

Physics 3112: Experimental & Computational Physics T1 2025

Never Stand Still

Science

School of Physics

Computational Physics: Simulation & Modeling

LECTURER (Weeks 5-9)

- Prof. Rajib Rahman (course convenor)
- Email: rajib.rahman@unsw.edu.au
- Course Resources: Moodle
- ➤ Administrative matters: Zofia Krawczyk

(z.krawczyk-bernotas@unsw.edu.au)



Objective (Weeks 5-9)

- Solving differential equations numerically
- Eigenvalue problems in physics
- Numerical linear algebra and solving large scale problems
- Monte Carlo simulations



Differential Equations in Physics

Electrostatics: Poisson's Equation (Charge to Potential)

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

$$V(r), \rho(r), \qquad r = x, y, z$$

$$\nabla^2 = \frac{d^2}{dx^2} + \frac{d^2}{dy^2} + \frac{d^2}{dz^2}$$
 (Laplacian)

Laplace Equation

$$\nabla^2 V = 0$$



Differential Equations in Physics

Thermal Physics: Heat Equation (2D)

$$\frac{\partial T}{\partial t} = \alpha \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + q$$

T: Temperature

q: Heat generation rate

 α : Thermal Expansion Coefficient

Q, Heat Flow
A, Area
I, Length

Statistical Mechanics: Motion of charge carriers

https://www.digitalengineering247.com/article/use-feathermal-analysis/



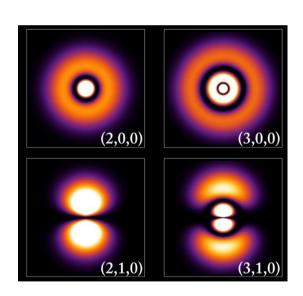
Hot End

Quantum Mechanics: Time independent Schrodinger Equation (Wavefunction)

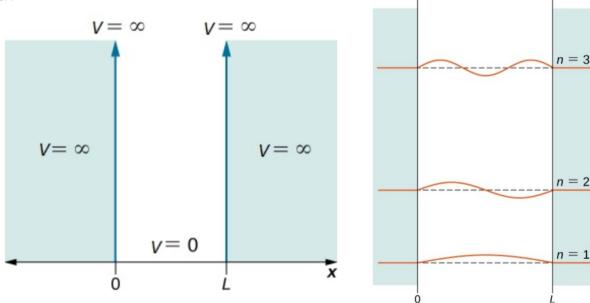
$$-\frac{\hbar^2}{2m}\nabla^2\psi + V\psi = E\psi \qquad V(r), \psi(r), \qquad r = x, y, z$$

Particle in a box wavefunction

Atoms: Solution with 1/r potential



https://en.wikipedia.org/wiki/Atomic_orbital#/media/File:Hydrogen_Density_Plots.png



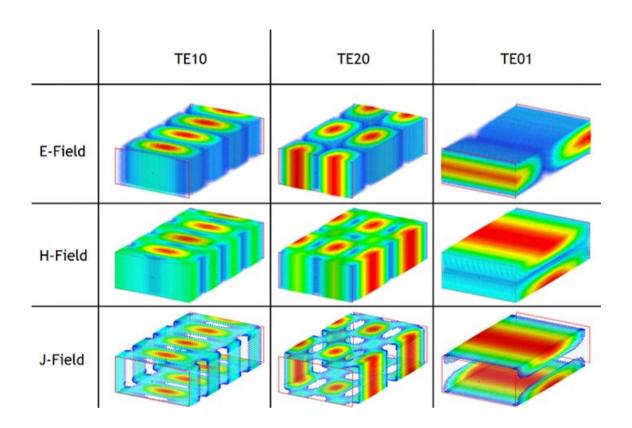
https://chem.libretexts.org/Bookshelves/Physical and Theoretical Chemistry Textbook Maps/Maps/3A Physical Chemistry (McQuarrie and Simon)/03%3A The Schr%C3%B6dinger Equation and a Particle in a Box/3.05%3A The Energy of a Particle in a Box is Quantized



Electro-magnetics: Maxwell's Equation (Electric & Magnetic Fields) Wave Guides

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{J}$$



https://www.technobyte.org/wp-content/uploads/2016/11/Rect-Hollow-WG-Modes.jpg

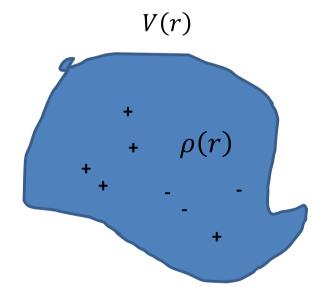


Differential Equations in Physics

Electrostatics: Poisson's Equation (Charge to Potential)

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

$$abla^2 V = -\frac{\rho}{\epsilon}$$



$$V(r), \rho(r), \qquad r = x, y, z$$

$$\nabla^2 = \frac{d^2}{dx^2} + \frac{d^2}{dy^2} + \frac{d^2}{dz^2}$$
 (Laplacian)

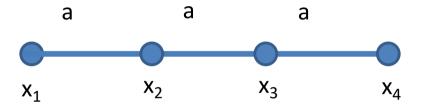


Numerical Solution of Differential Equations: The Finite Difference Method

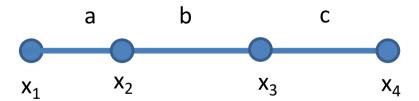
Example: 1D Poisson's Equation with FDM

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

Step 1: Mesh



Homogeneous mesh: Equal grid spacing



Inhomogeneous mesh: Unequal grid spacing

Needed for systems – when charge is very dense at one point and coarse in another.



Numerical Solution of Differential Equations: The Finite Difference Method

Example: 1D Poisson's Equation with FDM

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

Output: Potential V

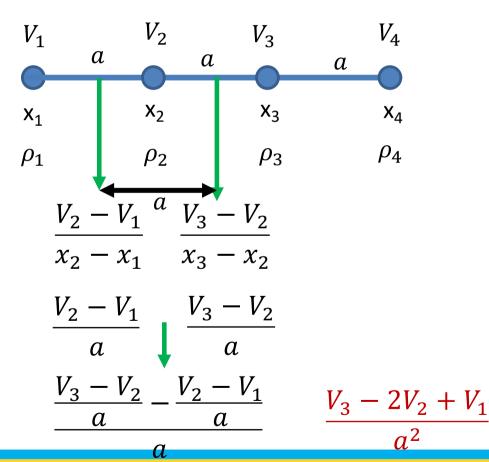
Step 2: Define operators

& Input / Output Variables

Input: Charge density
(Charge per unit length in 1D)

$$\nabla V = \frac{\Delta V}{\Delta x} =$$

$$\nabla^2 = \frac{d^2}{dx^2} \quad (at x_2)$$





Numerical Solution of Differential Equations: The Finite Difference Method

Example: 1D Poisson's Equation with FDM

 χ_1

 χ_2

 χ_3

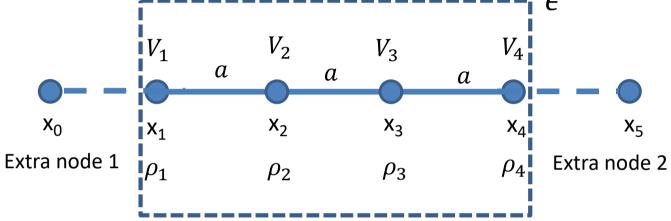
 χ_4

$$\nabla^2 V = -\frac{\rho}{2}$$

Step 3: Difference Equations

Attach two extra nodes (Boundary conditions)

$$\nabla^2 = \frac{d^2}{dx^2} = \frac{V_3 - 2V_2 + V_1}{a^2}$$



Simulation Domain

$$V_{2} - 2V_{1} + V_{0} = -\rho_{1}a^{2}/\epsilon$$

$$V_{3} - 2V_{2} + V_{1} = -\rho_{2}a^{2}/\epsilon$$

$$V_{4} - 2V_{3} + V_{2} = -\rho_{3}a^{2}/\epsilon$$

$$V_{5} - 2V_{4} + V_{3} = -\rho_{4}a^{2}/\epsilon$$

Boundary Conditions V=0 at the edges



Example: 1D Poisson's Equation with FDM

$$x_{1} V_{2} - 2V_{1} + V_{0} = -\rho_{1}\alpha^{2}/\epsilon$$

$$X_{2} V_{3} - 2V_{2} + V_{1} = -\rho_{2}\alpha^{2}/\epsilon$$

$$x_{3} V_{4} - 2V_{3} + V_{2} = -\rho_{3}\alpha^{2}/\epsilon$$

$$x_{4} V_{5} - 2V_{4} + V_{3} = -\rho_{4}\alpha^{2}/\epsilon$$

Boundary Conditions V=0 at the edges

Step 4: Matrix Representation

- Tri-diagonal matrix: 3
 diagonals, main + 2 off-diag.
- Applies to 1D case

$$AX = B$$

Linear (Matrix) Equation

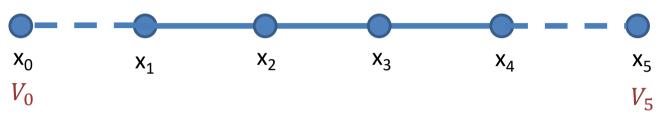
Solve for X (V in our case)



Numerical Solution of Differential Equations: The Finite Difference Method

Step 5: Boundary conditions (2 common types)

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$



1. Dirichlet Boundary condition (BC): Set Potential/Voltage to constant values at the edges.

$$x_1$$
 $V_2 - 2V_1 + V_0 = -\rho_1 a^2/\epsilon$ $V_2 - 2V_1 = -\rho_1 a^2/\epsilon - V_0$
 x_4 $V_5 - 2V_4 + V_3 = -\rho_4 a^2/\epsilon$ $-2V_4 + V_3 = -\rho_4 a^2/\epsilon - V_5$

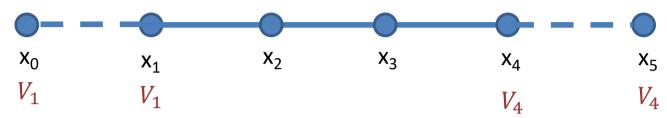
$$\begin{bmatrix} -2 & 1 & 0 & 0 \\ 1 & -2 & 1 & 0 \\ 0 & 1 & -2 & 1 \\ 0 & 0 & 1 & -2 \end{bmatrix} \begin{bmatrix} V_1 \\ V_2 \\ V_3 \\ V_4 \end{bmatrix} = \frac{a^2}{\epsilon} \begin{bmatrix} \rho_1 \\ \rho_2 \\ \rho_3 \\ \rho_4 \end{bmatrix} - \begin{bmatrix} V_0 \\ 0 \\ 0 \\ V_5 \end{bmatrix}$$



Numerical Solution of Differential Equations: The Finite Difference Method

Step 5: Boundary conditions (2 common types)

$$\nabla^2 V = -\frac{\rho}{\epsilon}$$



2. Neumann Boundary condition (BC): Set gradient of potential to constant at the edges.

$$E = -\nabla V$$
 Constant Example: Set E=0 at the edges

$$x_1$$
 $V_2 - 2V_1 + V_1 = -\rho_2 a^2/\epsilon$ $V_2 - V_1 = -\rho_1 a^2/\epsilon$ $v_4 - 2V_4 + V_3 = -\rho_4 a^2/\epsilon$ $v_4 - 2V_4 + V_3 = -\rho_4 a^2/\epsilon$

$$\begin{bmatrix} -1 & 1 & 0 & 0 \\ 1 & -2 & 1 & 0 \\ 0 & 1 & -2 & 1 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} V_1 \\ V_2 \\ V_3 \\ V_4 \end{bmatrix} = \frac{a^2}{\epsilon} \begin{bmatrix} \rho_1 \\ \rho_2 \\ \rho_3 \\ \rho_4 \end{bmatrix}$$



The Finite Difference Method: Error & Convergence

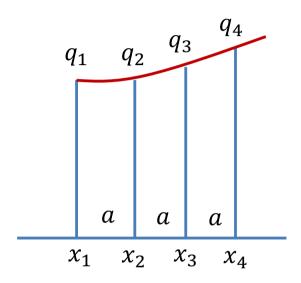
- Numerical solution is approximate.
- It approaches the exact solution in the limit grid spacing 'a' $\rightarrow 0$.
- Analytically unsolvable problems can be solved.

Numerical tolerance of solution

- How does the solution depend on the domain size (finite-size effects)?
- How does the solution depend on the grid spacing (meshing)?
- Order of accuracy in the definition of operators?
- Choice between accuracy and computational resources (memory and speed)?
- Define an acceptable tolerance level.
- Test the effect of various numerical parameters.
- Design simple cases to test against closed form analytic solutions.
- Build up confidence in your numerical solver.

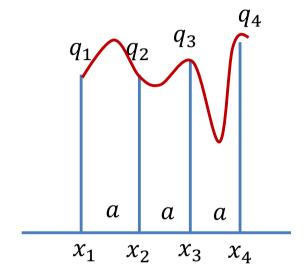


Mesh related errors



Numerical solution is good in this case.

- Curve is smooth between grid points.
- Approximately linear between grid points.

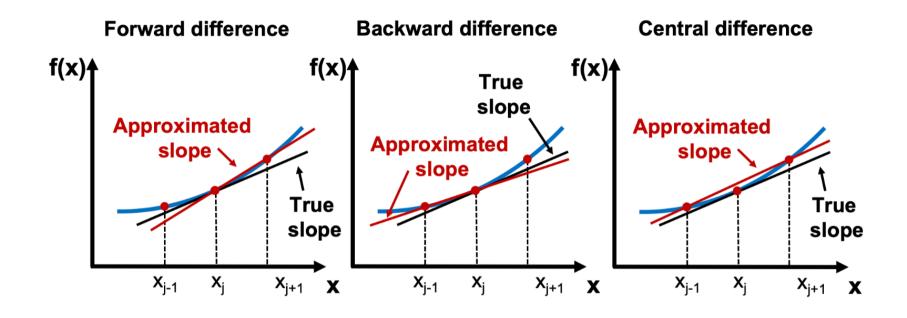


Not a good solution.

- Large variation on the scale of 'a'.
- Solution: Decrease 'a', Adaptive meshing.



Definition of operators



$$f'(x_j) = \frac{f(x_{j+1}) - f(x_j)}{x_{j+1} - x_j} \qquad f'(x_j) = \frac{f(x_j) - f(x_{j-1})}{x_j - x_{j-1}} \qquad f'(x_j) = \frac{f(x_{j+1}) - f(x_{j-1})}{x_{j+1} - x_{j-1}}$$

Source: https://pythonnumericalmethods.berkeley.edu/



Errors in approximating the first derivative (example)

Taylor series expansion of a function about x_i

$$f(x) = \frac{f(x_j)(x-x_j)^0}{0!} + \frac{f'(x_j)(x-x_j)^1}{1!} + \frac{f''(x_j)(x-x_j)^2}{2!} + \frac{f'''(x_j)(x-x_j)^3}{3!} + \cdots$$

Compute the derivative by forward difference (i.e. between x_j and x_{j+1}

$$f(x_{j+1}) = \frac{f(x_j)(x_{j+1}-x_j)^0}{0!} + \frac{f'(x_j)(x_{j+1}-x_j)^1}{1!} + \frac{f''(x_j)(x_{j+1}-x_j)^2}{2!} + \frac{f'''(x_j)(x_{j+1}-x_j)^3}{3!} + \cdots.$$

Grid spacing $h = x_{j+1} - x_j$

$$f'(x_j) = \frac{f(x_{j+1}) - f(x_j)}{h} + \left(-\frac{f''(x_j)h}{2!} - \frac{f'''(x_j)h^2}{3!} - \cdots\right).$$

$$f'(x_j) = \frac{f(x_{j+1}) - f(x_j)}{h} + O(h).$$



$O(h^p)$ p is the order of accuracy of the approximation

What is the order of accuracy of the central difference?

$$f(x) = \frac{f(x_j)(x - x_j)^0}{0!} + \frac{f'(x_j)(x - x_j)^1}{1!} + \frac{f''(x_j)(x - x_j)^2}{2!} + \frac{f'''(x_j)(x - x_j)^3}{3!} + \cdots$$

$$f(x_{j+1}) = f(x_j) + f'(x_j)h + \frac{1}{2}f''(x_j)h^2 + \frac{1}{6}f'''(x_j)h^3 + \cdots$$

$$f(x_{j-1}) = f(x_j) - f'(x_j)h + \frac{1}{2}f''(x_j)h^2 - \frac{1}{6}f'''(x_j)h^3 + \cdots$$

Subtract to find the derivative $f'(x_i)$

$$f'(x_j) \approx \frac{f(x_{j+1}) - f(x_{j-1})}{2h}.$$

What is the order of accuracy?



Summary: Finite Difference Method

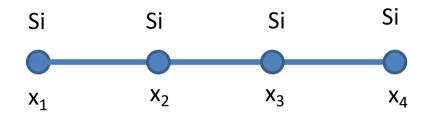
- Meshing
- Defining input / output variables
- Defining operators
- Representing as a matrix equation
- Boundary conditions
- Sources of errors
- Testing for convergence



Finite Difference Method: Generalization (Two Different Media)

Homogneous Medium: Si

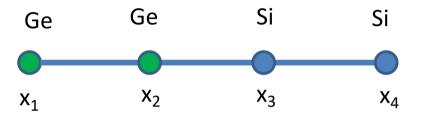
Dielectric constant: ϵ_{Si}



$$\nabla^2 V = -\frac{\rho}{\epsilon_{Si}}$$

Inhomogneous Medium: Ge, Si

Dielectric constant: ϵ_{Ge} , ϵ_{Si}



$$\nabla(\epsilon\nabla V) = -\rho$$

General form: $\epsilon(x)$



Finite Difference Method: Generalization (Two Different Media/Material)

Inhomogneous Medium: Ge, Si

Dielectric constant: ϵ_{Ge} , ϵ_{Si}

$$\nabla(\epsilon\nabla V) = -\rho$$

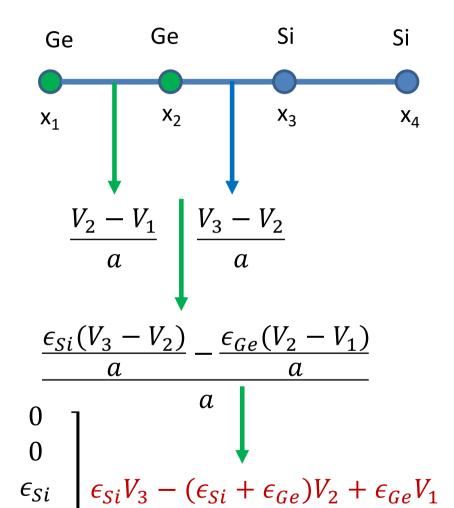
General form: $\epsilon(x)$

 (∇V)

Setting Dirichlet BC of V=0 at the boundaries

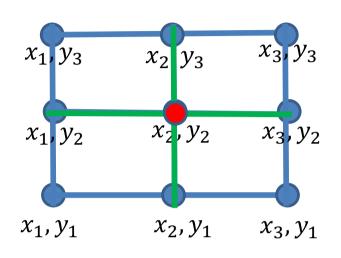
$$\nabla(\epsilon\nabla V)$$

$$\begin{bmatrix} -2\epsilon_{Ge} & \epsilon_{Ge} & 0 & 0 \\ \epsilon_{Ge} & -(\epsilon_{Ge} + \epsilon_{Si}) & \epsilon_{Si} & 0 \\ 0 & \epsilon_{Ge} & -(\epsilon_{Ge} + \epsilon_{Si}) & \epsilon_{Si} \\ 0 & 0 & \epsilon_{Si} & -2\epsilon_{Si} \end{bmatrix} \underbrace{\frac{\alpha}{\epsilon_{Si}V_3 - (\epsilon_{Si} + \epsilon_{Ge})V_2 + \epsilon_{Ge}V_1}_{\alpha^2}}_{a^2}$$





Finite Difference Method: Generalization (2D Problems)



$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

$$\nabla^2 = \frac{d^2}{dx^2} + \frac{d^2}{dy^2}$$

$$\frac{d^2}{dx^2}$$

$$x_2, y_2$$

$$V_{x_3,y_2} - 2V_{x_2,y_2} + V_{x_1,y_2}$$

Horizontally

$$\frac{d^2}{dy^2}$$

$$x_2, y_2$$

$$V_{x_2,y_3} - 2V_{x_2,y_2} + V_{x_2,y_1}$$

Vertically

$$\nabla^2$$

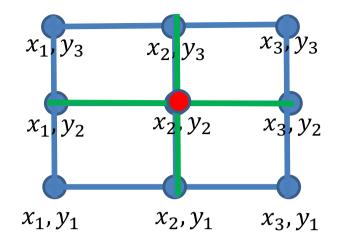
$$x_2, y_2$$

$$V_{x_3,y_2} + V_{x_2,y_3} - 4V_{x_2,y_2} + V_{x_1,y_2} + V_{x_2,y_1}$$

Diagonal elements (i,i) =-4, Off-diagonal elements 4 1's => 4 connections to nearest grid points



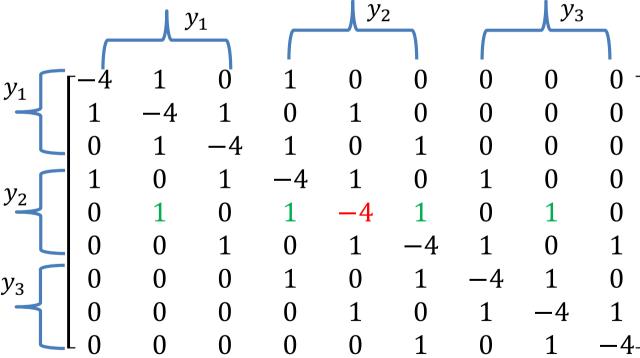
Finite Difference Method: Generalization (2D Problems)



$$\nabla^2 V = -\frac{\rho}{\epsilon}$$

Assuming BC: V=0 at the edges

$$\nabla^{2} \qquad x_{2}, y_{2}
V_{x_{3},y_{2}} + V_{x_{2},y_{3}}
- 4V_{x_{2},y_{2}} + V_{x_{1},y_{2}}
+ V_{x_{2},y_{1}}$$



Each 3x3 block is each horizontal row displaced in y.

Total Matrix: 3x3
(3 points in x & 3 points in y)



Generalization: Linear vs Non-linear Differential Equation

- In semiconductors, the left and right hand side of the Poisson equation are not independent.
- The charge density is a function of potential.
- Non-linear Poisson Equation

$$\nabla^2 V = -\frac{\rho(V)}{\epsilon}$$

Example:

$$\rho(V) \sim q N_C e^{(E_F - E_C + qV)/kT}$$

- Need to solve this iteratively with some initial guess for V.
- One method is Newton-Raphson.
- Define f(V) and find the root.

$$f(V) = \nabla^2 V + \frac{\rho(V)}{\epsilon}$$

Newton-Raphson technique of root finding

$$x_{n+1}=x_n-rac{f(x_n)}{f'(x_n)}$$



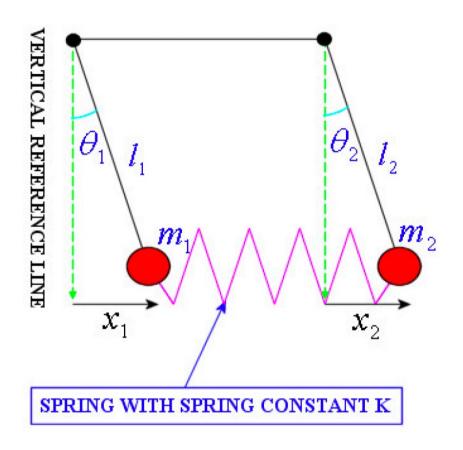
Generalization: Coupled Differential Equations

Example from your lab: Coupled Pendula

2 PDEs, solution of one depend on the other and vice versa.

$$\ddot{\phi}_1(t) + A^2 \phi_1(t) = -B^2 [\phi_2(t) - \phi_1(t)],$$

$$\ddot{\phi}_2(t) + A^2 \phi_2(t) = B^2 [\phi_2(t) - \phi_1(t)].$$

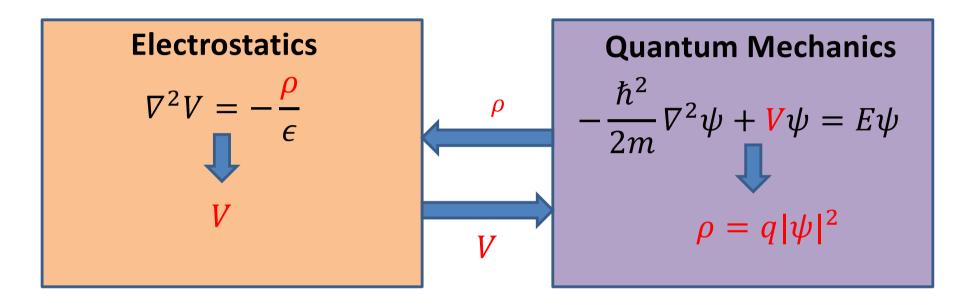


Where A and B are constants; $\phi_1(t)$ and $\phi_2(t)$ are functions of time.



Generalization: Coupled Differential Equations

Example: Schrodinger-Poisson



Self-consistent Solutions (Iterative)

Convergence: Both ρ and V are constant over iterations.



Summary: Generalizations

- Equations over heterogeneous domain (Poisson Eq. over two materials)
- 2 dimensions and more
- Non-linear differential equation (Non-linear Poisson)
- Coupled differential equations (Schrodinger-Poisson)

