

I. Background

- ▶ Einstein: 'locality' \implies Quantum Mechanics (QM) is incomplete [3].
- ▶ Bell: 'locality' $\implies \langle \hat{B} \rangle \leq 2$; For some $|\psi\rangle$, QM $\implies \langle \hat{B} \rangle = 2\sqrt{2}$ [1]. Verified experimentally (without loop holes in 2015)
- ▶ Comment: At roughly the same time, various physicists had produced proofs of the claim that one can't complete QM satisfactorily, that a sensible complete 'hidden variable' (HV) description of nature was impossible.
- ▶ Bohmian Mechanics (BM): a HV description, that (i) 'completes' QM in a simple, clear, precise but non-local manner, and (ii) is deterministic [2].
- ▶ Defn: *Deterministic* \equiv If in principle, the outcome of measuring each observable is predictable, given the HVs.
- ▶ Comment: Bell's inequality requires entanglement in some form, to prove Einstein's notion of locality incorrect. Recently, another peculiar feature of QM has been identified, namely contextuality.
- ▶ Impl Defn: *Context* \equiv If $[\hat{A}, \hat{B}] = 0$ and $[\hat{A}, \hat{C}] = 0$ but $[\hat{B}, \hat{C}] \neq 0$, then possible contexts are \hat{A}, \hat{B} and \hat{C} or \hat{A} and \hat{C} [5].
- ▶ Defn: *Non-contextual* \equiv Value an operator takes, depends only on the state (including 'hidden variables') and the choice of the operator \hat{A} (not it's context) [5].
- ▶ Defn: *Contextual* \equiv Value an operator takes, depends on it's context [5].
- ▶ Comment: This notion arises, at least in certain explicit constructions, where one is unable to assign values to operators, consistent with predictions of QM.
- ▶ Aim: Understand how a deterministic theory can be consistent with the notion of contextuality.

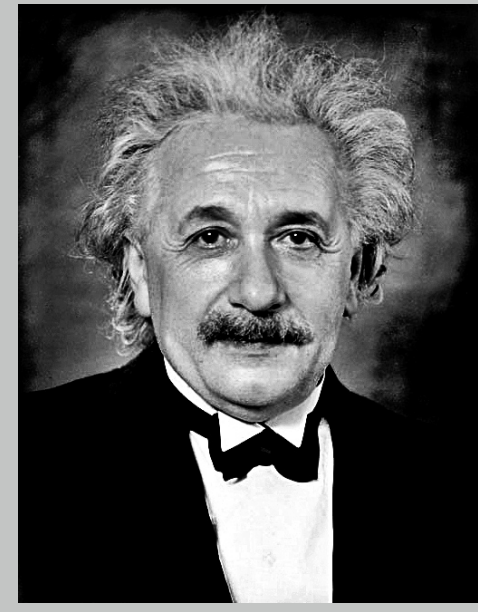


Figure 1: A. Einstein



Figure 2: J. Bell



Figure 3: D. Bohm

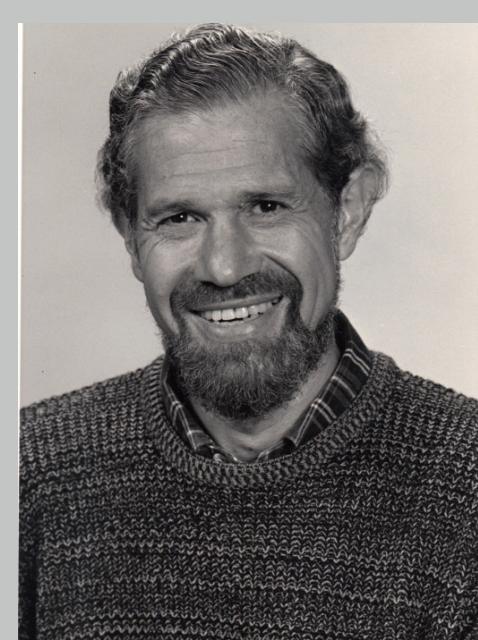


Figure 4: S. B. Kochen

II. Overview

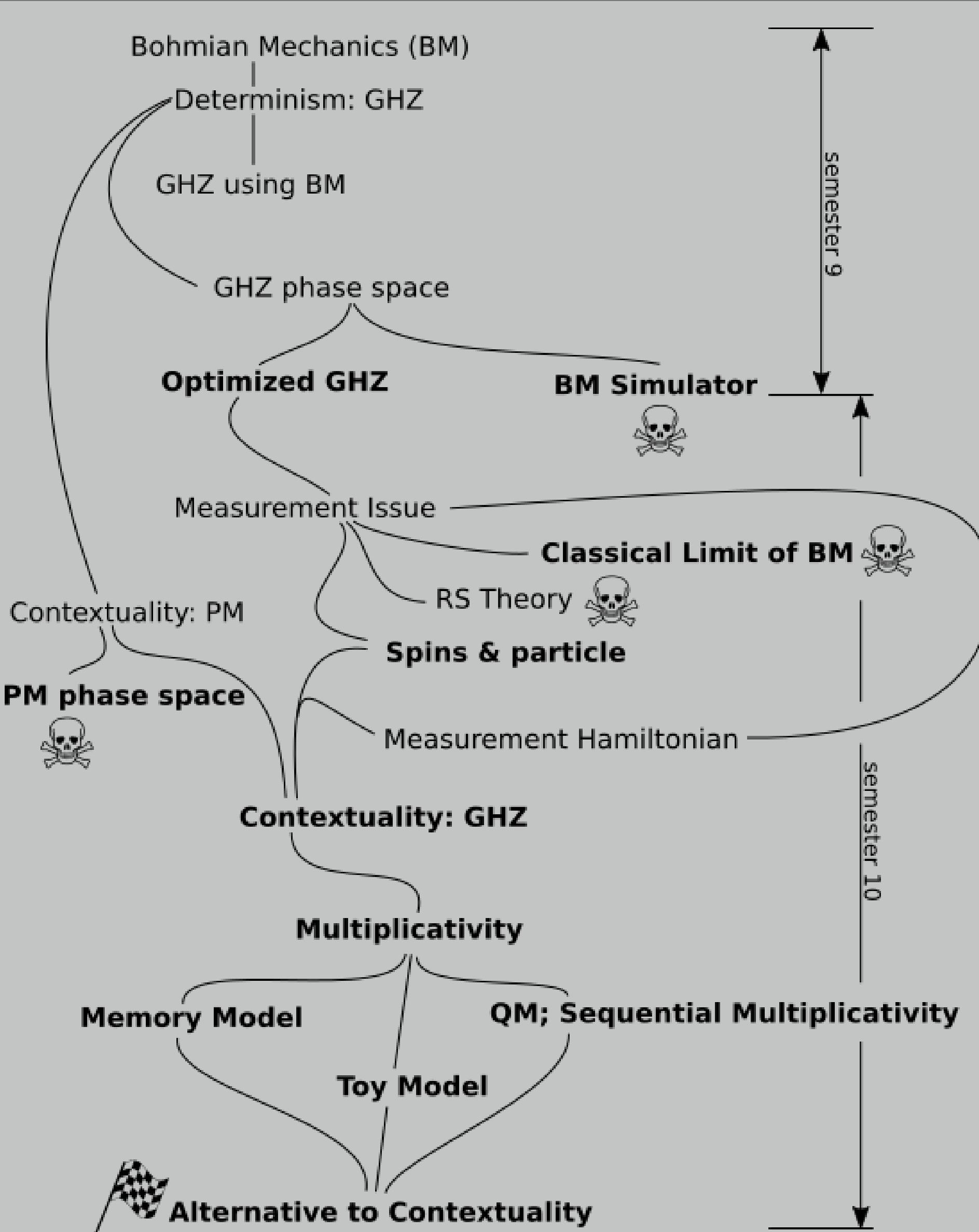


Figure 5: Exploration flow: Boldface titles represent new results

Early Acknowledgments

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III. Multiplicativity

- ▶ Defn: *Compatible operators* \equiv Two observables \hat{A} and \hat{B} are mutually compatible if, given that the system is prepared in a state s.t. measurement \hat{A} yields repeatable results, measurement of \hat{B} doesn't change the result of measuring \hat{A} . For projective measurements, its equivalent to $[\hat{A}, \hat{B}] = 0$.
- ▶ Defn: *Multiplicativity* \equiv For compatible operators \hat{B}_i , a model is multiplicative iff

$$f(m_1(\hat{B}_1), m_1(\hat{B}_2), \dots, m_1(\hat{B}_n)) = \quad (1)$$

$$m_1(f(\hat{B}_1, \hat{B}_2, \dots, \hat{B}_n)), \quad (2)$$

where $m_j(\hat{*})$ represents the assigned value of the operator, and j encodes the sequence of measurement. Note that this is an ontological statement and can't be experimentally tested.

- ▶ Defn: *Sequential Multiplicativity* \equiv For compatible operators \hat{B}_i , a model is sequentially multiplicative iff

$$f(m_{k_1}(\hat{B}_1), m_{k_2}(\hat{B}_2), \dots, m_{k_n}(\hat{B}_n)) = \quad (3)$$

$$m_1(f(\hat{B}_1, \hat{B}_2, \dots, \hat{B}_n)), \quad (4)$$

where $\mathbf{k} \equiv (k_1, k_2, \dots, k_n) \in ((1, 2, \dots, n) + \text{all possible permutations})$, $m_j(\hat{*})$ represents the assigned value of the operator, and j encodes the sequence of measurement.

- ▶ Example: $\hat{B}_1 = \hat{\sigma}_x \otimes \hat{\sigma}_y$, $\hat{B}_2 = \hat{\sigma}_y \otimes \hat{\sigma}_x$ so that $\hat{C} = \hat{B}_1 \hat{B}_2 = \hat{\sigma}_z \otimes \hat{\sigma}_z$. $|\psi\rangle = |00\rangle$, so that $m_1(\hat{C}) = 1$, while $m_1(\hat{B}_i) = \pm 1$. If say $m_1(\hat{B}_1) = -1$, then $\psi \rightarrow$ (figure this) so that entails $m_2(\hat{B}_2) = -1$ as well, consistent with $m_1(\hat{C}) = m_1(\hat{B}_1)m_2(\hat{B}_2)$.
- ▶ Claim: Quantum Mechanics is sequentially multiplicative.

The Toy Model — Example

Iteration	$i = 1$	$i = 2$	$i = 3$
$ \psi_{\text{init}}\rangle$	$ 00\rangle$	$ 00\rangle$	$\frac{ 00\rangle + 11\rangle}{\sqrt{2}}$
HV/Toss	$c = -1$	$c = -1$	$c = +1$
Predictions (Assignments)	$m_1(\hat{A}_{ij}) \doteq \begin{bmatrix} -1 & -1 & -1 \\ -1 & -1 & -1 \\ -1 & -1 & +1 \end{bmatrix}$	$m_2(\hat{A}_{ij}) \doteq \begin{bmatrix} -1 & -1 & -1 \\ -1 & -1 & -1 \\ -1 & -1 & +1 \end{bmatrix}$	$m_3(\hat{A}_{ij}) \doteq \begin{bmatrix} +1 & +1 & +1 \\ +1 & +1 & -1 \\ +1 & +1 & +1 \end{bmatrix}$
Operator Measured	$\hat{A}_{13} = \hat{\sigma}_z \otimes \hat{\sigma}_z$; $m_1(\hat{A}_{13}) = +1$	$\hat{A}_{23} = \hat{\sigma}_y \otimes \hat{\sigma}_y$; $m_2(\hat{A}_{23}) = -1$	$\hat{A}_{33} = \hat{\sigma}_x \otimes \hat{\sigma}_x$; $m_3(\hat{A}_{33}) = +1$
$ \psi_{\text{final}}\rangle$	$ 00\rangle$	$\frac{ 00\rangle + 11\rangle}{\sqrt{2}}$	$\frac{ 00\rangle + 11\rangle}{\sqrt{2}}$

(8)

V. Contextuality - Memory Model

An example of a contextual and non-multiplicative model; Sequential multiplicativity has been assumed.

- ▶ Initial: The assignment is as given in the first Mat in Eq. 9.
- ▶ Remark: The system is assumed to be capable of remembering the last three observables that were measured.
- ▶ Algorithm: Upon measurement of an observable, (i) yield the value as saved in the matrix, (ii) append the observable in the 3 element memory and (iii) update the matrix, once the context (set of commuting observables) is known, to satisfy the PM requirements.

$$m_1(\hat{A}_{ij}) = m_2(\hat{A}_{ij}) \doteq \begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix}, m_3(\hat{A}_{ij}) \doteq \begin{bmatrix} 1 & 1 & -1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix} \quad (9)$$

- ▶ For example:

Operator	Updated Array	Assignment	Value
\hat{A}_{33}	$\{*, *, \hat{A}_{33}\}$	$m_1(\hat{A}_{ij})$	1
\hat{A}_{23}	$\{*, \hat{A}_{33}, \hat{A}_{23}\}$	$m_2(\hat{A}_{ij})$	1
\hat{A}_{13}	$\{\hat{A}_{33}, \hat{A}_{23}, \hat{A}_{13}\}$	$m_3(\hat{A}_{ij})$	-1

Result: $m_1(\hat{C}_3) = -1$ as required.

VI. The Toy Model

An example of a non-contextual non-multiplicative model; Sequential multiplicativity is demonstrated.

- ▶ Initial: $|\psi\rangle$.
- ▶ 'hidden variable': Choose $c = +1$ for heads, $c = -1$ for tails, after a coin toss.
- ▶ Predictions/Assignments: For an operator $\hat{p}' \in \{\hat{A}_{ij}, \hat{R}_i, + \text{their products such as } \hat{C}_j (\forall i, j)\}$ check if $\exists \lambda$, s.t. $\hat{p}'|\psi\rangle = \lambda|\psi\rangle$. If $\exists \lambda$, then assign λ as the value. Else, assign c.
- ▶ Update: Say \hat{p} was observed. If \hat{p} is s.t. $\hat{p}|\psi\rangle = \lambda|\psi\rangle$, then leave the state unchanged. Else, find $|p_{\pm}\rangle$ (eigenkets of \hat{p}), s.t. $\hat{p}|p_{\pm}\rangle = \pm|p_{\pm}\rangle$ and update the state $|\psi\rangle \rightarrow |p_c\rangle$. NB: This would statistically agree with QM, for a few $|\psi\rangle$ s.

IV. Contextuality - PM Test

Kochen-Specker proved that non-contextual theories, are inconsistent with QM [6]. Mermin and Peres showed this for a four-level system [4].

- ▶ Simplified Proof: Consider the following operators.

$$\hat{A}_{ij} \doteq \begin{bmatrix} \hat{\sigma}_z \otimes \hat{\mathbb{I}} & \hat{\mathbb{I}} \otimes \hat{\sigma}_z & \hat{\sigma}_z \otimes \hat{\sigma}_z \\ \hat{\mathbb{I}} \otimes \hat{\sigma}_x & \hat{\sigma}_x \otimes \hat{\mathbb{I}} & \hat{\sigma}_x \otimes \hat{\sigma}_x \\ \hat{\sigma}_z \otimes \hat{\sigma}_x & \hat{\sigma}_x \otimes \hat{\sigma}_z & \hat{\sigma}_y \otimes \hat{\sigma}_y \end{bmatrix}$$

Note that operators along a given row (column) commute.

$$\hat{R}_i \equiv \prod_j \hat{A}_{ij} = \mathbb{I} \quad (5)$$

$$\hat{C}_j \equiv \prod_i \hat{A}_{ij} = \begin{cases} +\hat{\mathbb{I}} & (j \neq 3) \\ -\hat{\mathbb{I}} & (j = 3) \end{cases} \quad (6)$$

It entails that $\prod_{k=1,2,3} \hat{R}_k \hat{C}_k = -\hat{\mathbb{I}}$, whereas non-contextual models would yield $+1$.

NB: We also assumed multiplicativity.

To facilitate experimental validation, it has been shown that non-contextual models satisfy Eq. 7, while QM yields $\langle \hat{X}_{PM} \rangle = 6$.

$$\langle \hat{X}_{PM} \rangle = \langle \hat{R}_1 \rangle + \langle \hat{R}_2 \rangle + \langle \hat{R}_3 \rangle + \langle \hat{C}_1 \rangle + \langle \hat{C}_2 \rangle - \langle \hat{C}_3 \rangle \leq 4 \quad (7)$$

- ▶ Conclusion: Deterministic theories, that satisfy both (a) non-contextuality and (b) multiplicativity, are inconsistent with QM.

VII. Results and Conclusion

- ▶ Contextuality is not necessary.
 - ▶ The properties 'multiplicativity' and 'sequential multiplicativity' were identified, defined and proven where they hold.
 - ▶ Demonstrated that 'non-multiplicativity' is an alternative to 'contextuality', by constructing a 'non-contextual' theory, consistent with QM predictions.
 - ▶ Proposed a Minimalist HV theory; simplifies predictions.
- ▶ Tests of Determinism and Contextuality
 - ▶ Optimised phase-space GHZ
 - ▶ GHZ extension to a test of contextuality
 - ▶ PM extension to phase space (independently re-discovered)
- ▶ Measurements in Bohmian Mechanics
 - ▶ Generalised the Hamiltonian based measurement scheme to continuous variables
 - ▶ Analytic/graphical solution to measuring entangled spins using SG
 - ▶ Analytic/graphical proof of consistency of position measurements
 - ▶ Alternative proof of spins can't be associated with particles, only with wavefunctions

Bohmian Mechanics, being a deterministic and precise theory, has been successfully used to probe fundamental concepts in Quantum Mechanics and has radically clarified them (to the author at least).

References

- [1] John S. Bell. On the einstein-podolsky-rosen paradox. *Physics*, 1:195–200, 1964.
- [2] David Bohm. A suggested interpretation of the quantum theory in terms of "hidden" variables. i. *Phys. Rev.*, 85:166–179, Jan 1952.
- [3] A. Einstein, B. Podolsky, and N. Rosen. Can quantum-mechanical description of physical reality be considered complete? *Phys. Rev.*, 47:777–780, May 1935.
- [4] N. David Mermin. Simple unified form for the major no-hidden-variables theorems. *Phys. Rev. Lett.*, 65:3373–3376, Dec 1990.
- [5] A. Peres. *Quantum Theory: Concepts and Methods (Fundamental Theories of Physics)*. Springer, 1993.
- [6] E. Specker Simon Kochen. The problem of hidden variables in quantum mechanics. *Indiana Univ. Math. J.*, 17:59–87, 1968.