

UNIVERSITY OF ABERDEEN

SESSION 2018–19

EX3030

Degree Examination in EX3030 Heat, Mass, & Momentum Transfer

14th December 2018

9 am – 12 pm

PLEASE NOTE THE FOLLOWING

- (i) You **must not** have in your possession any material other than that expressly permitted in the rules appropriate to this examination. Where this is permitted, such material **must not** be amended, annotated or modified in any way.
- (ii) You **must not** have in your possession any material that could be determined as giving you an advantage in the examination.
- (iii) You **must not** attempt to communicate with any candidate during the exam, either orally or by passing written material, or by showing material to another candidate, nor must you attempt to view another candidate's work.
- (iv) You **must not** take to your examination desk any electronic devices such as mobile phones or other "smart" devices. The only exception to this rule is an approved calculator.

Failure to comply with the above will be regarded as cheating and may lead to disciplinary action as indicated in the Academic Quality Handbook.

Notes:

- (i) Candidates ARE permitted to use an approved calculator.*
- (ii) Candidates ARE NOT permitted to use the Engineering Mathematics Handbook.*
- (iii) Candidates ARE NOT permitted to use GREEN or RED pen in their exam booklet.*
- (iv) Data sheets are attached to the paper.*

Candidates must attempt ALL questions in BOTH parts. All questions in part A are worth 20 marks and should be answered in the standard examination booklet. Part B is multiple choice, your answers must be submitted on the multiple choice paper.

PART A: Answer using exam booklets

Question 1

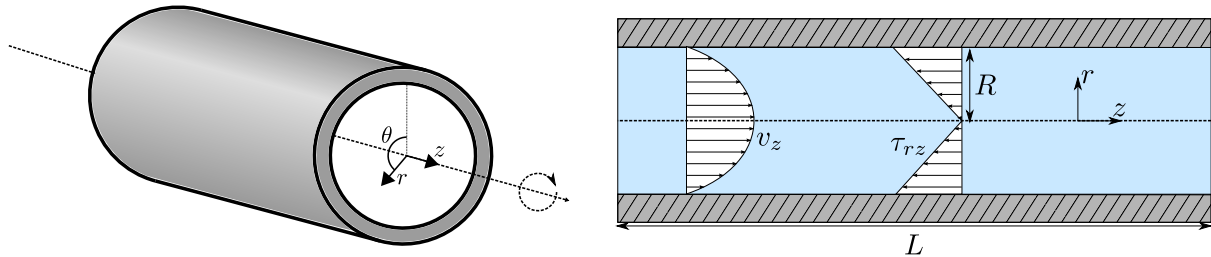


Figure 1: *An illustration of pipe flow.*

Consider pressure-driven flow along a horizontal pipe, as illustrated in Fig. 1.

- Simplify the continuity equation for this system, what does it tell you about the flow? Remember to make your assumptions and their effects clear. **[6 marks]**
- Derive the following differential equation from the Cauchy momentum equation.

$$\frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rz}) = - \frac{\partial p}{\partial z}$$

Remember to make your assumptions and their effects clear. **[7 marks]**

- Determine the following expression for the stress profile.

$$\tau_{rz} = - \frac{\Delta p}{2L} r$$

[3 marks]

- Demonstrate that the velocity profile is as given below.

$$v_z = \frac{\Delta p}{4\mu L} (r^2 - R^2)$$

[4 marks]

[Question total: 20 marks]

Question 2

A solid wire is being used to carry electrical current (see Fig. 2).

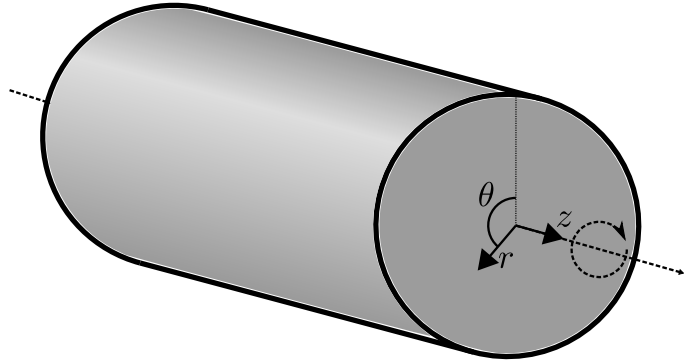


Figure 2: A representation of a solid wire (right) used as a high-power transmission line (left).

- a) You may assume that heat is generated constantly within the volume of the wire at the following rate,

$$\sigma_{\text{energy}}^{\text{current}} = \frac{I^2}{k_e}.$$

Simplify the differential energy balance equation for this system to the following form,

$$\frac{1}{r} \frac{\partial}{\partial r} (r q_r) = \frac{I^2}{k_e}$$

Ensure you clearly state any assumptions you make.

[6 marks]

- b) Derive the following expression for the heat flux within the wire,

[4 marks]

$$q_r = \frac{I^2}{2 k_e} r$$

- c) Demonstrate that the temperature profile has the following form,

[5 marks]

$$T - T_0 = \frac{I^2 R^2}{4 k_e k} \left(1 - \frac{r^2}{R^2} \right).$$

where T_0 arises from an assumption on the temperature at the surface of the wire.

- d) Discuss if the assumptions you have made are realistic.

[3 marks]

- e) How might the surface boundary condition be improved?

[2 marks]

[Question total: 20 marks]

Question 3

To maintain a pressure close to 1 atm, an industrial pipeline containing ammonia gas is vented to ambient air. Venting is achieved by tapping the pipe and inserting a 3 mm diameter tube, which extends for 20 m into the atmosphere. With the entire system operating at 25 °C and 1 bar, the ideal gas equation of state predicts a total molar concentration of 40.9 mol m^{-3} . Equimolar counter-diffusion can be assumed, and both the concentration of air in the pipeline and the concentration of ammonia in the atmosphere can be considered negligible. The diffusion coefficient of ammonia through air is approximately $2 \times 10^{-5} \text{ m}^2 \text{ s}^{-1}$.

- a) Determine the mass rate of ammonia (molar mass 17 g mol^{-1}) lost in to the atmosphere, in kg/h and the mass rate of contamination of the pipe with air (molar mass 29 g mol^{-1}) in the same units. **[12 marks]**
- b) A new high-tech membrane, which is impermeable to air, is installed at the bottom of the pipe to prevent air polluting the pipeline. The *air* within the tube is now **stationary** and the mole fraction of ammonia at the surface of the membrane is $x_A(z = 0) = 0.9$. Resolve the problem again to determine the flux of ammonia. **Note:** Stefan's law (in mole fractions for ideal gases) is given by the following

$$N_{A,z} = -D_{AB} \frac{C_T}{1 - x_A} \frac{\partial x_A}{\partial z}$$

[8 marks]**[Question total: 20 marks]**

Question 4

- a) A granite sphere of 15 cm in diameter and at uniform temperature of 120°C is suddenly placed in a controlled environment where temperature is kept at 30°C. Average convective heat transfer coefficient is $350 \text{ W m}^{-2}\text{K}^{-1}$. Calculate the temperature of the granite sphere at a radius of 4.5 cm after 21 minutes. Given properties of granite: $k = 3.2 \text{ W m}^{-1}\text{K}^{-1}$ and $\alpha = 13 \times 10^{-7} \text{ m}^2\text{s}^{-1}$. **[7 marks]**
- b) Saturated steam at 1 atm and 100°C is condensed in a shell-and-tube heat exchanger (one shell, two tube passes). Cooling water enters the tubes at 15°C with an average velocity of 3.5 m s^{-1} . The tubes are thin walled and made of copper with a diameter of $1.4 \times 10^{-2} \text{ m}$ and length of 0.5 m. The convective heat transfer coefficient for condensation on the outer surface of the tubes is $21.8 \text{ kW m}^{-2} \text{ K}^{-1}$. Determine:
- Rate of heat transfer (in W) in the heat exchanger: **[1 marks]**
 - Overall heat transfer coefficient (in $\text{kW m}^{-2} \text{ K}^{-1}$): **[3 marks]**
 - Number of tubes/pass required to condense 2.3 kg s^{-1} of steam and total surface area. If you have not solved (i) and/or (ii), you should assume $U = 7300 \text{ W m}^{-2} \text{ K}^{-1}$ and/or $\dot{Q} = 5 \text{ MW}$; **[5 marks]**
 - Outlet water temperature; **[2 marks]**
 - The maximum possible condensation rate that could be achieved with this heat exchanger using the same water flow rate and inlet temperature. **[2 marks]**

Given properties of:

- Saturated steam flow: $T_{\text{sat}} = 100^\circ\text{C}$ and $h_{fg} = 2257 \text{ kJ kg}^{-1}$;
- Cooling water: $\rho = 998 \text{ kg m}^{-3}$, $C_p = 4181 \text{ J kg}^{-1} \text{ K}^{-1}$, $\mu = 959 \times 10^{-6} \text{ N s m}^{-2}$, $k = 0.606 \text{ W m}^{-1} \text{ K}^{-1}$ and $\text{Pr} = 6.62$.

Also, cooling water convective heat transfer coefficient should be obtained from the Dittus-Boelter correlation:

$$Nu = 0.023 Re^{0.8} Pr^{1/3}.$$

Also, note that the heat capacity rate $C = \dot{m}C_p$ of a fluid condensing or evaporating in a heat exchanger is infinity.

[Question total: 20 marks]

PART B: Answer using the multiple choice sheet

All questions are worth 2 marks. Only one answer is required per question.

- 1) A vertical wall 3 m high is at a temperature of $T_w = 60^\circ\text{C}$ and ambient air is at $T_\infty = 10^\circ\text{C}$. The properties of air are given in the table below.

μ	$1.78 \times 10^{-5} \text{ Pa s}$	ρ	1.2 kg m^{-3}
k	$0.02685 \text{ W m}^{-1} \text{ K}^{-1}$	C_p	$1.005 \text{ kJ kg}^{-1} \text{ K}^{-1}$

What is the Grashof number for this flow (select the nearest value)? **[2 marks]**

- A) $\text{Gr} \approx 2 \times 10^3$
 B) $\text{Gr} \approx 2 \times 10^5$
 C) $\text{Gr} \approx 2 \times 10^9$
 D) $\text{Gr} \approx 2 \times 10^{11}$
 E) $\text{Gr} \approx 2 \times 10^{15}$
- 2) The James Webb space telescope uses a five-layer sunshield. To what extent is the radiative flux received from the sun reduced? **[2 marks]**
- A) All radiation is removed.
 B) 1/5th of the radiation passes through.
 C) 1/6th of the radiation passes through.
 D) 1/25th of the radiation passes through.
 E) 1/60th of the radiation passes through.
- 3) Which fluid types illustrated using the flow profiles in Fig. 3 cannot be fitted using a power-law model? **[2 marks]**

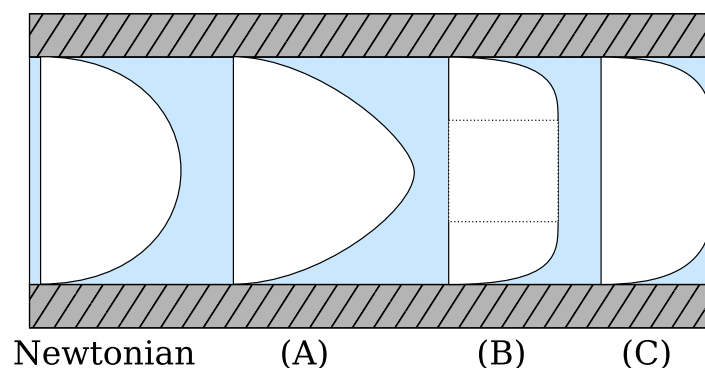


Figure 3: Non-newtonian flow profiles compared against the newtonian flow profile.

4) If a fluid has a flow index of $n = 1$, what type of fluid is it? **[2 marks]**

- A) Shear thickening.
- B) Shear thinning.
- C) Viscoplastic.
- D) Newtonian.
- E) Thixotropic.

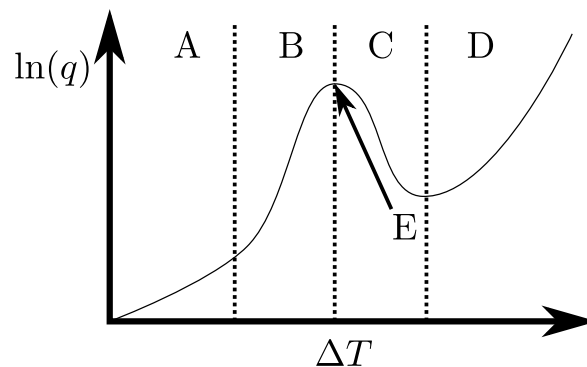


Figure 4: A typical boiling heat flux versus driving temperature difference curve.

5) In what region or point of Fig. 4 is radiation the dominant heat transfer mechanism for boiling **[2 marks]**

6) Which region or point should a boiler be operated in? **[2 marks]**

7) A wall, composed of plasterboard, brick, and insulation, is both radiating and convecting heat on one of its sides. The other side is at a constant temperature. What is the correct expression for its overall heat transfer coefficient? **[2 marks]**,

- A) $R_{total} = R_{brick} + R_{plasterboard} + R_{insulation} + (R_{radiative}^{-1} + R_{convective}^{-1})^{-1}$
- B) $R_{total}^{-1} = R_{brick}^{-1} + R_{plasterboard}^{-1} + R_{insulation}^{-1} + (R_{radiative} + R_{convective})^{-1}$
- C) $R_{total} = R_{brick} + R_{plasterboard} + R_{insulation} + R_{radiative} + R_{convective}$
- D) $R_{total}^{-1} = R_{radiative} + R_{convective} + (R_{brick}^{-1} + R_{plasterboard}^{-1} + R_{insulation}^{-1})^{-1}$

8) The walls of your house have a overall heat transfer coefficient of $U \approx 0.5 \text{ W m}^{-2} \text{ K}^{-1}$. If the temperature outside is 5°C , and inside is 23°C , what is the heat flux? **[2 marks]**

- A) 0.009 kW m^{-2}
- B) 9 kW m^{-2}
- C) 18 W m^{-2}
- D) 0.18 kW m^{-2}

- 9) What is not a valid boundary condition for an air-water interface? **[2 marks]**
- A) Stress in each phase is equal at the interface.
 - B) No-slip between the two phases at the interface.
 - C) Approximate that there is no stress at the interface.
 - D) The velocity is zero at the interface.
- 10) Consider the inside of an annulus (the zone between two concentric pipes) where the inner radius is 20% of the outer radius. What fraction of radiation emitted from the outer surface falls on the outer surface? **[2 marks]**
- A) 0.8
 - B) 0.2
 - C) 1.0
 - D) 0.0

END OF PAPER

DATASHEET

General balance equations:

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \rho \mathbf{v} \quad (\text{Mass/Continuity}) \quad (1)$$

$$\frac{\partial C_A}{\partial t} = -\nabla \cdot \mathbf{N}_A + \sigma_A \quad (\text{Species}) \quad (2)$$

$$\rho \frac{\partial \mathbf{v}}{\partial t} = -\rho \mathbf{v} \cdot \nabla \mathbf{v} - \nabla \cdot \boldsymbol{\tau} - \nabla p + \rho \mathbf{g} \quad (\text{Momentum}) \quad (3)$$

$$\rho C_p \frac{\partial T}{\partial t} = -\rho C_p \mathbf{v} \cdot \nabla T - \nabla \cdot \mathbf{q} - \boldsymbol{\tau} : \nabla \mathbf{v} - p \nabla \cdot \mathbf{v} + \sigma_{\text{energy}} \quad (\text{Heat/Energy}) \quad (4)$$

In Cartesian coordinate systems, ∇ can be treated as a vector of derivatives. In curve-linear coordinate systems, the directions $\hat{\mathbf{r}}$, $\hat{\boldsymbol{\theta}}$, and $\hat{\boldsymbol{\phi}}$ depend on the position. For convenience in these systems, look-up tables are provided for common terms involving ∇ .

Cartesian coordinates (with index notation examples)

where s is a scalar, \mathbf{v} is a vector, and $\boldsymbol{\tau}$ is a tensor.

$$\begin{aligned} \nabla s &= \nabla_i s = \left[\frac{\partial s}{\partial x}, \frac{\partial s}{\partial y}, \frac{\partial s}{\partial z} \right] \\ \nabla^2 s &= \nabla_i \nabla_i s = \frac{\partial^2 s}{\partial x^2} + \frac{\partial^2 s}{\partial y^2} + \frac{\partial^2 s}{\partial z^2} \\ \nabla \cdot \mathbf{v} &= \nabla_i v_i = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \\ \nabla \cdot \boldsymbol{\tau} &= \nabla_i \tau_{ij} \\ [\nabla \cdot \boldsymbol{\tau}]_x &= \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} \\ [\nabla \cdot \boldsymbol{\tau}]_y &= \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} \\ [\nabla \cdot \boldsymbol{\tau}]_z &= \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} \\ \mathbf{v} \cdot \nabla \mathbf{v} &= v_i \nabla_i v_j \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_x &= v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} + v_z \frac{\partial v_x}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_y &= v_x \frac{\partial v_y}{\partial x} + v_y \frac{\partial v_y}{\partial y} + v_z \frac{\partial v_y}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_z &= v_x \frac{\partial v_z}{\partial x} + v_y \frac{\partial v_z}{\partial y} + v_z \frac{\partial v_z}{\partial z} \end{aligned}$$

Cylindrical coordinates

where s is a scalar, \mathbf{v} is a vector, and τ is a tensor. All expressions involving τ are for symmetrical τ only.

$$\begin{aligned}\nabla s &= \left[\frac{\partial s}{\partial r}, \frac{1}{r} \frac{\partial s}{\partial \theta}, \frac{\partial s}{\partial z} \right] \\ \nabla^2 s &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial s}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 s}{\partial \theta^2} + \frac{\partial^2 s}{\partial z^2} \\ \nabla \cdot \mathbf{v} &= \frac{1}{r} \frac{\partial}{\partial r} (r v_r) + \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{\partial v_z}{\partial z} \\ [\nabla \cdot \tau]_r &= \frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rr}) + \frac{1}{r} \frac{\partial \tau_{r\theta}}{\partial \theta} - \frac{1}{r} \tau_{\theta\theta} + \frac{\partial \tau_{rz}}{\partial z} \\ [\nabla \cdot \tau]_\theta &= \frac{1}{r} \frac{\partial \tau_{\theta\theta}}{\partial \theta} + \frac{\partial \tau_{r\theta}}{\partial r} + \frac{2}{r} \tau_{r\theta} + \frac{\partial \tau_{\theta z}}{\partial z} \\ [\nabla \cdot \tau]_z &= \frac{1}{r} \frac{\partial}{\partial r} (r \tau_{rz}) + \frac{1}{r} \frac{\partial \tau_{\theta z}}{\partial \theta} + \frac{\partial \tau_{zz}}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_r &= v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta^2}{r} + v_z \frac{\partial v_r}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_\theta &= v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r v_\theta}{r} + v_z \frac{\partial v_\theta}{\partial z} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_z &= v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z}\end{aligned}$$

Spherical coordinates

where s is a scalar, \mathbf{v} is a vector, and τ is a tensor. All expressions involving τ are for symmetrical τ only.

$$\begin{aligned}\nabla s &= \left[\frac{\partial s}{\partial r}, \frac{1}{r} \frac{\partial s}{\partial \theta}, \frac{1}{r \sin \theta} \frac{\partial s}{\partial \phi} \right] \\ \nabla^2 s &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial s}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial s}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 s}{\partial \phi^2} \\ \nabla \cdot \mathbf{v} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (v_\theta \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} \\ [\nabla \cdot \tau]_r &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{rr}) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\tau_{r\theta} \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial \tau_{r\phi}}{\partial \phi} - \frac{\tau_{\theta\theta} + \tau_{\phi\phi}}{r} \\ [\nabla \cdot \tau]_\theta &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{r\theta}) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\tau_{\theta\theta} \sin \theta) + \frac{1}{r \sin \theta} \frac{\partial \tau_{\theta\phi}}{\partial \phi} + \frac{\tau_{r\theta}}{r} - \frac{\cot \theta}{r} \tau_{\phi\phi} \\ [\nabla \cdot \tau]_\phi &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \tau_{r\phi}) + \frac{1}{r} \frac{\partial \tau_{\theta\phi}}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial \tau_{\phi\phi}}{\partial \phi} + \frac{\tau_{r\theta}}{r} + \frac{2 \cot \theta}{r} \tau_{\theta\phi} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_r &= v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{v_\theta^2 + v_\phi^2}{r} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_\theta &= v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial v_\theta}{\partial \phi} + \frac{v_r v_\theta - v_\phi^2 \cot \theta}{r} \\ [\mathbf{v} \cdot \nabla \mathbf{v}]_\phi &= v_r \frac{\partial v_\phi}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\phi}{\partial \theta} + \frac{v_\phi}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} + \frac{v_r v_\phi + v_\theta v_\phi \cot \theta}{r}\end{aligned}$$

Rectangular		Cylindrical		Spherical	
q_x	$-k \frac{\partial T}{\partial x}$	q_r	$-k \frac{\partial T}{\partial r}$	q_r	$-k \frac{\partial T}{\partial r}$
q_y	$-k \frac{\partial T}{\partial y}$	q_θ	$-k \frac{1}{r} \frac{\partial T}{\partial \theta}$	q_θ	$-k \frac{1}{r} \frac{\partial T}{\partial \theta}$
q_z	$-k \frac{\partial T}{\partial z}$	q_z	$-k \frac{\partial T}{\partial z}$	q_ϕ	$-k \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi}$
τ_{xx}	$-2\mu \frac{\partial v_x}{\partial x} + \mu^B \nabla \cdot \mathbf{v}$	τ_{rr}	$-2\mu \frac{\partial v_r}{\partial r} + \mu^B \nabla \cdot \mathbf{v}$	τ_{rr}	$-2\mu \frac{\partial v_r}{\partial r} + \mu^B \nabla \cdot \mathbf{v}$
τ_{yy}	$-2\mu \frac{\partial v_y}{\partial y} + \mu^B \nabla \cdot \mathbf{v}$	$\tau_{\theta\theta}$	$-2\mu \left(\frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r}{r} \right) + \mu^B \nabla \cdot \mathbf{v}$	$\tau_{\theta\theta}$	$-2\mu \left(\frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r}{r} \right) + \mu^B \nabla \cdot \mathbf{v}$
τ_{zz}	$-2\mu \frac{\partial v_z}{\partial z} + \mu^B \nabla \cdot \mathbf{v}$	τ_{zz}	$-2\mu \frac{\partial v_z}{\partial z} + \mu^B \nabla \cdot \mathbf{v}$	$\tau_{\phi\phi}$	$-2\mu \left(\frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi} + \frac{v_r + v_\theta \cot \theta}{r} \right) + \mu^B \nabla \cdot \mathbf{v}$
τ_{xy}	$-\mu \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right)$	$\tau_{r\theta}$	$-\mu \left(r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_r}{\partial \theta} \right)$	$\tau_{r\theta}$	$-\mu \left(r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_r}{\partial \theta} \right)$
τ_{yz}	$-\mu \left(\frac{\partial v_y}{\partial z} + \frac{\partial v_z}{\partial y} \right)$	$\tau_{\theta z}$	$-\mu \left(\frac{1}{r} \frac{\partial v_z}{\partial \theta} + \frac{\partial v_\theta}{\partial z} \right)$	$\tau_{\theta\phi}$	$-\mu \left(\frac{\sin \theta}{r} \frac{\partial}{\partial \theta} \left(\frac{v_\phi}{\sin \theta} \right) + \frac{1}{r \sin \theta} \frac{\partial v_\theta}{\partial \phi} \right)$
τ_{xz}	$-\mu \left(\frac{\partial v_x}{\partial z} + \frac{\partial v_z}{\partial x} \right)$	τ_{zr}	$-\mu \left(\frac{\partial v_r}{\partial z} + \frac{\partial v_z}{\partial r} \right)$	$\tau_{\phi r}$	$-\mu \left(\frac{1}{r \sin \theta} \frac{\partial v_r}{\partial \phi} + r \frac{\partial}{\partial r} \left(\frac{v_\phi}{r} \right) \right)$

Table 1: Fourier's law for the heat flux and Newton's law for the stress in several coordinate systems. Please remember that the stress is symmetric, so $\tau_{ij} = \tau_{ji}$.

Viscous models:

Power-Law Fluid:

$$|\tau_{xy}| = k \left| \frac{\partial v_x}{\partial y} \right|^n \quad (5)$$

Bingham-Plastic Fluid:

$$\frac{\partial v_x}{\partial y} = \begin{cases} -\mu^{-1} (\tau_{xy} - \tau_0) & \text{if } \tau_{xy} > \tau_0 \\ 0 & \text{if } \tau_{xy} \leq \tau_0 \end{cases}$$

Dimensionless Numbers

$$\text{Re} = \frac{\rho \langle v \rangle D}{\mu} \quad \text{Re}_H = \frac{\rho \langle v \rangle D_H}{\mu} \quad \text{Re}_{MR} = -\frac{16 L \rho \langle v \rangle^2}{R \Delta p} \quad (6)$$

The hydraulic diameter is defined as $D_H = 4 A / P_w$.

Single phase pressure drop calculations in pipes:

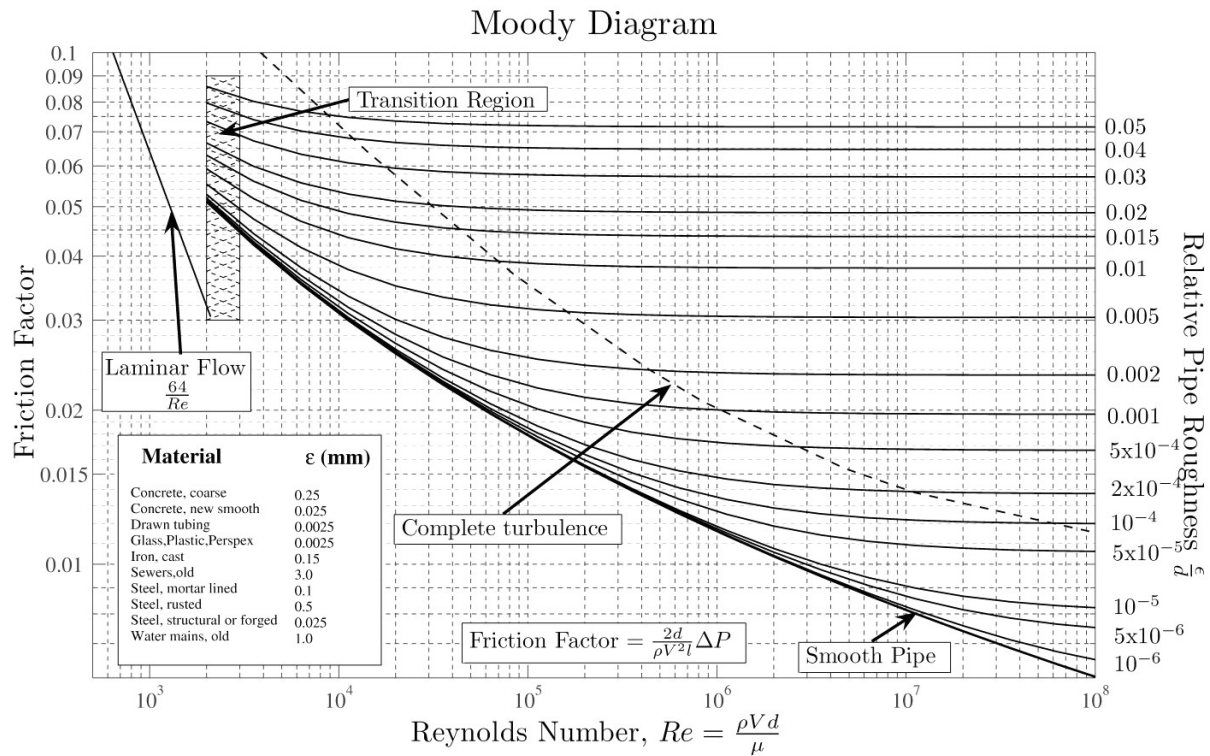
Darcy-Weisbach equation:

$$\frac{\Delta p}{L} = -\frac{C_f \rho \langle v \rangle^2}{R} \quad (7)$$

where $C_f = 16/Re$ for laminar Newtonian flow. For turbulent flow of Newtonian fluids in smooth pipes, we have the Blasius correlation:

$$C_f = 0.079 Re^{-1/4} \quad \text{for } 2.5 \times 10^3 < Re < 10^5 \text{ and smooth pipes.}$$

Otherwise, you may refer to the Moody diagram.



Laminar Power-Law fluid:

$$\dot{V} = \frac{n \pi R^3}{3n+1} \left(\frac{R}{2k} \right)^{\frac{1}{n}} \left(-\frac{\Delta p}{L} \right)^{\frac{1}{n}}$$

Two-Phase Flow:

Lockhart-Martinelli parameter:

$$X^2 = \frac{\Delta p_{liq.-only}}{\Delta p_{gas-only}}$$

Pressure drop calculation:

$$\Delta p_{two-phase} = \Phi_{liq.}^2 \Delta p_{liq.-only} = \Phi_{gas}^2 \Delta p_{gas-only}$$

Chisholm's relation:

$$\Phi_{gas}^2 = 1 + cX + X^2$$

$$\Phi_{liq.}^2 = 1 + \frac{c}{X} + \frac{1}{X^2}$$

$$c = \begin{cases} 20 & \text{turbulent liquid \& turbulent gas} \\ 12 & \text{laminar liquid \& turbulent gas} \\ 10 & \text{turbulent liquid \& laminar gas} \\ 5 & \text{laminar liquid \& laminar gas} \end{cases}$$

Farooqi and Richardson expression for liquid hold-up in co-current flows of Newtonian fluids and air in horizontal pipes:

$$h = \begin{cases} 0.186 + 0.0191 X & 1 < X < 5 \\ 0.143 X^{0.42} & 5 < X < 50 \\ 1 / (0.97 + 19/X) & 50 < X < 500 \end{cases}$$

Heat Transfer Dimensionless numbers:

$$\text{Nu} = \frac{hL}{k} \quad \text{Pr} = \frac{\mu C_p}{k} \quad \text{Gr} = \frac{g \beta (T_w - T_\infty) L^3}{\nu^2}$$

where $\beta = V^{-1}(\partial V / \partial T)$.

Heat transfer: Resistances

$$Q = U_T A_T \Delta T = R_T^{-1} \Delta T$$

	Conduction Shell Resistances			Radiation
	Rect.	Cyl.	Sph.	
R	$\frac{X}{kA}$	$\frac{\ln(R_{outer}/R_{inner})}{2\pi Lk}$	$\frac{R_{inner}^{-1} - R_{outer}^{-1}}{4\pi k}$	$\left[A \varepsilon \sigma (T_j^2 + T_i^2) (T_j + T_i) \right]^{-1}$

Radiation Heat Transfer:

Stefan-Boltzmann constant $\sigma = 5.6703 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$.

Summation relationship, $\sum_j F_{i \rightarrow j} = 1$, and reciprocity relationship, $F_{i \rightarrow j} A_i = F_{j \rightarrow i} A_j$.

Radiation shielding factor $1/(N+1)$.

$$Q_{rad., i \rightarrow j} = \sigma \varepsilon F_{i \rightarrow j} A_i (T_j^4 - T_i^4) = h_{rad.} A (T_\infty - T_w)$$

Natural Convection

$Ra = Gr Pr$	C	m
$< 10^4$	1.36	1/5
$10^4 - 10^9$	0.59	1/4
$> 10^9$	0.13	1/3

Table 2: Natural convection coefficients for isothermal vertical plates in the empirical relation $Nu \approx C (Gr Pr)^m$.

For isothermal vertical cylinders, the above expressions for isothermal vertical plates may be used but must be scaled by a factor, F (i.e., $Nu_{v.cyl.} = F Nu_{v.plate}$):

$$F = \begin{cases} 1 & \text{for } (D/H) \geq 35 Gr_H^{-1/4} \\ 1.3 [H D^{-1} Gr_D^{-1}]^{1/4} + 1 & \text{for } (D/H) < 35 Gr_H^{-1/4} \end{cases}$$

where D is the diameter and H is the height of the cylinder. The subscript on Gr indicates which length is to be used as the critical length to calculate the Grashof number.

Churchill and Chu expression for natural convection from a horizontal pipe:

$$Nu^{1/2} = 0.6 + 0.387 \left\{ \frac{Gr Pr}{[1 + (0.559/Pr)^{9/16}]^{16/9}} \right\}^{1/6} \quad \text{for } 10^{-5} < Gr Pr < 10^{12}$$

Forced Convection:

Laminar flows:

$$Nu \approx 0.332 Re^{1/2} Pr^{1/3}$$

Well-Developed turbulent flows in smooth pipes:

$$Nu \approx \frac{(C_f/2) Re Pr}{1.07 + 12.7(C_f/2)^{1/2} (Pr^{2/3} - 1)} \left(\frac{\mu_b}{\mu_w} \right)^{0.14}$$

Boiling:

Forster-Zuber pool-boiling coefficient:

$$h_{nb} = 0.00122 \frac{k_L^{0.79} C_{p,L}^{0.45} \rho_L^{0.49}}{\gamma^{0.5} \mu_L^{0.29} h_{fg}^{0.24} \rho_G^{0.24}} (T_w - T_{sat})^{0.24} (p_w - p_{sat})^{0.75}$$

Mostinski correlations:

$$h_{nb} = 0.104 p_c^{0.69} q^{0.7} \left[1.8 \left(\frac{p}{p_c} \right)^{0.17} + 4 \left(\frac{p}{p_c} \right)^{1.2} + 10 \left(\frac{p}{p_c} \right)^{10} \right]$$

$$q_c = 3.67 \times 10^4 p_c \left(\frac{p}{p_c} \right)^{0.35} \left[1 - \frac{p}{p_c} \right]^{0.9}$$

(Note: for the Mostinski correlations, the pressures are in units of bar)

Condensing:

Horizontal pipes

$$h = 0.72 \left(\frac{k^3 \rho^2 g_x E_{latent}}{D \mu (T_w - T_\infty)} \right)^{1/4}$$

Lumped capacitance method:

$$\text{Bi} = \frac{h L_c}{\kappa}$$

$$L_c = V/A \quad \text{for Bi} < 0.1$$

$$\frac{T(t) - T_\infty}{T_0 - T_\infty} = e^{-bt} \quad b = \frac{hA_s}{\rho V C_p}$$

1-D Transient Heat Conduction:

$$Fo = \frac{\alpha \Delta t}{(\Delta x)^2} = \tau, \quad \alpha = \kappa (\rho C_p)^{-1}$$

$$\theta_{\text{wall}} = \frac{T(x, t) - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau} \cos\left(\frac{\lambda_1 x}{L}\right), \quad \theta_{\text{cyl}} = \frac{T(r, t) - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau} \mathbf{J}_0\left(\frac{\lambda_1 r}{r_0}\right)$$

$$\theta_{\text{sph}} = \frac{T(r, t) - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau} \frac{\sin\left(\frac{\lambda_1 r}{r_0}\right)}{\frac{\lambda_1 r}{r_0}}$$

$$\theta_{0,\text{wall}} = \theta_{0,\text{cyl}} = \theta_{0,\text{sph}} = \frac{T_0 - T_\infty}{T_i - T_\infty} = A_1 e^{-\lambda_1^2 \tau}$$

$$\left(\frac{Q}{Q_{\text{max}}}\right)_{\text{wall}} = 1 - \theta_{0,\text{wall}} \frac{\sin \lambda_1}{\lambda_1}, \quad \left(\frac{Q}{Q_{\text{max}}}\right)_{\text{cyl}} = 1 - 2\theta_{0,\text{cyl}} \frac{\mathbf{J}_1(\lambda_1)}{\lambda_1}$$

$$\left(\frac{Q}{Q_{\text{max}}}\right)_{\text{sph}} = 1 - 3\theta_{0,\text{sph}} \frac{\sin \lambda_1 - \lambda_1 \cos \lambda_1}{\lambda_1^3}$$

Finite-Difference Method:

$$\frac{\partial}{\partial t}(\rho \phi) + \frac{\partial}{\partial x}(\rho \mathbf{v} \phi) = \nabla \cdot (\Gamma \nabla \phi) + \mathcal{S} \quad (1D \text{ transport equation})$$

$$\left(\frac{d\phi}{dx}\right)_i = \frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x} \quad \text{and} \quad \left(\frac{d^2\phi}{dx^2}\right)_i = \frac{\phi_{i-1} + \phi_{i+1} - 2\phi_i}{(\Delta x)^2}$$

$$T_i^{j+1} = (1 - 2\tau) T_i^j + \tau (T_{i+1}^j + T_{i-1}^j) + \frac{\tau (\Delta x)^2}{\kappa} S_i^j$$

Overall Heat Transfer Coefficient:

$$\dot{Q} = \frac{\Delta T}{\mathcal{R}} = UA\Delta T = U_i A_i \Delta T = U_o A_o \Delta T$$

$$\mathcal{R} = R_i + R_{\text{wall}} + R_o = \frac{1}{h_i A_i} + \frac{\ln D_o/D_i}{2\pi\kappa L} + \frac{1}{h_o A_o}$$

Fouling Factor:

$$\mathcal{R} = \frac{1}{h_i A_i} + \frac{R_{f,i}}{A_i} + R_{\text{wall}} + \frac{R_{f,o}}{A_o} + \frac{1}{h_o A_o}$$

LMTD Method:

$$\dot{Q} = UA_s \Delta T_{\text{lm}} \quad \text{with} \quad \Delta T_{\text{lm}} = \frac{\Delta T_2 - \Delta T_1}{\ln \frac{\Delta T_2}{\Delta T_1}} = \frac{\Delta T_1 - \Delta T_2}{\ln \frac{\Delta T_1}{\Delta T_2}}$$

$$\text{Parallel flows: } \begin{cases} \Delta T_1 = T_{\text{hot,in}} - T_{\text{cold,in}} \\ \Delta T_2 = T_{\text{hot,out}} - T_{\text{cold,out}} \end{cases}$$

$$\text{Counter flows: } \begin{cases} \Delta T_1 = T_{\text{hot,in}} - T_{\text{cold,out}} \\ \Delta T_2 = T_{\text{hot,out}} - T_{\text{cold,in}} \end{cases}$$

ε-NTU Method:

$$\epsilon = \frac{\dot{Q}}{\dot{Q}_{\text{max}}}, \quad \text{with } \dot{Q}_{\text{max}} = C_{\text{min}} (T_{\text{hot,in}} - T_{\text{cold,in}}) \quad \text{and} \quad C_{\text{min}} = \text{Min} \{ \dot{m}_{\text{hot}} C_{p,\text{hot}}, \dot{m}_{\text{cold}} C_{p,\text{cold}} \}$$

$$\text{NTU} = \frac{UA_s}{C_{\text{min}}}$$

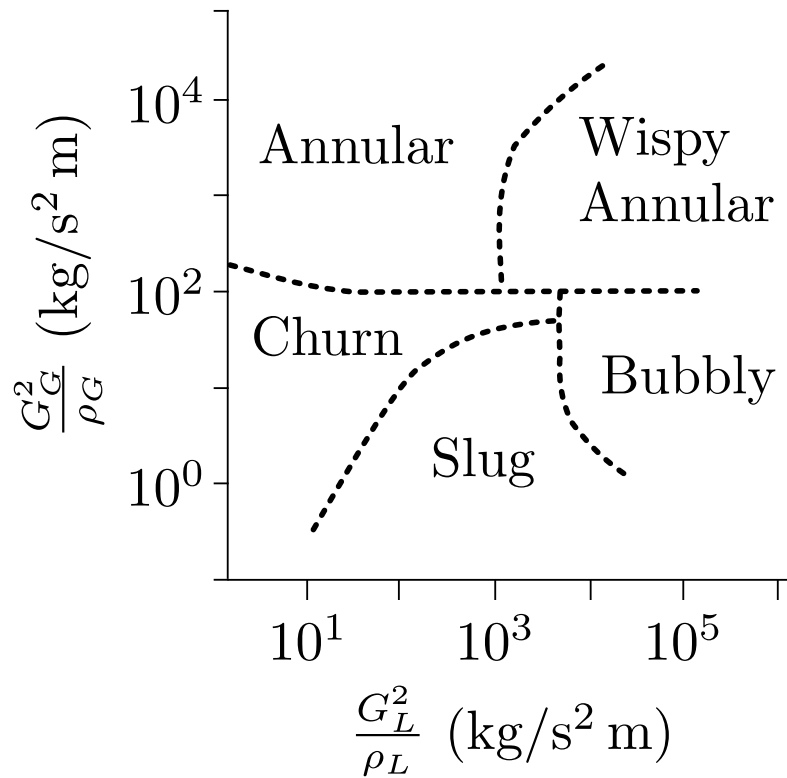


Figure 5: *Hewitt-Taylor flow pattern map for multiphase flows in vertical pipes.*

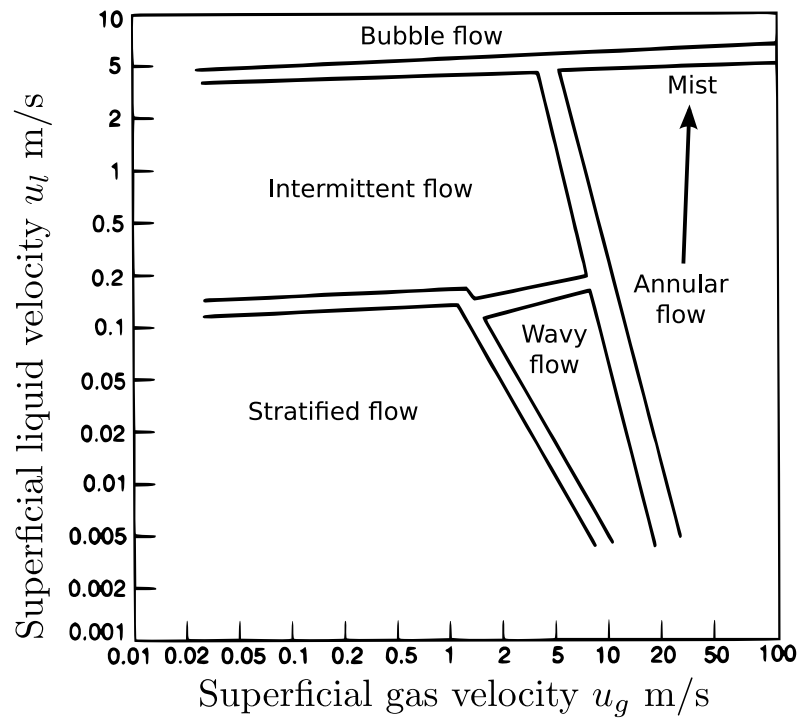


Figure 6: *Chhabra and Richardson flow pattern map for horizontal pipes.*

TABLE 4–2

Coefficients used in the one-term approximate solution of transient one-dimensional heat conduction in plane walls, cylinders, and spheres ($Bi = hL/k$ for a plane wall of thickness $2L$, and $Bi = hr_o/k$ for a cylinder or sphere of radius r_o)

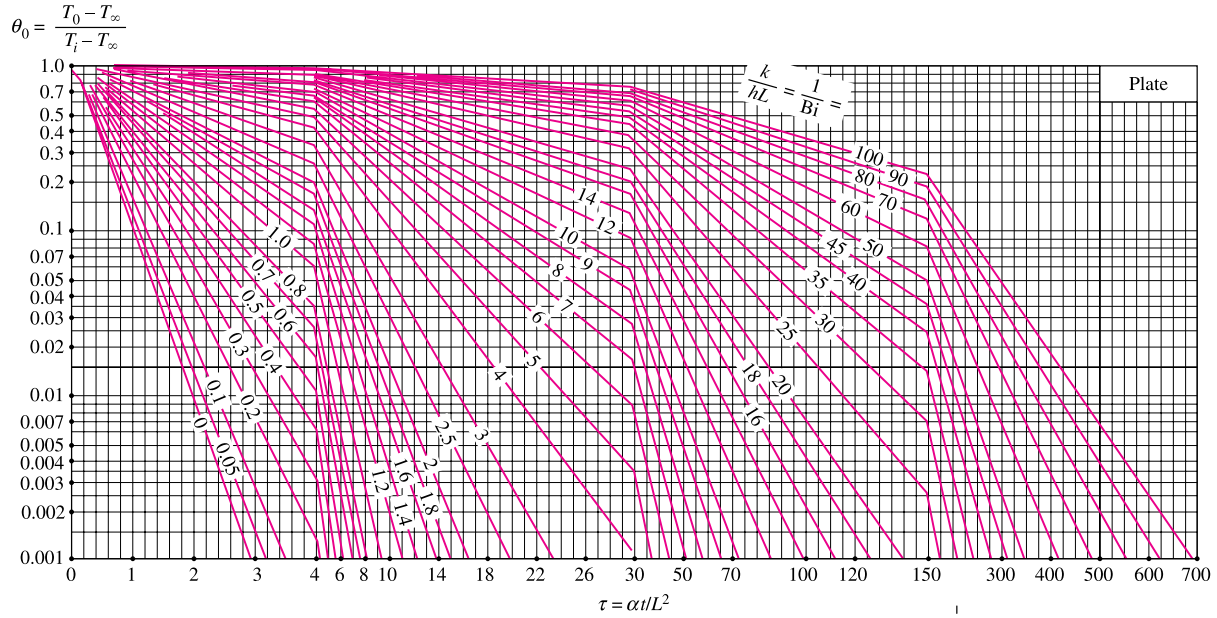
Bi	Plane Wall		Cylinder		Sphere	
	λ_1	A_1	λ_1	A_1	λ_1	A_1
0.01	0.0998	1.0017	0.1412	1.0025	0.1730	1.0030
0.02	0.1410	1.0033	0.1995	1.0050	0.2445	1.0060
0.04	0.1987	1.0066	0.2814	1.0099	0.3450	1.0120
0.06	0.2425	1.0098	0.3438	1.0148	0.4217	1.0179
0.08	0.2791	1.0130	0.3960	1.0197	0.4860	1.0239
0.1	0.3111	1.0161	0.4417	1.0246	0.5423	1.0298
0.2	0.4328	1.0311	0.6170	1.0483	0.7593	1.0592
0.3	0.5218	1.0450	0.7465	1.0712	0.9208	1.0880
0.4	0.5932	1.0580	0.8516	1.0931	1.0528	1.1164
0.5	0.6533	1.0701	0.9408	1.1143	1.1656	1.1441
0.6	0.7051	1.0814	1.0184	1.1345	1.2644	1.1713
0.7	0.7506	1.0918	1.0873	1.1539	1.3525	1.1978
0.8	0.7910	1.1016	1.1490	1.1724	1.4320	1.2236
0.9	0.8274	1.1107	1.2048	1.1902	1.5044	1.2488
1.0	0.8603	1.1191	1.2558	1.2071	1.5708	1.2732
2.0	1.0769	1.1785	1.5995	1.3384	2.0288	1.4793
3.0	1.1925	1.2102	1.7887	1.4191	2.2889	1.6227
4.0	1.2646	1.2287	1.9081	1.4698	2.4556	1.7202
5.0	1.3138	1.2403	1.9898	1.5029	2.5704	1.7870
6.0	1.3496	1.2479	2.0490	1.5253	2.6537	1.8338
7.0	1.3766	1.2532	2.0937	1.5411	2.7165	1.8673
8.0	1.3978	1.2570	2.1286	1.5526	2.7654	1.8920
9.0	1.4149	1.2598	2.1566	1.5611	2.8044	1.9106
10.0	1.4289	1.2620	2.1795	1.5677	2.8363	1.9249
20.0	1.4961	1.2699	2.2880	1.5919	2.9857	1.9781
30.0	1.5202	1.2717	2.3261	1.5973	3.0372	1.9898
40.0	1.5325	1.2723	2.3455	1.5993	3.0632	1.9942
50.0	1.5400	1.2727	2.3572	1.6002	3.0788	1.9962
100.0	1.5552	1.2731	2.3809	1.6015	3.1102	1.9990
∞	1.5708	1.2732	2.4048	1.6021	3.1416	2.0000

TABLE 4–3

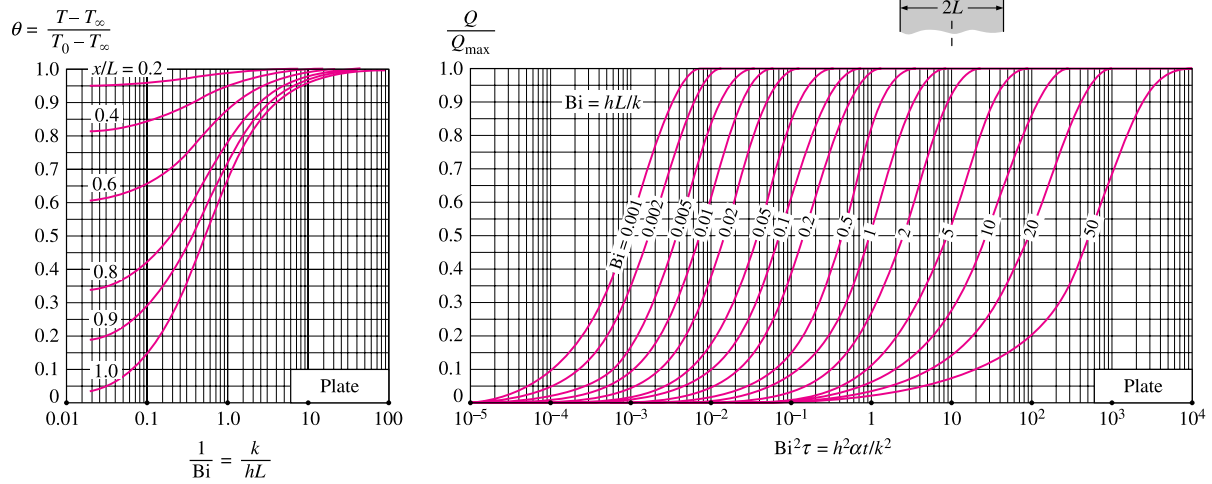
The zeroth- and first-order Bessel functions of the first kind

η	$J_0(\eta)$	$J_1(\eta)$
0.0	1.0000	0.0000
0.1	0.9975	0.0499
0.2	0.9900	0.0995
0.3	0.9776	0.1483
0.4	0.9604	0.1960
0.5	0.9385	0.2423
0.6	0.9120	0.2867
0.7	0.8812	0.3290
0.8	0.8463	0.3688
0.9	0.8075	0.4059
1.0	0.7652	0.4400
1.1	0.7196	0.4709
1.2	0.6711	0.4983
1.3	0.6201	0.5220
1.4	0.5669	0.5419
1.5	0.5118	0.5579
1.6	0.4554	0.5699
1.7	0.3980	0.5778
1.8	0.3400	0.5815
1.9	0.2818	0.5812
2.0	0.2239	0.5767
2.1	0.1666	0.5683
2.2	0.1104	0.5560
2.3	0.0555	0.5399
2.4	0.0025	0.5202
2.6	−0.0968	−0.4708
2.8	−0.1850	−0.4097
3.0	−0.2601	−0.3391
3.2	−0.3202	−0.2613

Figure 7: Coefficients for the 1D transient equations.



(a) Midplane temperature (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)

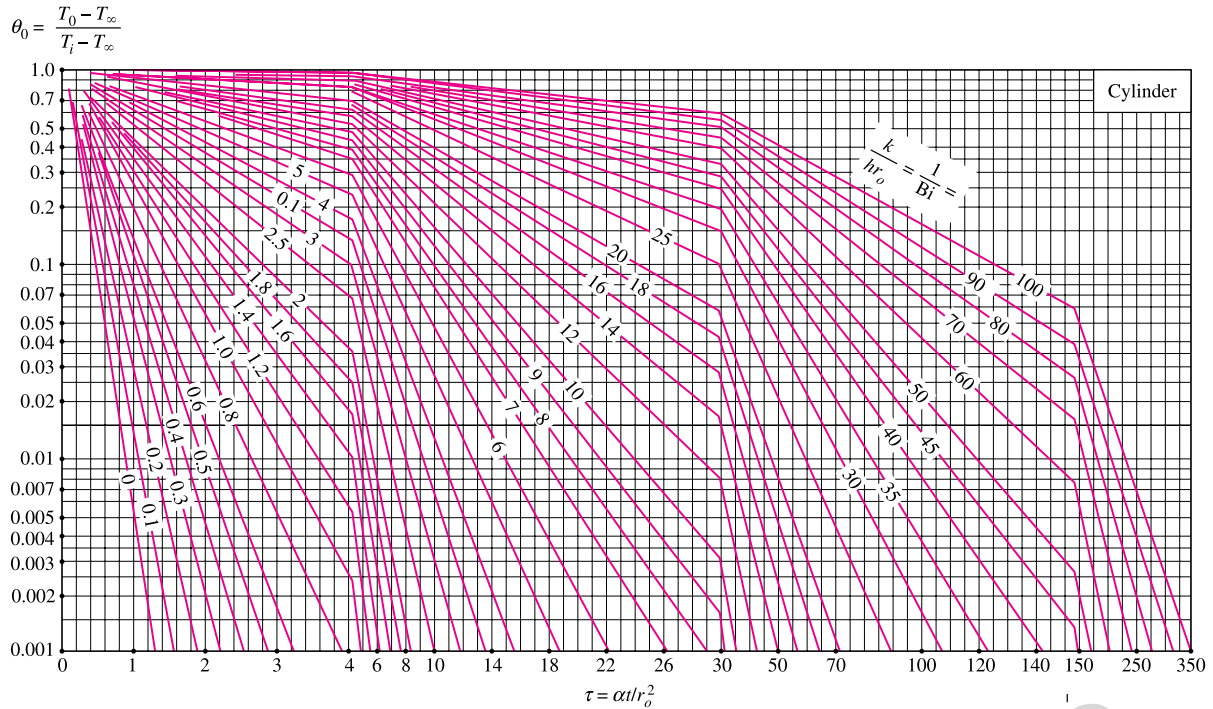


(b) Temperature distribution (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)

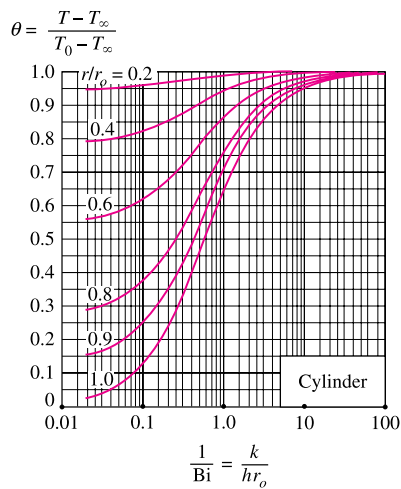
(c) Heat transfer (from H. Gröber et al.)

Transient temperature and heat transfer charts for a plane wall of thickness $2L$ initially at a uniform temperature T_i subjected to convection from both sides to an environment at temperature T_∞ with a convection coefficient of h .

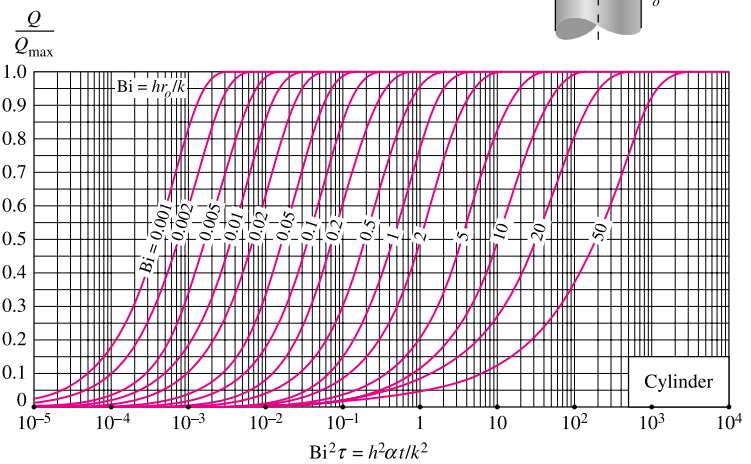
Figure 8:



(a) Centerline temperature (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)



(b) Temperature distribution (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)

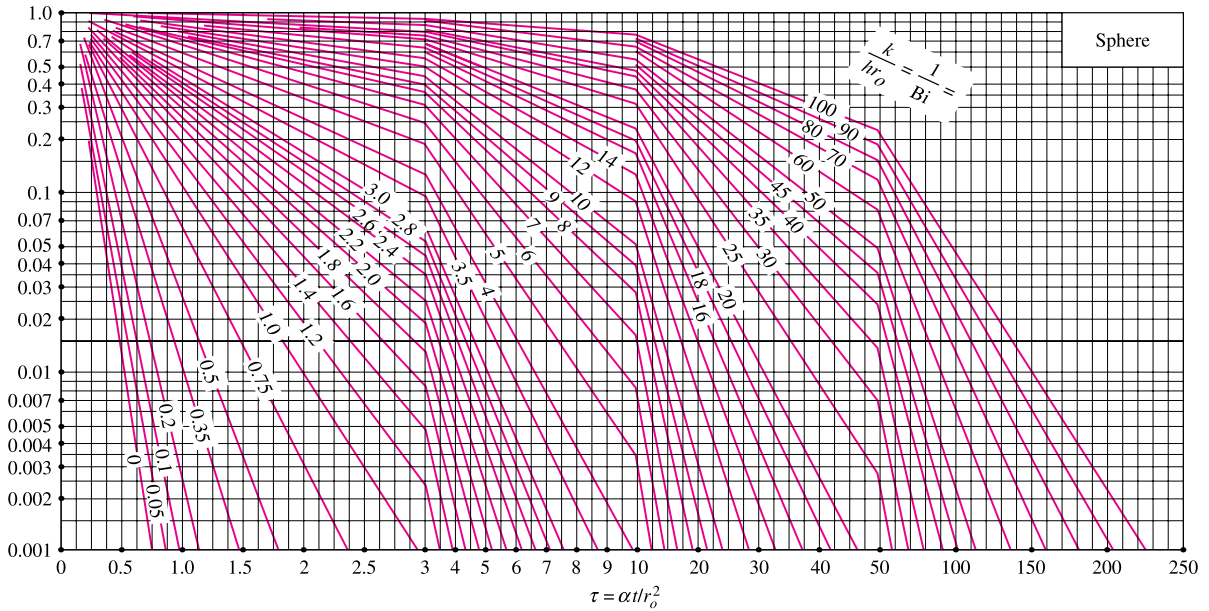


(c) Heat transfer (from H. Gröber et al.)

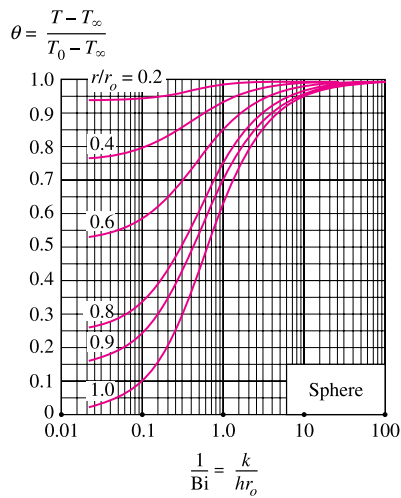
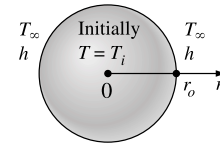
Transient temperature and heat transfer charts for a long cylinder of radius r_o initially at a uniform temperature T_i subjected to convection from all sides to an environment at temperature T_∞ with a convection coefficient of h .

Figure 9:

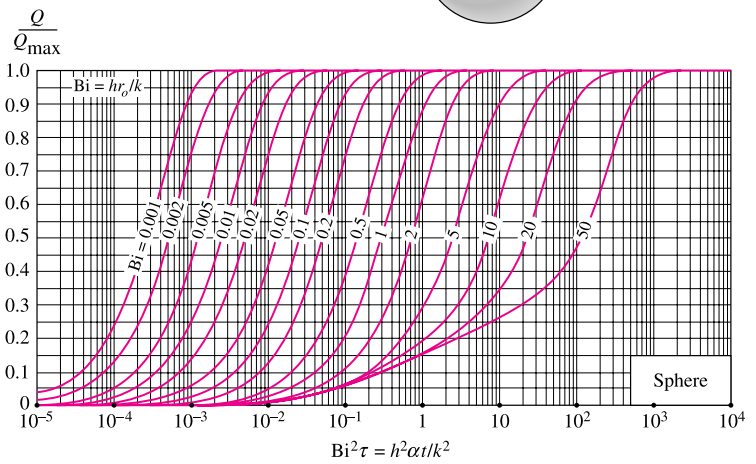
$$\theta_0 = \frac{T_0 - T_\infty}{T_i - T_\infty}$$



(a) Midpoint temperature (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947, pp. 227–36. Reprinted by permission of ASME International.)



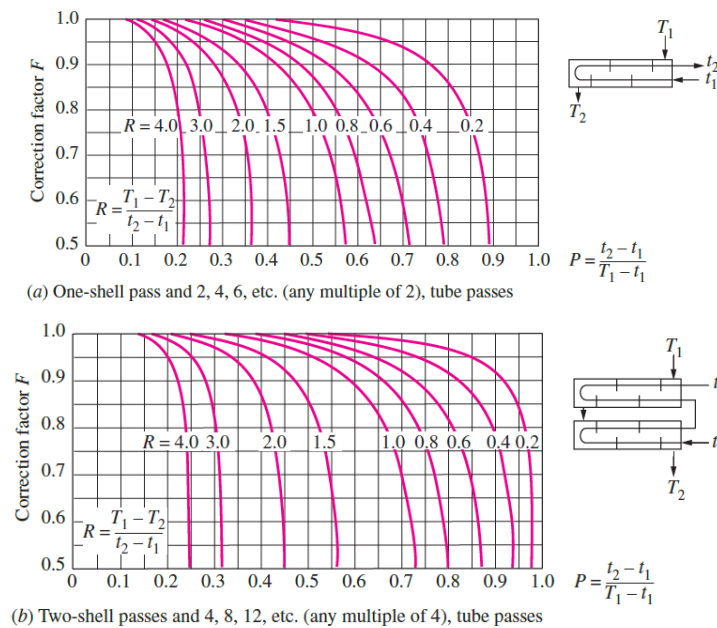
(b) Temperature distribution (from M. P. Heisler, "Temperature Charts for Induction and Constant Temperature Heating," *Trans. ASME* 69, 1947,



(c) Heat transfer (from H. Gröber et al.)

Transient temperature and heat transfer charts for a sphere of radius r_o initially at a uniform temperature T_i subjected to convection from all sides to an environment at temperature T_∞ with a convection coefficient of h .

Figure 10:

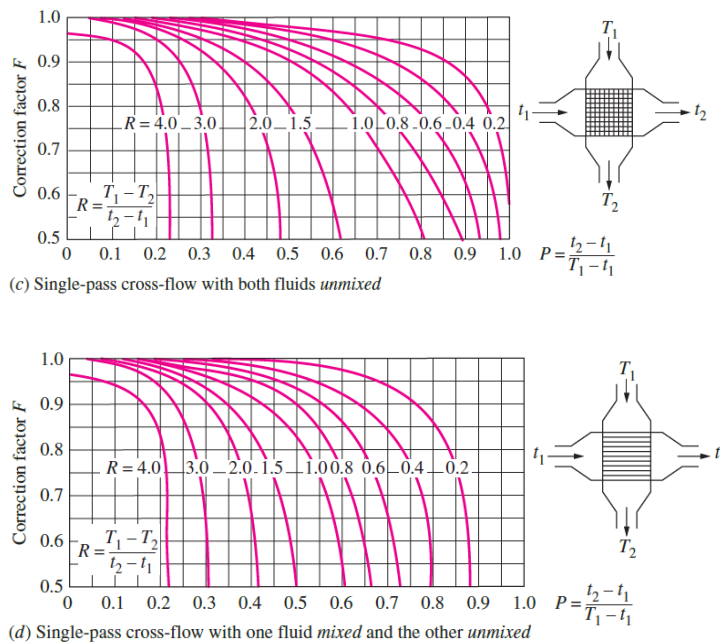


Extracted from Y.A. Cengel, "Heat Transfer: A Practical Approach", 2nd Edition.

Figure 10.8

Correction factor F charts for common shell-and-tube and cross-flow heat exchangers (from Bowman, Mueller, and Nagle, Ref. 2).

Figure 11: Correction-factors for LMTD Method, extracted from Y. A. Cengel, "Heat transfer:A practical approach", 2nd Ed.



Extracted from Y.A. Cengel, "Heat Transfer: A Practical Approach", 2nd Edition.

Figure 10.8

Correction factor F charts for common shell-and-tube and cross-flow heat exchangers (from Bowman, Mueller, and Nagle, Ref. 2).

Figure 12: Correction-factors for LMTD Method, extracted from Y. A. Cengel, "Heat transfer:A practical approach", 2nd Ed.

Effectiveness relations for heat exchangers: $NTU = UA_s/C_{\min}$ and $c = C_{\min}/C_{\max} = (\dot{m}C_p)_{\min}/(\dot{m}C_p)_{\max}$ (Kays and London, Ref. 5.)

Heat exchanger type	Effectiveness relation
1 Double-pipe: Parallel-flow	$\varepsilon = \frac{1 - \exp[-NTU(1+c)]}{1+c}$
Counter-flow	$\varepsilon = \frac{1 - \exp[-NTU(1-c)]}{1-c \exp[-NTU(1-c)]}$
2 Shell and tube: One-shell pass 2, 4, ... tube passes	$\varepsilon = 2 \left\{ 1 + c + \sqrt{1+c^2} \frac{1 + \exp[-NTU\sqrt{1+c^2}]}{1 - \exp[-NTU\sqrt{1+c^2}]} \right\}^{-1}$
3 Cross-flow (single-pass) Both fluids unmixed	$\varepsilon = 1 - \exp \left\{ \frac{NTU^{0.22}}{c} [\exp(-c NTU^{0.78}) - 1] \right\}$
C_{\max} mixed, C_{\min} unmixed	$\varepsilon = \frac{1}{c} (1 - \exp[1 - c(1 - \exp(-NTU))])$
C_{\min} mixed, C_{\max} unmixed	$\varepsilon = 1 - \exp \left\{ -\frac{1}{c} [1 - \exp(-c NTU)] \right\}$
4 All heat exchangers with $c = 0$	$\varepsilon = 1 - \exp(-NTU)$

NTU relations for heat exchangers $NTU = UA_s/C_{\min}$ and $c = C_{\min}/C_{\max} = (\dot{m}C_p)_{\min}/(\dot{m}C_p)_{\max}$ (Kays and London, Ref. 5.)

Heat exchanger type	NTU relation
1 Double-pipe: Parallel-flow	$NTU = -\frac{\ln[1 - \varepsilon(1+c)]}{1+c}$
Counter-flow	$NTU = \frac{1}{c-1} \ln \left(\frac{\varepsilon-1}{\varepsilon c-1} \right)$
2 Shell and tube: One-shell pass 2, 4, ... tube passes	$NTU = -\frac{1}{\sqrt{1+c^2}} \ln \left(\frac{2/\varepsilon - 1 - c - \sqrt{1+c^2}}{2/\varepsilon - 1 - c + \sqrt{1+c^2}} \right)$
3 Cross-flow (single-pass) C_{\max} mixed, C_{\min} unmixed	$NTU = -\ln \left[1 + \frac{\ln(1-\varepsilon c)}{c} \right]$
C_{\min} mixed, C_{\max} unmixed	$NTU = \frac{-\ln[c \ln(1-\varepsilon) + 1]}{c}$
4 All heat exchangers with $c = 0$	$NTU = -\ln(1-\varepsilon)$

Figure 13: NTU relations extracted from Y. A. Cengel, "Heat transfer: A practical approach", 2nd Ed.

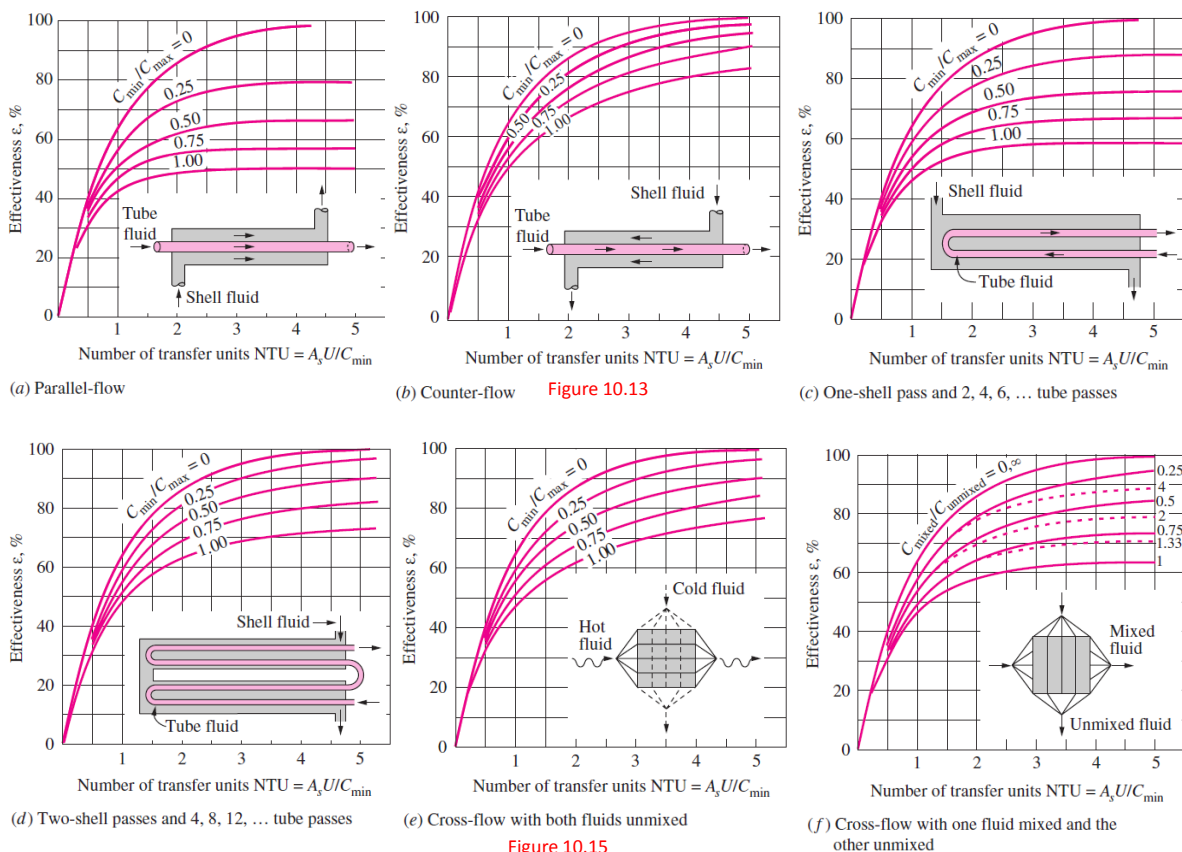


Figure 10.15

Extracted from Y.A. Cengel, "Heat Transfer: A Practical Approach", 2nd Edition.

Figure 14: NTU plots extracted from Y. A. Cengel, "Heat transfer: A practical approach", 2nd Ed.

Diffusion Dimensionless Numbers

$$Sc = \frac{\mu}{\rho D_{AB}}$$

$$Le = \frac{k}{\rho C_p D_{AB}}$$

Diffusion

General expression for the flux:

$$\mathbf{N}_A = \mathbf{J}_A + x_A \sum_B \mathbf{N}_B$$

Fick's law:

$$\mathbf{J}_A = -D_{AB} \nabla C_A$$

Stefan's law:

$$N_{s,r} = -D \frac{c}{1-x} \frac{\partial x}{\partial r}$$

Ideal Gas

$$P V = n R T$$

$$R \approx 8.314598 \text{ J K}^{-1} \text{ mol}^{-1}$$

Geometry

$$P_{\text{circle}} = 2 \pi r$$

$$A_{\text{circle}} = \pi r^2$$

$$A_{\text{sphere}} = 4 \pi r^2$$

$$V_{\text{sphere}} = \frac{4}{3} \pi r^3$$

$$A_{\text{cylinder}} = P_{\text{circle}} L$$

$$V_{\text{cylinder}} = A_{\text{circle}} L$$