

Introduction to general relativity: prerelativity gravitation

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Text: Wald, "General Relativity"

Other sources:

- Misner, Thorne, & Wheeler
- Weinberg

Structure: Mathematical methods \Rightarrow Physics

Differential geometry
↑ ↓ ↗

Quantum gravity & string theory.

Videos: Lenny Susskind

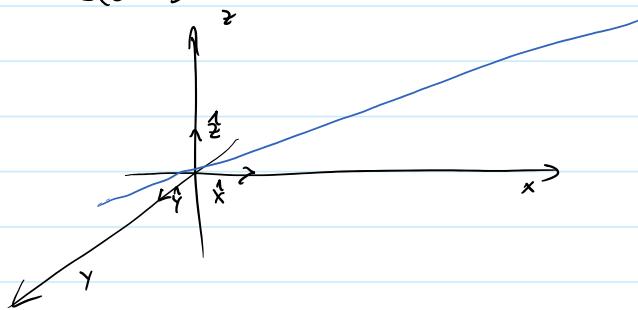
Assume:

- * Special relativity (Theoretische Physik A)
- * Calculus
- * Linear algebra
- * Combinatorics

Lecture 1: prerelativity gravitation

Assumption:

Newtonian physics: existence of a preferred reference system & clocks.



A free test mass (ie. no force acts) moves along a Straight trajectory at a constant rate uniformly in time.

\Rightarrow Inertial reference frame

Newton's 2nd law is written w.r.t. an inertial reference system

$$\underline{F}(\underline{x}(t), t) = m_I \ddot{\underline{x}}(t)$$

↓ ↓ ↘
 Force Inertial mass acceleration

W.r.t non-inertial reference frame ?

Define three unit vectors $\hat{x}'(t)$, $\hat{y}'(t)$, $\hat{z}'(t)$

In reference^{frame}, with respect to old basis

$$\underline{w}(t) = w_x(t) \hat{x} + w_y(t) \hat{y} + w_z(t) \hat{z}$$

coordinates

trajectory

$$= w'_x \hat{x}'(t) + w'_y \hat{y}'(t) + w'_z \hat{z}'(t)$$

Notation

a \in denotes vector

M \Leftarrow matrix

$$\underline{M} \times = \underline{x}$$

Tensors : won't use underline notation

Index notation :

$$\text{or } [\underline{a}]_x = a_x \quad \text{=} \quad x \text{ component}$$

$$[M]_{xy} = x-y \text{ Component}$$

$$= M_{xy}$$

Newton's

Physics in non-inertial frames (as captured by 2nd law)
gives rise to additional force terms - inertial forces -
E.g. centrifugal, coriolis, euler, etc..

H.W. Distinguish "true forces" from such

apparent forces". In particular in Newtonian physics Gravity is a free force (not be true in GR)
general relativity

Mass: (at least) 3 conceptually distinct quantities

- Inertial mass: constant of proportionality in Newton's 2nd law

$$\boxed{F = m_I a}$$

- "active" gravitational mass: this is the mass that produces the gravitational field that other masses respond to

$$g(x) \equiv -G \frac{m^{(a)} g}{\|x - x'\|^3}$$

x' = location of generating mass and G → gravitational constant

$$\boxed{G = 6.67430 (15) \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}}$$

- "passive" gravitational mass: the mass that responds to an external gravitational field by accelerating.

$$F_g = m^{(p)} g$$

It is assumed that $m_I, m^{(a)}, m^{(p)} \geq 0$

Relationships between these three masses.

Exercise: using Newton's 3rd law argue
that

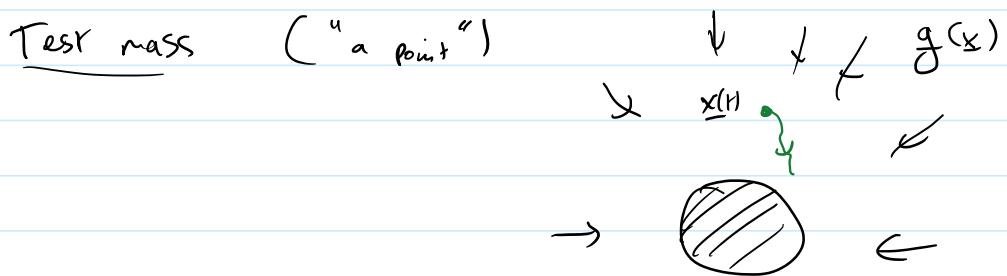
$$\frac{m^{(a)}}{m^{(p)}} = \text{universal constant}$$

Choosing units: $m^{(a)} = m^{(p)} = m_g$

Relationship between m_g & m_I

Extensive experimental evidence suggests $m g = m \ddot{x}$

Universality of Free Fall (UFF): In a gravitational field g the orbit $\underline{x}(t)$ of a test mass only depends on initial conditions (IC). Same ICs \Rightarrow same trajectories & test masses.



$$m g g(\underline{x}(t)) = m \ddot{x}(t)$$

or

$$\ddot{x}(t) = \frac{m g}{m_I} g(\underline{x}(t))$$

\downarrow A universal constant

Choose units so $m g = m_I$.

Remark: (1) UFF is not a consequence of Newtonian theory! It is a falsifiable hypothesis (2) UFF relates to test masses (subtle concept)

From Newtonian gravity to GR

→ Newton gravity as a field theory

Electrodynamics and SR are compatible
Newtonian gravity \leftrightarrow Electrostatics

Highlight similarities between Newton gravity & Electrodynamics

Gravitational field:

$$g : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}^3$$

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$$\downarrow \quad \downarrow$$

$$(t, \mathbf{x}) \rightarrow g(t, \mathbf{x})$$

Sources: (gravitational) mass density

$$j : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}^+$$

with

$$\boxed{\nabla \cdot g = -4\pi G \rho}$$

Just like Coulomb's law!

We have $\nabla \times g = 0$, so g is conservative

Thus by Poincaré's lemma \exists

$$\phi : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}$$

s.t.

$$\boxed{g = -\nabla \phi} \quad \text{convention!}$$

Just as for electrostatics

$$\boxed{\Delta \phi = 4\pi G \rho}$$

Poisson's equation

Equation of motion for test masses

$$\boxed{\ddot{\mathbf{x}}(t) = -\nabla \phi(t, \mathbf{x}(t))}$$

\overline{F} Assumed here is that

$$\begin{aligned} \text{supp}(\rho) &= \overline{\{ \mathbf{x} \in \mathbb{R}^3 \mid \rho(\mathbf{x}) \neq 0 \}} \\ &= \text{compact} \quad (\text{Hectic borel} = \text{closed \& bounded}) \end{aligned}$$

↙ closure

With boundary conditions

$$\phi(t, \|\mathbf{x}\| \rightarrow \infty) \rightarrow 0$$

unique solution

$$\phi(t, \underline{x}) = -\zeta \int_{\mathbb{R}^3} d^3x' \frac{\rho(t, \underline{x}')}{\|\underline{x} - \underline{x}'\|}$$

Thus we have an analogue "gravito statics"

How to add dynamics to grav. field?

- (i) Assume instantaneous propagation?
- (ii) Assume scalar wave equation for ϕ ?
- (iii) Complete $\phi \rightarrow$ to a 4-vector, (e.g. electrodynamics)
- (iv) Complete $\phi \rightarrow$ a tensor (symmetric)

MTW chapter 7

- (i) obvious causality problems \times
- (ii) No bending of light
- (iii) Perihelion precession problems (negative energy waves?)
- (iv) This route leads to linearized Einstein field equations \Rightarrow GR