

Communication Beyond Information: Manifold Expansion via High-Dimensional Coupling

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Abstract Many information-theoretic bounds assume a fixed statistical model class. We identify a regime where this assumption fails: when high-dimensional systems couple, the identifiable parameter set changes because the image rank of the dynamics-to-distribution map changes. We formalize this as *manifold expansion*: coupling-induced increase in Fisher rank. Our main theorem provides checkable criteria for rank transitions: the *transversality criterion* (coupling moves the accessible family off a constraint submanifold) and the *symmetry-breaking criterion* (coupling breaks a group invariance). We prove Fisher rank equals Jacobian rank, so transitions are coordinate-invariant. Two examples demonstrate superadditive complexity: coupled Ornstein–Uhlenbeck processes (transversality, $\kappa_c = 0$) and Kuramoto oscillators (symmetry breaking, $\kappa_c > 0$). The framework predicts Fisher eigenvalues emerging from zero as coupling crosses critical thresholds.

Keywords Fisher information metric · Parameter identifiability · Superadditivity · Coupled dynamical systems · Statistical manifolds

1 Introduction

Information theory provides powerful bounds on communication: channel capacity limits transmission rates, mutual information bounds learning, and data processing inequalities constrain inference (Cover and Thomas 2006). Many commonly used bounds—PAC generalization, channel coding theorems, Cramér–Rao bounds—assume a fixed model class, channel law, or hypothesis space.

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This assumption is usually invisible. When two agents exchange messages, we take for granted that the receiver’s model class is fixed: decoding maps incoming symbols to a pre-existing internal space. Complexity growth is then bounded by how much information the channel can carry.

But what happens when the communicating systems are themselves high-dimensional dynamical objects? When a protocell network coordinates, or when coupled oscillators synchronize, or when interacting agents learn together—the systems don’t merely exchange symbols. They *reshape each other’s dynamics*.

This paper identifies a regime where the fixed-manifold assumption need not hold, so that bounds derived for fixed model classes may underpredict complexity growth:

Core claim: When high-dimensional coherent systems couple, they exchange constraints rather than tokens. This can create new collective coordinates, increasing effective dimensionality beyond the sum of parts.

We call this *manifold expansion*: the accessible statistical manifold grows under coupling. This is not a violation of information theory—it is a regime where information theory’s foundational assumption (fixed model class) need not hold.

The implications are significant:

- Model complexity (identifiable dimensions) can grow faster than fixed-class bounds predict—because coupling activates new coordinates, not because information bounds are violated
- Communication becomes manifold deformation, not message passing
- The mechanism is scale-free: it applies wherever high-D systems interact

1.1 The fixed-manifold assumption in practice

To be precise about what we are *not* claiming: we do not argue that data-processing inequalities are violated, or that channel capacity theorems fail. These results hold within their stated assumptions.

What we identify is a regime where the *model class itself* changes under coupling, so that bounds derived for a fixed class can underpredict emergent complexity. The fixed-manifold assumption appears in several standard settings:

1. **Generalization bounds:** PAC-learning and VC-dimension arguments assume a fixed hypothesis class \mathcal{H} . Sample complexity bounds scale with $\dim(\mathcal{H})$ (Cover and Thomas 2006).
2. **Channel capacity:** Shannon’s theorem assumes a fixed channel law $p(y|x)$. Capacity is computed over this fixed conditional distribution.
3. **Data processing inequalities:** DPI holds for fixed processing maps. If the map itself changes under interaction, DPI applies to each fixed map but not to the trajectory across maps.

4. **Fisher information bounds:** Cramér-Rao and related bounds assume a fixed parametric family. If coupling changes which parameters are identifiable, the bounds shift.

Our claim: *coupling between high-dimensional systems can change the identifiable parameter set, so that applying bounds derived for a fixed family may underpredict the achievable complexity.*

To be explicit: we are tracking changes in the *image dimension* of the parameter-to-distribution map, not claiming violations of information inequalities. The inequalities hold; what changes is the dimension of the family to which they apply.

1.2 Coherence vs. information: the key distinction

We distinguish two conceptual modes of inter-system interaction that motivate our geometric framework:

Information transfer: System A transmits discrete, addressable, copyable tokens to system B . The receiver decodes tokens into a pre-existing internal representation. Complexity growth is bounded by channel capacity. This is the regime where standard information-theoretic bounds apply.

Constraint exchange: System A couples to system B via shared dynamical modes. The coupling reshapes both systems' accessible state spaces. Complexity growth depends on geometric properties of the coupling, not channel capacity. This is the regime we study.

High-dimensional coherent systems (our focus) are those with: (i) many coupled internal degrees of freedom, (ii) long-lived collective modes or metastable states, and (iii) dynamics that maintain internal correlations over timescales longer than the coupling timescale. Examples include oscillator networks, neural populations, and coupled reaction systems near bifurcations.

The mathematical content of this paper concerns the constraint-exchange regime. We prove that coupling can increase the Fisher rank of an accessible family—a phenomenon invisible to channel-capacity arguments that assume a fixed model class.

1.3 Paper outline

This paper provides:

1. **Section 2:** Background on information geometry and the fixed-manifold assumption
2. **Section 3:** Formal setup—accessible families, collective coordinates, and the main theorem on coupling-induced rank transitions (Theorem 1)
3. **Section 4:** Consequences for complexity growth, including the complexity acceleration conjecture (Conjecture 1)
4. **Section 5:** Discussion, testable predictions, and scope

2 Background: Information geometry

Information geometry studies the differential geometry of statistical manifolds (Amari 2016; Ay et al. 2017; Nielsen 2020). A statistical manifold \mathcal{M} is a space of probability distributions, equipped with the Fisher information metric:

$$g_{ij}(\theta) = \mathbb{E} \left[\frac{\partial \log p(x|\theta)}{\partial \theta_i} \frac{\partial \log p(x|\theta)}{\partial \theta_j} \right] \quad (1)$$

The Fisher metric induces a Riemannian structure on \mathcal{M} . A key feature of information geometry is the existence of *dual affine connections* ($\nabla^{(e)}$ and $\nabla^{(m)}$) corresponding to exponential and mixture geodesics. In dually flat manifolds (including exponential families), KL divergence is a canonical divergence whose second-order expansion recovers the Fisher metric, and information projections satisfy a Pythagorean theorem (Amari 2016).

Communication and learning are naturally described as motion on this manifold:

- Parameter estimation moves along geodesics (e- or m-geodesics depending on the estimation criterion)
- KL divergence locally approximates squared Fisher–Rao distance to second order
- Natural gradient descent follows the manifold’s intrinsic geometry
- Information projections onto submanifolds minimize divergence

The hidden assumption: The manifold \mathcal{M} is fixed. The parameterization θ may change, but the model class—the set of distributions considered—does not.

This assumption is reasonable when:

- Systems have fixed, known structure
- Communication channels are well-defined
- Complexity is measured by localization on a fixed space

It can be inapplicable when:

- Interacting systems are high-dimensional and adaptive
- Coupling creates new collective modes
- The model class itself is shaped by interaction

2.1 Relation to hierarchical information geometry

Amari (2001)’s hierarchical decomposition of probability distributions provides the natural framework for our analysis. For a joint distribution $p(x_1, x_2)$, the log-linear decomposition separates marginal and interaction terms:

$$\log p(x_1, x_2) = \theta_1(x_1) + \theta_2(x_2) + \eta(x_1, x_2) + \psi \quad (2)$$

where θ_1, θ_2 are marginal parameters and η captures interactions. The *independence submanifold* \mathcal{M}_{ind} is the e-flat submanifold where $\eta = 0$.

Amari (2001) showed that interaction coordinates η are always *structurally present* in the ambient model—they parameterize directions transverse to independence. Our contribution is orthogonal: we study when these coordinates become *dynamically accessible*.

The distinction is crucial. In standard information geometry, one chooses a model class \mathcal{M} and studies its geometry. The interaction coordinates either belong to \mathcal{M} or they don’t—this is a modeling choice. We consider a different situation: the dynamics determine which distributions are reachable, and coupling can change the reachable set. The “manifold expansion” we study is not a change in the chosen model, but a change in what the dynamics can produce.

This connects to recent work on sufficient statistics and statistical morphisms (Ay et al. 2015). The map from physical parameters to observable distributions is a statistical transformation; our rank-change theorems characterize when this transformation becomes more “expressive” under coupling.

Relation to singular models. Our focus on Fisher rank degeneracy connects to singular learning theory, where the Fisher information matrix is degenerate at certain parameter values (Watanabe 2009). In that literature, singularities are typically fixed features of a chosen model class. Here, the novelty is that degeneracy can be *lifted by coupling*: parameter directions that are unidentifiable (in $\ker(dF_0)$) at zero coupling become identifiable (leave the kernel) at positive coupling. This “identifiability activation” is a dynamic phenomenon, not a property of a fixed model.

3 Manifold expansion under coupling

3.1 Setup: three distinct objects

Consider two dynamical systems with state spaces $X_1 \subset \mathbb{R}^{n_1}$ and $X_2 \subset \mathbb{R}^{n_2}$. We distinguish three objects that are often conflated:

1. **Ambient manifold** $\mathcal{P}(X_1 \times X_2)$: The space of all probability distributions over the joint state space. This is fixed.
2. **Chosen model class** $\mathcal{M} \subset \mathcal{P}$: A parametric family the analyst uses for inference (e.g., “all bivariate Gaussians”). Standard bounds assume this is fixed.
3. **Accessible family** $\mathcal{M}_{\text{acc}}(\kappa)$: The distributions that the *dynamics* can actually produce at coupling strength κ . This may change with κ .

The independence submanifold: A key example of a submodel is $\mathcal{M}_{\text{ind}} = \{p_1 \otimes p_2 : p_i \in \mathcal{M}_i\}$ —the product distributions with zero interaction. In information geometry terminology, this is an *e-flat* submanifold corresponding to zero interaction coordinates in the hierarchical log-linear decomposition (Amari 2016).

The fixed-model assumption: Many learning and communication bounds assume the accessible family is contained in a fixed model class. We do *not*

claim this assumption is wrong in general—only that it need not hold in the regime we study.

Formal setup: Let $\phi_\kappa : X_1 \times X_2 \rightarrow X_1 \times X_2$ denote coupled dynamics with coupling strength $\kappa \geq 0$, and let $h : X_1 \times X_2 \rightarrow \mathbb{R}^m$ be an observation map. Let \mathcal{B} be a smooth d -dimensional manifold of *physical parameters* β (e.g., damping rates, noise intensities, frequencies).

Key convention. The coupling strength κ is an *index* selecting a dynamical system, not a coordinate in \mathcal{B} . This prevents trivial rank increases from simply adding κ as a new parameter. Rank changes must arise from how coupling restructures the parameter→distribution map, not from expanding the parameter space.

For each fixed κ , we have a map:

$$F_\kappa : \mathcal{B} \rightarrow \mathcal{P}(\mathbb{R}^m), \quad \beta \mapsto Q_{(\kappa, \beta)} = (h)_\# P_{(\kappa, \beta)} \quad (3)$$

where $P_{(\kappa, \beta)}$ is the stationary distribution of ϕ_κ at parameters β , and $(h)_\#$ denotes the pushforward through h . The *accessible family* at coupling κ is the image:

$$\mathcal{M}_{\text{acc}}(\kappa) = \text{Image}(F_\kappa) \subset \mathcal{P}(\mathbb{R}^m)$$

Role of the observation map. The observation map $h : X_1 \times X_2 \rightarrow \mathbb{R}^m$ determines which features of the joint state are statistically accessible. When h is the identity (full state observation), identifiability depends only on dynamics. When h is a low-dimensional projection (partial observation), some coordinates may remain unidentifiable even if the dynamics distinguish them. Our rank-change theorems concern the map $F_\kappa = (h)_\# \circ P_{(\kappa, \cdot)}$; both dynamics and observation contribute to the accessible family. In the examples, we specify h explicitly and note when identifiability depends on observation richness.

Standing assumptions. Throughout, we assume:

- (A1) For each (κ, β) , the dynamics ϕ_κ admit a (locally) unique stationary distribution $P_{(\kappa, \beta)}$.
- (A2) The map $\beta \mapsto P_{(\kappa, \beta)}$ (and hence F_κ) is smooth.
- (A3) The observation map h is smooth with sufficient regularity that $Q = (h)_\# P$ has a smooth density on its support.

These assumptions exclude bifurcation points where stationary distributions split or vanish. Our results apply to parameter regions where (A1)–(A3) hold; analysis at bifurcation points requires additional care.

Our claim: For high-dimensional coherent systems, the rank of dF_κ can increase as κ crosses a threshold—the image manifold gains dimension.

Definition 1 (Manifold expansion) Define the *generic rank* of F_κ as:

$$r(\kappa) := \sup_{\beta \in \mathcal{B}} \text{rank } dF_\kappa(\beta) \quad (4)$$

By lower semicontinuity of rank (the sets $\{\beta : \text{rank } dF_\kappa(\beta) \geq r\}$ are open), the set where rank equals the supremum is open. We call $r(\kappa)$ the “generic rank” when this open set is nonempty; in applications we restrict to parameter regions where the supremum is attained.

Manifold expansion occurs at κ_c if the generic rank increases as κ crosses κ_c :

$$r(\kappa_c^+) > r(\kappa_c^-) \quad (5)$$

Equivalently, the dimension of $\mathcal{M}_{\text{acc}}(\kappa)$ increases: a parameter direction in \mathcal{B} that was generically in the kernel of dF_κ becomes generically non-degenerate under coupling.

In Fisher-metric terms: manifold expansion corresponds to an increase in Fisher rank—the number of generically linearly independent score directions.

This definition is coordinate-invariant: rank is preserved under reparameterization. It captures “new identifiable directions appear” in an intrinsic way.

The novelty is not that interaction parameters exist in the full parametric family—Amari (2001)’s hierarchical decomposition already accounts for this. The novelty is that **dynamics and observation determine which interaction coordinates are identifiable**. Coupling can *activate* coordinates that were structurally present but dynamically inaccessible—a phenomenon we term **coupling-induced identifiability activation**.

3.2 Mechanism: collective coordinates

Let $\phi : X_1 \times X_2 \rightarrow X_1 \times X_2$ be the coupled dynamics. Under coupling, new *collective coordinates* can become accessible:

Definition 2 (Collective coordinate) A collective coordinate is a function $\psi : X_1 \times X_2 \rightarrow \mathbb{R}$ that is:

1. Not reducible to functions of X_1 or X_2 alone
2. *Dynamically stable*: slow manifold, metastable basin, or conserved quantity
3. *Statistically identifiable*: non-degenerate Fisher information along ψ under observation h

Examples:

- Phase difference between coupled oscillators (Strogatz 2000; Acebrón et al. 2005)
- Synchronization manifold coordinates (Pikovsky et al. 2001)
- Order parameters of collective states (e.g., chimera states) (Panaggio and Abrams 2015)
- Interface modes at boundaries between systems

These coordinates represent genuinely new degrees of freedom: they were not accessible to either system in isolation.

3.3 Main theorem: coupling-induced rank transitions

We now state and prove our main result, which provides checkable criteria for when coupling induces Fisher-rank transitions.

Structure and novelty. The theorem has three parts. Part I recalls the standard pullback characterization of Fisher information; we include it for completeness and to fix notation. **The novel content is Parts II and III**, which provide checkable criteria—in terms of kernel activation—for when coupling *increases* the Fisher rank.

Theorem 1 (Coupling-induced rank transitions) *Let \mathcal{B} be a smooth d -dimensional parameter manifold, let \mathcal{P} be a finite-dimensional statistical manifold on observation space Y equipped with the Fisher information metric g , and let $F_\kappa : \mathcal{B} \rightarrow \mathcal{P}$ be a smooth family of maps indexed by coupling strength $\kappa \geq 0$, where $F_\kappa(\beta) = Q_{(\kappa, \beta)}$ is the distribution of observations under dynamics with parameters (κ, β) .*

Part I (Pullback identity). *The Fisher information matrix on \mathcal{B} induced by the accessible family $\mathcal{M}_{\text{acc}}(\kappa) = \text{Image}(F_\kappa)$ is the pullback of the Fisher metric:*

$$I_{\mathcal{B}}(\kappa) = F_\kappa^* g \quad (6)$$

Consequently:

$$\text{rank } I_{\mathcal{B}}(\kappa) = \text{rank } dF_\kappa \quad (7)$$

Part II (Constraint-release criterion). *Let $\mathcal{M}_0 \subset \mathcal{P}$ be a smooth submanifold of codimension c (e.g., the independence submanifold). Choose local coordinates on \mathcal{P} near \mathcal{M}_0 so that $\mathcal{M}_0 = \{\eta = 0\}$ locally, and write $F_\kappa(\beta) = (f_\kappa(\beta), \eta_\kappa(\beta))$ where f are coordinates along \mathcal{M}_0 and η are transverse coordinates. Let $\beta_0 \in \mathcal{B}$ achieve the maximal rank $r(0) = \text{rank } dF_0(\beta_0)$. Suppose:*

- (C1) $F_0(\mathcal{B}) \subseteq \mathcal{M}_0$, i.e., $\eta_0(\beta) \equiv 0$ (dynamics constrain to submanifold at $\kappa = 0$)
- (C2) For $\kappa > \kappa_c$, there exists β such that $\eta_\kappa(\beta) \neq 0$ (constraint released)
- (C3) **Kernel activation:** There exists $v \in \ker(df_0(\beta_0))$ such that $d\eta_\kappa(v) \neq 0$ for $\kappa > \kappa_c$

Then for $\kappa > \kappa_c$:

$$r(\kappa) \geq r(0) + 1 \quad (8)$$

Part III (Symmetry-breaking criterion). *Let G be a k -dimensional Lie group acting smoothly on \mathcal{P} , and let $\text{Fix}(G) = \{p \in \mathcal{P} : \gamma \cdot p = p \text{ for all } \gamma \in G\}$ be the fixed-point set. Let $\beta_0 \in \mathcal{B}$ achieve the maximal rank $r(0) = \text{rank } dF_0(\beta_0)$. Suppose:*

- (S1) $F_0(\mathcal{B}) \subseteq \text{Fix}(G)$ (distributions are G -invariant at $\kappa = 0$)
- (S2) For $\kappa > \kappa_c$, there exists β such that $F_\kappa(\beta) \notin \text{Fix}(G)$ (symmetry broken)
- (S3) The orbit $G \cdot F_\kappa(\beta)$ has dimension k (free action locally)

(S4) **Kernel-to-orbit activation:** *There exist k linearly independent vectors $v_1, \dots, v_k \in \ker(dF_0(\beta_0))$ such that $dF_\kappa(v_1), \dots, dF_\kappa(v_k)$ project to linearly independent vectors in $T_p(G \cdot p)$ for $\kappa > \kappa_c$*

Then for $\kappa > \kappa_c$:

$$r(\kappa) \geq r(0) + k \quad (9)$$

Proof Part I. Let $\beta(t)$ be a smooth curve in \mathcal{B} with $\beta(0) = \beta_0$ and $\dot{\beta}(0) = v \in T_{\beta_0}\mathcal{B}$. The image curve $Q(t) = F_\kappa(\beta(t))$ has tangent vector $dF_\kappa(v) \in T_{Q_0}\mathcal{P}$.

The Fisher metric on \mathcal{P} evaluates tangent vectors via score functions. By the chain rule:

$$\frac{\partial \log q(y; \kappa, \beta)}{\partial \beta^i} = \sum_{\alpha} \frac{\partial \log q}{\partial \theta^\alpha} \frac{\partial \theta^\alpha}{\partial \beta^i} \quad (10)$$

where θ are coordinates on \mathcal{P} . The pullback metric is:

$$(F_\kappa^*g)_{ij} = g\left(\frac{\partial F_\kappa}{\partial \beta^i}, \frac{\partial F_\kappa}{\partial \beta^j}\right) = (I_{\mathcal{B}}(\kappa))_{ij} \quad (11)$$

This establishes the pullback identity.

For the rank equality: the pullback metric F_κ^*g is degenerate precisely along the kernel of dF_κ . A vector $v \in T_{\beta}\mathcal{B}$ satisfies $(F_\kappa^*g)(v, v) = 0$ if and only if $dF_\kappa(v) = 0$ (since g is positive-definite on $T_Q\mathcal{P}$). Thus $\ker(I_{\mathcal{B}}(\kappa)) = \ker(dF_\kappa)$ and $\text{rank}(I_{\mathcal{B}}(\kappa)) = \text{rank } dF_\kappa$.

Part II. In local coordinates with $\mathcal{M}_0 = \{\eta = 0\}$, write $F_\kappa(\beta) = (f_\kappa(\beta), \eta_\kappa(\beta))$. At $\kappa = 0$, condition (C1) gives $\eta_0 \equiv 0$, so the image of dF_0 lies entirely in the f -directions with $\dim(\text{Image}(dF_0(\beta_0))) = r(0)$.

Let $v \in \ker(df_0(\beta_0))$ be as in (C3). At $\kappa = 0$, $dF_0(v) = (df_0(v), d\eta_0(v)) = (0, 0)$, so $v \in \ker(dF_0(\beta_0))$. By (C3), for $\kappa > \kappa_c$, $d\eta_\kappa(v) \neq 0$, so $dF_\kappa(v)$ has a nonzero transverse component.

The key observation is that $T_p\mathcal{P} \cong T_p\mathcal{M}_0 \oplus N_p\mathcal{M}_0$. At $\kappa = 0$, the image of dF_0 lies entirely in $T\mathcal{M}_0$. The vector $dF_\kappa(v)$ has a nonzero component in $N\mathcal{M}_0$, so it cannot lie in the span of the old $r(0)$ tangential directions. By continuity of the Gram determinant, the old $r(0)$ directions persist, so $r(\kappa) \geq r(0) + 1$.

Part III. At $\kappa = 0$, (S1) implies $F_0(\mathcal{B}) \subseteq \text{Fix}(G)$, so $\text{Image}(dF_0) \subseteq T(\text{Fix}(G))$ (no orbit directions).

Let $v_1, \dots, v_k \in \ker(dF_0(\beta_0))$ be as in (S4). At $\kappa = 0$, $dF_0(v_i) = 0$ for all i . For $\kappa > \kappa_c$, condition (S4) states that $dF_\kappa(v_1), \dots, dF_\kappa(v_k)$ project to k linearly independent vectors in the orbit tangent space $T_p(G \cdot p)$.

At points $p \notin \text{Fix}(G)$, the tangent space decomposes as $T_p\mathcal{P} \cong T_p(G \cdot p) \oplus (T_p(G \cdot p))^\perp$. The vectors $dF_\kappa(v_i)$ have nonzero projections onto $T_p(G \cdot p)$, so they cannot lie in the span of the old $r(0)$ directions. Therefore $r(\kappa) \geq r(0) + k$.

3.4 Superadditive accessible complexity

The preceding theorem establishes conditions for rank increase under coupling. We now formalize “the whole exceeds the sum of its parts” as a precise inequality.

Definition 3 (Component accessible families) For each subsystem $i = 1, \dots, n$, let \mathcal{B}_i be its parameter manifold and let $F_i : \mathcal{B}_i \rightarrow \mathcal{P}_i$ be its dynamics-to-distribution map. Define the component accessible family $\mathcal{M}_i := \text{Image}(F_i)$ and its generic rank $r_i := \max_{\beta_i \in \mathcal{B}_i} \text{rank } dF_i(\beta_i)$.

Definition 4 (Superadditive accessible complexity) Let $F_\kappa : \mathcal{B} \rightarrow \mathcal{P}$ be the coupled dynamics-to-distribution map with generic rank $r(\kappa) := \max_{\beta \in \mathcal{B}} \text{rank } dF_\kappa(\beta)$. The coupled system exhibits *superadditive accessible complexity* at coupling κ if

$$r(\kappa) > r_\oplus = \sum_i r_i.$$

Lemma 1 (Additivity under independence) For the independent composite map F_\oplus , $r_\oplus = \sum_{i=1}^n r_i$.

Proof The differential dF_\oplus is block-diagonal with blocks dF_i on $T_{\beta_i} \mathcal{B}_i$. Hence $\text{rank } dF_\oplus(\beta) = \sum_i \text{rank } dF_i(\beta_i)$. Maximizing over β yields $r_\oplus = \sum_i r_i$.

Proposition 1 (Sufficient condition via non-factorizable coordinates)

Suppose there exists a smooth map $\pi : \mathcal{P} \rightarrow \mathbb{R}^k$ such that:

- (A1) $\pi \circ F_\oplus$ is locally constant on some open $U \subset \mathcal{B}$;
- (A2) $\pi \circ F_\kappa$ has Jacobian rank k on U .

Then $r(\kappa) \geq r_\oplus + k$, so the coupled system is superadditive whenever $k \geq 1$.

3.5 Example 1: coupled Ornstein–Uhlenbeck (transversality)

This example demonstrates Theorem 1 Part II (transversality criterion) with $\kappa_c = 0$.

Dynamics. Consider two coupled Ornstein–Uhlenbeck processes:

$$dX_1 = -\gamma_1 X_1 dt + \kappa(X_2 - X_1) dt + \sigma_1 dW_1 \quad (12)$$

$$dX_2 = -\gamma_2 X_2 dt + \kappa(X_1 - X_2) dt + \sigma_2 dW_2 \quad (13)$$

where $\gamma_i > 0$ are mean-reversion rates, $\sigma_i > 0$ are noise intensities, and $\kappa \geq 0$ is the coupling strength.

Stationary distribution. The stationary distribution is bivariate Gaussian with zero mean and covariance matrix Σ satisfying the Lyapunov equation $A\Sigma + \Sigma A^T + D = 0$. For the symmetric case $\gamma_1 = \gamma_2 = \gamma$ and $\sigma_1 = \sigma_2 = \sigma$:

$$\Sigma_{11} = \Sigma_{22} = \frac{\sigma^2}{2\gamma}, \quad \Sigma_{12} = \frac{\kappa\sigma^2}{2\gamma(\gamma + 2\kappa)} \quad (14)$$

giving correlation $\rho = \kappa/(\gamma + 2\kappa)$.

This confirms $\rho(\kappa = 0) = 0$ and shows the rank transition explicitly. Computing $\partial\rho/\partial\gamma$:

$$\frac{\partial\rho}{\partial\gamma} = -\frac{\kappa}{(\gamma + 2\kappa)^2} \neq 0 \quad \text{for } \kappa > 0 \quad (15)$$

This nonzero derivative demonstrates kernel activation: a direction that left ρ unchanged at $\kappa = 0$ now varies ρ at $\kappa > 0$.

Verification of Theorem 1 Part II. The ambient manifold \mathcal{P} of centered bivariate Gaussians is 3-dimensional, parameterized by (v_1, v_2, ρ) . The independence submanifold is $\mathcal{M}_0 = \{(v_1, v_2, 0)\}$, which has codimension $c = 1$.

At $\kappa = 0$: $\rho = 0$ for all β , so $F_0(\mathcal{B}) \subseteq \mathcal{M}_0$ and $r(0) = 2$.

At $\kappa > 0$: $\rho > 0$, so $F_\kappa(\mathcal{B})$ escapes \mathcal{M}_0 and $r(\kappa) = 3$.

Superadditivity. Each OU process in isolation has a 1-dimensional accessible family (variance only), so $r_1 = r_2 = 1$ and $r_\oplus = 2$. The coupled system achieves $r(\kappa) = 3 > r_\oplus = 2$: *strict superadditivity*.

3.6 Example 2: Kuramoto synchronization (symmetry breaking)

This example demonstrates Theorem 1 Part III (symmetry-breaking criterion) with $\kappa_c > 0$. We work in the **mean-field limit** ($N \rightarrow \infty$) with added noise, where the stationary density is unique up to symmetry (Strogatz 2000; Acebrón et al. 2005).

Mean-field dynamics. Let $\rho(\theta, \omega, t)$ be the density of oscillators with natural frequency ω at phase θ . With phase diffusion of strength $D > 0$, the density evolves according to the nonlinear Fokker–Planck equation:

$$\frac{\partial\rho}{\partial t} + \frac{\partial}{\partial\theta} [\rho(\omega + Kr \sin(\Psi - \theta))] = D \frac{\partial^2\rho}{\partial\theta^2} \quad (16)$$

where the order parameter $re^{i\Psi} = \int e^{i\theta} \rho(\theta, \omega, t) \nu(\omega) d\omega d\theta$ is determined self-consistently. The system has S^1 symmetry.

Stationary solutions. For $K < K_c$, the stable stationary solution is the incoherent S^1 -invariant density with $r = 0$. For $K > K_c$, there exists an S^1 -orbit of synchronized stationary solutions with $r > 0$.

Verification of Theorem 1 Part III.

- (S1): For $K < K_c$, the image $F_K(\mathcal{B})$ is the single incoherent distribution, which is S^1 -invariant.
- (S2): For $K > K_c$, the synchronized stationary solutions break S^1 symmetry ($r > 0$).
- (S3): The S^1 -orbit of synchronized solutions is 1-dimensional.
- (S4): For $K < K_c$, the map $F_K : \Psi \mapsto \rho_{K,\Psi}$ is constant (all Ψ give the same incoherent distribution), so $dF_K = 0$. For $K > K_c$, $\partial F_K / \partial \Psi \neq 0$ and maps onto the orbit direction (kernel-to-orbit activation).

The theorem predicts: $r(K) \geq r(0)+1$ for $K > K_c$. Since $r(0) = 0$ (constant map), we get $r(K) \geq 1$.

Key difference from OU. In the OU example (Part II), any $\kappa > 0$ immediately releases the constraint: $\kappa_c = 0$. In Kuramoto (Part III), the symmetry persists for $0 < K < K_c$; manifold expansion occurs at a *genuine* critical threshold $K_c > 0$.

4 Consequences for complexity growth

4.1 Faster than fixed-class bounds predict

Standard bounds on complexity growth assume fixed parameterization, fixed sufficient statistics, and fixed dimensionality.

When coupling creates new coordinates:

- The Fisher geometry changes because the model class changes—new significant Fisher directions appear
- KL-based learning rates underpredict adaptation
- Channel capacity arguments don’t apply (the “channel” itself is being re-structured)

Definition 5 (Model complexity) We define *complexity* C of a statistical family \mathcal{M} as the Fisher rank: the dimension of the identifiable tangent space. This is coordinate-invariant and intrinsic to the manifold geometry.

Conjecture 1 (Eigenvalue emergence under coupling) In regimes where the accessible family $\mathcal{M}_{\text{acc}}(\kappa)$ can be embedded in a fixed parameterization across κ , eigenvalue emergence provides a concrete signature. As coupling strength κ crosses a critical threshold κ_c , one or more Fisher eigenvalues emerge from zero:

$$\lambda_k(\kappa_c^-) < \epsilon \quad \text{and} \quad \lambda_k(\kappa_c^+) > \epsilon \quad (17)$$

This implies rank increase as coupling activates previously degenerate interaction coordinates.

4.2 Heuristic scaling: colony dynamics

Consider N high-D systems in a “colony” (weakly coupled network). The *ambient space* of possible interaction parameters grows rapidly: with N subsystems of internal dimension n , pairwise cross-covariances alone contribute $O(N^2 \cdot n^2)$ potential coordinates. However, the *accessible* submanifold dimension is constrained by coupling parameters, symmetry constraints, attractor structure, and observation map.

A rough upper bound, if each pairwise coupling can create k new identifiable coordinates:

$$\text{rank } I_{\text{colony}} \leq N \cdot \text{rank } I_{\text{single}} + |E| \cdot k \quad (18)$$

where $|E|$ is the number of coupling edges. For dense networks $|E| \sim N^2$; for sparse networks $|E| \sim N$.

5 Discussion

5.1 Summary of contributions

This paper provides:

1. A formal definition of **manifold expansion** as Fisher-rank increase under coupling (Definition 1)
2. **Theorem 1**: Coupling-induced rank transitions, with three parts:
 - Part I: Pullback identity (Fisher rank = Jacobian rank)
 - Part II: Transversality criterion for rank increase
 - Part III: Symmetry-breaking criterion for rank increase
3. Worked examples verifying the theorem: coupled OU (Part II, $\kappa_c = 0$) and Kuramoto oscillators (Part III, $\kappa_c > 0$)

The core technical contribution is identifying **coupling-induced identifiability activation**: interaction coordinates that are structurally present but dynamically inaccessible can become identifiable under coupling, increasing the Fisher rank of the observable family.

5.2 Testable predictions

The framework makes specific predictions:

1. **Fisher rank increase**: For coupled dynamical systems, estimate Fisher information from time-series data at varying coupling strengths κ . Predict: $\text{rank } I(\kappa)$ increases as κ crosses critical thresholds.
2. **Eigenvalue emergence**: New Fisher eigenvalues should emerge continuously from zero as coupling increases.
3. **Superlinear scaling**: For colonies of N coupled subsystems, the identifiable parameter count should grow faster than N in the manifold-expansion regime.

5.3 Scope and limitations

Our analysis assumes that the accessible family forms a smooth parametric model. In practice, the set of stationary distributions induced by dynamics may have irregular structure; our results apply where smooth-family regularity holds at least locally.

The mechanism is geometric, not substrate-specific. It applies wherever high-dimensional coherent systems couple and satisfy the regularity conditions of Theorem 1.

6 Conclusion

Information geometry has been remarkably successful because the fixed-manifold assumption usually holds. We have identified a regime where it does not: when high-dimensional coherent systems couple, the identifiable statistical family can expand.

This is not a rejection of information-theoretic bounds. It is recognition that such bounds apply to fixed model classes, and coupling can change which parameters are identifiable. The key contribution is formalizing this as **manifold expansion**—defined via Fisher-rank increase—and establishing conditions under which superadditive dimensionality occurs.

In the regime described here, a useful description is constraint exchange rather than token exchange: high-dimensional systems reshape each other's accessible state spaces rather than transmitting discrete symbols through a fixed channel.

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