

# Quantum Measurements and Relativity

Kishor Bharti

Centre for Quantum Technologies. Singapore October 23, 2016

· Quantum mechanics not compatible with Relativity

- · Quantum mechanics not compatible with Relativity
- · What if it does?

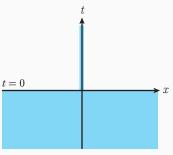
- · Quantum mechanics not compatible with Relativity
- What if it does?  $\implies$  Paradox

- · Quantum mechanics not compatible with Relativity
- What if it does?  $\implies$  Paradox
  - · What is the paradox?

 $\cdot$  According to QM, measurement  $\implies$  instantaneous collapse

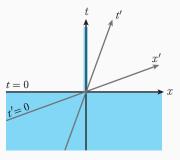
- $\cdot$  According to QM, measurement  $\implies$  instantaneous collapse
- E.g.: Particle in momentum eigenstate

- $\cdot$  According to QM, measurement  $\implies$  instantaneous collapse
- E.g.: Particle in momentum eigenstate



• In another frame of reference?

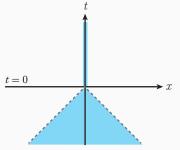
• In another frame of reference?



 $\boldsymbol{\cdot}$  What if particle confined in past light cone?

- · What if particle confined in past light cone?
  - No longer momentum eigenstate

- What if particle confined in past light cone?
  - · No longer momentum eigenstate



$$\boxed{|\Psi_{-}\rangle = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle\right)}$$

• E.g.: 2 distinguishable particle singlet state, space-like separated

$$\boxed{\ket{\Psi_-}=rac{1}{\sqrt{2}}\left(\ket{\uparrow\downarrow}-\ket{\downarrow\uparrow}
ight)}$$

• Who collapses state first?

$$\boxed{|\Psi_{-}\rangle = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle\right)}$$

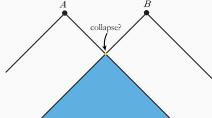
- · Who collapses state first?
  - If Alice measures first, ∃ intertial frame s.t. Bob measures first

$$\boxed{|\Psi_{-}\rangle = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle\right)}$$

- · Who collapses state first?
  - $\cdot$  If Alice measures first,  $\exists$  intertial frame s.t. Bob measures first
- Collapse at intersection of light cone?

$$\boxed{|\Psi_{-}\rangle = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle\right)}$$

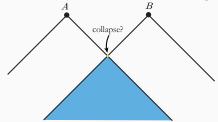
- · Who collapses state first?
  - $\cdot$  If Alice measures first,  $\exists$  intertial frame s.t. Bob measures first
- Collapse at intersection of light cone?



• E.g.: 2 distinguishable particle singlet state, space-like separated

$$\boxed{|\Psi_{-}\rangle = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle\right)}$$

- · Who collapses state first?
  - $\cdot$  If Alice measures first,  $\exists$  intertial frame s.t. Bob measures first
- · Collapse at intersection of light cone?



Collapse happens in the past ⇒ A, B measures product state ⇒ Obeys Bell's inequality ⇒ Paradox

Landau and Peierls:

Theorem (Relativistic Uncertainty Relation)

Measurement of momentum cannot be instantaneous. Momentum with accuracy  $\Delta p$  cannot take time less than  $\Delta t$  such that

$$\Delta p \, \Delta t \geq \frac{\hbar}{c} \, .$$

Landau and Peierls:

Theorem (Relativistic Uncertainty Relation)

Measurement of momentum cannot be instantaneous. Momentum with accuracy  $\Delta p$  cannot take time less than  $\Delta t$  such that

$$\Delta p \, \Delta t \geq \frac{\hbar}{c} \, .$$

· Measurement of non-local operator cannot be instantaneous.

$$\left|\Psi_{\alpha\beta}\right\rangle = \alpha\left|\uparrow\downarrow\right\rangle + \beta\left|\downarrow\uparrow\right\rangle \;.$$

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle$$
.

• Suppose A and B can measure  $S^2$  instantaneously.

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle$$
.

- Suppose A and B can measure  $S^2$  instantaneously.
- · Rewrite state in total spin basis:

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$
,

where

$$|0,0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$
 and  $|2,0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$ .

8

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

Measurement 1:

9

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- Measurement 1:
  - 1. A and B measure  $S^2$

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- · Measurement 1:
  - 1. A and B measure  $S^2$
  - 2. A measures  $S_z \implies \frac{1}{2}$  probability get  $\uparrow$  or  $\downarrow$

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- · Measurement 1:
  - 1. A and B measure  $S^2$
  - 2. A measures  $S_z \implies \frac{1}{2}$  probability get  $\uparrow$  or  $\downarrow$
- · Measurement 2:

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- · Measurement 1:
  - 1. A and B measure  $S^2$
  - 2. A measures  $S_z \implies \frac{1}{2}$  probability get  $\uparrow$  or  $\downarrow$
- · Measurement 2:
  - 1. B flips his spin

$$\left|\Psi'_{\alpha\beta}\right\rangle = \alpha\left|2,1\right\rangle + \beta\left|2,-1\right\rangle$$

9

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- · Measurement 1:
  - 1. A and B measure  $S^2$
  - 2. A measures  $S_z \implies \frac{1}{2}$  probability get  $\uparrow$  or  $\downarrow$
- · Measurement 2:
  - 1. B flips his spin

$$|\Psi'_{\alpha\beta}\rangle = \alpha |2,1\rangle + \beta |2,-1\rangle$$

2. A and B measure  $S^2$ 

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- · Measurement 1:
  - 1. A and B measure  $S^2$
  - 2. A measures  $S_z \implies \frac{1}{2}$  probability get  $\uparrow$  or  $\downarrow$
- · Measurement 2:
  - 1. B flips his spin

$$|\Psi'_{\alpha\beta}\rangle = \alpha |2,1\rangle + \beta |2,-1\rangle$$

- 2. A and B measure  $S^2$
- 3. A measures  $S_z \implies |\alpha|^2$  probability get  $\uparrow$ ,  $|\beta|^2$  get  $\downarrow$

g

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

E.g.: Measure total spin of 2 localised spin- $\frac{1}{2}$  particles with state

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

· B can send superluminal message by flipping in his lab!

$$|\Psi_{\alpha\beta}\rangle = \frac{\alpha+\beta}{\sqrt{2}}|2,0\rangle + \frac{\alpha-\beta}{\sqrt{2}}|0,0\rangle$$

- · B can send superluminal message by flipping in his lab!
- Unless  $|\alpha|^2 = |\beta|^2 = \frac{1}{2}$ .

#### Non-local Measurements

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\implies$  state-specific measurement

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\Longrightarrow$  state-specific measurement

· Measurement interaction

$$\mathcal{H}_{int} = g(t) \left[ S_z^A P_z^A + S_z^B P_z^B \right]$$

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\Longrightarrow$  state-specific measurement

Measurement interaction

$$\mathcal{H}_{int} = g(t) \left[ S_z^A P_z^A + S_z^B P_z^B \right]$$

where

$$g(t) = \begin{cases} \frac{1}{T}, & 0 \le t \le T \\ 0, & \text{otherwise} \end{cases}$$

What if  $|\alpha|^2 = |\beta|^2 = \frac{1}{2}$ ?  $\implies$  state-specific measurement

Measurement interaction

$$\mathcal{H}_{int} = g(t) \left[ S_z^A P_z^A + S_z^B P_z^B \right]$$

where

$$g(t) = \begin{cases} \frac{1}{T}, & 0 \le t \le T \\ 0, & \text{otherwise} \end{cases}$$

· Measuring device state

$$Q_z^A + Q_z^B = 0 = P_z^A - P_z^B$$
.

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\Longrightarrow$  state-specific measurement

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\implies$  state-specific measurement

• Measures 
$$Q_z^A + Q_z^B \implies$$
 Measures  $S_z = S_z^A + S_z^B$ 

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\implies$  state-specific measurement

- Measures  $Q_z^A + Q_z^B \implies$  Measures  $S_z = S_z^A + S_z^B$
- Do this for  $S_x$  and  $S_y$  too.

What if 
$$|\alpha|^2 = |\beta|^2 = \frac{1}{2}$$
?  $\Longrightarrow$  state-specific measurement

- Measures  $Q_z^A + Q_z^B \implies$  Measures  $S_z = S_z^A + S_z^B$
- Do this for  $S_x$  and  $S_y$  too.
- If  $S_x = S_y = S_z$ , then  $S^2 = 0 \implies \text{verified } |0,0\rangle$

 $\exists$  operator measurable instantaneously?  $\Longrightarrow$  Operator-specific measurement

 $\exists$  operator measurable instantaneously?  $\Longrightarrow$  Operator-specific measurement

• Suppose ∃ such non-degenerate operator *W* with eigenstate

$$|\Psi_{\alpha\beta}\rangle = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle$$
 with  $\alpha, \beta \in \mathbb{R}$ .

 $\exists$  operator measurable instantaneously?  $\Longrightarrow$  Operator-specific measurement

• Suppose ∃ such non-degenerate operator W with eigenstate

$$|\Psi_{\alpha\beta}\rangle = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle \quad \text{with} \quad \alpha, \beta \in \mathbb{R}.$$

• The other 3 eigenstates are  $|W_1\rangle$ ,  $|W_2\rangle$  and  $|W_3\rangle$ 

 $\exists$  operator measurable instantaneously?  $\Longrightarrow$  Operator-specific measurement

· Suppose ∃ such non-degenerate operator W with eigenstate

$$|\Psi_{\alpha\beta}\rangle = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle$$
 with  $\alpha, \beta \in \mathbb{R}$ .

- The other 3 eigenstates are  $|W_1\rangle$ ,  $|W_2\rangle$  and  $|W_3\rangle$
- · The space is spanned by bases

$$|\Psi_{\alpha\beta}\rangle = \alpha |\uparrow\downarrow\rangle + \beta |\downarrow\uparrow\rangle; \quad |\Psi_{\alpha\beta}^{\perp}\rangle = \beta |\uparrow\downarrow\rangle - \alpha |\downarrow\uparrow\rangle; \quad |\uparrow\uparrow\rangle; \quad |\downarrow\downarrow\rangle.$$

• Measurement: Prepare state 
$$|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} \, |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} \, |\Psi_{\alpha\beta}^{\perp}\rangle$$

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} \, |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} \, |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} \, |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} \, |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously
    - · Probability  $\frac{|lpha-eta|^2}{2}$ , getting  $|\Psi_{lphaeta}
      angle$

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously
    - Probability  $\frac{|\alpha-\beta|^2}{2}$ , getting  $|\Psi_{\alpha\beta}\rangle$
    - Probability  $\frac{|\alpha+\beta|^2}{2}$ , getting any of the other eigenstates  $W_i$

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously
    - Probability  $\frac{|\alpha-\beta|^2}{2}$ , getting  $|\Psi_{\alpha\beta}\rangle$
    - · Probability  $\frac{|\alpha+\beta|^2}{2}$ , getting any of the other eigenstates  $W_i$
  - 2. A measures  $S_z^A$

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously
    - Probability  $\frac{|\alpha-\beta|^2}{2}$ , getting  $|\Psi_{\alpha\beta}\rangle$
    - · Probability  $\frac{|\alpha+\beta|^2}{2}$ , getting any of the other eigenstates  $W_i$
  - 2. A measures  $S_z^A$
- · For Alice, probability of getting up

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} \left| \Psi_{\alpha\beta} \right\rangle + \frac{\alpha+\beta}{\sqrt{2}} \left| \Psi_{\alpha\beta}^{\perp} \right\rangle$ 
  - 1. A and B measure W instantaneously
    - · Probability  $\frac{|\alpha-\beta|^2}{2}$ , getting  $|\Psi_{\alpha\beta}\rangle$
    - · Probability  $\frac{|\alpha+\beta|^2}{2}$ , getting any of the other eigenstates  $W_i$
  - 2. A measures  $S_z^A$
- · For Alice, probability of getting up

$$P(\uparrow_{A}) = \frac{1 - 2\alpha\beta}{2}\alpha^{2} + \frac{1 + 2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \beta^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right]$$

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously
    - Probability  $\frac{|\alpha-\beta|^2}{2}$ , getting  $|\Psi_{\alpha\beta}\rangle$
    - Probability  $\frac{|\alpha+\beta|^2}{2}$ , getting any of the other eigenstates  $W_i$
  - 2. A measures  $S_z^A$
- · For Alice, probability of getting up

$$P(\uparrow_{A}) = \frac{1 - 2\alpha\beta}{2}\alpha^{2} + \frac{1 + 2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \beta^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right]$$

For Bob probability of getting up

- Measurement: Prepare state  $|0,0\rangle = \frac{\alpha-\beta}{\sqrt{2}} |\Psi_{\alpha\beta}\rangle + \frac{\alpha+\beta}{\sqrt{2}} |\Psi_{\alpha\beta}^{\perp}\rangle$ 
  - 1. A and B measure W instantaneously
    - · Probability  $\frac{|\alpha-\beta|^2}{2}$ , getting  $|\Psi_{\alpha\beta}\rangle$
    - Probability  $\frac{|\alpha+\beta|^2}{2}$ , getting any of the other eigenstates  $W_i$
  - 2. A measures  $S_z^A$
- · For Alice, probability of getting up

$$P(\uparrow_{A}) = \frac{1 - 2\alpha\beta}{2}\alpha^{2} + \frac{1 + 2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \beta^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right]$$

For Bob probability of getting up

$$P(\uparrow_B) = \frac{1 - 2\alpha\beta}{2}\beta^2 + \frac{1 + 2\alpha\beta}{2}\sum_i \left| \left\langle W_i \middle| \Psi_{\alpha\beta}^{\perp} \right\rangle \right|^2 \left[ \left| \left\langle W_i \middle| \uparrow \uparrow \right\rangle \right|^2 + \alpha^2 \left| \left\langle W_i \middle| \Psi_{\alpha\beta}^{\perp} \right\rangle \right|^2 \right]$$

What if Alice and Bob apply local unitary transformations (change relative phase)

What if Alice and Bob apply local unitary transformations (change relative phase)

$$|0,0\rangle \mapsto |2,0\rangle$$

What if Alice and Bob apply local unitary transformations (change relative phase)

$$|0,0\rangle \mapsto |2,0\rangle$$

Repeat measurement:

What if Alice and Bob apply local unitary transformations (change relative phase)

$$|0,0\rangle\mapsto|2,0\rangle$$

· Repeat measurement:

$$\begin{split} P(\uparrow_{A}) &= \frac{1+2\alpha\beta}{2}\alpha^{2} + \frac{1-2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \beta^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right] \\ P(\uparrow_{B}) &= \frac{1+2\alpha\beta}{2}\beta^{2} + \frac{1-2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \alpha^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right] \end{split}$$

What if Alice and Bob apply local unitary transformations (change relative phase)

$$|0,0\rangle\mapsto|2,0\rangle$$

· Repeat measurement:

$$P(\uparrow_{A}) = \frac{1 + 2\alpha\beta}{2}\alpha^{2} + \frac{1 - 2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \beta^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right]$$

$$P(\uparrow_{B}) = \frac{1 + 2\alpha\beta}{2}\beta^{2} + \frac{1 - 2\alpha\beta}{2}\sum_{i}\left|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\left[\left|\left\langle W_{i}\middle|\uparrow\uparrow\rangle\right|^{2} + \alpha^{2}\middle|\left\langle W_{i}\middle|\Psi_{\alpha\beta}^{\perp}\right\rangle\right|^{2}\right]$$

Demand probabilities to be the same

$$\implies \alpha\beta = 0$$
 or  $\alpha = \pm \beta$ 

Conclusion:

#### Conclusion:

A and B cannot measure W instantaneously unless each eigenstate of W is  $|\Psi_{\alpha\beta}\rangle$  with  $\alpha\beta=0$  or  $\alpha=\pm\beta$ , up to local unitary transformation.

#### Conclusion:

A and B cannot measure W instantaneously unless each eigenstate of W is  $|\Psi_{\alpha\beta}\rangle$  with  $\alpha\beta=0$  or  $\alpha=\pm\beta$ , up to local unitary transformation.

• This is necessary but not sufficient. Why? Think of  $S^2$ .

• What about an operator with an eigenvector that is not entangled ( $\alpha\beta=0$ )?

- What about an operator with an eigenvector that is not entangled ( $\alpha\beta=0$ )?
- · Can A and B simultaneously measure such an operator?

- What about an operator with an eigenvector that is not entangled ( $\alpha\beta=0$ )?
- · Can A and B simultaneously measure such an operator?
- Take two eigenvectors to be  $\frac{1}{\sqrt{2}}\left(|\uparrow\downarrow\rangle-|\downarrow\uparrow\rangle\right)$  and  $|\uparrow\uparrow\rangle$

- What about an operator with an eigenvector that is not entangled ( $\alpha\beta=0$ )?
- · Can A and B simultaneously measure such an operator?
- Take two eigenvectors to be  $\frac{1}{\sqrt{2}} \left( |\uparrow\downarrow\rangle |\downarrow\uparrow\rangle \right)$  and  $|\uparrow\uparrow\rangle$
- Can Alice and bob instantaneously measure an operator with these two states as eigenstates?

· If yes

- $\cdot$  If yes  $\Longrightarrow$ 
  - A and B prepare their spins in the state  $\left|\uparrow\uparrow\right\rangle$

- $\cdot$  If yes  $\Longrightarrow$ 
  - A and B prepare their spins in the state  $|\!\!\uparrow\uparrow\rangle$
  - $\boldsymbol{\cdot}$  A and B measure instantaneously

- $\cdot$  If yes  $\Longrightarrow$ 
  - A and B prepare their spins in the state  $|\uparrow\uparrow\rangle$
  - · A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z=1$

- $\cdot$  If yes  $\Longrightarrow$ 
  - · A and B prepare their spins in the state  $|\uparrow\uparrow\rangle$
  - · A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?

- $\cdot$  If yes  $\Longrightarrow$ 
  - $\cdot$  A and B prepare their spins in the state  $|\uparrow\uparrow\rangle$
  - A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?
  - $\cdot$   $|\uparrow\uparrow\rangle$  changes to  $|\uparrow\downarrow\rangle$

- $\cdot$  If yes  $\Longrightarrow$ 
  - $\cdot$  A and B prepare their spins in the state  $|\uparrow\uparrow\rangle$
  - A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?
  - $\cdot \mid \uparrow \uparrow \rangle$  changes to  $\mid \uparrow \downarrow \rangle$
  - This is not an eigenstate of the chosen non-local operator!

- $\cdot$  If yes  $\Longrightarrow$ 
  - A and B prepare their spins in the state  $|\!\!\uparrow\uparrow\rangle$
  - A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?
  - $\cdot \mid \uparrow \uparrow \rangle$  changes to  $\mid \uparrow \downarrow \rangle$
  - · This is not an eigenstate of the chosen non-local operator!
  - The non-local operator has a chance of leaving the spins in the state  $\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle |\downarrow\uparrow\rangle)$

- $\cdot$  If yes  $\Longrightarrow$ 
  - · A and B prepare their spins in the state  $|\uparrow\uparrow\rangle$
  - · A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?
  - $\cdot \mid \uparrow \uparrow \rangle$  changes to  $\mid \uparrow \downarrow \rangle$
  - · This is not an eigenstate of the chosen non-local operator!
  - The non-local operator has a chance of leaving the spins in the state  $\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle |\downarrow\uparrow\rangle)$
  - $\cdot$  Now Alice has a chance of obtaining -1 when she measures  $\sigma_{z}$

- $\cdot$  If yes  $\Longrightarrow$ 
  - A and B prepare their spins in the state  $|\!\!\uparrow\uparrow\rangle$
  - A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?
  - $\cdot \mid \uparrow \uparrow \rangle$  changes to  $\mid \uparrow \downarrow \rangle$
  - · This is not an eigenstate of the chosen non-local operator!
  - The non-local operator has a chance of leaving the spins in the state  $\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle |\downarrow\uparrow\rangle)$
  - · Now Alice has a chance of obtaining -1 when she measures  $\sigma_z$
  - Bob can send superluminal signal to Alice!

- $\cdot$  If yes  $\Longrightarrow$ 
  - $\cdot$  A and B prepare their spins in the state  $|\uparrow\uparrow\rangle$
  - · A and B measure instantaneously
  - Alice measures  $\sigma_z$  on her spin and she is certain to obtain  $\sigma_z = 1$
  - But, what if just before the measuring the non-local operator, Bob flips his spin?
  - $\cdot \mid \uparrow \uparrow \rangle$  changes to  $\mid \uparrow \downarrow \rangle$
  - · This is not an eigenstate of the chosen non-local operator!
  - The non-local operator has a chance of leaving the spins in the state  $\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle |\downarrow\uparrow\rangle)$
  - · Now Alice has a chance of obtaining -1 when she measures  $\sigma_z$
  - Bob can send superluminal signal to Alice! ⇒ Alice and Bob can't measure such an operator instantaneously!

#### Conclusion:

A and B cannot measure W instantaneously unless each eigenstate of W is  $|\Psi_{\alpha\beta}\rangle$  with  $\alpha=\pm\beta$ , up to local unitary transformation.

 $\cdot$  Collapse is not Lorentz invariant  $\implies$ 

 $\cdot$  Collapse is not Lorentz invariant  $\implies$  Observers in different frames disagree about collapse

- $\cdot$  Collapse is not Lorentz invariant  $\implies$  Observers in different frames disagree about collapse
- So what?

- $\cdot$  Collapse is not Lorentz invariant  $\implies$  Observers in different frames disagree about collapse
- · So what?
- · This is not a contradiction

- $\cdot$  Collapse is not Lorentz invariant  $\implies$  Observers in different frames disagree about collapse
- · So what?
- · This is not a contradiction
- We have a contradiction only if observers in different frames verify incompatible account of collapse

- $\cdot$  Collapse is not Lorentz invariant  $\implies$  Observers in different frames disagree about collapse
- · So what?
- · This is not a contradiction
- We have a contradiction only if observers in different frames verify incompatible account of collapse
- · Which they can't!

 State specific verifications disturb one another and so verifications are incompatible.

- State specific verifications disturb one another and so verifications are incompatible.
- · No contradiction over collapse!

- State specific verifications disturb one another and so verifications are incompatible.
- · No contradiction over collapse!
- · In fact, disagreement is fine!

- State specific verifications disturb one another and so verifications are incompatible.
- · No contradiction over collapse!
- · In fact, disagreement is fine!
- Observers in different frames disagree over length, temporal order etc.

· So what?

- · So what?
- Is everything okay?

- · So what?
- · Is everything okay?
- · What about measurements?

- · So what?
- Is everything okay?
- · What about measurements?
- Relativistic causality forbids an instantaneous collapse of almost all non-local operators!

Then what happens in our experiments?

