

Dirac and Majorana Mass

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Overview of the Talk

Outline

Introduction

Prerequisites

Quantum theory of fields

Dirac and Majorana Mass

Physical Relevance

Closing Remarks

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Motivation | Mass

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Towards a quantum theory of fields

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- Conclusion: Parameter m is mass

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- To be Klein Gordon, $\{\gamma^\mu, \gamma^\nu\} = -2\delta^{\mu\nu}$

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- To be Klein Gordon, $\{\gamma^\mu, \gamma^\nu\} = -2\delta^{\mu\nu}$

- Claim

γ^μ are 4×4 matrices and ψ then is a 4-component object, called a Dirac spinor.

The Dirac Equation

- Hat:

$$\begin{aligned}
 \sigma^\mu &\equiv (1, \vec{\sigma}) & \bar{\sigma}^\mu &\equiv (1, -\vec{\sigma}) \\
 \gamma^\mu &\equiv \begin{pmatrix} 0 & \sigma^\mu \\ \bar{\sigma}^\mu & 0 \end{pmatrix}
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$$m\bar{\psi}\psi = m \begin{pmatrix} \psi_L^\dagger & \psi_R^\dagger \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix} = m(\psi_R^\dagger \psi_L + \psi_L^\dagger \psi_R)$$

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- Mass term mixes the left and right spinors
- Explore: mass term that doesn't mix

Projectors, mixing of mass terms

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- In this notation also, there's mixing

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 - I can write C as a product of γ matrices as

$$C = i \begin{pmatrix} \sigma^2 & 0 \\ 0 & -\sigma^2 \end{pmatrix} = -i\gamma^2\gamma^0$$

which is easy to verify.

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- To ensure reality, we add $-m\psi^\dagger C\psi^*$

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- There’re alternatives, such as ‘see-saw’ model

The End

?

References

- An Introduction to Quantum Field Theory
M. E. Peskin, D. V. Schroeder
Addison-Wesley Publishing Company
- PHY659 Lectures
Prof. C. S. Aulakh
Spring 2015, IISER Mohali
- Quantum Field Theory, Second Edition
L. H. Ryder
Cambridge University Press

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