

Linear 2d \rightarrow Non-linear
 fixed pt: classify?
 closed orbits: stable? flow
 $\vec{x}(t+T) = \vec{x}(t)$

Chaos 11.4
 (in 11.3, 11.2 was lab)

$\dot{x}_1 = f_1(x_1, x_2)$	$\dot{x} = x + e^{-y}$	$\dot{y} = -y$	$y = 0 \Rightarrow y^* = 0$	Nullclines
$\dot{x}_2 = f_2(x_1, x_2)$	$\dot{x} = x + e^{-y}$	$\dot{y} = -y$	$x = 0 \Rightarrow$	$\dot{x} = 0$
$f_1(x_1^*, x_2^*) = 0$	$y = -y$	$y(t) = y_0 e^{-t}$	$x + e^{-0} = 0$	$y = 0$
$f_2(x_1^*, x_2^*) = 0$	$y(t) = y_0 e^{-t}$	$x^* = -1$	$x \approx x_1$	$\dot{y} = 0$
	$x \approx x_1$			

saddles
 spirals (stable)
 nodes
 Centre; $\text{Re}(\lambda) = 0$
 From Hartman-Borg man

(u) $= \begin{pmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$
 remember: at x^*, y^*

$\dot{u} = u \frac{\partial f}{\partial x} + v \frac{\partial f}{\partial y}$
 $(\because f(x^*, y^*) = 0)$
 similarly
 $\dot{v} = u \frac{\partial g}{\partial x} + v \frac{\partial g}{\partial y}$

stability (local)
 Non-linear \rightarrow linearize
 $\dot{x} = f(x, y) \quad x^*, y^*$
 $\dot{u} = g(x, y) \quad u = x - x^*$
 $f(x^*, y^*) = 0 \quad v = y - y^*$
 $g(x^*, y^*) = 0$

$\dot{x} = -x + x^3 = f$
 $\dot{y} = -2y = g$
 $\dot{y} = 0$

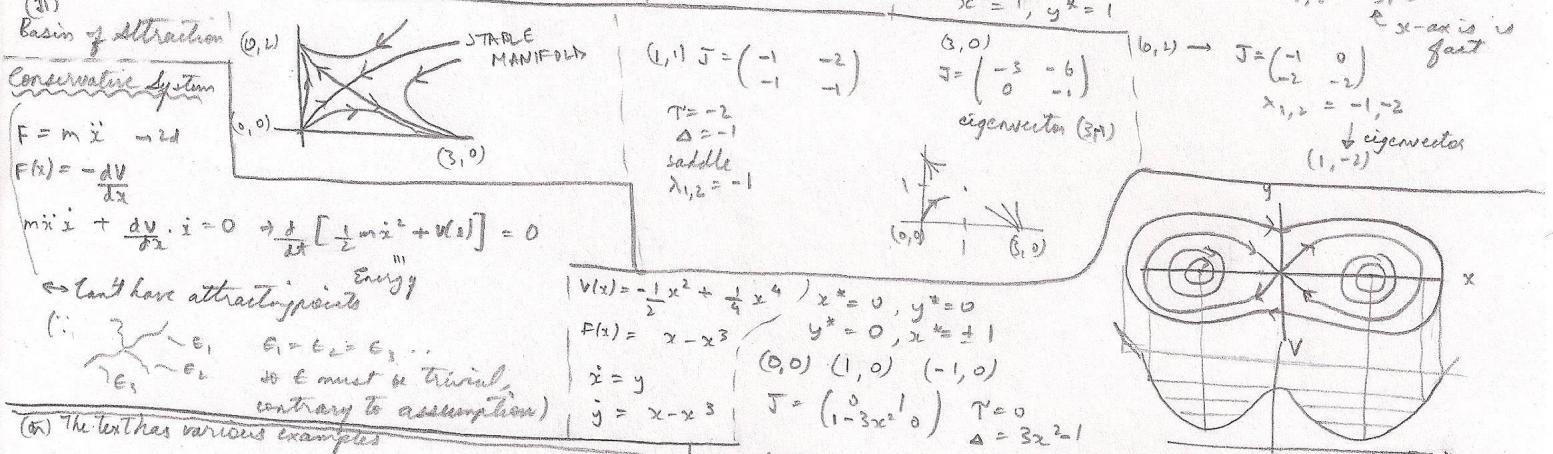
$x = 0, +1, -1$
 $(0, 0), (+1, 0), (-1, 0)$

$J = \begin{pmatrix} -1 + 3x^2 & 0 \\ 0 & -2 \end{pmatrix}$

$\text{at } x = 0: A = \begin{pmatrix} -1 & 0 \\ 0 & -2 \end{pmatrix}$ STABLE NODE
 $\text{at } x = \pm 1: A = \begin{pmatrix} 2 & 0 \\ 0 & -2 \end{pmatrix}$ SADDLES

$\dot{x}_i = f_i(x_1, \dots, x_n)$ 1) Fixed Points; solve a simultaneous eq.
 $\dot{x}_n = f_n(x_1, \dots, x_n)$ 2) Stability of fixed points
 $\frac{\partial f_i}{\partial x_j}$
 (J)
 Poincaré, eigenvalue/eigenvectors
 (a) $\det(\lambda_i) \neq 0$; robust
 3) Nullclines $\dot{x}_i = 0 \rightarrow$ flow across nullclines

Sadesha 12-1
 (2) Poincaré-Bendixson (2)
 $2d < FP$ closed orbit.
 LOTKA-VOLTERRA MODEL
 2 species - competition
 $\dot{x} = x(3-x-2y)$
 $\dot{y} = y(2-3x-y)$
 2) $x^* = 0, y^* = 2$ for $(0,0), (0,2), (3,0), (1,1)$
 we find J
 3) $y^* = 0, x^* = 3$ $(0,0) \rightarrow J = \begin{pmatrix} 3 & 0 \\ 0 & 2 \end{pmatrix}$
 4) Simultaneous $x^* = 1, y^* = 1$
 $x_{1,2} = 3, 2$
 x-axis is fast



$\dot{x} = f(x, y)$
 $\dot{y} = g(x, y)$
 $u = x - x^*$
 $v = y - y^*$

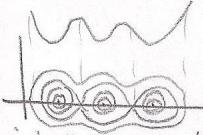
$u = \dot{u}$
 $= f(x^* + u, y^* + v)$
 $= f(x^*, y^*) + u \frac{\partial f}{\partial x} + v \frac{\partial f}{\partial y} + O(u^2, v^2, uv)$
 $= u \frac{\partial f}{\partial x} + v \frac{\partial f}{\partial y}$
 $v = u \frac{\partial g}{\partial x} + v \frac{\partial g}{\partial y}$
 $\begin{pmatrix} \dot{u} \\ \dot{v} \end{pmatrix} = \begin{pmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$

Def: Robust case: Repellers, Attractors, Saddles
 Marginal case: Centres (+1 more)
 Criterion: At least one λ s.t. $R(\lambda) = 0$
 Def": Hyperbolic fixed points $\Rightarrow \text{Re}(\lambda) \neq 0$ for both λ .

Hartman-Grobman Thm: For hyperbolic pts, linearization faithfully captures the phase portrait.

(3) Given an attracting fixed point x^* , we define its basin of attraction as the set of init conditions x_0 , s.t. $x(t) \rightarrow x^*$ as $t \rightarrow \infty$.
 Basin Boundary = stable manifold of saddle.

Conservative
Centre
Saddle

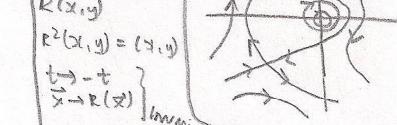


TODO: Homoclinic Orbit
(in) Reversible Time Reversible
 $t \rightarrow -t$

$$\frac{dx}{dt} : \text{even sign } \dot{x} = y \\ \frac{d^2x}{dt^2} : \text{invariant } y, \dot{y} = f(x) \\ t \rightarrow -t \quad \begin{cases} x \rightarrow -x \\ y \rightarrow -y \end{cases}$$

$\dot{x} = f(x, y)$? $f(x, -y) = -f(x, y)$ odd
 $\dot{y} = g(x, y)$ $g(-x, -y) = g(x, y)$ even
you want $x \rightarrow -x$ leaves the system invariant.

Syndesma 12.2
 $\begin{cases} \dot{x} = y - y^2 & \dot{x} = 0 \quad (0, 0) \text{ its centre} \\ \dot{y} = -x - y^2 & \dot{y} = 0 \quad \therefore t \mapsto -t \end{cases}$
 $J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \Delta = 1 \quad \text{centre becomes unchanged}$
 $(1, 1) \quad \text{more fixed points}$
 $(-1, -1) \quad \text{fixed points}$
claim: reversible \Rightarrow robust



$$R(x, y) \\ R^2(x, y) = (x, y) \\ \begin{cases} t \mapsto -t \\ x \mapsto R(x) \end{cases} \quad \text{invariant} \\ \begin{cases} \dot{x} = -2\cos x - \cos y \\ \dot{y} = -2\sin y - \cos x \end{cases} \\ t \mapsto -t \quad \text{leaves invariant} \\ x \mapsto -x \\ y \mapsto -y \quad \text{de eq's.}$$

Pendulum

$$\frac{d^2\theta}{dt^2} + \frac{L}{m} \sin \theta = 0$$

$$T = wt$$

$$E_{\text{kin}} = \frac{1}{2} w^2$$

$$\frac{d^2\theta}{dt^2} + \sin \theta = 0$$

$$\dot{\theta} = v$$

$$v = -\sin \theta$$

$$(\theta^*, v^*) = (k\pi, 0)$$

$$(0, 0), (\pi, 0) \xrightarrow{T} \left(\begin{matrix} 0 \\ 1 \end{matrix} \right) \quad \pi \approx 0, \Delta = -1$$

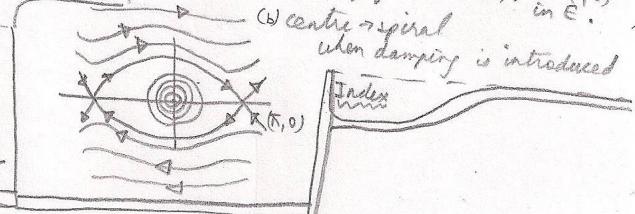
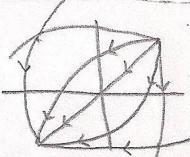
$$\lambda_{1,2} = -1, 1 \quad \text{invariant under } t \mapsto -t$$

$$v \mapsto -v \quad \text{saddle.}$$

$$E = \frac{1}{2} v^2 - \omega_0^2$$

$$E = \frac{1}{2} (v^2 + \theta^2) - 1 \quad \text{circle}$$

(a) Energy of separatrix: out $(1, 0)$ in $(-1, 0)$
(b) centre \Rightarrow spiral when damping is introduced



(in) Homoclinic orbits: trajectories that start and end at the same fixed point

N.B.: Even though centres are predicted by linearization, they're indeed reliable: the system is conservative.

Heteroclinic trajectories

\Rightarrow saddle connections = trajectories joining two saddle points

Thm: Consider $\dot{\vec{x}} = \vec{F}(\vec{x})$ with $\vec{x} = (x, y) \in \mathbb{R}^2$,

f is continuously differentiable.

Suppose \exists a conserved quantity $E(\vec{x})$

& that \vec{x}^* is an isolated fixed pt.

(no fixed pt. in a small neighbourhood)

If \vec{x}^* is also a local minimum (or maximum)

of E , then all trajectories sufficiently

close to \vec{x}^* are closed.

Thm: Suppose that origin $\vec{x}^* = 0$ is a linear centre for the continuously differentiable system

$$\dot{x} = f(x, y)$$

$$\dot{y} = g(x, y)$$

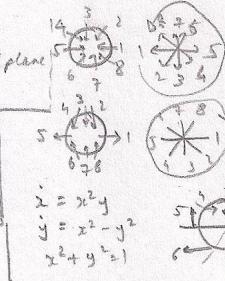
& suppose that the system is reversible. Then sufficiently close to the origin, all trajectories are closed curves.

Index Theory

$$\frac{1}{\pi} = f(z)$$

$$\phi = \tan^{-1}\left(\frac{y}{x}\right)$$

- $\int \phi dz \Rightarrow I_c = \frac{1}{2\pi} [\phi]_c$
- 1) Shouldn't sit on a fixed pt.
 - 2) Shouldn't intersect itself.



$$\begin{aligned} x &= x^2 y \\ y &= x^2 - y^2 \\ x^2 + y^2 &= 1 \end{aligned}$$

No fixed points

No closed orbit

Next: Limit Cycles
Relaxation Oscillators

$$\partial L = I_c$$

$$c' \rightarrow c$$

No fp crossed

Proof: Pumps value is 1 if c is, try, $I_c = 1$
int continuity

$$-\pi + 2\pi - \pi = 0$$

stable unstable node saddle

$$I_c = 1$$

stable unstable node saddle

$$I_c = -1$$

stable unstable node saddle

$$I_c = 0$$

closed curve

$$\Rightarrow I_c = 1$$

fixed points exist

$$x = x(2 - x - 2y)$$

$$y = y(2 - x - y)$$

$$x^2 + y^2 = 1$$

No fixed pt

saddle, $I_c = -1$

Index theory isn't ok,

but $x=0 \& y=0$ are fixed pts,

you can't cross 'em.

(b) if C doesn't enclose a f.p.

$\Rightarrow I_c = 0$ ' then you shrink to zero.

(c) Kernel arrows - index is invariant

$\phi \rightarrow \phi + \pi \Rightarrow$ Index doesn't change

int continuity

$\frac{1}{2\pi} [\phi]_c = I_c$

$\Rightarrow \frac{1}{2\pi} \sum [\phi]_{c,p.}$

This is 'n' isolated fixed pt.

$\Sigma^+, \Sigma^-, \dots$