

UNIVERSITY OF ABERDEEN

SESSION 2019–20

EX3030

Degree Examination in EX3030 Heat, Mass, & Momentum Transfer

5th December 2019

9 am – 12 pm

PLEASE NOTE THE FOLLOWING

- (i) You **must not** have in your possession any material other than that expressly permitted in the rules appropriate to this examination. Where this is permitted, such material **must not** be amended, annotated or modified in any way.
- (ii) You **must not** have in your possession any material that could be determined as giving you an advantage in the examination.
- (iii) You **must not** attempt to communicate with any candidate during the exam, either orally or by passing written material, or by showing material to another candidate, nor must you attempt to view another candidate's work.
- (iv) You **must not** take to your examination desk any electronic devices such as mobile phones or other "smart" devices. The only exception to this rule is an approved calculator.

Failure to comply with the above will be regarded as cheating and may lead to disciplinary action as indicated in the Academic Quality Handbook.

Notes:

- (i) Candidates ARE ONLY permitted to use APPROVED calculators.*
- (ii) Candidates ARE NOT permitted to use the Engineering Mathematics Handbook.*
- (iii) Candidates ARE NOT permitted to use GREEN or RED pen in their exam booklet.*
- (iv) Data sheets are attached to the paper.*

Candidates must attempt ALL questions in BOTH parts. All questions in part A are worth 20 marks and should be answered in the standard examination booklet. Part B is multiple choice, your answers must be submitted on the multiple choice paper.

PART A: Answer using exam booklets**Question 1**

- a) A copper sphere 12.7 mm in diameter is initially at 66°C and is placed in a airstream at 27°C. In order to determine the heat transfer coefficient of air flowing over the sphere, a time series of the surface temperature is used. A sensor placed at the outer surface of the sphere indicates a temperature of 55°C at 69 seconds after the sphere was introduced into the airstream. Calculate the heat transfer coefficient. Describe the assumptions made to solve this problem. Given $\rho = 8933 \text{ kg m}^{-3}$, $C_p = 389 \text{ J kg}^{-1} \text{ K}^{-1}$, and $\kappa = 398 \text{ W m}^{-1} \text{ K}^{-1}$. **[9 marks]**
- b) A long cylindrical wood log ($\kappa = 0.17 \text{ W m}^{-1} \text{ K}^{-1}$ and $\alpha = 1.28 \times 10^{-7} \text{ m}^2 \text{ s}^{-1}$) is 10 cm in diameter and is initially at a uniform temperature of 10°C. It is exposed to hot gases at 500°C in a fireplace with a heat transfer coefficient of $13.6 \text{ W m}^{-2} \text{ K}^{-1}$ on the surface. If the ignition temperature of the wood is 420°C, determine how long it will be before the log ignites. **[11 marks]**

[Question total: 20 marks]

Question 2

Water is overflowing a dam and down an inclined slope (see Fig. 1). The surface of the dam can be idealised as a rectangular plane which is symmetric in the z -direction, and (for now) only laminar flow is being considered.

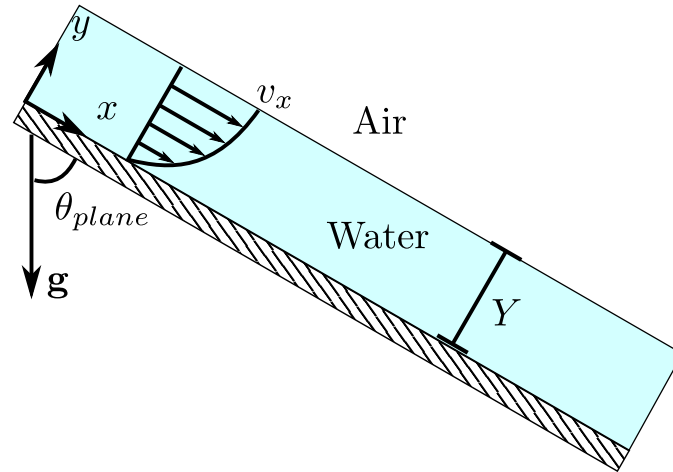


Figure 1: Water flowing down an inclined plane.

- a) Simplify the continuity equation for this system and state any assumptions you make. **[4 marks]**
- b) Derive the following results from the Cauchy momentum equation and the general form of Newton's law of viscosity: **[6 marks]**

$$\frac{\partial \tau_{yx}}{\partial y} = \rho g_x \qquad \tau_{yx} = -\mu \frac{\partial v_x}{\partial y}.$$

- c) Define your boundary conditions and derive the following expression for the velocity profile, **[5 marks]**

$$v_x = \frac{\rho g_x}{\mu} \left(Y y - \frac{y^2}{2} \right)$$

- d) Use an integration of the velocity over the flow area to determine the following expression for the volumetric flow rate, **[3 marks]**

$$\dot{V}_x = \frac{\rho g_x Y^3 Z}{3 \mu},$$

where Z is the width of the plate in the z -direction.

- e) Provide an expression for the maximum flow velocity. **[2 marks]**

[Question total: 20 marks]

Question 3

The following integrated expressions for heat transfer in a plate and a pipe are available:

$$Q_x = \frac{k}{X} A (T_{inner} - T_{outer}) \quad Q_r = \frac{2 \pi L k}{\ln(R_{outer}/R_{inner})} (T_{in} - T_{out})$$

An equivalent equation is required for spherical geometries.

- a) What single assumption is made in the derivation energy balance equation below? **[2 marks]**

$$\rho C_p \frac{\partial T}{\partial t} = -\rho C_p \mathbf{v} \cdot \nabla T - \nabla \cdot \mathbf{q} - \boldsymbol{\tau} : \nabla \mathbf{v} - p \nabla \cdot \mathbf{v} + \sigma_{energy}$$

- b) Simplify the energy balance equation to the following expression,

$$\frac{\partial}{\partial r} r^2 q_r = 0.$$

Clearly state any assumptions you make along the way. **[7 marks]**

- c) Derive the following equation for the heat flux in spherical shells. **[6 marks]**

$$q_r = \frac{k}{r^2 (R_{inner}^{-1} - R_{outer}^{-1})} (T_{inner} - T_{outer})$$

- d) Demonstrate that the resistance to heat transfer, for a spherical shell is given by the following expression:

$$R = \frac{1}{UA} = \frac{R_{inner}^{-1} - R_{outer}^{-1}}{4 \pi k}$$

Note: You will need to derive the expression for the overall heat transfer rate, Q_r , and then isolate the $R = (UA)^{-1}$ term. **[5 marks]**

[Question total: 20 marks]

Question 4

Helium gas ($\approx 4 \text{ g mol}^{-1}$) at 100 bar and 20°C is stored in a 140 mm outer-diameter vessel with a pyrex wall 4 mm thick and a height of 850 mm. The molar concentration of helium in the pyrex is 35 mol m^{-3} at the inner surface and negligible at the outer surface, while the diffusion coefficient of helium in pyrex is approximately $0.2 \times 10^{-12} \text{ m}^2 \text{ s}^{-1}$.

- a) Assuming the curvature of the tank is negligible (you can use rectangular coordinates) and steady-state, one-dimensional conditions, show that the molar flux of helium is constant through the wall. **[7 marks]**

$$N_{\text{He},z} = N_{\text{He},0}$$

- b) The concentration of helium in the pyrex wall is very low $x_{\text{He}} \ll 1$, allowing the use of the simple form of Fick's law. Determine the concentration profile of helium in the wall. **[7 marks]**
- c) Calculate the total mass flow-rate of helium transported through the side walls of the vessel (consider just the cylindrical sides). **[3 marks]**
- d) Fick's law is often modified to the following form:

$$N_{A,x} = - (D_{AB} + E_D) \frac{\partial C_A}{\partial x}$$

What is the parameter E_D and what does it represent? **[3 marks]**

[Question total: 20 marks]

PART B: Answer using the multiple choice sheet

- 1) Thick solid suspensions, like toothpaste or mud, are best modelled at low shear stress using which fluid type? Note that this is at very low shear stress and time-dependent effects should be ignored. **[2 marks]**

- A) Shear-thinning.
- B) Shear-thickening.
- C) Viscoplastic.
- D) Newtonian.

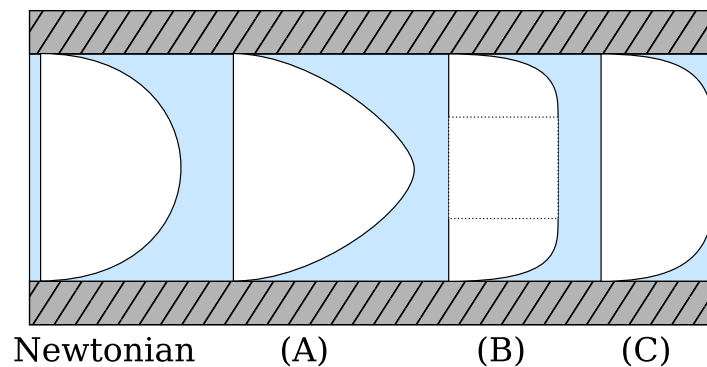


Figure 2: *Non-newtonian flow profiles compared against the newtonian flow profile.*

- 2) Which flow profile in Fig. 2 best describes toothpaste flowing through a tube? **[2 marks]**
- 3) What are the value(s) of the flow index n in the Power-law model for a shear-thinning fluid? **[2 marks]**
- A) $n < 0$
 - B) $n < 1$
 - C) $n = 1$
 - D) $n > 1$

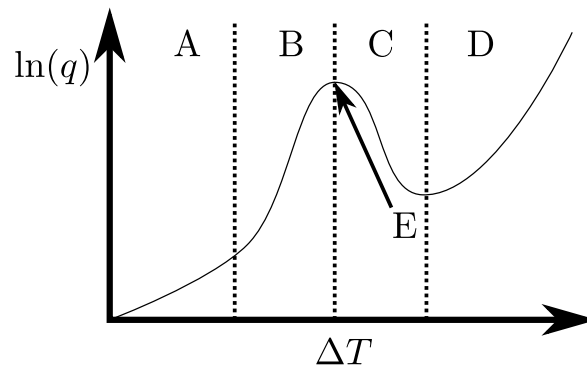


Figure 3: A typical boiling heat flux versus driving temperature difference curve.

- 4) Which region of the boiling curve in Fig. 3 is the nucleate boiling regime? **[2 marks]**
- 5) Which region of the boiling curve in Fig. 3 is the operation of the boiler unstable? **[2 marks]**
- 6) At what temperature should the properties used in the Prandtl number be evaluated for isothermal vertical walls surrounded by a fluid which has another different temperature far from the wall? **[2 marks]**
 - A) Inlet temperature.
 - B) Average of the wall and far fluid temperature.
 - C) Wall temperature.
 - D) Centerline temperature.
 - E) Average of inlet and outlet temperature.
- 7) The Schmidt number is a ratio of which two properties? **[2 marks]**
 - A) Viscosity and density.
 - B) Momentum diffusivity and molar diffusivity.
 - C) Thermal diffusivity and mass diffusivity.
 - D) Viscosity and dimensional density.
- 8) Two liquids are in contact with a species diffusing between them. Which boundary condition is inappropriate? **[2 marks]**
 - A) Temperature in each phase is equal at the interface.
 - B) No-slip between the two phases at the interface.
 - C) The chemical potential of the diffusing species in each phase is equal at the interface.
 - D) There is no mass flux at the interface.

- 9) How many shield layers are required to lower incoming thermal radiation by a factor of $1/4$? **[2 marks]**
- A) 1
 - B) 2
 - C) 3
 - D) 4
- 10) Consider the inside of an annulus (the zone between two concentric pipes) where the inner radius is 10% of the outer radius. What fraction of radiation emitted from the outer surface falls on the outer surface? **[2 marks]**
- A) 0.8
 - B) 0.2
 - C) 1.0
 - D) 0.9

END OF PAPER

DATASHEET

General balance equations:

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \rho \mathbf{v} \quad (\text{Mass/Continuity}) \quad (1)$$

$$\frac{\partial C_A}{\partial t} = -\nabla \cdot \mathbf{N}_A + \sigma_A \quad (\text{Species}) \quad (2)$$

$$\rho \frac{\partial \mathbf{v}}{\partial t} = -\rho \mathbf{v} \cdot \nabla \mathbf{v} - \nabla \cdot \boldsymbol{\tau} - \nabla p + \rho \mathbf{g} \quad (\text{Momentum}) \quad (3)$$

$$\rho C_p \frac{\partial T}{\partial t} = -\rho C_p \mathbf{v} \cdot \nabla T - \nabla \cdot \mathbf{q} - \boldsymbol{\tau} : \nabla \mathbf{v} - p \nabla \cdot \mathbf{v} + \sigma_{\text{energy}} \quad (\text{Heat/Energy}) \quad (4)$$

In Cartesian coordinate systems, ∇ can be treated as a vector of derivatives. In curve-linear coordinate systems, the directions $\hat{\mathbf{r}}$, $\hat{\boldsymbol{\theta}}$, and $\hat{\boldsymbol{\phi}}$ depend on the position. For convenience in these systems, look-up tables are provided for common terms involving ∇ .

Cartesian coordinates (with index notation examples)

where s is a scalar, \mathbf{v} is a vector, and $\boldsymbol{\tau}$ is a tensor.

$$\begin{aligned} \nabla s &= \nabla_i s = \left[\frac{\partial s}{\partial x}, \frac{\partial s}{\partial y}, \frac{\partial s}{\partial z} \right] \\ \nabla^2 s &= \nabla_i \nabla_i s = \frac{\partial^2 s}{\partial x^2} + \frac{\partial^2 s}{\partial y^2} + \frac{\partial^2 s}{\partial z^2} \\ \nabla \cdot \mathbf{v} &= \nabla_i v_i = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \\ \nabla \cdot \boldsymbol{\tau} &= \nabla_i \tau_{ij} \\ [\nabla \cdot \boldsymbol{\tau}]_x &= \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} \\ [\nabla \cdot \boldsymbol{\tau}]_y &= \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} \\ [\nabla \cdot \boldsymbol{\tau}]_z &= \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} \\ \mathbf{v} \cdot \nabla \mathbf{v} &= v_i \nabla_i v_j \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_x &= v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} + v_z \frac{\partial v_x}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_y &= v_x \frac{\partial v_y}{\partial x} + v_y \frac{\partial v_y}{\partial y} + v_z \frac{\partial v_y}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_z &= v_x \frac{\partial v_z}{\partial x} + v_y \frac{\partial v_z}{\partial y} + v_z \frac{\partial v_z}{\partial z} \end{aligned}$$

Cylindrical coordinates

where s is a scalar, \mathbf{v} is a vector, and τ is a tensor. All expressions involving τ are for symmetrical τ only.

$$\begin{aligned}\nabla s &= \left[\frac{\partial s}{\partial r}, \frac{1}{r} \frac{\partial s}{\partial \theta}, \frac{\partial s}{\partial z} \right] \\ \nabla^2 s &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial s}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 s}{\partial \theta^2} + \frac{\partial^2 s}{\partial z^2} \\ \nabla \cdot \mathbf{v} &= \frac{1}{r} \frac{\partial}{\partial r} (r v_r) + \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{\partial v_z}{\partial z} \\ [\nabla \cdot \tau]_r &= \frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rr}) + \frac{1}{r} \frac{\partial \tau_{r\theta}}{\partial \theta} - \frac{1}{r} \tau_{\theta\theta} + \frac{\partial \tau_{rz}}{\partial z} \\ [\nabla \cdot \tau]_\theta &= \frac{1}{r} \frac{\partial \tau_{\theta\theta}}{\partial \theta} + \frac{\partial \tau_{r\theta}}{\partial r} + \frac{2}{r} \tau_{r\theta} + \frac{\partial \tau_{\theta z}}{\partial z} \\ [\nabla \cdot \tau]_z &= \frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rz}) + \frac{1}{r} \frac{\partial \tau_{\theta z}}{\partial \theta} + \frac{\partial \tau_{zz}}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_r &= v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta^2}{r} + v_z \frac{\partial v_r}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_\theta &= v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r v_\theta}{r} + v_z \frac{\partial v_\theta}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_z &= v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z}\end{aligned}$$

Spherical coordinates

where s is a scalar, \mathbf{v} is a vector, and τ is a tensor. All expressions involving τ are for symmetrical τ only.

$$\begin{aligned}\nabla s &= \left[\frac{\partial s}{\partial r}, \frac{1}{r} \frac{\partial s}{\partial \theta}, \frac{1}{r \sin \theta} \frac{\partial s}{\partial \phi} \right] \\ \nabla^2 s &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial s}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial s}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 s}{\partial \phi^2} \\ \nabla \cdot \mathbf{v} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (v_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} \\ [\nabla \cdot \tau]_r &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{rr}) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\tau_{r\theta} \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial \tau_{r\phi}}{\partial \phi} - \frac{\tau_{\theta\theta} + \tau_{\phi\phi}}{r} \\ [\nabla \cdot \tau]_\theta &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{r\theta}) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\tau_{\theta\theta} \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial \tau_{\theta\phi}}{\partial \phi} + \frac{\tau_{r\theta}}{r} - \frac{\cot \theta}{r} \tau_{\phi\phi} \\ [\nabla \cdot \tau]_\phi &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{r\phi}) + \frac{1}{r} \frac{\partial \tau_{\theta\phi}}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial \tau_{\phi\phi}}{\partial \phi} + \frac{\tau_{r\theta}}{r} + \frac{2 \cot \theta}{r} \tau_{\theta\phi} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_r &= v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{v_\theta^2 + v_\phi^2}{r} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_\theta &= v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial v_\theta}{\partial \phi} + \frac{v_r v_\theta - v_\phi^2 \cot \theta}{r} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_\phi &= v_r \frac{\partial v_\phi}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\phi}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} + \frac{v_r v_\phi + v_\theta v_\phi \cot \theta}{r}\end{aligned}$$

Rectangular		Cylindrical		Spherical	
q_x	$-k \frac{\partial T}{\partial x}$	q_r	$-k \frac{\partial T}{\partial r}$	q_r	$-k \frac{\partial T}{\partial r}$
q_y	$-k \frac{\partial T}{\partial y}$	q_θ	$-k \frac{1}{r} \frac{\partial T}{\partial \theta}$	q_θ	$-k \frac{1}{r} \frac{\partial T}{\partial \theta}$
q_z	$-k \frac{\partial T}{\partial z}$	q_z	$-k \frac{\partial T}{\partial z}$	q_ϕ	$-k \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi}$
τ_{xx}	$-2\mu \frac{\partial v_x}{\partial x} + \mu^B \nabla \cdot \mathbf{v}$	τ_{rr}	$-2\mu \frac{\partial v_r}{\partial r} + \mu^B \nabla \cdot \mathbf{v}$	τ_{rr}	$-2\mu \frac{\partial v_r}{\partial r} + \mu^B \nabla \cdot \mathbf{v}$
τ_{yy}	$-2\mu \frac{\partial v_y}{\partial y} + \mu^B \nabla \cdot \mathbf{v}$	$\tau_{\theta\theta}$	$-2\mu \left(\frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r}{r} \right) + \mu^B \nabla \cdot \mathbf{v}$	$\tau_{\theta\theta}$	$-2\mu \left(\frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r}{r} \right) + \mu^B \nabla \cdot \mathbf{v}$
τ_{zz}	$-2\mu \frac{\partial v_z}{\partial z} + \mu^B \nabla \cdot \mathbf{v}$	τ_{zz}	$-2\mu \frac{\partial v_z}{\partial z} + \mu^B \nabla \cdot \mathbf{v}$	$\tau_{\phi\phi}$	$-2\mu \left(\frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} + \frac{v_r + v_\theta \cot \theta}{r} \right) + \mu^B \nabla \cdot \mathbf{v}$
τ_{xy}	$-\mu \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right)$	$\tau_{r\theta}$	$-\mu \left(r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_r}{\partial \theta} \right)$	$\tau_{r\theta}$	$-\mu \left(r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_r}{\partial \theta} \right)$
τ_{yz}	$-\mu \left(\frac{\partial v_y}{\partial z} + \frac{\partial v_z}{\partial y} \right)$	$\tau_{\theta z}$	$-\mu \left(\frac{1}{r} \frac{\partial v_z}{\partial \theta} + \frac{\partial v_\theta}{\partial z} \right)$	$\tau_{\theta\phi}$	$-\mu \left(\frac{\sin \theta}{r} \frac{\partial}{\partial \theta} \left(\frac{v_\phi}{\sin \theta} \right) + \frac{1}{r \sin \theta} \frac{\partial v_\theta}{\partial \phi} \right)$
τ_{xz}	$-\mu \left(\frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} \right)$	τ_{zr}	$-\mu \left(\frac{\partial v_r}{\partial z} + \frac{\partial v_z}{\partial r} \right)$	$\tau_{\phi r}$	$-\mu \left(\frac{1}{r \sin \theta} \frac{\partial v_r}{\partial \phi} + r \frac{\partial}{\partial r} \left(\frac{v_\phi}{r} \right) \right)$

Table 1: Fourier's law for the heat flux and Newton's law for the stress in several coordinate systems. Please remember that the stress is symmetric, so $\tau_{ij} = \tau_{ji}$.

Viscous models:

Power-Law Fluid:

$$|\tau_{xy}| = k \left| \frac{\partial v_x}{\partial y} \right|^n \quad (5)$$

Bingham-Plastic Fluid:

$$\frac{\partial v_x}{\partial y} = \begin{cases} -\mu^{-1} (\tau_{xy} - \tau_0) & \text{if } \tau_{xy} > \tau_0 \\ 0 & \text{if } \tau_{xy} \leq \tau_0 \end{cases}$$

Dimensionless Numbers

$$\text{Re} = \frac{\rho \langle v \rangle D}{\mu} \quad \text{Re}_H = \frac{\rho \langle v \rangle D_H}{\mu} \quad \text{Re}_{MR} = -\frac{16 L \rho \langle v \rangle^2}{R \Delta p} \quad (6)$$

The hydraulic diameter is defined as $D_H = 4 A / P_w$.

Single phase pressure drop calculations in pipes:

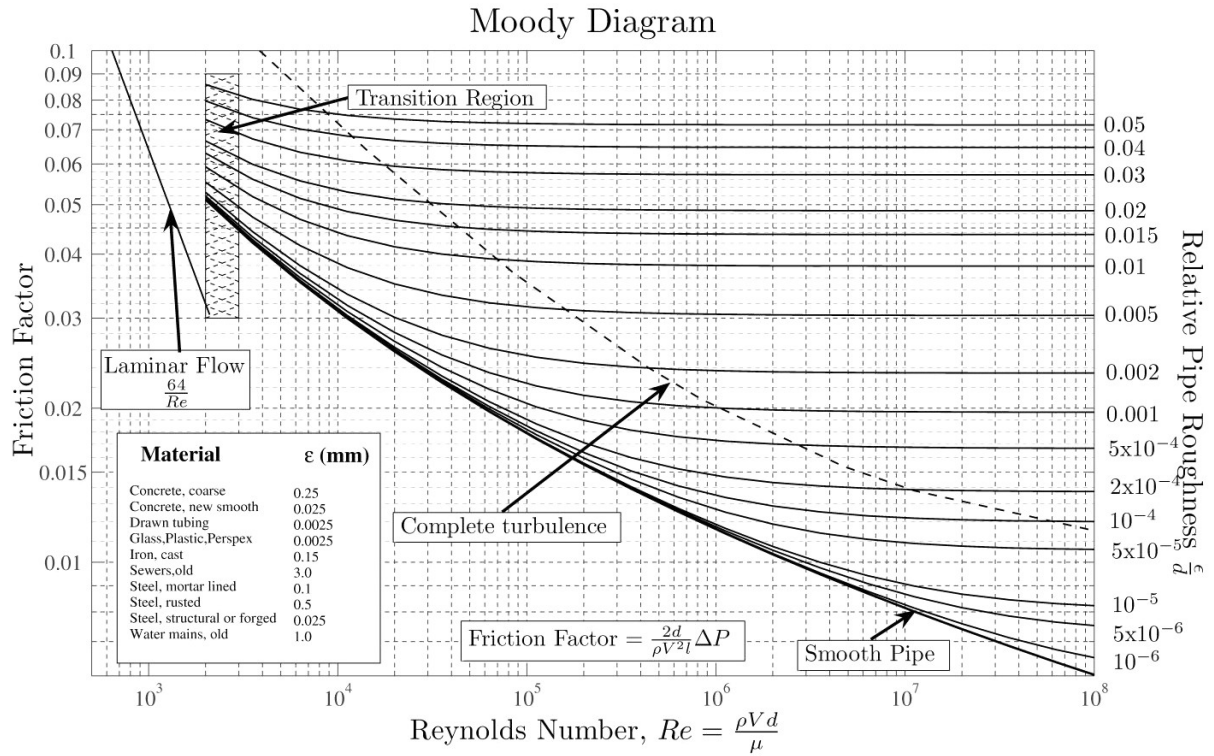
Darcy-Weisbach equation:

$$\frac{\Delta p}{L} = -\frac{C_f \rho \langle v \rangle^2}{R} \quad (7)$$

where $C_f = 16/Re$ for laminar Newtonian flow. For turbulent flow of Newtonian fluids in smooth pipes, we have the Blasius correlation:

$$C_f = 0.079 Re^{-1/4} \quad \text{for } 2.5 \times 10^3 < Re < 10^5 \text{ and smooth pipes.}$$

Otherwise, you may refer to the Moody diagram.



Laminar Power-Law fluid:

$$\dot{V} = \frac{n \pi R^3}{3n+1} \left(\frac{R}{2k} \right)^{\frac{1}{n}} \left(-\frac{\Delta p}{L} \right)^{\frac{1}{n}}$$

Two-Phase Flow:

Lockhart-Martinelli parameter:

$$X^2 = \frac{\Delta p_{liq.-only}}{\Delta p_{gas-only}}$$

Pressure drop calculation:

$$\Delta p_{two-phase} = \Phi_{liq.}^2 \Delta p_{liq.-only} = \Phi_{gas}^2 \Delta p_{gas-only}$$

Chisholm's relation:

$$\Phi_{gas}^2 = 1 + cX + X^2$$

$$\Phi_{liq.}^2 = 1 + \frac{c}{X} + \frac{1}{X^2}$$

$$c = \begin{cases} 20 & \text{turbulent liquid \& turbulent gas} \\ 12 & \text{laminar liquid \& turbulent gas} \\ 10 & \text{turbulent liquid \& laminar gas} \\ 5 & \text{laminar liquid \& laminar gas} \end{cases}$$

Farooqi and Richardson expression for liquid hold-up in co-current flows of Newtonian fluids and air in horizontal pipes:

$$h = \begin{cases} 0.186 + 0.0191 X & 1 < X < 5 \\ 0.143 X^{0.42} & 5 < X < 50 \\ 1 / (0.97 + 19/X) & 50 < X < 500 \end{cases}$$

Heat Transfer Dimensionless numbers:

$$\text{Nu} = \frac{hL}{k} \quad \text{Pr} = \frac{\mu C_p}{k} \quad \text{Gr} = \frac{g \beta (T_w - T_\infty) L^3}{\nu^2}$$

where $\beta = V^{-1}(\partial V / \partial T)$.

Heat transfer: Resistances

$$Q = U_T A_T \Delta T = R_T^{-1} \Delta T$$

	Conduction Shell Resistances			Radiation
	Rect.	Cyl.	Sph.	
R	$\frac{X}{kA}$	$\frac{\ln(R_{outer}/R_{inner})}{2\pi Lk}$	$\frac{R_{inner}^{-1} - R_{outer}^{-1}}{4\pi k}$	$[A \varepsilon \sigma (T_j^2 + T_i^2) (T_j + T_i)]^{-1}$

Radiation Heat Transfer:

Stefan-Boltzmann constant $\sigma = 5.6703 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$.

Summation relationship, $\sum_j F_{i \rightarrow j} = 1$, and reciprocity relationship, $F_{i \rightarrow j} A_i = F_{j \rightarrow i} A_j$.

Radiation shielding factor $1/(N+1)$.

$$Q_{rad., i \rightarrow j} = \sigma \varepsilon F_{i \rightarrow j} A_i (T_j^4 - T_i^4) = h_{rad.} A (T_\infty - T_w)$$

Natural Convection

$\text{Ra} = \text{Gr Pr}$	C	m
$< 10^4$	1.36	1/5
$10^4 - 10^9$	0.59	1/4
$> 10^9$	0.13	1/3

Table 2: Natural convection coefficients for isothermal vertical plates in the empirical relation $\text{Nu} \approx C (\text{Gr Pr})^m$.

For isothermal vertical cylinders, the above expressions for isothermal vertical plates may be used but must be scaled by a factor, F (i.e., $Nu_{v.cyl.} = F Nu_{v.plate}$):

$$F = \begin{cases} 1 & \text{for } (D/H) \geq 35 Gr_H^{-1/4} \\ 1.3 [H D^{-1} Gr_D^{-1}]^{1/4} + 1 & \text{for } (D/H) < 35 Gr_H^{-1/4} \end{cases}$$

where D is the diameter and H is the height of the cylinder. The subscript on Gr indicates which length is to be used as the critical length to calculate the Grashof number.

Churchill and Chu expression for natural convection from a horizontal pipe:

$$Nu^{1/2} = 0.6 + 0.387 \left\{ \frac{Gr Pr}{[1 + (0.559/Pr)^{9/16}]^{16/9}} \right\}^{1/6} \quad \text{for } 10^{-5} < Gr Pr < 10^{12}$$

Forced Convection:

Laminar flows:

$$Nu \approx 0.332 Re^{1/2} Pr^{1/3}$$

Well-Developed turbulent flows in smooth pipes:

$$Nu \approx \frac{(C_f/2) Re Pr}{1.07 + 12.7(C_f/2)^{1/2} (Pr^{2/3} - 1)} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Boiling:

Forster-Zuber pool-boiling coefficient:

$$h_{nb} = 0.00122 \frac{k_L^{0.79} C_{p,L}^{0.45} \rho_L^{0.49}}{\gamma^{0.5} \mu_L^{0.29} h_{fg}^{0.24} \rho_G^{0.24}} (T_w - T_{sat})^{0.24} (p_w - p_{sat})^{0.75}$$

Mostinski correlations:

$$h_{nb} = 0.104 p_c^{0.69} q^{0.7} \left[1.8 \left(\frac{p}{p_c} \right)^{0.17} + 4 \left(\frac{p}{p_c} \right)^{1.2} + 10 \left(\frac{p}{p_c} \right)^{10} \right]$$

$$q_c = 3.67 \times 10^4 p_c \left(\frac{p}{p_c} \right)^{0.35} \left[1 - \frac{p}{p_c} \right]^{0.9}$$

(Note: for the Mostinski correlations, the pressures are in units of bar)

Condensing:

Horizontal pipes

$$h = 0.72 \left(\frac{k^3 \rho^2 g_x E_{latent}}{D \mu (T_w - T_\infty)} \right)^{1/4}$$

Lumped capacitance method:

$$\text{Bi} = \frac{h L_c}{\kappa}$$

$$L_c = V/A \quad \text{for Bi} < 0.1$$

$$\frac{T(t) - T_\infty}{T_0 - T_\infty} = e^{-bt} \quad b = \frac{hA_s}{\rho V C_p}$$

1-D Transient Heat Conduction:

$$Fo = \frac{\alpha \Delta t}{(\Delta x)^2} = \tau, \quad \alpha = \kappa (\rho C_p)^{-1}$$

$$\theta_{\text{wall}} = \frac{T(x, t) - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau} \cos\left(\frac{\lambda_1 x}{L}\right), \quad \theta_{\text{cyl}} = \frac{T(r, t) - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau} \mathbf{J}_0\left(\frac{\lambda_1 r}{r_0}\right)$$

$$\theta_{\text{sph}} = \frac{T(r, t) - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau} \frac{\sin\left(\frac{\lambda_1 r}{r_0}\right)}{\frac{\lambda_1 r}{r_0}}$$

$$\theta_{0,\text{wall}} = \theta_{0,\text{cyl}} = \theta_{0,\text{sph}} = \frac{T_0 - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau}$$

$$\left(\frac{Q}{Q_{\text{max}}}\right)_{\text{wall}} = 1 - \theta_{0,\text{wall}} \frac{\sin \lambda_1}{\lambda_1}, \quad \left(\frac{Q}{Q_{\text{max}}}\right)_{\text{cyl}} = 1 - 2\theta_{0,\text{cyl}} \frac{\mathbf{J}_1(\lambda_1)}{\lambda_1}$$

$$\left(\frac{Q}{Q_{\text{max}}}\right)_{\text{sph}} = 1 - 3\theta_{0,\text{sph}} \frac{\sin \lambda_1 - \lambda_1 \cos \lambda_1}{\lambda_1^3}$$

Finite-Difference Method:

$$\frac{\partial}{\partial t}(\rho \phi) + \frac{\partial}{\partial x}(\rho \mathbf{v} \phi) = \nabla \cdot (\Gamma \nabla \phi) + \mathcal{S} \quad (1D \text{ transport equation})$$

$$\left(\frac{d\phi}{dx}\right)_i = \frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x} \quad \text{and} \quad \left(\frac{d^2\phi}{dx^2}\right)_i = \frac{\phi_{i-1} + \phi_{i+1} - 2\phi_i}{(\Delta x)^2}$$

$$T_i^{j+1} = (1 - 2\tau) T_i^j + \tau (T_{i+1}^j + T_{i-1}^j) + \frac{\tau (\Delta x)^2}{\kappa} S_i^j$$

Overall Heat Transfer Coefficient:

$$\dot{Q} = \frac{\Delta T}{\mathcal{R}} = UA\Delta T = U_i A_i \Delta T = U_o A_o \Delta T$$

$$\mathcal{R} = R_i + R_{\text{wall}} + R_o = \frac{1}{h_i A_i} + \frac{\ln D_o/D_i}{2\pi\kappa L} + \frac{1}{h_o A_o}$$

Fouling Factor:

$$\mathcal{R} = \frac{1}{h_i A_i} + \frac{R_{f,i}}{A_i} + R_{\text{wall}} + \frac{R_{f,o}}{A_o} + \frac{1}{h_o A_o}$$

LMTD Method:

$$\dot{Q} = UA_s \Delta T_{\text{lm}} \quad \text{with} \quad \Delta T_{\text{lm}} = \frac{\Delta T_2 - \Delta T_1}{\ln \frac{\Delta T_2}{\Delta T_1}} = \frac{\Delta T_1 - \Delta T_2}{\ln \frac{\Delta T_1}{\Delta T_2}}$$

$$\text{Parallel flows: } \begin{cases} \Delta T_1 = T_{\text{hot,in}} - T_{\text{cold,in}} \\ \Delta T_2 = T_{\text{hot,out}} - T_{\text{cold,out}} \end{cases}$$

$$\text{Counter flows: } \begin{cases} \Delta T_1 = T_{\text{hot,in}} - T_{\text{cold,out}} \\ \Delta T_2 = T_{\text{hot,out}} - T_{\text{cold,in}} \end{cases}$$

ϵ -NTU Method:

$$\epsilon = \frac{\dot{Q}}{\dot{Q}_{\text{max}}}, \quad \text{with } \dot{Q}_{\text{max}} = C_{\min} (T_{\text{hot,in}} - T_{\text{cold,in}}) \quad \text{and} \quad C_{\min} = \text{Min} \{ \dot{m}_{\text{hot}} C_{p,\text{hot}}, \dot{m}_{\text{cold}} C_{p,\text{cold}} \}$$

$$\text{NTU} = \frac{UA_s}{C_{\min}}$$

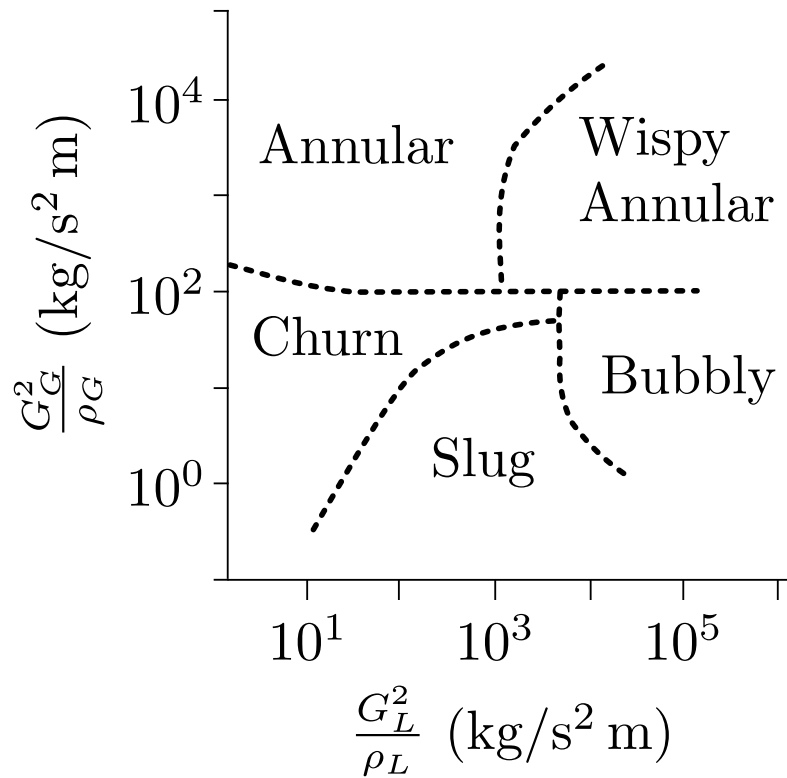


Figure 4: Hewitt-Taylor flow pattern map for multiphase flows in vertical pipes.

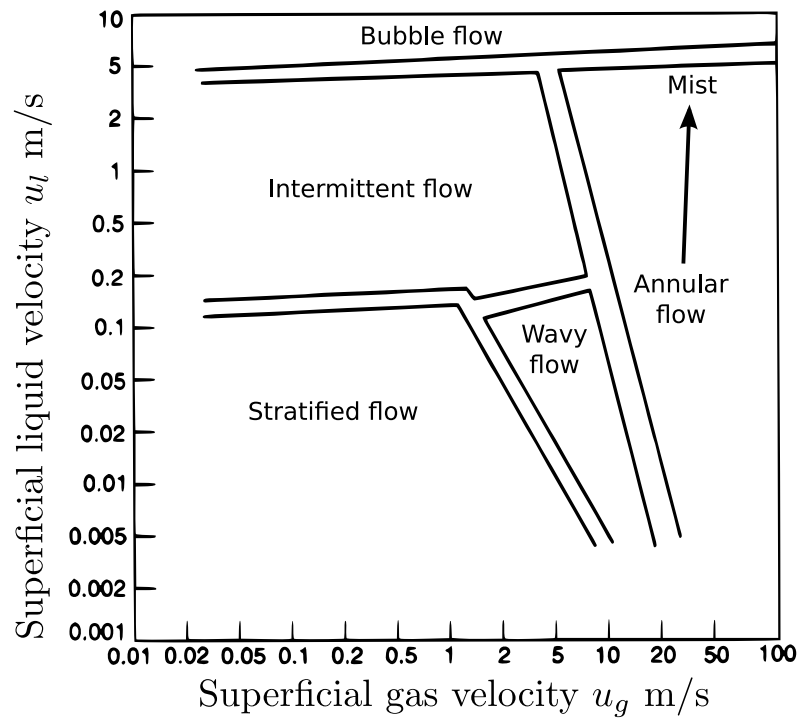


Figure 5: Chhabra and Richardson flow pattern map for horizontal pipes.

TABLE 4–2

Coefficients used in the one-term approximate solution of transient one-dimensional heat conduction in plane walls, cylinders, and spheres ($Bi = hL/k$ for a plane wall of thickness $2L$, and $Bi = hr_o/k$ for a cylinder or sphere of radius r_o)

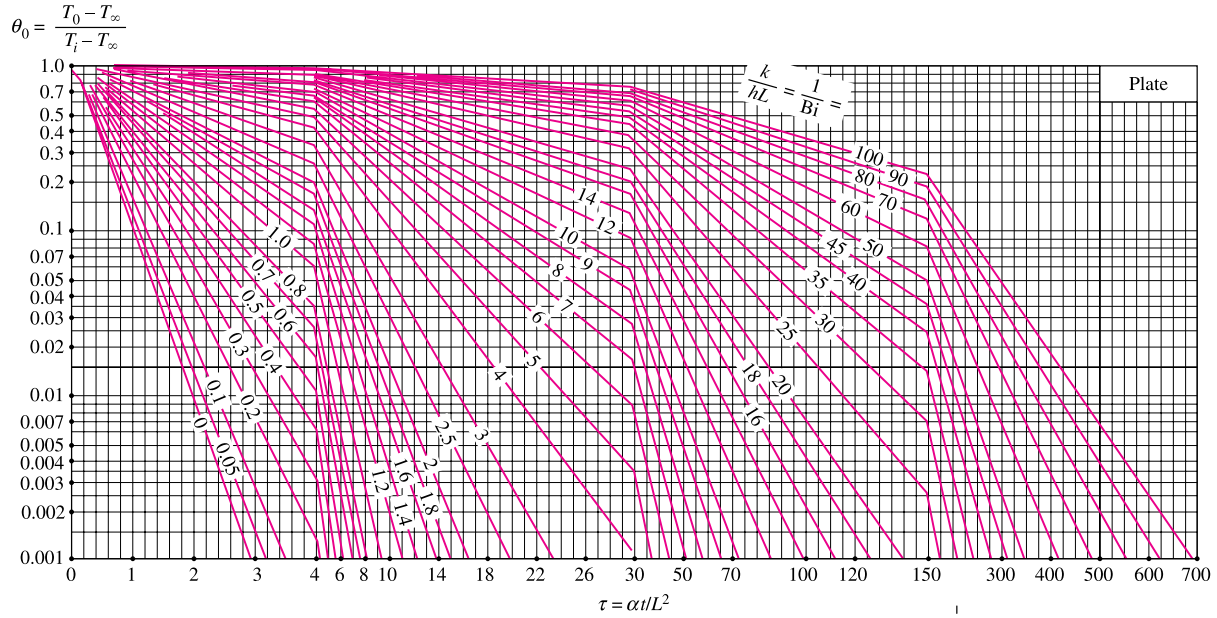
Bi	Plane Wall		Cylinder		Sphere	
	λ_1	A_1	λ_1	A_1	λ_1	A_1
0.01	0.0998	1.0017	0.1412	1.0025	0.1730	1.0030
0.02	0.1410	1.0033	0.1995	1.0050	0.2445	1.0060
0.04	0.1987	1.0066	0.2814	1.0099	0.3450	1.0120
0.06	0.2425	1.0098	0.3438	1.0148	0.4217	1.0179
0.08	0.2791	1.0130	0.3960	1.0197	0.4860	1.0239
0.1	0.3111	1.0161	0.4417	1.0246	0.5423	1.0298
0.2	0.4328	1.0311	0.6170	1.0483	0.7593	1.0592
0.3	0.5218	1.0450	0.7465	1.0712	0.9208	1.0880
0.4	0.5932	1.0580	0.8516	1.0931	1.0528	1.1164
0.5	0.6533	1.0701	0.9408	1.1143	1.1656	1.1441
0.6	0.7051	1.0814	1.0184	1.1345	1.2644	1.1713
0.7	0.7506	1.0918	1.0873	1.1539	1.3525	1.1978
0.8	0.7910	1.1016	1.1490	1.1724	1.4320	1.2236
0.9	0.8274	1.1107	1.2048	1.1902	1.5044	1.2488
1.0	0.8603	1.1191	1.2558	1.2071	1.5708	1.2732
2.0	1.0769	1.1785	1.5995	1.3384	2.0288	1.4793
3.0	1.1925	1.2102	1.7887	1.4191	2.2889	1.6227
4.0	1.2646	1.2287	1.9081	1.4698	2.4556	1.7202
5.0	1.3138	1.2403	1.9898	1.5029	2.5704	1.7870
6.0	1.3496	1.2479	2.0490	1.5253	2.6537	1.8338
7.0	1.3766	1.2532	2.0937	1.5411	2.7165	1.8673
8.0	1.3978	1.2570	2.1286	1.5526	2.7654	1.8920
9.0	1.4149	1.2598	2.1566	1.5611	2.8044	1.9106
10.0	1.4289	1.2620	2.1795	1.5677	2.8363	1.9249
20.0	1.4961	1.2699	2.2880	1.5919	2.9857	1.9781
30.0	1.5202	1.2717	2.3261	1.5973	3.0372	1.9898
40.0	1.5325	1.2723	2.3455	1.5993	3.0632	1.9942
50.0	1.5400	1.2727	2.3572	1.6002	3.0788	1.9962
100.0	1.5552	1.2731	2.3809	1.6015	3.1102	1.9990
∞	1.5708	1.2732	2.4048	1.6021	3.1416	2.0000

TABLE 4–3

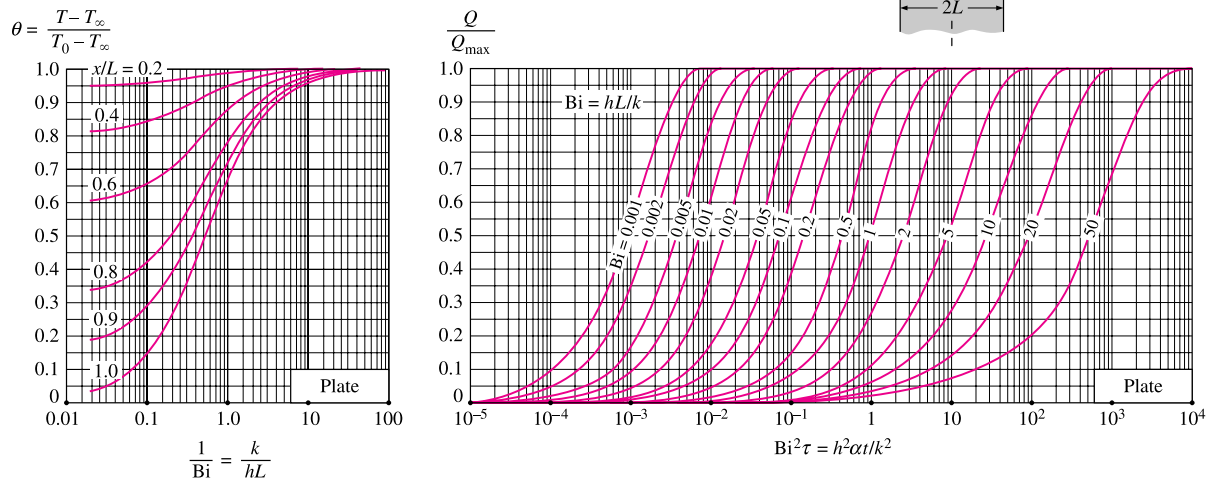
The zeroth- and first-order Bessel functions of the first kind

η	$J_0(\eta)$	$J_1(\eta)$
0.0	1.0000	0.0000
0.1	0.9975	0.0499
0.2	0.9900	0.0995
0.3	0.9776	0.1483
0.4	0.9604	0.1960
0.5	0.9385	0.2423
0.6	0.9120	0.2867
0.7	0.8812	0.3290
0.8	0.8463	0.3688
0.9	0.8075	0.4059
1.0	0.7652	0.4400
1.1	0.7196	0.4709
1.2	0.6711	0.4983
1.3	0.6201	0.5220
1.4	0.5669	0.5419
1.5	0.5118	0.5579
1.6	0.4554	0.5699
1.7	0.3980	0.5778
1.8	0.3400	0.5815
1.9	0.2818	0.5812
2.0	0.2239	0.5767
2.1	0.1666	0.5683
2.2	0.1104	0.5560
2.3	0.0555	0.5399
2.4	0.0025	0.5202
2.6	−0.0968	−0.4708
2.8	−0.1850	−0.4097
3.0	−0.2601	−0.3391
3.2	−0.3202	−0.2613

Figure 6: Coefficients for the 1D transient equations.



(a) Midplane temperature (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227-36. Reprinted by permission of ASME International.)



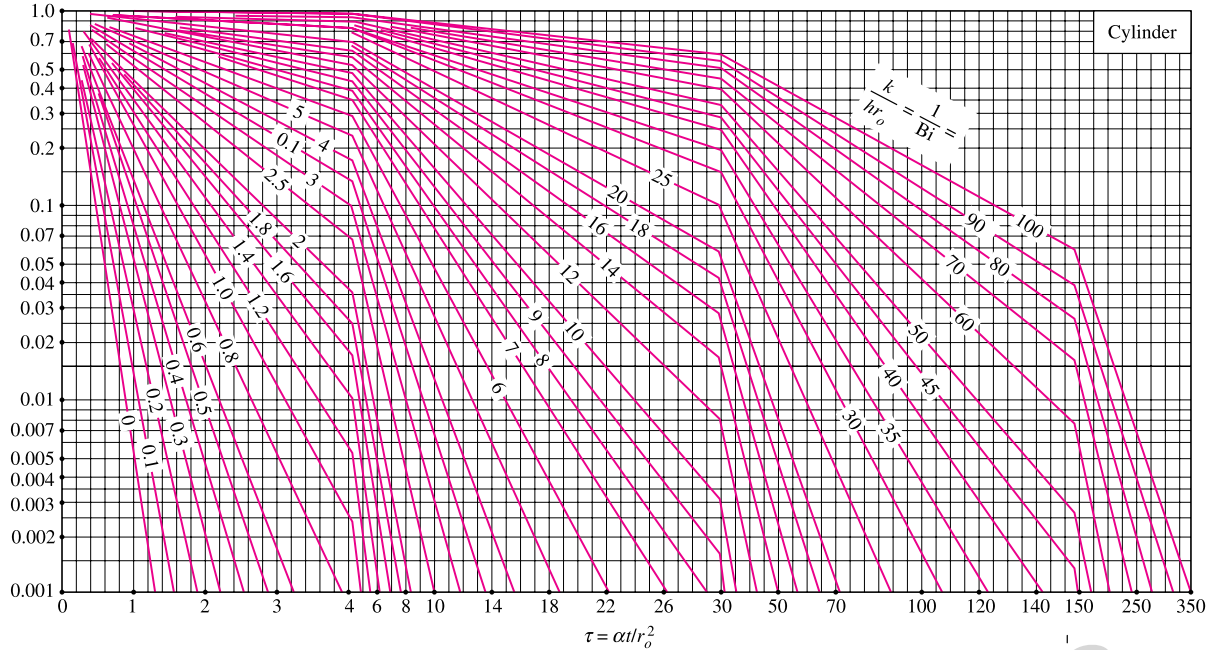
(b) Temperature distribution (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227-36. Reprinted by permission of ASME International.)

(c) Heat transfer (from H. Gröber et al.)

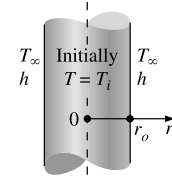
Transient temperature and heat transfer charts for a plane wall of thickness $2L$ initially at a uniform temperature T_i subjected to convection from both sides to an environment at temperature T_∞ with a convection coefficient of h .

Figure 7:

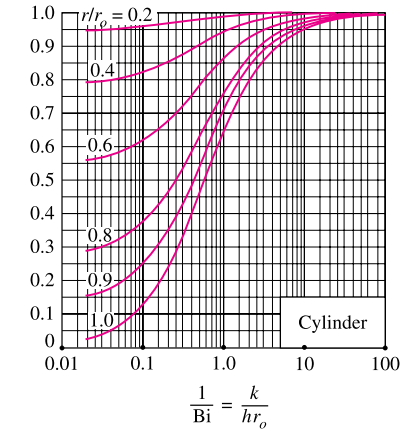
$$\theta_0 = \frac{T_0 - T_\infty}{T_i - T_\infty}$$



(a) Centerline temperature (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)

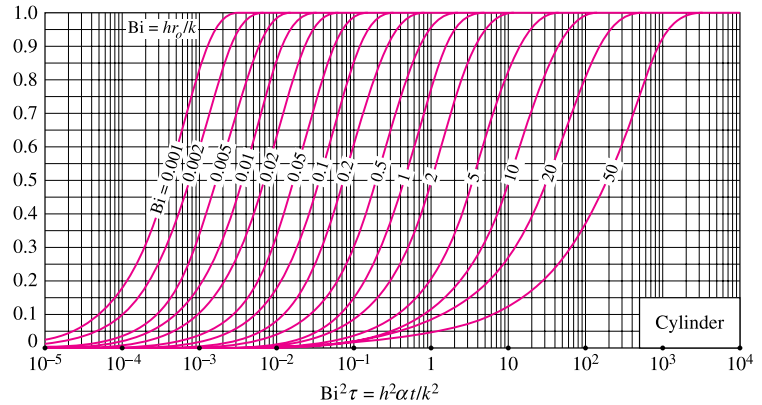


$$\theta = \frac{T - T_\infty}{T_0 - T_\infty}$$



(b) Temperature distribution (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)

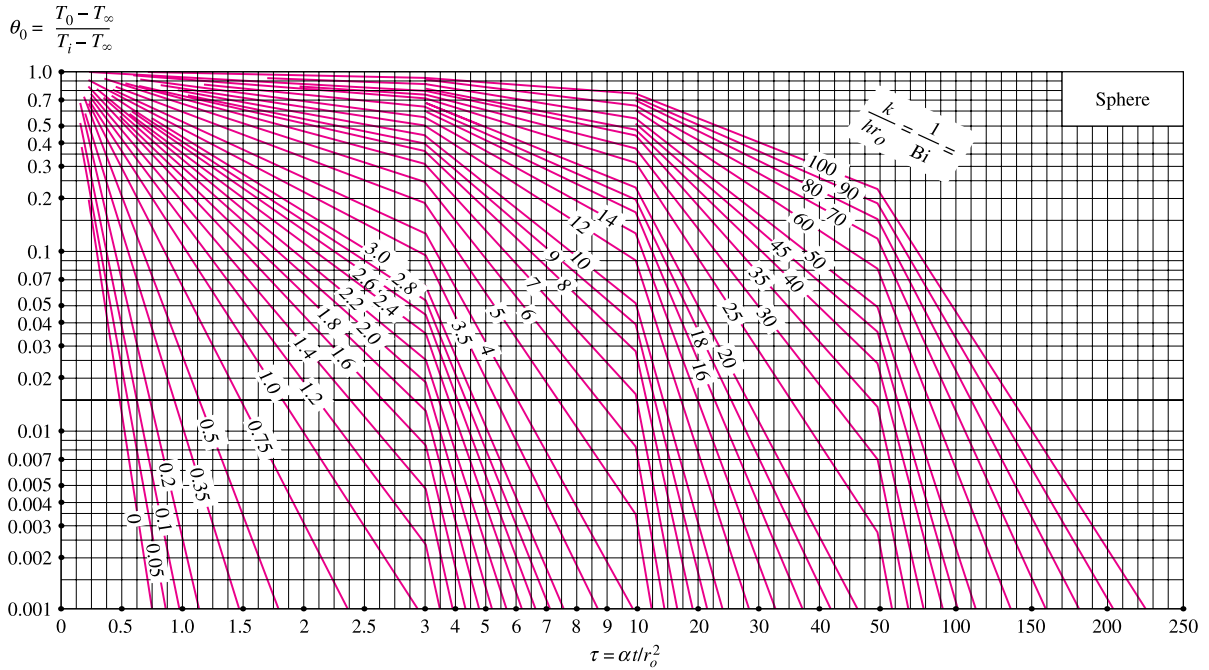
$$\frac{Q}{Q_{\max}}$$



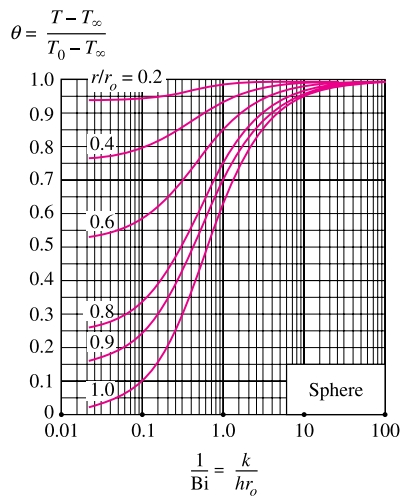
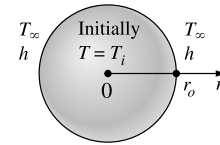
(c) Heat transfer (from H. Gröber et al.)

Transient temperature and heat transfer charts for a long cylinder of radius r_o initially at a uniform temperature T_i subjected to convection from all sides to an environment at temperature T_∞ with a convection coefficient of h .

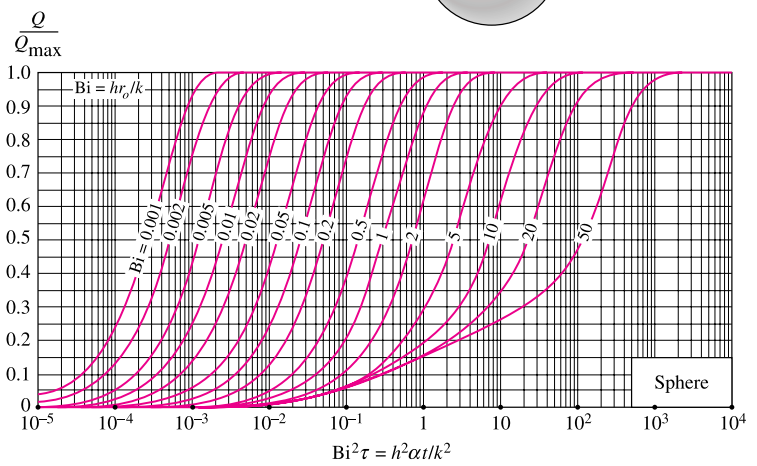
Figure 8:



(a) Midpoint temperature (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)



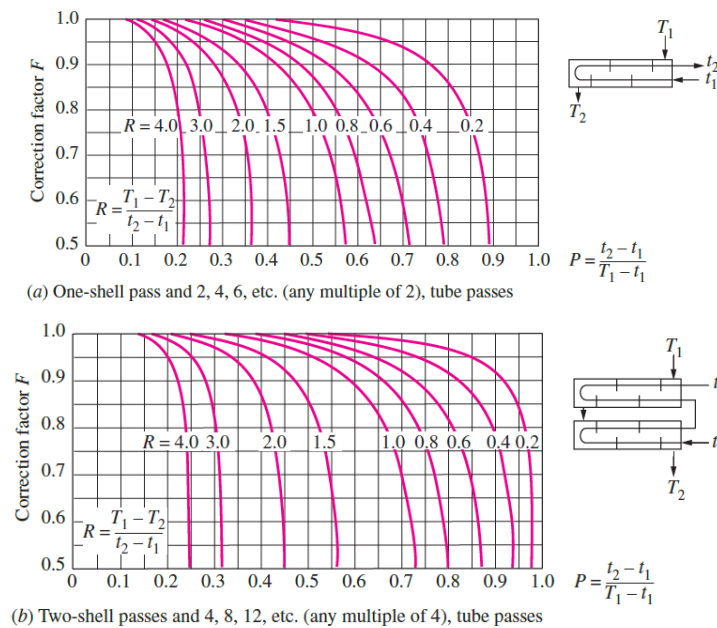
(b) Temperature distribution (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947,



(c) Heat transfer (from H. Gröber et al.)

Transient temperature and heat transfer charts for a sphere of radius r_o initially at a uniform temperature T_i subjected to convection from all sides to an environment at temperature T_∞ with a convection coefficient of h .

Figure 9:

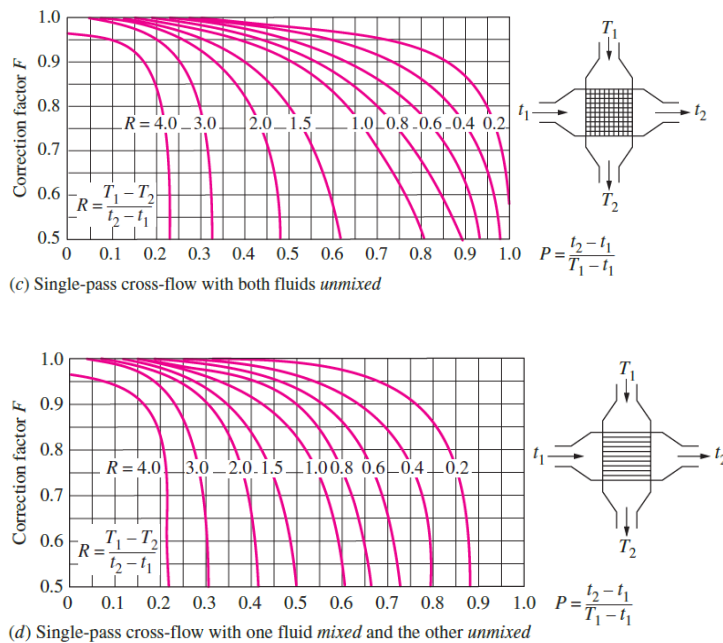


Extracted from Y.A. Cengel, "Heat Transfer: A Practical Approach", 2nd Edition.

Figure 10.8

Correction factor F charts for common shell-and-tube and cross-flow heat exchangers (from Bowman, Mueller, and Nagle, Ref. 2).

Figure 10: Correction-factors for LMTD Method, extracted from Y. A. Cengel, "Heat transfer:A practical approach", 2nd Ed.



Extracted from Y.A. Cengel, "Heat Transfer: A Practical Approach", 2nd Edition.

Figure 10.8

Correction factor F charts for common shell-and-tube and cross-flow heat exchangers (from Bowman, Mueller, and Nagle, Ref. 2).

Figure 11: Correction-factors for LMTD Method, extracted from Y. A. Cengel, "Heat transfer:A practical approach", 2nd Ed.

Effectiveness relations for heat exchangers: $NTU = UA_s/C_{\min}$ and $c = C_{\min}/C_{\max} = (\dot{m}C_p)_{\min}/(\dot{m}C_p)_{\max}$ (Kays and London, Ref. 5.)

Heat exchanger type	Effectiveness relation
1 Double-pipe: Parallel-flow	$\varepsilon = \frac{1 - \exp[-NTU(1 + c)]}{1 + c}$
Counter-flow	$\varepsilon = \frac{1 - \exp[-NTU(1 - c)]}{1 - c \exp[-NTU(1 - c)]}$
2 Shell and tube: One-shell pass 2, 4, ... tube passes	$\varepsilon = 2 \left\{ 1 + c + \sqrt{1 + c^2} \frac{1 + \exp[-NTU\sqrt{1 + c^2}]}{1 - \exp[-NTU\sqrt{1 + c^2}]} \right\}^{-1}$
3 Cross-flow (single-pass) Both fluids unmixed	$\varepsilon = 1 - \exp \left\{ \frac{NTU^{0.22}}{c} [\exp(-c NTU^{0.78}) - 1] \right\}$
C_{\max} mixed, C_{\min} unmixed	$\varepsilon = \frac{1}{c} (1 - \exp[1 - c(1 - \exp(-NTU))])$
C_{\min} mixed, C_{\max} unmixed	$\varepsilon = 1 - \exp \left\{ -\frac{1}{c} [1 - \exp(-c NTU)] \right\}$
4 All heat exchangers with $c = 0$	$\varepsilon = 1 - \exp(-NTU)$

NTU relations for heat exchangers $NTU = UA_s/C_{\min}$ and $c = C_{\min}/C_{\max} = (\dot{m}C_p)_{\min}/(\dot{m}C_p)_{\max}$ (Kays and London, Ref. 5.)

Heat exchanger type	NTU relation
1 Double-pipe: Parallel-flow	$NTU = -\frac{\ln[1 - \varepsilon(1 + c)]}{1 + c}$
Counter-flow	$NTU = \frac{1}{c - 1} \ln \left(\frac{\varepsilon - 1}{\varepsilon c - 1} \right)$
2 Shell and tube: One-shell pass 2, 4, ... tube passes	$NTU = -\frac{1}{\sqrt{1 + c^2}} \ln \left(\frac{2/\varepsilon - 1 - c - \sqrt{1 + c^2}}{2/\varepsilon - 1 - c + \sqrt{1 + c^2}} \right)$
3 Cross-flow (single-pass) C_{\max} mixed, C_{\min} unmixed	$NTU = -\ln \left[1 + \frac{\ln(1 - \varepsilon c)}{c} \right]$
C_{\min} mixed, C_{\max} unmixed	$NTU = \frac{\ln[c \ln(1 - \varepsilon) + 1]}{c}$
4 All heat exchangers with $c = 0$	$NTU = -\ln(1 - \varepsilon)$

Figure 12: NTU relations extracted from Y. A. Cengel, "Heat transfer: A practical approach", 2nd Ed.

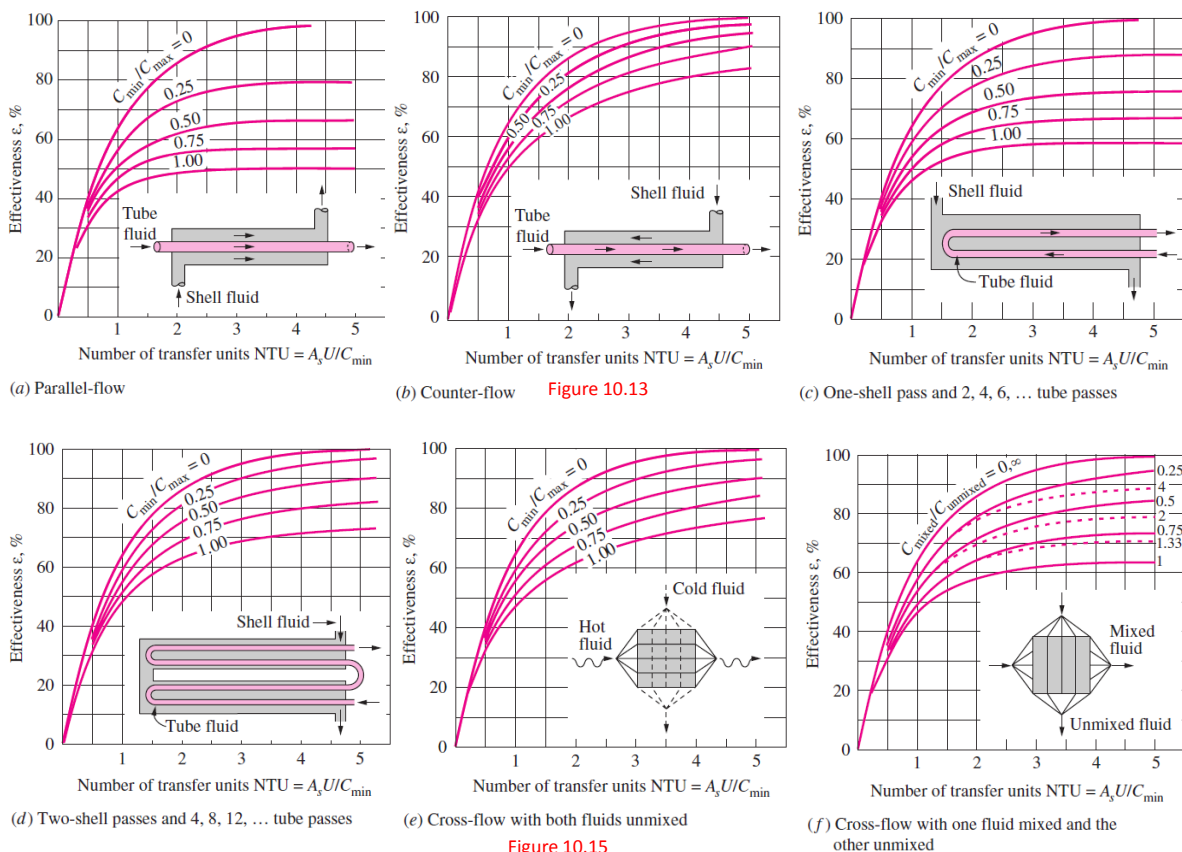


Figure 10.13

Extracted from Y.A. Cengel, "Heat Transfer: A Practical Approach", 2nd Edition.

Figure 13: NTU plots extracted from Y. A. Cengel, "Heat transfer: A practical approach", 2nd Ed.

Diffusion Dimensionless Numbers

$$Sc = \frac{\mu}{\rho D_{AB}}$$

$$Le = \frac{k}{\rho C_p D_{AB}}$$

Diffusion

General expression for the flux:

$$\mathbf{N}_A = \mathbf{J}_A + x_A \sum_B \mathbf{N}_B$$

Fick's law:

$$\mathbf{J}_A = -D_{AB} \nabla C_A$$

Stefan's law:

$$N_{s,r} = -D \frac{c}{1-x} \frac{\partial x}{\partial r}$$

Ideal Gas

$$P V = n R T$$

$$R \approx 8.314598 \text{ J K}^{-1} \text{ mol}^{-1}$$

Geometry

$$P_{\text{circle}} = 2 \pi r$$

$$A_{\text{circle}} = \pi r^2$$

$$A_{\text{sphere}} = 4 \pi r^2$$

$$V_{\text{sphere}} = \frac{4}{3} \pi r^3$$

$$A_{\text{cylinder}} = P_{\text{circle}} L$$

$$V_{\text{cylinder}} = A_{\text{circle}} L$$