

Ex 1. we provide a different method which can be generalized to $n \geq 3$. See remark.

For $t \in [0, T]$, let $Z_t^z = \mathbb{1}_{\tau_t} e^{\int_0^t \mathbb{1}_{\tau_s} \lambda_s ds}$ and define

$$\frac{d\alpha^2}{d\alpha} = Z_T^2 \quad \text{on } (\Sigma, g_T)$$

Then, $H_t^Z = \mathbb{1}_{Z \leq t}$, $t \in [0, T]$, is 0, Q^Z -a.e.

Für $T > U \geq 0$, $\mathcal{Q}(\tau_T, \bar{\tau}_U)$

$$= \mathcal{Q} \left(Z_T^2 e^{-\int_0^T \mathbb{1}_{\tau > s} \lambda_s ds} \mathbb{1}_{\bar{\tau} > T} \right)$$

$$= Q(Z_T^Z e^{-\int_0^T \lambda_s ds} \mathbb{1}_{\bar{Z}_U})$$

$$= Q^2 (e^{-\int_0^T a_1 + a_2 \mathbb{1}_{\bar{z} \leq s} ds} \mathbb{1}_{\bar{z} > 0})$$

Since under \mathcal{Q}^2 , $\bar{X}_s = \bar{a}_1 + \bar{a}_2 \mathbb{1}_{\tau \leq s} = \bar{a}_1$ for $s \in [0, \tau]$, \mathcal{Q}^2 -a.e.

it follows that $\mathcal{Q}(\mathcal{C}_T, \mathcal{C}_U)$

$$= e^{-a_1 T} \int_{-\infty}^T e^{-a_2 (T-x)} \bar{a}_1 e^{-\bar{a}_1 x} dx + \int_T^{\infty} \bar{a}_1 e^{-\bar{a}_1 x} dx$$

In turn,

$$\frac{\partial}{\partial T} \frac{\partial}{\partial U} Q(\omega_T, \bar{z}, U) = \frac{\partial}{\partial T} \frac{\partial}{\partial U} \left(e^{-a_1 T} \int_{-1}^1 e^{-a_2(T-x)} \bar{a}_1 e^{-\bar{a}_1 x} dx \right)$$

$$= \frac{\partial}{\partial T} e^{-a_1 T} \cdot (-e^{-a_2 (T+U)} a_1 e^{-a_1 U})$$

$$= \bar{a}_1 (a_1 + a_2) e^{-(a_1 + a_2)T - (\bar{a}_1 - a_1)U} \quad T > U \geq 0$$

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Remark: For $n=3$, consider z^1, z^2, z^3 with their corresponding intensities

$$\lambda'_s = a + b \mathbb{1}_{c_3 \leq t}$$

$$\lambda_s^2 = a + b \mathbb{1}_{\tau_3 \leq t}$$

$$\lambda_s^3 = c + d \mathbb{1}_{\tau_1 \leq t} + d \mathbb{1}_{\tau_2 \leq t}.$$



Case 1. $T^3 T^1 T^2 \geq 0$. let $Z_t^j = \frac{1}{\tau_{j,t}^3} e^{\int_0^t \tau_{j,s}^3 \lambda_s^3 ds}$ and define

$$\frac{dQ^3}{dQ} = Z_T^3 \ln(\ln g_T^3)$$

7/0. $11^3 \pm 1$, i.e. -37 is $10^3 - 40$

$$\frac{d\tilde{\alpha}}{d\alpha} = Z_{T^3} \text{ on } (\Omega, \mathcal{G}_{T^3})$$

Then $H_t^3 \triangleq \mathbb{1}_{t^3 \leq t}$, $t \in [0, T^3]$, is 0 \mathcal{Q}^3 -a.e.

$$\begin{aligned} \mathcal{Q}(\tau^i > T^i, 1 \leq i \leq 3) &= \mathcal{Q}(Z_{T^3}^3 e^{-\int_0^{T^3} \mathbb{1}_{\tau^3 > s} \lambda_s^3 ds} \mathbb{1}_{\tau^1 > T^1} \mathbb{1}_{\tau^2 > T^2}) \\ &= \mathcal{Q}(Z_{T^3}^3 e^{-\int_0^{T^3} \lambda_s^3 ds} \mathbb{1}_{\tau^1 > T^1} \mathbb{1}_{\tau^2 > T^2}) \\ &= \mathcal{Q}^3(e^{-\int_0^{T^3} c + d \mathbb{1}_{\tau^3 > s} + d \mathbb{1}_{\tau^2 \leq s} ds} \mathbb{1}_{\tau^1 > T^1} \mathbb{1}_{\tau^2 > T^2}) \end{aligned}$$

Since under \mathcal{Q}^3 , $\lambda^1 = \lambda^2 = a$ for $s \in [0, T^3]$, \mathcal{Q}^3 -a.e.

it follows that $\mathcal{Q}(\tau^i > T^i, 1 \leq i \leq 3)$

$$\begin{aligned} &= e^{-cT^3} \left(\int_{T^1}^{T^3} dx_1 \int_{T^2}^{T^3} dx_2 e^{-d(T^3-x_1)} e^{-d(T^3-x_2)} f_{\tau^1}(x_1) f_{\tau^2}(x_2) \right. \\ &\quad + \int_{T^1}^{T^3} dx_1 \int_{T^3}^{\infty} dx_2 e^{-d(T^3-x_1)} f_{\tau^1}(x_1) f_{\tau^2}(x_2) \\ &\quad + \int_{T^3}^{\infty} dx_1 \int_{T^2}^{T^3} dx_2 e^{-d(T^3-x_1)} f_{\tau^1}(x_1) f_{\tau^2}(x_2) \\ &\quad \left. + \int_{T^3}^{\infty} dx_1 \int_{T^3}^{\infty} dx_2 f_{\tau^1}(x_1) f_{\tau^2}(x_2) \right) \text{ with } f_{\tau^i}(x) = ae^{-ax} \end{aligned}$$

In turn, $\frac{\partial}{\partial T^1} \frac{\partial}{\partial T^2} \frac{\partial}{\partial T^3} \mathcal{Q}(\tau^i > T^i, 1 \leq i \leq 3)$

$$\begin{aligned} &= \frac{\partial}{\partial T^1} \frac{\partial}{\partial T^2} \frac{\partial}{\partial T^3} \left(e^{-cT^3} \int_{T^1}^{T^3} dx_1 \int_{T^2}^{T^3} dx_2 e^{-d(T^3-x_1)} e^{-d(T^3-x_2)} ae^{-ax_1} ae^{-ax_2} \right) \\ &= \frac{\partial}{\partial T^3} e^{-cT^3} \cdot (-e^{-d(T^3-T^1)} ae^{-aT^1}) \cdot (-e^{-d(T^3-T^2)} ae^{-aT^2}) \\ &= \frac{\partial}{\partial T^3} (-1)^2 a^2 e^{-(c+2d)T^3} e^{-(a-d)(T^1+T^2)} \\ &= (-1)^3 a^2 (c+2d) e^{-(c+2d)T^3 - (a-d)(T^1+T^2)} \text{ for } T^3 > T^1 \vee T^2 \geq 0 \end{aligned}$$

Case 2 $T^1 \vee T^2 > T^3 \geq 0$. Without loss of generality, consider $T^1 > T^2$.

Let $Z_t^1 = \mathbb{1}_{\tau^1 > t} e^{-\int_0^t \mathbb{1}_{\tau^1 > s} \lambda_s^1 ds}$ and define

$$\frac{d\alpha'}{d\alpha} = Z_{T^1}^1 \text{ on } (\Omega, \mathcal{G}_{T^1})$$

Th. 4.1.1. $\mathbb{1}_{\tau^1 > t}$ is 0 \mathcal{Q}^1 -a.e.

$$\frac{d\alpha}{dQ} = 2T'_1 \text{ on } (\Omega, \mathcal{F}_{T'_1})$$

Then $H'_t = \mathbb{1}_{\tau^1 \leq t}$, $t \in [0, T'_1]$, is 0 Q' -a.e.

Note that under Q' : $\lambda_s^2 = a + b \mathbb{1}_{\tau^3 \leq s}$

$$\lambda_s^3 = c + d \mathbb{1}_{\tau^2 \leq s}$$

$$\text{Hence, } Q(\tau^i > T^i, 1 \leq i \leq 3) = Q'(e^{-\int_0^{T'_1} a + b \mathbb{1}_{\tau^3 \leq s} ds} \mathbb{1}_{\tau^2 > T^2} \mathbb{1}_{\tau^3 > T^3})$$

$$= e^{-aT'} \left(\int_{T^2}^{\infty} dx_2 \int_{T^3}^{T'_1} dx_3 e^{-b(T'_1 - x_3)} f_{23}(x_2, x_3) \right.$$

$$\left. + \int_{T^2}^{\infty} dx_2 \int_{T^1}^{\infty} dx_3 f_{23}(x_2, x_3) \right)$$

with $f_{23}(x_2, x_3)$ joint density of (τ^2, τ^3) derived in Ex 1.

$$\text{In turn, } \frac{\partial}{\partial T_1} \frac{\partial}{\partial T_2} \frac{\partial}{\partial T_3} Q(\tau^i > T^i, 1 \leq i \leq 3)$$

$$= \frac{\partial}{\partial T_1} \frac{\partial}{\partial T_2} \frac{\partial}{\partial T_3} \left(e^{-aT'} \int_{T^2}^{\infty} dx_2 \int_{T^3}^{T'_1} dx_3 e^{-b(T'_1 - x_3)} f_{23}(x_2, x_3) \right)$$

$$= \frac{\partial}{\partial T_1} \frac{\partial}{\partial T_3} e^{-aT'} \cdot \left(- \int_{T^3}^{T'_1} dx_3 e^{-b(T'_1 - x_3)} f_{23}(T^2, x_3) \right)$$

$$= \frac{\partial}{\partial T_1} e^{-aT'} \cdot e^{-b(T'_1 - T^3)} f_{23}(T^2, T^3)$$

$$= (-1)(a+b) e^{-(a+b)T'_1 + bT^3} f_{23}(T^2, T^3) \quad T' > T^2 \vee T^3 \geq 0$$

Similarly, for $T^2 \geq T'$,

$$\frac{\partial}{\partial T_1} \frac{\partial}{\partial T_2} \frac{\partial}{\partial T_3} Q(\tau^i > T^i, 1 \leq i \leq 3) = (-1)(a+b) e^{-(a+b)T^2 + bT^3} f_{13}(T', T^3)$$

$$T^2 > T' \vee T^3 \geq 0.$$

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