

# Introduction to general relativity: prerelativity gravitation

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Text: Wald, "General Relativity"

Other sources:

- Misner, Thorne, & Wheeler
- Weinberg

Structure: Mathematical methods  $\rightarrow$  Physics

$\uparrow$   
Differential geometry  
 $\swarrow \downarrow \searrow$

Quantum gravity & string theory.

Videos: Lenny Susskind

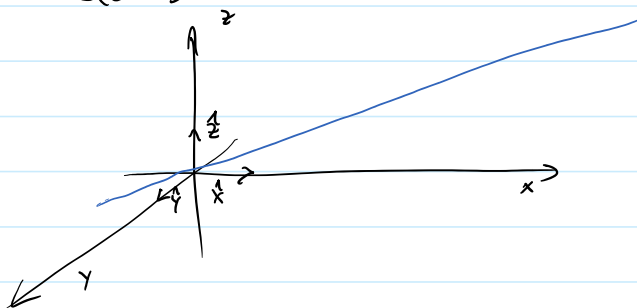
Assume:

- \* Special relativity (Theoretische Physik A)
- \* Calculus
- \* linear algebra
- \* combinatorics

## Lecture 1: prerelativity gravitation

Assumption

Newtonian physics: existence of a preferred reference system & clocks.



A free test mass (i.e. no force acts) moves along a straight trajectory at a constant rate uniformly in time.

$\Rightarrow$  Inertial reference frame

Newton's 2nd law is written w.r.t. an inertial reference system

$$\boxed{\underline{F}(x(t), t) = m_I \ddot{\underline{x}}(t)}$$

$\downarrow$  Force                       $\downarrow$  Inertial mass                      acceleration

W.r.t non-inertial reference frame?

Define three unit vectors  $\hat{x}'(t), \hat{y}'(t), \hat{z}'(t)$

In reference <sup>frame</sup>, with respect to old basis

trajectories  $\rightarrow$

$$\underline{w}(t) = w_x(t) \hat{x} + w_y(t) \hat{y} + w_z(t) \hat{z}$$

coordinates  $\rightarrow$

$$= w'_x(t) \hat{x}'(t) + w'_y(t) \hat{y}'(t) + w'_z(t) \hat{z}'(t)$$

Notation:

$\underline{a} \Leftarrow$  denotes vector

$\underline{\underline{M}} \Leftarrow$  matrix

$$\underline{\underline{M}} \underline{x} = \underline{y}$$

Tensors: Don't use underline notation

Index notation:  $[\underline{a}]_x \equiv x \text{ component}$   
 or  $\equiv a_x$

$$[\underline{\underline{M}}]_{xy} = x-y \text{ component}$$

$$\equiv M_{xy}$$

Physics in non-inertial frames (as captured by 2nd law) gives rise to additional force terms - inertial forces -  
 Eg. centrifugal, Coriolis, Euler, etc..

HW Distinguish "true forces" from such

apparent forces. In particular in Newtonian physics  
Gravity is a true force (not be true in GR)  
general relativity

Mass: (at least) 3 conceptually distinct quantities

- Inertial mass: constant of proportionality in Newton's 2nd law

$$\underline{F} = m_I \underline{a}$$

- "active" gravitational mass: this is the mass that produces the gravitational field that other masses respond to

$$\underline{g}(\underline{x}) \equiv -G \underbrace{m_g^{(a)}}_{\text{active}} \frac{\underline{x} - \underline{x}'}{\|\underline{x} - \underline{x}'\|^3}$$

$\underline{x}' \equiv$  location of generating mass and  $G \rightarrow$  gravitational constant

$$G = 6.67430(15) \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$$

- "passive" gravitational mass: the mass that responds to an external gravitational field by acceleration.

$$\underline{F}_g = \underbrace{m_g^{(p)}}_{\text{passive}} \cdot \underline{g}$$

It is assumed that  $m_I, m_g^{(a)}, m_g^{(p)} \geq 0$

Relationships between these three masses.

Exercise: using Newton's 3rd law argue that

$$\frac{m_g^{(a)}}{m_g^{(p)}} = \text{universal constant}$$

Choosing units:

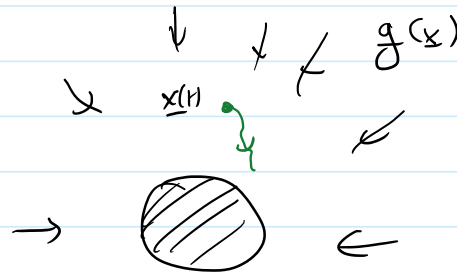
$$m_g^{(a)} = m_g^{(p)} = m_g$$

Relationship between  $m_g$  &  $m_I$

Extensive experimental evidence suggests  $m_g = m_I$

Universality of Free Fall (UFF): In a gravitational field  $g$  the orbit  $\underline{x}(t)$  of a test mass only depends on initial conditions (IC). Same ICs  $\Rightarrow$  same trajectories  $\forall$  test masses

Test mass ("a point")



$$m_g g(\underline{x}(t)) = m_I \ddot{\underline{x}}(t)$$

or

$$\ddot{\underline{x}}(t) = \underbrace{\frac{m_g}{m_I}}_{\downarrow} g(\underline{x}(t))$$

Choose units so  $m_g = m_I$ .  
A universal constant

Remark: (1) UFF is not a consequence of Newtonian theory! It is a falsifiable hypothesis (2) UFF relates to test masses (subtle concept)

From Newtonian gravity to GR

$\rightarrow$  Newton gravity as a field theory

Electrodynamics and SR are compatible  
Newtonian gravity  $\nleftrightarrow$  Electrostatics

Highlight similarities between Newton gravity & Electrostatics

Gravitational field:

$$g : \underset{1.}{\mathbb{R}} \times \underset{1.}{\mathbb{R}^3} \rightarrow \mathbb{R}^3$$

$$g: \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}^3$$

$$\downarrow \quad \downarrow$$

$$(t, x) \rightarrow g(t, x)$$

Sources: (gravitational) mass densities

$$\rho: \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}^+$$

with

$$\nabla \cdot g = -4\pi G \rho$$

Just like Coulomb's law!

We have  $\nabla \times g = 0$ , so  $g$  is conservative

Thus by Poincaré's lemma  $\exists$

$$\phi: \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}$$

s.t

$$g = -\nabla \phi$$

convention!

Just as for electrostatics

$$\Delta \phi = 4\pi G \rho$$

Poisson's equation

Equation of motion for test masses

$$\ddot{x}(t) = -\nabla \phi(t, x(t))$$

Assumed here is that

$$\text{supp}(\rho) = \overline{\{x \in \mathbb{R}^3 \mid \rho(x) \neq 0\}}$$

↙ closure

$$= \text{compact} \quad (\text{Metric space} = \text{closed \& bounded})$$

With boundary conditions

$$\phi(t, \|x\| \rightarrow \infty) \rightarrow 0$$

unique solution

$$\phi(t, \underline{x}) = - \int_{\mathbb{R}^3} d^3 \underline{x}' \frac{\rho(t, \underline{x}')}{\|\underline{x} - \underline{x}'\|}$$

Thus we have an analogue "gravito statics"

How to add dynamics to grav. field?

- (0) Assume instantaneous propagation?
- (1) Assume scalar wave equation for  $\phi$ ?
- (2) Complete  $\phi \rightarrow$  to a 4-vector, (eg. electrodynamics)
- (3) Complete  $\phi \rightarrow$  a tensor (symmetric)

MTW chapter 7

- (0) obvious causality problems X
- (1) No bending of light
- (2) Perichthon precession problems (negative energy waves?)
- (3) This route leads to linearised Einstein field equations  $\Rightarrow$  GR