

$$C^{i_1 \dots i_p}_{j_1 \dots j_q}(\vec{x}') =$$

$$\text{Rank} \quad \partial_{k_1} x'^{i_1} \dots \partial_{k_p} x'^{i_p} \cdot \partial_{g_1} x'^{j_1} \dots \partial_{g_q} x'^{j_q} C^{i_1 \dots i_p}_{j_1 \dots j_q}(\vec{x}') \quad \text{--- This transforms as a } (p+q, q) \text{ tensor}$$

$$(0,0) \text{ scalar} \quad \phi'(\vec{x}') = \phi(\vec{x})$$

$$(1,0) \text{ contravariant vector}$$

$$(0,1) \text{ covariant vector}$$

$$A^{i_1 \dots i_p}_{j_1 \dots j_q} B^{k_1 \dots k_r}_{l_1 \dots l_s} \quad \text{--- This transforms as a } (p+r, q+s) \text{ tensor}$$

$$A^{i_1 \dots i_p}_{j_1 \dots j_q}$$

$$A^{i_1 \dots i_p}_{j_1 \dots j_q}(\vec{x}') = \partial_{k_1} x'^{i_1} \dots \partial_{k_p} x'^{i_p} \partial_{l_1} x'^{j_1} \dots \partial_{l_q} x'^{j_q} A^{k_1 \dots k_p}_{l_1 \dots l_q}(\vec{x})$$

$$\frac{\partial x'^{i_1}}{\partial x^{k_1}} \dots \frac{\partial x'^{i_p}}{\partial x^{k_p}} \frac{\partial x'^{j_1}}{\partial x^{l_1}} \dots \frac{\partial x'^{j_q}}{\partial x^{l_q}} = \frac{\partial x'^{l_1}}{\partial x^{k_1}} = \delta^{l_1}_{k_1}$$

$$= A^{k_1 \dots k_p}_{l_1 \dots l_q}(\vec{x}) \quad \text{which is consistent}$$

Claim: Take some point p on the manifold. \vec{x}_0 is the coordinate of p in some coordinate system.

Then \exists a coordinate transf $\vec{x} \rightarrow \vec{x}'$ s.t. $g'_{ij}(\vec{x}_0) = \text{diag}(\pm 1, \pm 1, \dots, \pm 1)$

NB: For continuous metrics, the number of $+$ & $-$ entries will be constant.

$$= f^i(\vec{x}') \Rightarrow x'_0 \equiv f^i(\vec{x}'_0)$$

$$x^i = x^i_0 + A^i_k (x'^k - x'^k_0) + B^i_{kl} (x'^k - x'^k_0)(x'^l - x'^l_0) + \dots$$

Here we expand about x_0 .

$$g'_{kl}(\vec{x}'_0) = \partial'_k x^i \partial'_l x^j g_{ij} \Big|_{\vec{x}' = \vec{x}'_0}$$

$$= A^i_k A^j_l g_{ij}(\vec{x}_0) \quad (\text{Take } \partial' \text{ & should be clear})$$

$$= (A^T g A)_{kl} \quad (\text{in matrix notation})$$

Notation: Regard g_{ij} as matrix; i as row index, j as column index

- Recall: S is symmetric; Thus can be diagonalized with orthogonal matrices.

$$\text{viz. } g = S^T g_d S \quad \text{s.t. } S^T S = 1$$

$$g_d = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix}$$

Non-singular metric only, i.e. eigenvalues are non-zero.

$$\text{Defn: } R \equiv R^T \equiv \begin{pmatrix} \sqrt{|\lambda_1|} & & 0 \\ & \ddots & \\ 0 & & \sqrt{|\lambda_n|} \end{pmatrix}; \Rightarrow R^T \eta R = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix} \quad \text{where } \eta = \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{pmatrix}$$

$$\text{Thus, } g = S^T g_d S = S^T R^T \eta R S$$

$$g'_{kl} = (A^T g A)_{kl} = (A^T S^T R^T \eta R A)_{kl} \Rightarrow \text{for } A = (R S)^{-1} \text{ we'll get } g'_{kl}(\vec{x}'_0) = \eta_{kl}$$

This is not unique; eg. $\eta = 1$, then $A = (URS)^{-1}$ will get the right form, where

Signature:
 When all +, then euclidean
 when one -, then lorentzian

- NB: (1) Locally then, there's no more information
 (2) This will not distinguish b/w a plane & the surface of a sphere.
 (3) This is more than diagonalization; A is more than just a PTR like matrix.

Defⁿ: $K_{ijk} = \partial_i g_{jk}$ $K \rightarrow K'$

$$K'_{i_1 i_2 i_3} = \partial'_{i_1} (g'_{i_2 i_3}(\vec{x}'))$$

$$= \partial'_{i_1} (\partial'_{i_2} x^{j_2} \partial'_{i_3} x^{j_3} g_{j_2 j_3}(\vec{x}'))$$

$$= \partial'_{i_1} \partial'_{i_2} x^{j_2} \partial'_{i_3} x^{j_3} g_{j_2 j_3} + \partial'_{i_1} x^{j_2} \partial'_{i_2} \partial'_{i_3} x^{j_3} g_{j_2 j_3} + \partial'_{i_1} x^{j_2} \partial'_{i_2} x^{j_3} \partial'_{i_3} g_{j_2 j_3}$$

$$\underline{\partial'_{i_1} \partial'_{i_2} x^{j_2} \partial'_{i_3} x^{j_3} g_{j_2 j_3}}$$

This is the term we want (why? we'll see soon)

chain rule

$$= \partial'_{i_2} x^{j_2} \partial'_{i_3} x^{j_3} \partial'_{i_1} x^{i_1} \partial_{i_1} g_{j_2 j_3}$$

$$\left(\because \frac{\partial x^{j_1}}{\partial x'^{i_1}} \frac{\partial g_{j_2 j_3}}{\partial x^{i_1}} \right)$$

$$= \frac{\partial g_{j_2 j_3}}{\partial x'^{i_1}}$$

rearrange

$$= \underline{\partial'_{i_1} x^{j_1} \partial'_{i_2} x^{j_2} \partial'_{i_3} x^{j_3} \partial_{i_1} g_{j_2 j_3}} \rightarrow K_{j_2 j_3}$$

This is how K_{ijk} should transform if K_{ijk} was a tensor.

However, since the other terms are non-zero, $\Rightarrow K_{ijk}$ is not a tensor under general coordinate transformation.

Idea: Take a linear combination of K_{ijk} to construct a tensor.

: This is not possible. (even polynomials of K_{ijk})

Proof sketch: Assume we can show that given a point, \exists a coordinate transformation (CT) s.t. K vanishes. Therefore, it must vanish in all coordinates (\because tensors transform in that way). Therefore it's impossible to find such a tensor (it'll be zero essentially), (if this argument is repeated for all points)

Recall: $x^i = x^i_0 + A^i_j (x'^j - x^j_0) + B^i_{jk} (x'^j - x^j_0)(x'^k - x^k_0) + \dots$

We need 2 derivatives $\Rightarrow B^i_{jk}$ will be important

We want $K'_{ijk} = 0$, so $\frac{n(n+1)}{2} \cdot n$ eq's are required to be zero.

constants is also $\frac{n(n+1)}{2} \cdot n$ ($\because B^i_{jk}$ also has similar symmetric aspects as K_{ijk})

B^i_{jk} will be unique. ∞

Exercise: Find B^i_{jk} explicitly, to make $K'^i_{j,k} = 0$

Idea: Take 2 derivatives of g_{ij}

$$S_{ijkl} = \partial_i \partial_j g_{kl}$$

(Claim: This itself is not a tensor.)

: Use the old argument. See if it can be shown to be zero in some coordinate sys.

At \vec{x}_0 , # of components of $S_{ijkl} = \frac{n(n+1)}{2} \cdot \frac{n(n+1)}{2}$

Recall: $x^i = x^i_0 + A^i_j(\dots)^j + B^i_{jk}(\dots)^{jk} + C^i_{jkl}(\dots)^{jkl}$

Now # of parameters in $C^i_{jkl} = n \cdot \frac{n(n+1)(n+2)}{3!}$ & \leq (for large N)

\Rightarrow There're more constraints than parameters. This means one can't in general get

$$S_{ijkl} = 0$$

Result: So $\frac{n^2(n+1)^2}{4} - \frac{n^2(n+1)(n+2)}{6} = \frac{1}{12} n^2(n^2-1)$ — 1 for $n=2$
6 for $n=3$
70 for $n=4$

which means the tensor so constructed will have the aforesaid # of free parameters.

Defⁿ: $\Gamma^i_{jk} = \frac{1}{2} g^{il} (\partial_j g_{lk} + \partial_k g_{lj} - \partial_l g_{jk}) :=$ Christoffel symbol / connection

$$R^i_{jkl} = \partial_l \Gamma^i_{jk} - \partial_k \Gamma^i_{jl} + \Gamma^m_{jk} \Gamma^i_{lm} - \Gamma^m_{jl} \Gamma^i_{km} := \text{Riemann tensor}$$

Exercise: $\Gamma'^l_{mn}(\vec{x}) = \partial_i x'^l \partial'_m x^i \partial'_n x^k \Gamma^i_{jk}(\vec{x}) + \partial_n x^l \partial'_m \partial'_n x^k$

NB: This is not a tensor \because of the second term.

: R^i_{jkl} is a tensor of rank (1,3)

: R^i_{jkl} has $\frac{1}{12} n^2(n^2-1)$ independent parameters.