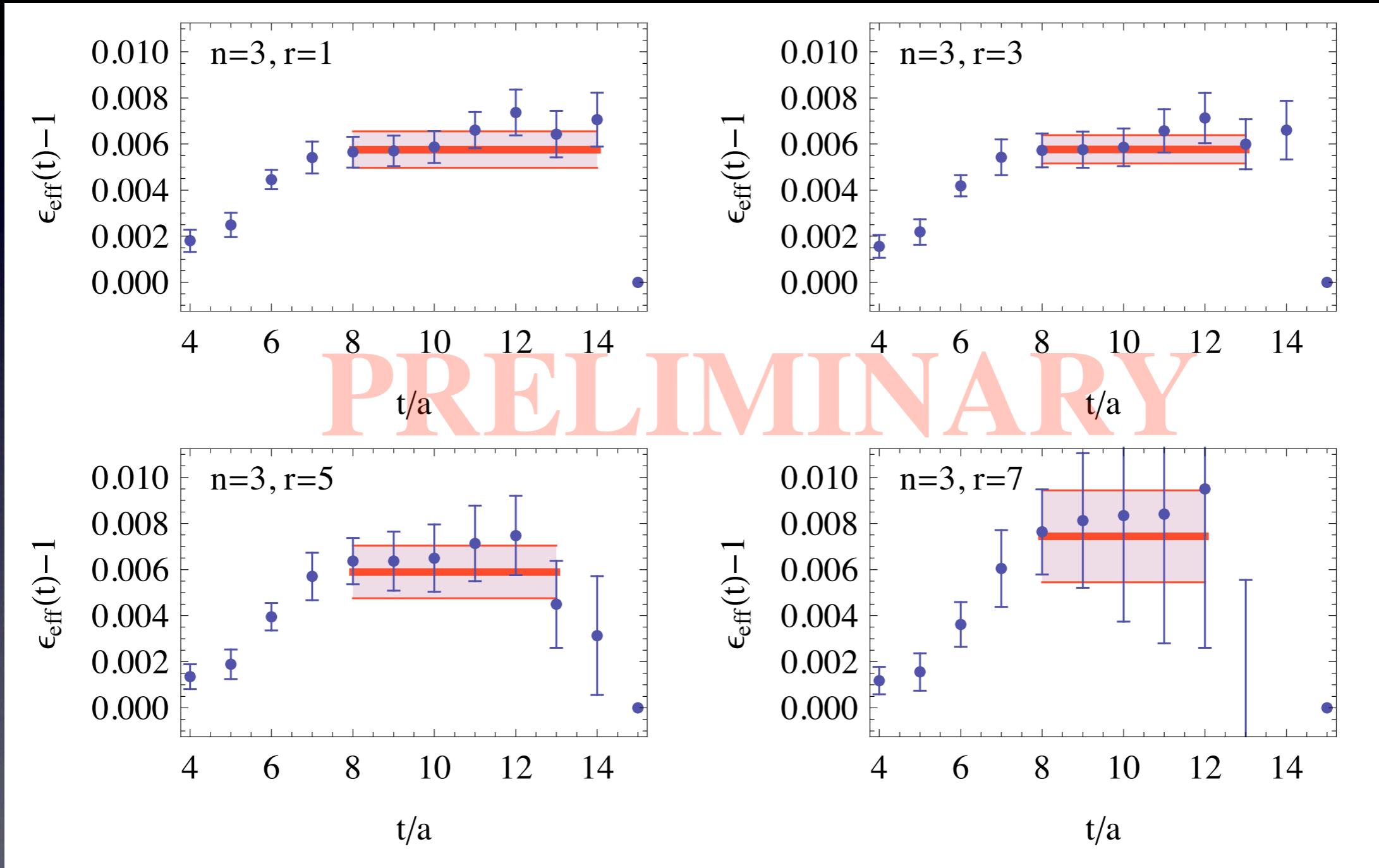
Dielectric “Constant”

- Effects of medium shown by

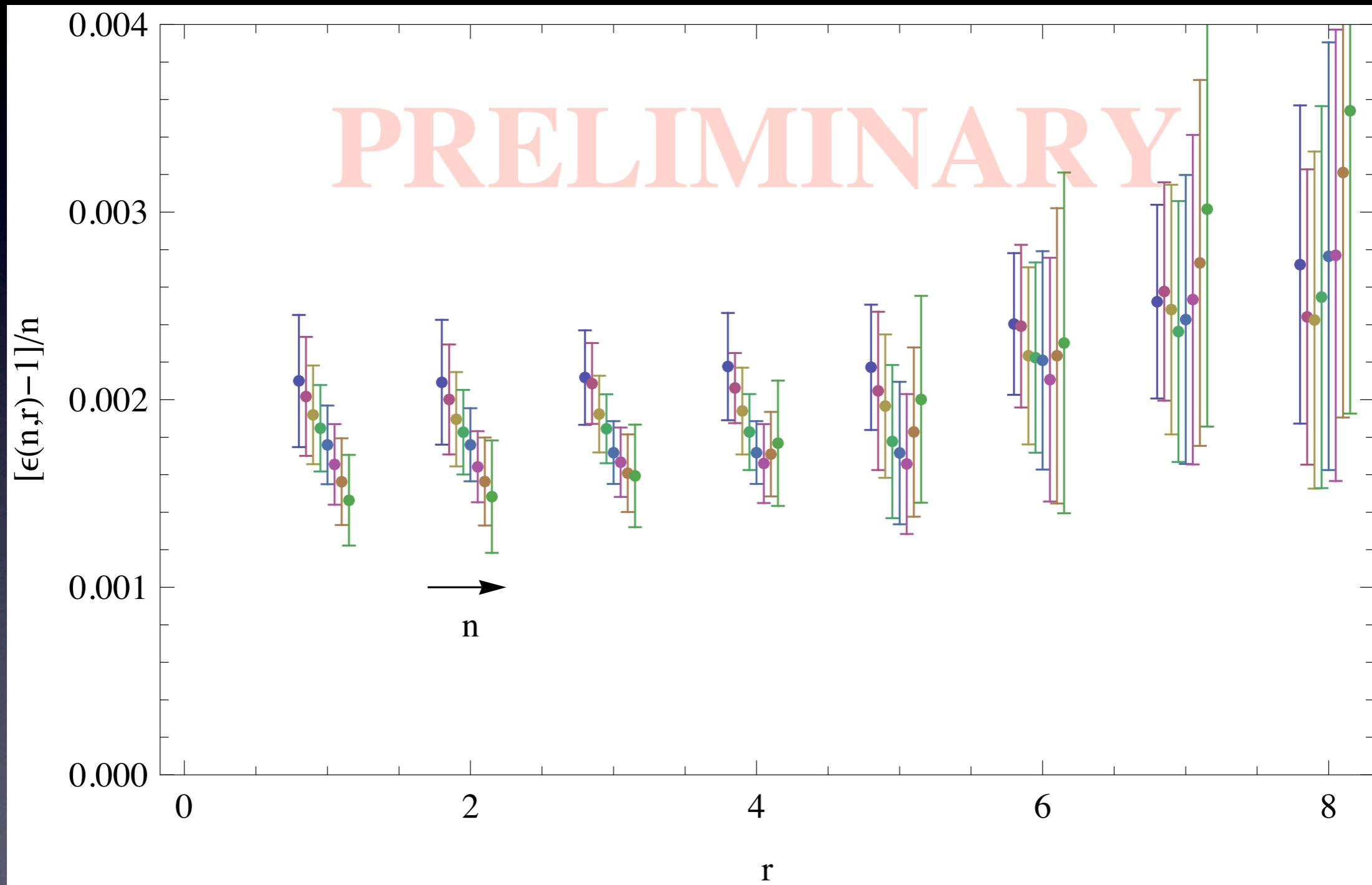
$$\epsilon(r, n) \equiv \left[1 + \frac{\delta V(r, n)}{V(r)} \right]^{-1}$$

- Would be constant for a dielectric medium

Dielectric “Constant”

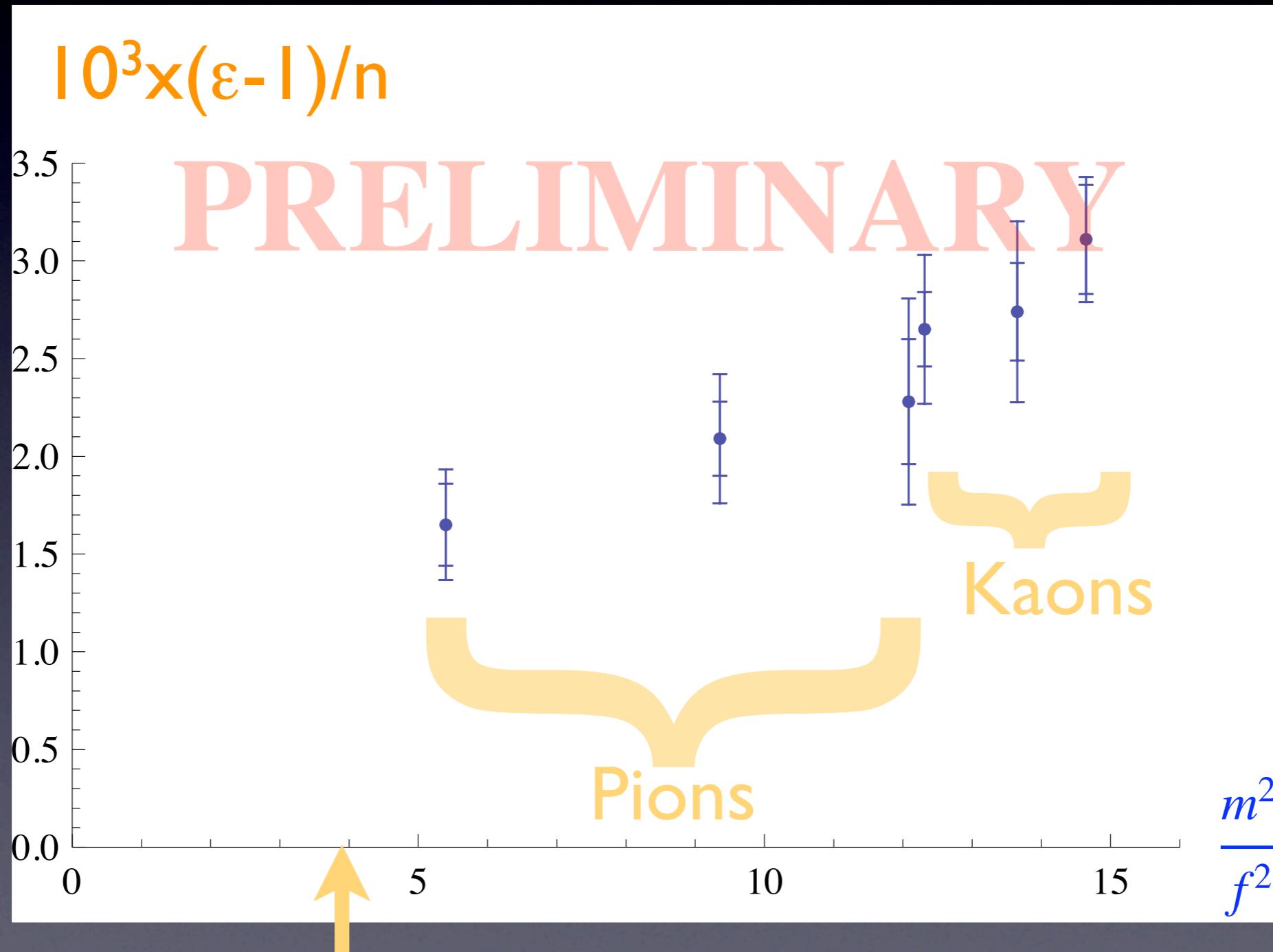


Dielectric “Constant”



$$\epsilon(n,r) \sim 1 + A_n ??$$

Dielectric “Constant”



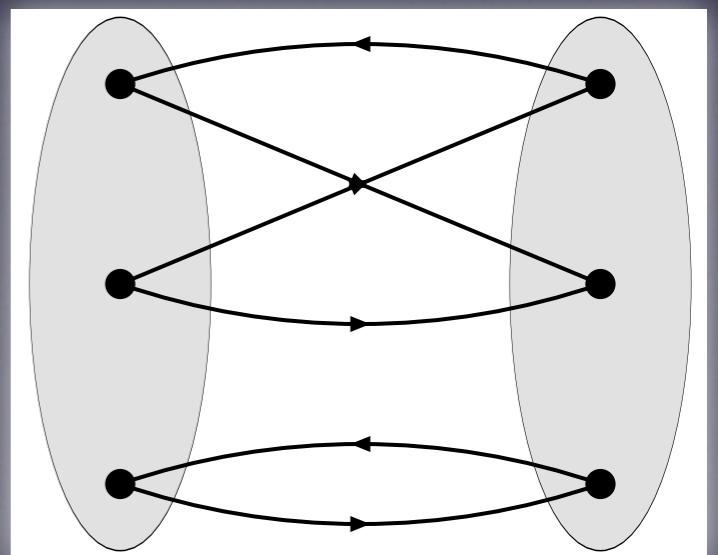
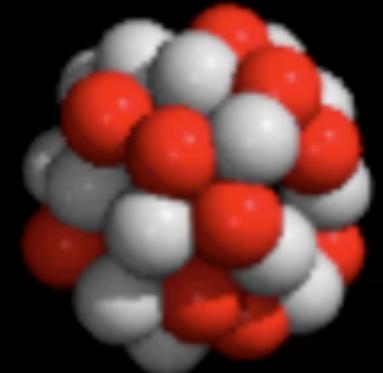
Dielectric “Constant”

- Work in progress
- Well described as a dielectric medium
 - Small vs large R
 - *Effects of gauge field smearing TBD*
- J/ ψ - π scattering
- Relevance to experiment?

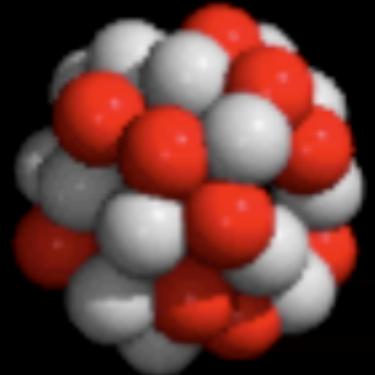
Summary

- Multi-hadron systems
 - Two and three meson interactions
 - Equation of state of meson gases
 - Hadronic screening of static quark potential
 - More to come...

Contractions and Interpolators



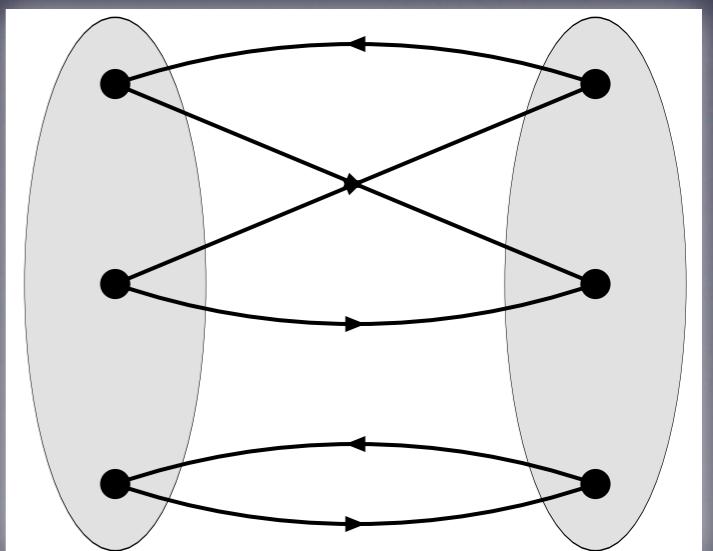
Contractions and Interpolators



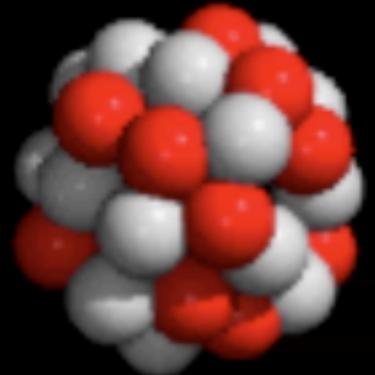
- Number of contractions for ${}^A X_{xyz}$:

$$(A+Z)!(2A-Z)!$$

- Pb: 10^{1289}
- evaluation with fewer operations
- Need many propagators
- What interpolating operators?



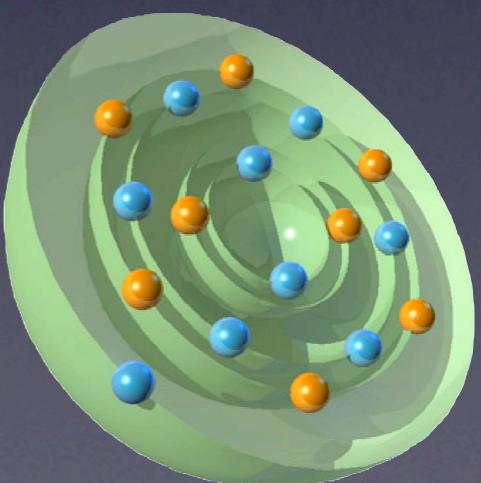
Contractions and Interpolators



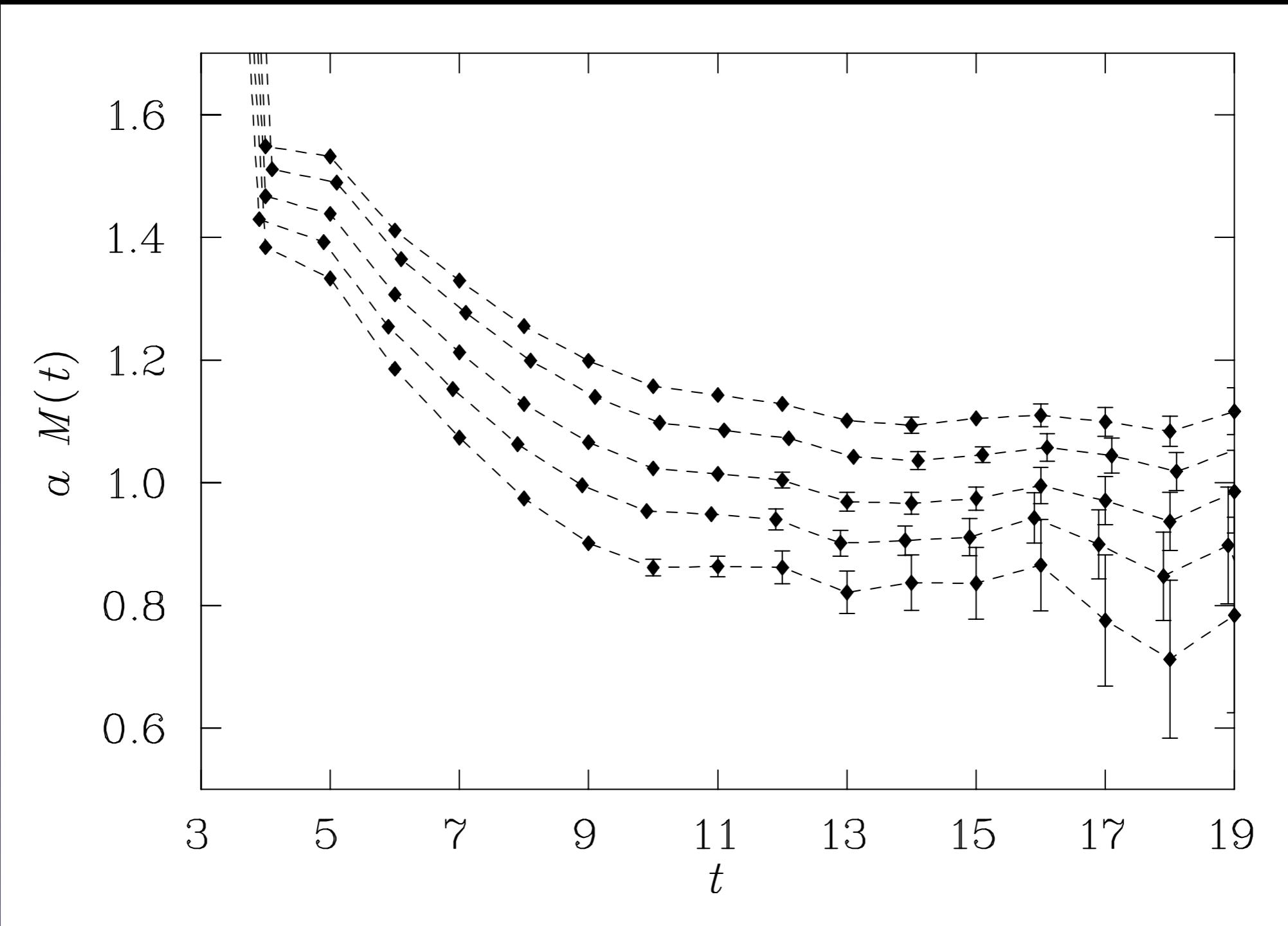
- Number of contractions for ${}^A\text{Xyz}$:

$$(A+Z)!(2A-Z)!$$

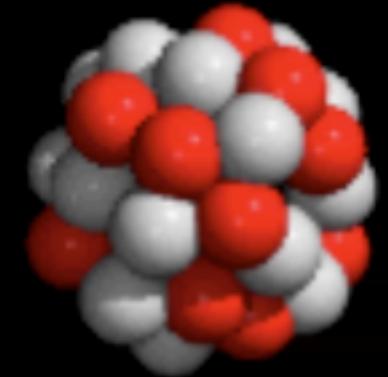
- Pb: 10^{1289}
- evaluation with fewer operations
- Need many propagators
- What interpolating operators?
- The shell model+!



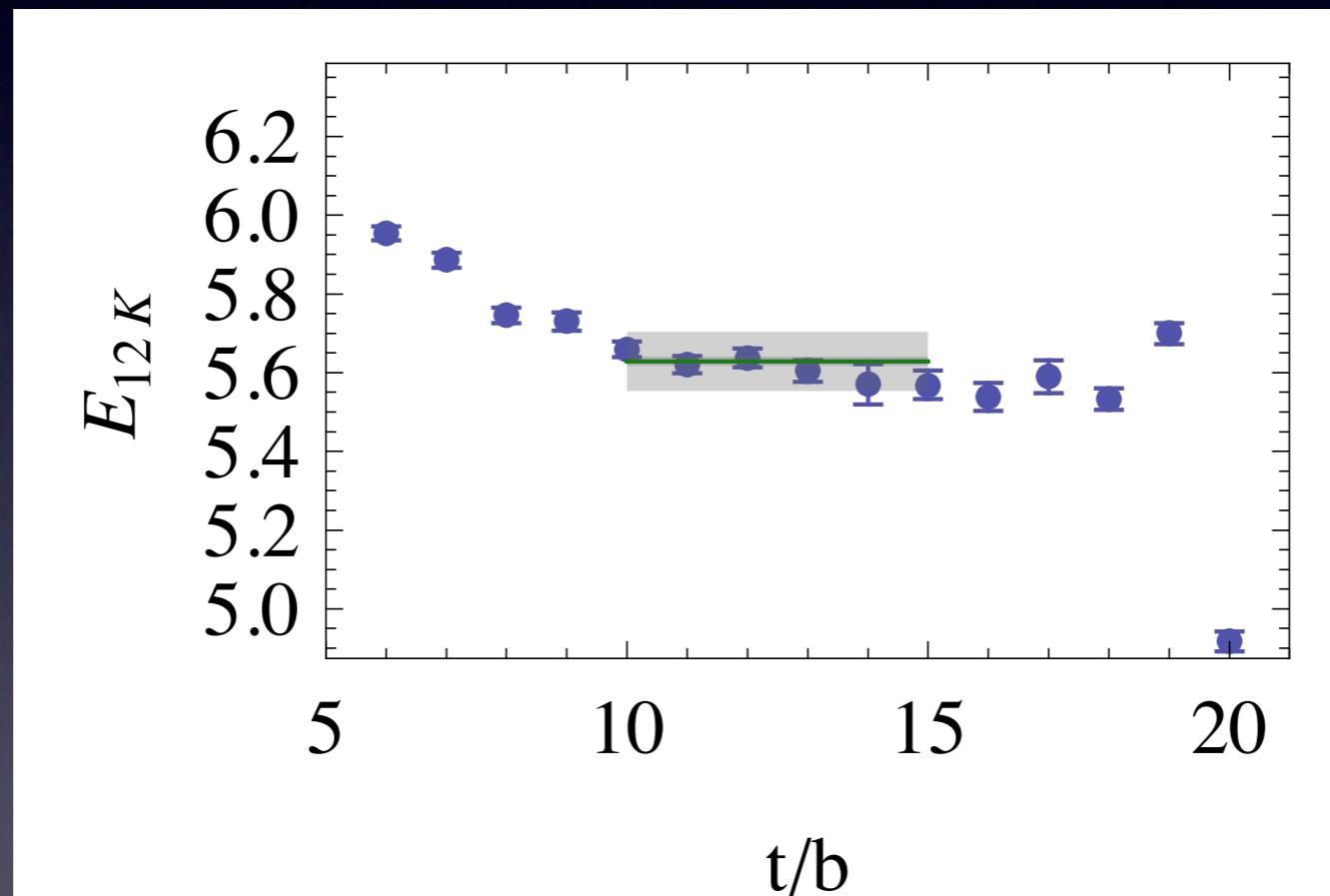
An example



Massive states

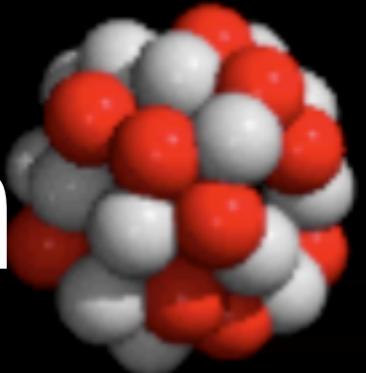


- Nuclear masses \gg typical lattice scales

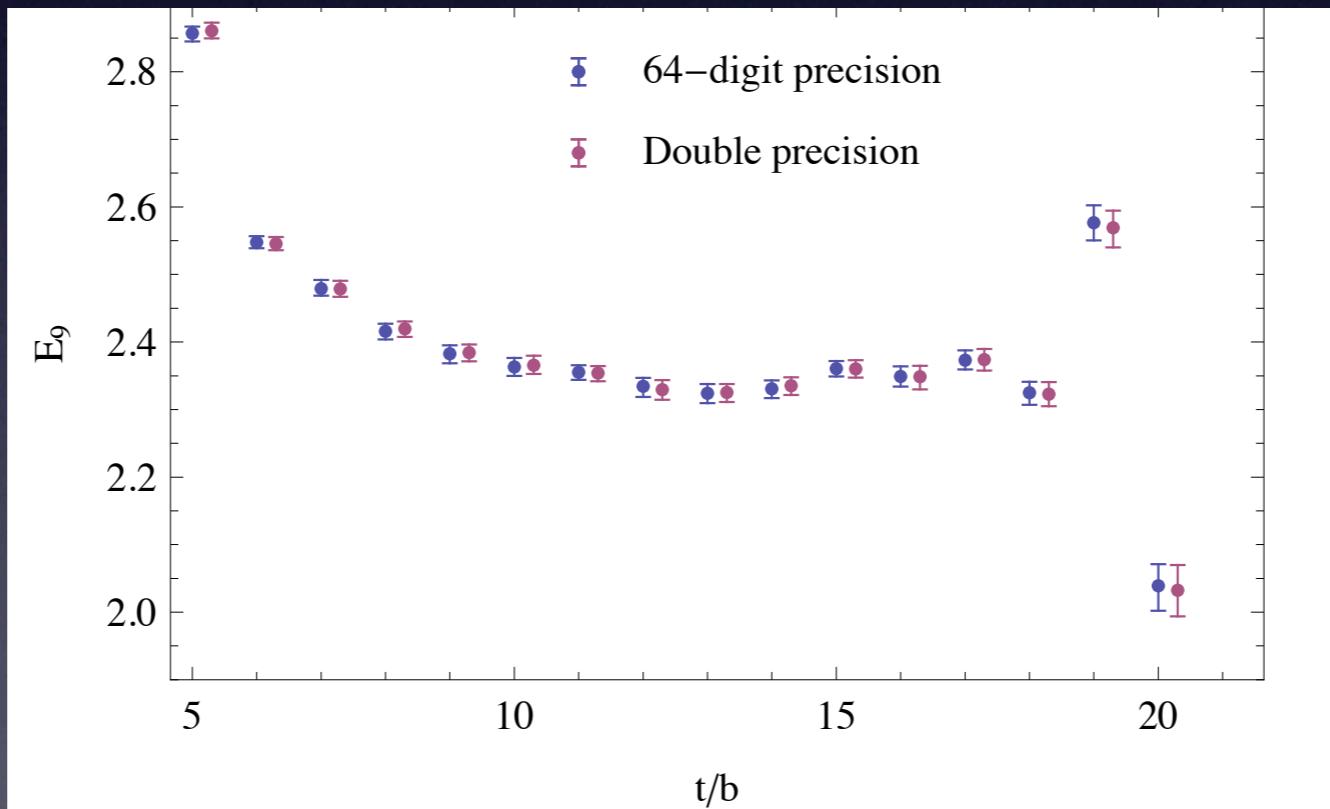


- Twelve kaon energy: $5.6/a \sim 8.8$ GeV

Numerical precision

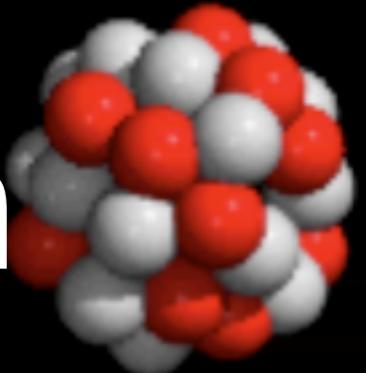


- Double precision is not enough: not gauge invariant
- Need very high precision contraction code

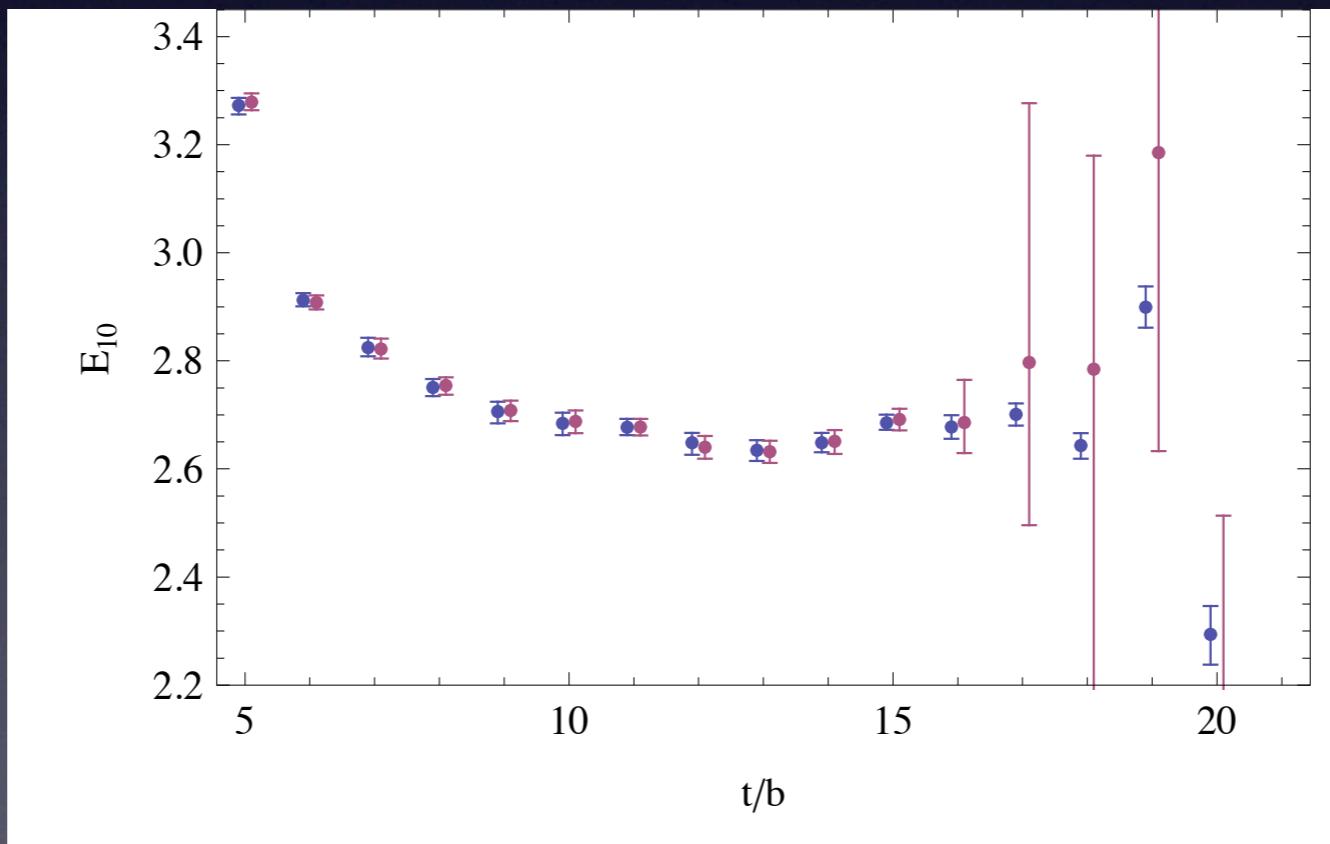


- Propagator precision??

Numerical precision

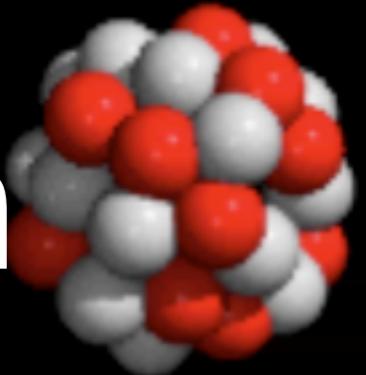


- Double precision is not enough: not gauge invariant
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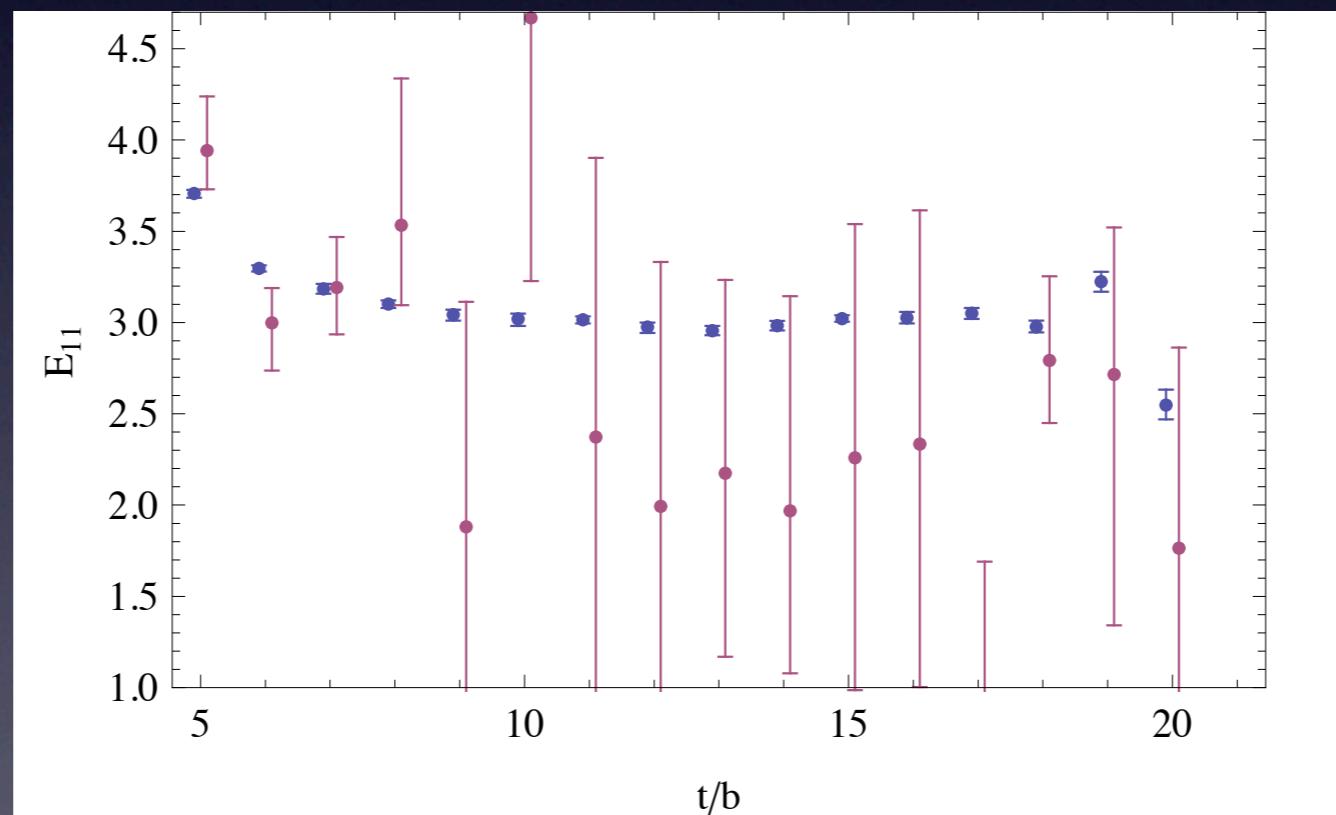


- Propagator precision??

Numerical precision



- Double precision is not enough: not gauge invariant
- Need very high precision contraction code



- Propagator precision??

$n=13$ & matrix identities

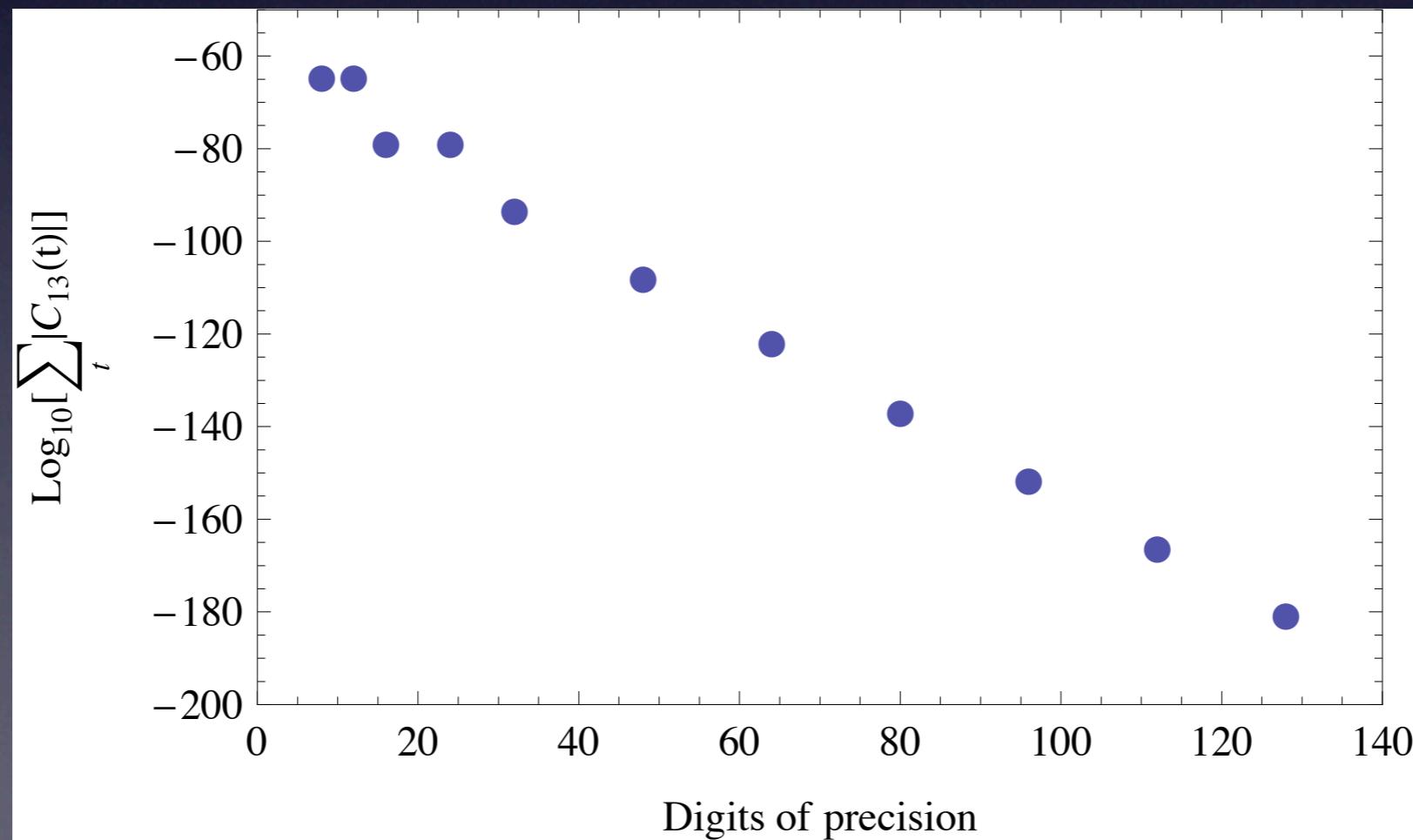
- Pauli principle: $C_{13}(t)$ correlator vanishes
- Matrix Identity: vanishes $\forall 12 \times 12$ matrices

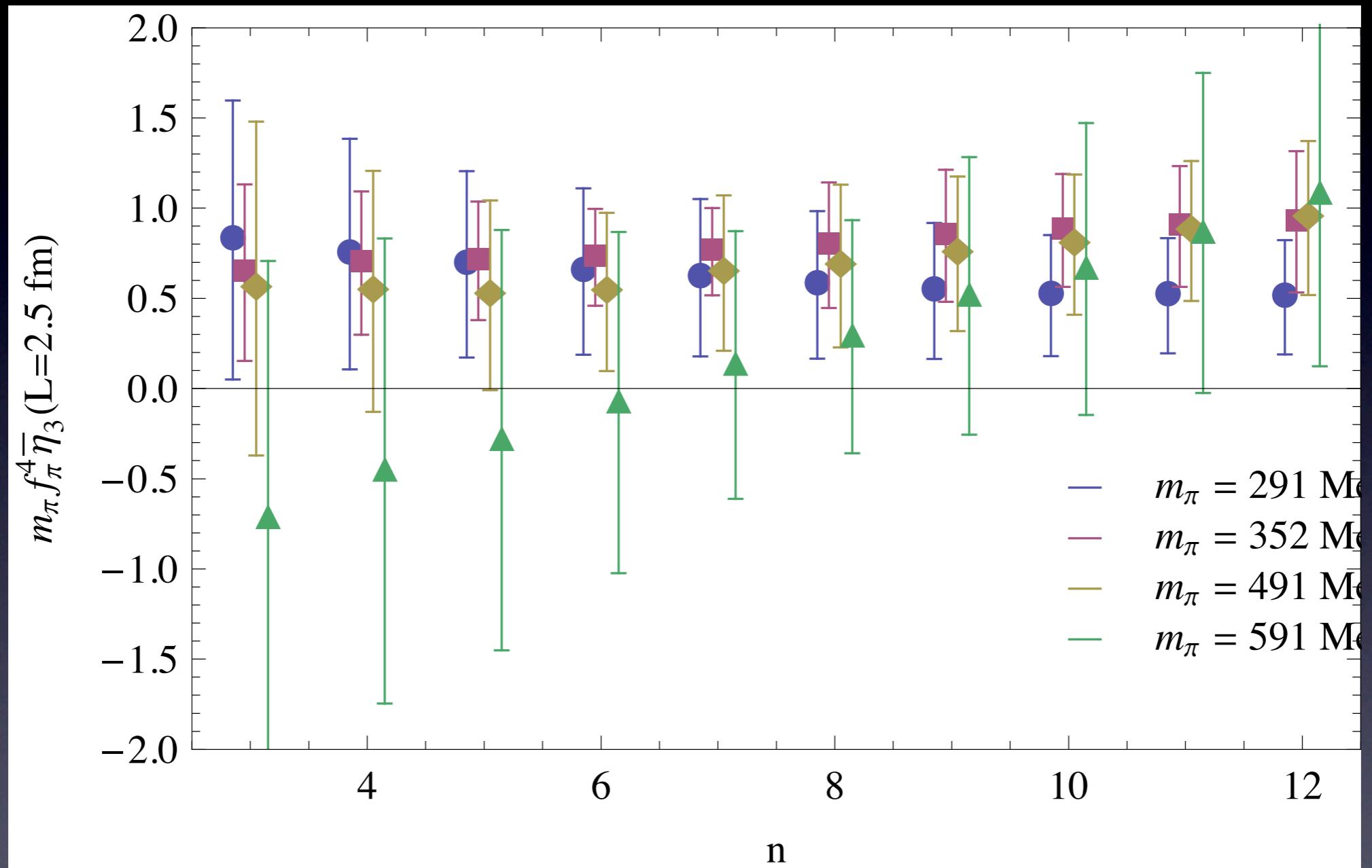
$$\begin{aligned}
C_{13}(t) = & T_1^{13} - 78T_2T_1^{11} + 572T_3T_1^{10} + 2145T_2^2T_1^9 - 4290T_4T_1^9 - 25740T_2T_3T_1^8 + 30888T_5T_1^8 \\
& - 25740T_2^3T_1^7 + 68640T_3^2T_1^7 + 154440T_2T_4T_1^7 - 205920T_6T_1^7 + 360360T_2^2T_3T_1^6 \\
& - 720720T_3T_4T_1^6 - 864864T_2T_5T_1^6 + 1235520T_7T_1^6 + 135135T_2^4T_1^5 - 1441440T_2T_3^2T_1^5 \\
& + 1621620T_4^2T_1^5 - 1621620T_2^2T_4T_1^5 + 3459456T_3T_5T_1^5 + 4324320T_2T_6T_1^5 - 6486480T_8T_1^5 \\
& + 1601600T_3^3T_1^4 - 1801800T_2^3T_3T_1^4 + 10810800T_2T_3T_4T_1^4 + 6486480T_2^2T_5T_1^4 - 12972960T_4T_5T_1^4 \\
& - 14414400T_3T_6T_1^4 - 18532800T_2T_7T_1^4 + 28828800T_9T_1^4 - 270270T_2^5T_1^3 + 7207200T_2^2T_3^2T_1^3 \\
& - 16216200T_2T_4^2T_1^3 + 20756736T_5^2T_1^3 + 5405400T_2^3T_4T_1^3 - 14414400T_3^2T_4T_1^3 - 34594560T_2T_3T_5T_1^3 \\
& - 21621600T_2^2T_6T_1^3 + 43243200T_4T_6T_1^3 + 49420800T_3T_7T_1^3 + 64864800T_2T_8T_1^3 - 103783680T_{10}T_1^3 \\
& - 9609600T_2T_3^3T_1^2 + 32432400T_3T_4^2T_1^2 + 2702700T_2^4T_3T_1^2 - 32432400T_2^2T_3T_4T_1^2 \\
& - 12972960T_2^3T_5T_1^2 + 34594560T_3^2T_5T_1^2 + 77837760T_2T_4T_5T_1^2 + 86486400T_2T_3T_6T_1^2 \\
& - 103783680T_5T_6T_1^2 + 55598400T_2^2T_7T_1^2 - 111196800T_4T_7T_1^2 - 129729600T_3T_8T_1^2 \\
& - 172972800T_2T_9T_1^2 + 283046400T_{11}T_1^2 + 135135T_2^6T_1 + 3203200T_3^4T_1 - 16216200T_4^3T_1 \\
& - 7207200T_2^3T_3^2T_1 + 24324300T_2^2T_4^2T_1 - 62270208T_2T_5^2T_1 + 86486400T_6^2T_1
\end{aligned}$$

$T_i^j = \text{tr}[X^i]^j$

$n=13$ & matrix identities

- Pauli principle: $C_{13}(t)$ correlator vanishes
- Matrix Identity: vanishes $\forall 12 \times 12$ matrices
- Numerically





Analytic details

- Hamiltonian formulation = pionless EFT

$$\begin{aligned} H = & \sum_{\mathbf{k}} h_{\mathbf{k}}^\dagger h_{\mathbf{k}} \left(\frac{|\mathbf{k}|^2}{2M} - \frac{|\mathbf{k}|^4}{8M^3} \right) \\ & + \frac{1}{(2!)^2} \sum_{\mathbf{Q}, \mathbf{k}, \mathbf{p}} h_{\frac{\mathbf{Q}}{2} + \mathbf{k}}^\dagger h_{\frac{\mathbf{Q}}{2} - \mathbf{k}}^\dagger h_{\frac{\mathbf{Q}}{2} + \mathbf{p}} h_{\frac{\mathbf{Q}}{2} - \mathbf{p}} \left(\frac{4\pi a}{M} + \frac{\pi a}{M} \left(ar - \frac{1}{2M^2} \right) (|\mathbf{k}|^2 + |\mathbf{p}|^2) \right) \\ & + \frac{\eta_3(\mu)}{(3!)^2} \sum_{\mathbf{Q}, \mathbf{k}, \mathbf{p}, \mathbf{r}, \mathbf{s}} h_{\frac{\mathbf{Q}}{3} + \mathbf{k}}^\dagger h_{\frac{\mathbf{Q}}{3} + \mathbf{p}}^\dagger h_{\frac{\mathbf{Q}}{3} - \mathbf{k} - \mathbf{p}}^\dagger h_{\frac{\mathbf{Q}}{3} + \mathbf{r}} h_{\frac{\mathbf{Q}}{3} + \mathbf{s}} h_{\frac{\mathbf{Q}}{3} - \mathbf{r} - \mathbf{s}} , \end{aligned}$$

- Calculate to 5th order
- Three loop diagrams: 9d integer sums