

# Introduction to general relativity: homogeneity and isotropy cont.

14 June 2021 09:33

Dynamics of homogeneous & isotropic universe  
Idealise galaxies  $\rightarrow$  "grains of dust"; assume galaxies move on worldlines of isotropic observers. Fluctuations of galaxy velocities is small  $\rightarrow$  pressure is neglected. Thus stress-energy-tensor of galaxy content

$$T_{ab} = \rho u_a u_b$$

$\rho$  is average density of matter.

To model background thermal radiation (@ 3K) we use perfect-fluid stress-energy tensor

$$T_{ab} = P_{\text{rad}} (g_{ab} + u_a u_b) + \rho_{\text{rad}} u_a u_b$$

where  $P_{\text{rad}} = \rho_{\text{rad}}/3$ . Thus most general stress-energy tensor consistent with homogeneity & isotropy is (ex.)

$$T_{ab} = \rho u_a u_b + P (g_{ab} + u_a u_b)$$

We aim to solve Einstein's equation

(\*)

$$G_{ab} + \Lambda g_{ab} = 8\pi T_{ab}$$

where we have added  $\Lambda$ , the cosmological constant, which is necessary to account for cosmic background radiation and dark energy surveys.

In principle we need to solve 10 unknowns & 10 equations

$$\begin{pmatrix} \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix}$$

Because of the symmetries of homogeneity & isotropy only two equations are independent: Firstly note, the  $G^{ab} u_b$  &  $T^{ab} u_b$  cannot have spatial components (isotropy). Similarly, restricted to  $\Sigma_t$ , the same arguments we made for  ${}^{(3)}R^{ab\,cd}$  implies  $G_{ab}^{(3)}$  is proportional to identity. Off-diagonal terms of  $(k)$  vanish. And diagonal space-space components are equal. Due to homogeneity implies space-time components vanish. Thus independent components of  $(k)$  are

$$\begin{aligned} G_{\tau\tau} + \Lambda g_{\tau\tau} &= 8\pi\rho \\ G_{\mu\mu} + \Lambda g_{\mu\mu} &= 8\pi p \end{aligned}$$

Here  $G_{\tau\tau} \equiv G_{ab} u^a u^b$  and  $G_{\mu\mu} = G_{ab} s^a s^b$  with  $s^a u_a = 0$ ;  $s^a s_a = 1$ .

To carry out subsequent calculations we rewrite the FLRW metric as (ex)

$$ds^2 = -d\tau^2 + a^2(\tau) \left[ \frac{dr^2}{1-kr^2} + r^2 (d\theta^2 + \sin^2\theta d\phi^2) \right]$$

where  $k = +1 \equiv$  the sphere;  $k = 0$ : flat space;  $k = -1 \equiv$  hyperboloid.