# An Alternative to Contextuality

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#### I. Background

- ► Einstein: 'locality' ⇒ Quantum Mechanics (QM) is incomplete [3].
- ▶ Bell: 'locality'  $\Longrightarrow \langle \hat{B} \rangle \leq 2$ ; For some  $|\psi\rangle$ , QM  $\Longrightarrow$   $\langle \hat{B} \rangle = 2\sqrt{2}$ [1]. Verified experimentally (without loop holes in 2015)
- ► Comment: At roughly the same time, various physicists had produced proofs of the claim that one can't complete QM satisfactorily, that a sensible complete 'hidden variable' (HV) description of nature was impossible.
- Bohmian Mechanics (BM): a HV description, that (i) 'completes' QM in a simple, clear, precise but non-local manner, and (ii) is deterministic [2].
- ightharpoonup Defn: Deterministic  $\equiv$  If in principle, the outcome of measuring each observable is predictable, given the HVs.
- ► Comment: Bell's inequality requires entanglement in some form, to prove Einstein's notion of locality incorrect. Recently, another peculiar feature of QM has been identified, namely contextuality.
- ▶ Impl Defn: Context  $\equiv$  If  $[\hat{A}, \hat{B}] = 0$ and [A, C] = 0 but  $[B, C] \neq 0$ , then possible contexts are A, A and B or A and C [5].
- ightharpoonup Defn: Non-contextual  $\equiv$  Value an operator takes, depends only on the state (including 'hidden variables') and the choice of the operator A (not it's context) [5].
- ightharpoonup Defn: Contextual  $\equiv$  Value an operator takes, depends on it's context [5].
- ► Comment: This notion arises, atleast in certain explicit constructions, where one is unable to assign values to operators, consistent with predictions of QM.
- ► Aim: Understand how a deterministic theory can be consistent with the notion of contextuality.

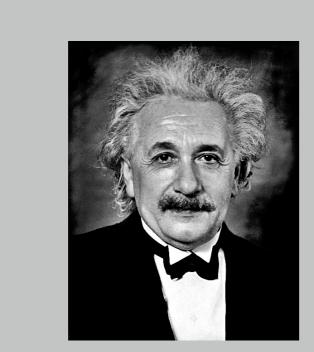


Figure 1: A. Einstein

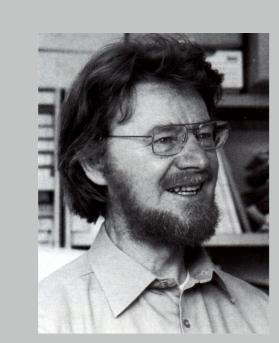


Figure 2: J. Bell



Figure 3: D. Bohm

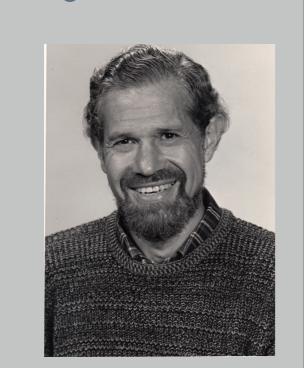


Figure 4: S. B. Kochen

#### III. Multiplicativity

- ightharpoonup Defn: Compatible operators  $\equiv$  Two observables  $\hat{A}$  and  $\hat{B}$  are mutually compatible if, given that the system is prepared in a state s.t. measurement  $\hat{A}$  yields repeatable results, measurement of  $\hat{B}$  doesn't change the result of measuring  $\hat{A}$ . For projective measurements, its equivalent to  $[\hat{A}, \hat{B}] = 0$ .
- Defn: Multiplicativity  $\equiv$  For compatible operators  $\hat{B}_i$ , a model is multiplicative iff

$$f(m_1(\hat{B}_1), m_1(\hat{B}_2), \dots, m_1(\hat{B}_n)) = (1)$$

$$m_1(f(\hat{B}_1, \hat{B}_2, \dots, \hat{B}_n), (2)$$

where  $m_i(\hat{*})$  represents the assigned value of the operator, and j encodes the sequence of measurement. Note that this is an ontological statement and can't be experimentally tested.

 $\triangleright$  Defn: Sequential Multiplicativity  $\equiv$  For compatible operators  $B_i$ , a model is sequentially multiplicative iff

$$f(m_{k_1}(\hat{B}_1), m_{k_2}(\hat{B}_2), \dots, m_{k_n}(\hat{B}_n)) = (3)$$

$$m_1(f(\hat{B}_1, \hat{B}_2, \dots, \hat{B}_n), (4)$$

where  $\mathbf{k} \equiv (k_1, k_2, \dots, k_n) \in ((1, 2, \dots, n) + \text{all possible})$ permutations),  $m_i(\hat{*})$  represents the assigned value of the operator, and j encodes the sequence of measurement.

- $\blacktriangleright$  Example:  $\hat{B}_1 = \hat{\sigma}_{\scriptscriptstyle X} \otimes \hat{\sigma}_{\scriptscriptstyle Y}, \ \hat{B}_2 = \hat{\sigma}_{\scriptscriptstyle Y} \otimes \hat{\sigma}_{\scriptscriptstyle X}$  so that  $\hat{C}=\hat{B}_1\hat{B}_2=\hat{\sigma}_z\otimes\hat{\sigma}_z.\,\,|\psi
  angle=|00
  angle$ , so that  $m_1(\hat{C})=1$ , while  $m_1(\hat{B}_i) = \pm 1$ . If say  $m_1(\hat{B}_1) = -1$ , then  $\psi \to \text{(figure this) so that entails } m_2(\hat{B}_2) = -1 \text{ as well,}$ consistent with  $m_1(\hat{C}) = m_1(\hat{B}_1)m_2(\hat{B}_2)$ .
- ► Claim: Quantum Mechanics is sequentially multiplicative.

#### IV. Contextuality - PM Test

Kochen-Specker proved that non-contextual theories, are inconsistent with QM [6]. Mermin and Peres showed this for a four-level system [4].

Simplified Proof: Consider the following operators.

$$A_{ij} \doteq \begin{bmatrix} \sigma_z \otimes \mathbb{I} & \mathbb{I} \otimes \sigma_z & \sigma_z \otimes \sigma_z \\ \mathbb{I} \otimes \sigma_x & \sigma_x \otimes \mathbb{I} & \sigma_x \otimes \sigma_x \\ \sigma_z \otimes \sigma_x & \sigma_x \otimes \sigma_z & \sigma_y \otimes \sigma_y \end{bmatrix}$$

Note that operators along a given row (column) commute.

$$R_i \equiv \prod_i A_{ij} = \mathbb{I}$$
 (

$$R_{i} \equiv \prod_{j} A_{ij} = \mathbb{I}$$

$$C_{j} \equiv \prod_{i} A_{ij} = \begin{cases} +\mathbb{I} & (j \neq 3) \\ -\mathbb{I} & (j = 3) \end{cases}$$
(6)

It entails that  $\prod_{k=1,2,3} R_k C_k = -\mathbb{I}$ , whereas non-contextual models would yield +1.

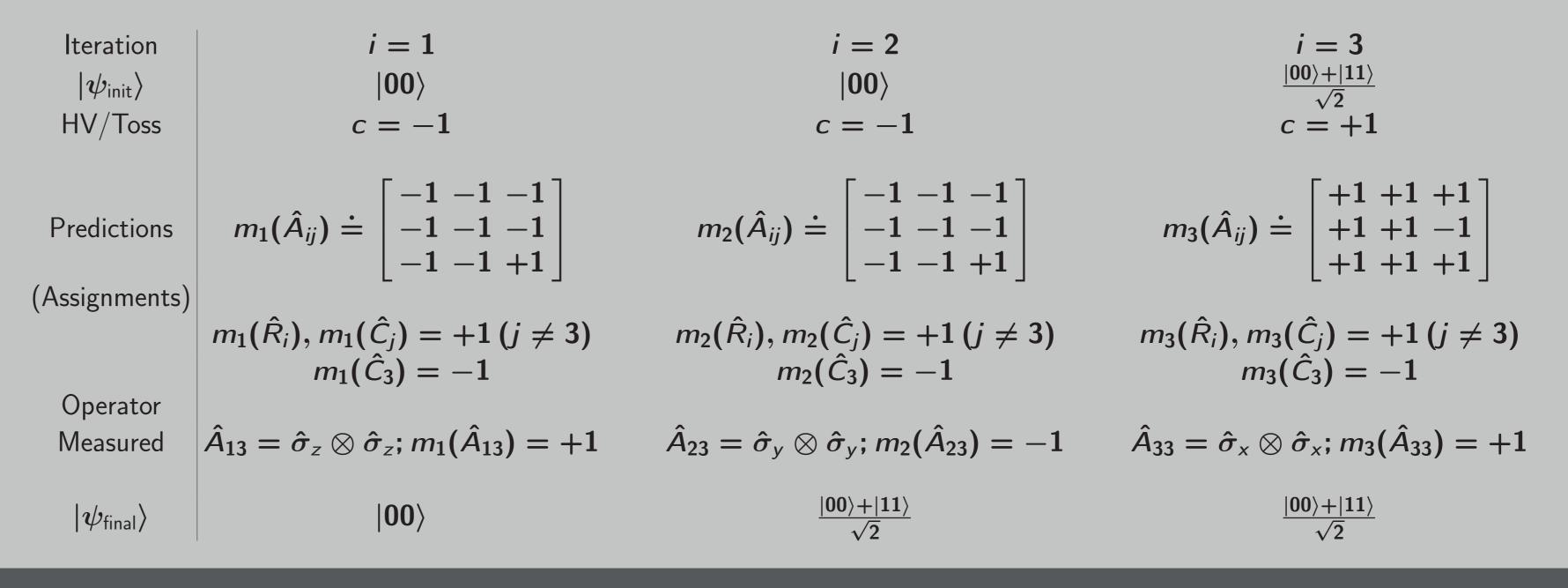
NB: We also assumed multiplicativity.

To facilitate experimental validation, it has been shown that non-contextual models satisfy Eq. 7, while QM yields  $\langle \chi_{PM} \rangle = 6.$ 

$$\langle \chi_{PM} \rangle = \langle R_1 \rangle + \langle R_2 \rangle + \langle R_3 \rangle + \langle C_1 \rangle + \langle C_2 \rangle - \langle C_3 \rangle \le 4$$
(7)

► Conclusion: Deterministic theories, that satisfy both (a) non-contextuality and (b) multiplicativity, are inconsistent with QM.

## The Toy Model — Example



### II. Overview

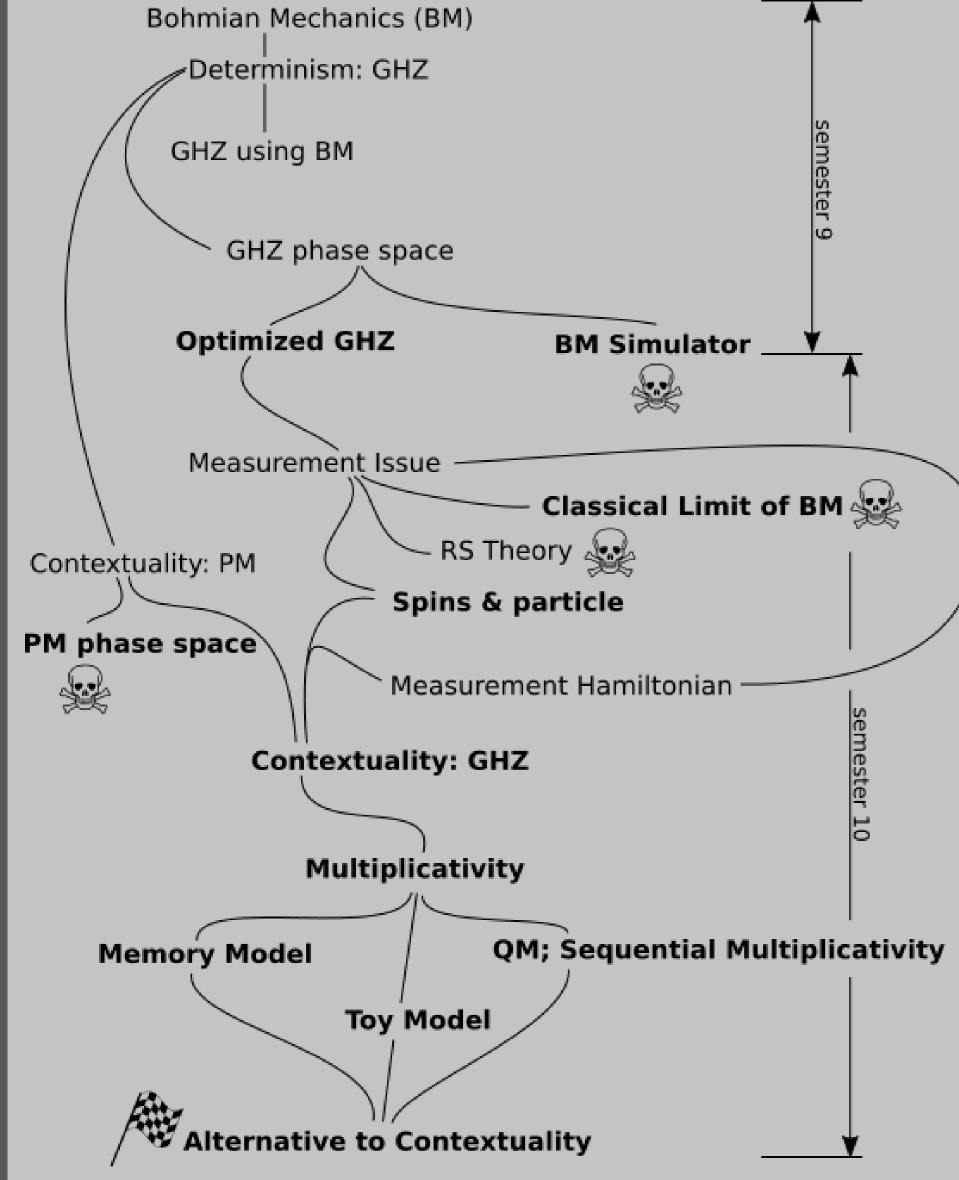


Figure 5: Exploration flow: Boldface titles represent new results

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### V. Contextuality - Memory Model

An example of a contextual and non-multiplicative model; Sequential multiplicativity has been assumed.

- ▶ Initial: The assignment is as given in the first Mat in Eq. 9.
- ► Remark: The system is assumed to be capable of remembering the last three observables that were measured.
- ► Algorithm: Upon measurement of an observable, (i) yield the value as saved in the matrix, (ii) append the observable in the 3 element memory and (iii) update the matrix, once the context (set of commuting observables) is known, to satisfy the PM requirements.

$$m_1(\hat{A}_{ij}) = m_2(\hat{A}_{ij}) \doteq \begin{bmatrix} 1 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix}, m_3(\hat{A}_{ij}) \doteq \begin{bmatrix} 1 & 1 & -1 \\ 1 & 1 & 1 \\ 1 & 1 & 1 \end{bmatrix}$$

► For example:

Operator
 Updated Array
 Assignment Value

 
$$\hat{A}_{33}$$
 $\{*, *, \hat{A}_{33}\}$ 
 $m_1(\hat{A}_{ij})$ 
 1

  $\hat{A}_{23}$ 
 $\{*, \hat{A}_{33}, \hat{A}_{23}\}$ 
 $m_2(\hat{A}_{ij})$ 
 1

  $\hat{A}_{13}$ 
 $\{\hat{A}_{33}, \hat{A}_{23}, \hat{A}_{13}\}$ 
 $m_3(\hat{A}_{ij})$ 
 -1

Result:  $m_1(C_3) = -1$  as required.

### VI. The Toy Model

An example of a non-contextual non-multiplicative model; Sequential multiplicativity is demonstrated.

- ightharpoonup Initial:  $|\psi\rangle$ .
- ightharpoonup 'hidden variable': Choose c=+1 for heads, c=-1 for tails, after a coin toss.
- ightharpoonup Predictions/Assignments: For an operator  $\hat{p}' \in \{\hat{A}_{ii}, \hat{R}_i, + \}$ their products such as  $\hat{C}_i$  ( $\forall i, j$ ) check if  $\exists$  a  $\lambda$ , s.t.  $\hat{p}' | \psi \rangle = \lambda | \psi \rangle$ . If  $\exists$  a  $\lambda$ , then assign  $\lambda$  as the value. Else, assign c.
- lacksquare Update: Say  $\hat{p}$  was observed. If  $\hat{p}$  is s.t.  $\hat{p} | \psi \rangle = \lambda | \psi \rangle$ , then leave the state unchanged. Else, find  $|p_{\pm}\rangle$  (eigenkets of  $\hat{p}$ ), s.t.  $\hat{p} | p_{\pm} \rangle = \pm | p_{\pm} \rangle$  and update the state  $| \psi \rangle \to | p_c \rangle$ . NB: This would statistically agree with QM, for a few  $|\psi\rangle$ s.

### VII. Results and Conclusion

- Contextuality is not necessary.
  - ▶ The properties 'multiplicativity' and 'sequential multiplicativity' were identified, defined and proven where they hold.
- Demonstrated that 'non-multiplicativity' is an alternative to 'contextuality', by constructing a 'non-contextual' theory, consistent with QM predictions.
- Proposed a Minimalistic HV theory; simplifies predictions.
- ► Tests of Determinism and Contextuality
  - Optimized phase-space GHZ
  - □ GHZ extension to a test of contextuality
  - ▶ PM extension to phase space (independently re-discovered)
- ► Measurements in Bohmain Mechanics
- ▶ Generalized the Hamiltonian based measurement scheme to continuous variables
- Analytic/graphical solution to measuring entangled spins using SG
- Analytic/graphical proof of consistency of position measurements
- ▶ Alternative proof of spins can't be associated with particles, only with wavefunctions

Bohmian Mechanics, being a deterministic and precise theory, has been successfully used to probe fundamental concepts in Quantum Mechanics and has radically clarified them (to the author atleast).

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