

UNIVERSITY OF WATERLOO

PHYS460B

Radio Frequency Electronics

EXPERIMENT #10

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A Voltage Controlled Oscillator

This part is to identify voltage-frequency response of ZX95-850+ voltage controlled oscillator.

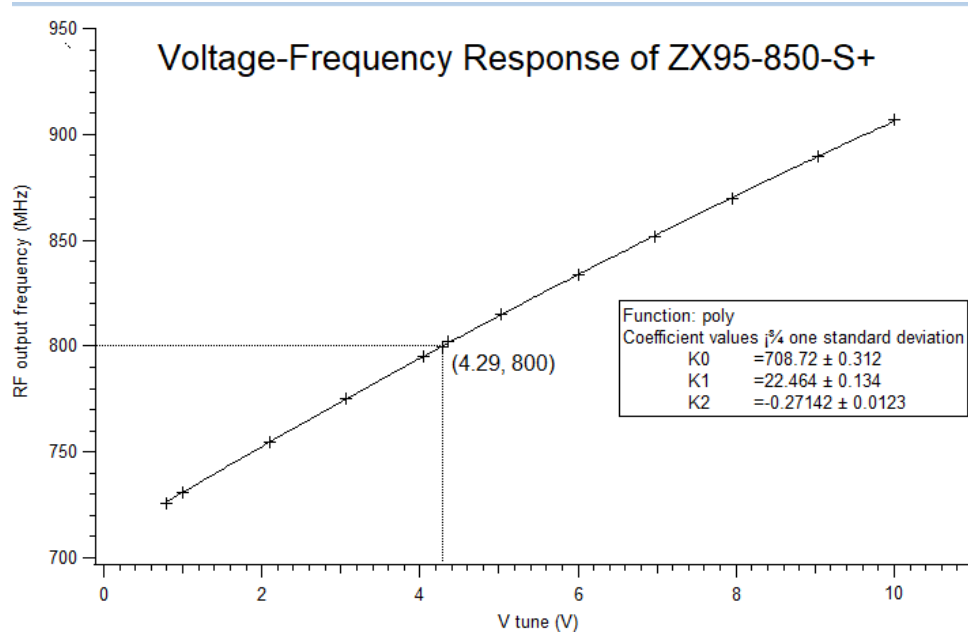


Figure 1: Voltage-frequency response of ZX95-850-S+. The error in measurements was too small to be displayed in the graph. ($\pm 0.01 MHz$)

The voltage-frequency response of the VCO is shown above. It is close to a linear graph, but it is a parabolic graph with the best fit equation below:

$$Frequency(MHz) = -0.27142V_{tune}^2 + 22.464V_{tune} + 708.72 \quad (1)$$

With above equation, $f = 800 MHz$ when $V_{tune} \cong 4.29V$. To determine the slope at $800 MHz$, we just differentiated above equation and substituted $V_{tune} \cong 4.29V$.

$$slope(800MHz) = -0.54284V_{tune} + 22.464 = (20.135 \pm 0.922) MHz/V \quad (2)$$

Error analysis

$$\Delta slope = slope \cdot \sqrt{\left(\frac{0.0246}{0.54280}\right)^2 + \left(\frac{0.134}{22.464}\right)^2 + \left(\frac{0.01}{4.29}\right)^2} = 0.922 \quad (3)$$

A.1 Additional Questions

1. Comparison with the known data from the Minicircuits web-site. The equation of the parabolic best fit function is :

$$Frequency(MHz) = -0.27043V_{tune}^2 + 22.766V_{tune} + 712.04 \quad (4)$$

Substituting $f = 800 MHz$, $V_{tune} \cong 4.06V$. The slope with the differentiated equation is:

$$slope(800MHz) = -0.54086V_{tune} + 22.766 = (20.570 \pm 1.108) MHz/V \quad (5)$$

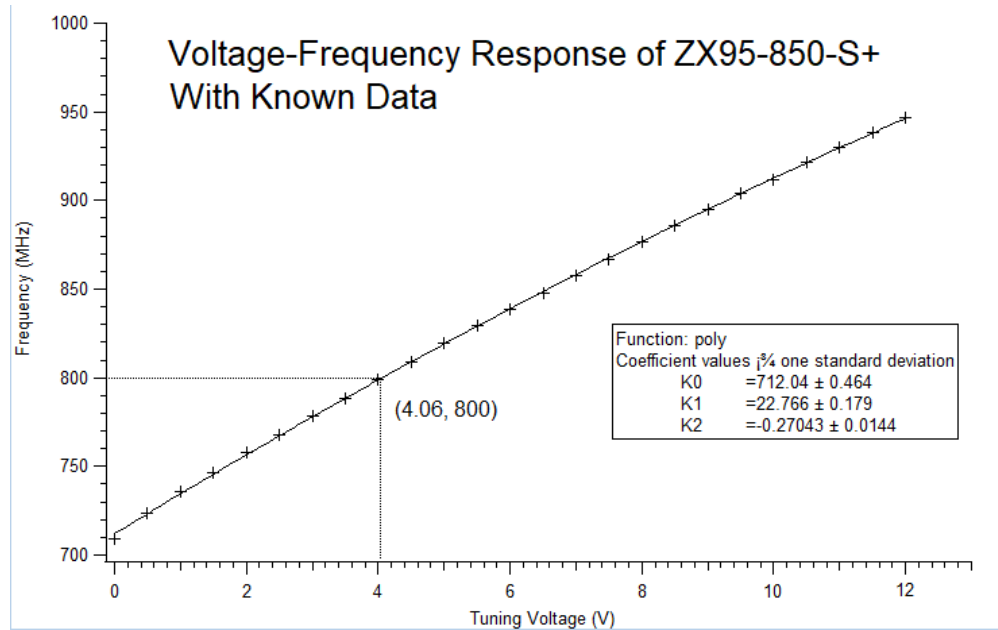


Figure 2: Voltage-frequency response of ZX95-850-S+ when $T=25^{\circ}\text{C}$ with given data from the Minicircuits web-site.

Error analysis

$$\Delta \text{slope} = \text{slope} \cdot \sqrt{\left(\frac{0.0288}{0.54086}\right)^2 + \left(\frac{0.179}{22.766}\right)^2 + \left(\frac{0.01}{4.06}\right)^2} = 1.108 \quad (6)$$

$$\frac{\text{slope}_{\text{known}} - \text{slope}_{\text{measured}}}{\text{slope}_{\text{known}}} = 0.021 \quad (7)$$

$$\left| \frac{V_{800\text{MHz}}^{\text{known}} - V_{800\text{MHz}}^{\text{measured}}}{V_{800\text{MHz}}^{\text{known}}} \right| = \left| \frac{4.06 - 4.29}{4.06} \right| \cong 5.7\text{percents} \quad (8)$$

The slope at $f = 800\text{MHz}$ of the known data is only 2.1 percents off from the slope of the measured data, and they agree to each other within the error range calculated. The constants of parabolic best fits agrees to each other within 1 percents of error. The difference between V_{tune} values when $f = 800\text{MHz}$ is 5.7 percents which is higher than others, but this is due to the voltage drawn from the extra components of the circuit (ie. wire). The extra voltage drawn caused x-shift in right direction of the graph, but it did not affect on the slope when the frequency is 800MHz.

There should be an error due to the temperature of the laboratory. Although the temperature was not explicitly measured, it was lower than 25°C , which is the temperature that the known data collected. However, change in voltage-frequency response of VCO due to the temperature is negligible. In Minicircuits website, data collected with $T=-55^{\circ}\text{C}$ and $T=85^{\circ}\text{C}$ have no mentionable difference with data with $T=25^{\circ}\text{C}$. Therefore, error due to the temperature can be ignored.

2.The isolator is an equipment that transmits radio frequency in one direction only. In this circuit the isolator is used to reduce phase noise of the VCO.

B Cavity Coupling and Quality Factor

This portion of the experiment is concerned with the cavity, which is a quarter wavelength coaxial resonator.

We note that the cavity only resonates at a particular frequency, at all other frequencies we will have a reflection of the voltage signal.

B.1 Cavity Reflection

First we test this by ramping through the V_{tune} voltages, which equates to ramping through the frequencies as shown in Part A more specifically in equation 4.

An examples of such a graph is shown in Figure 3. We note that a raw V_{tune} ramp is noisy as the oscilloscope does not output the required accuracy, so we take a linear fit of the V_{tune} . From this point onward unless otherwise stated, we will always present the V_{tune} as frequency which has undergone a linear voltage fit and gone through the conversion shown in equation 4.

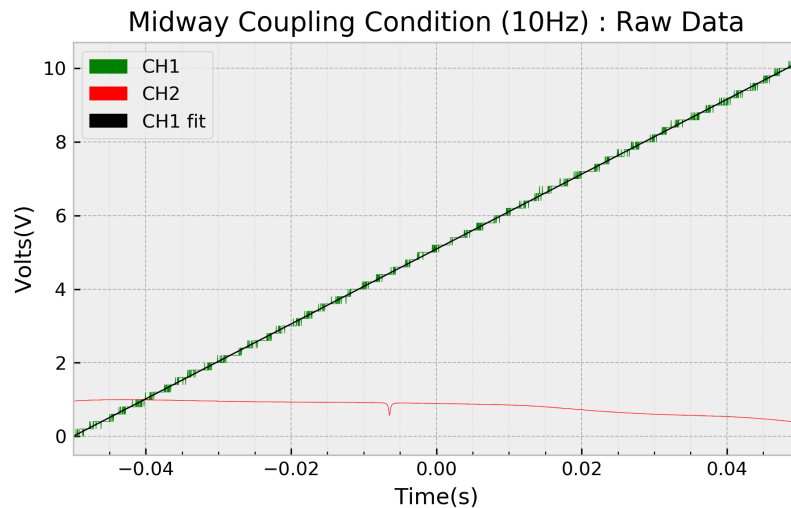
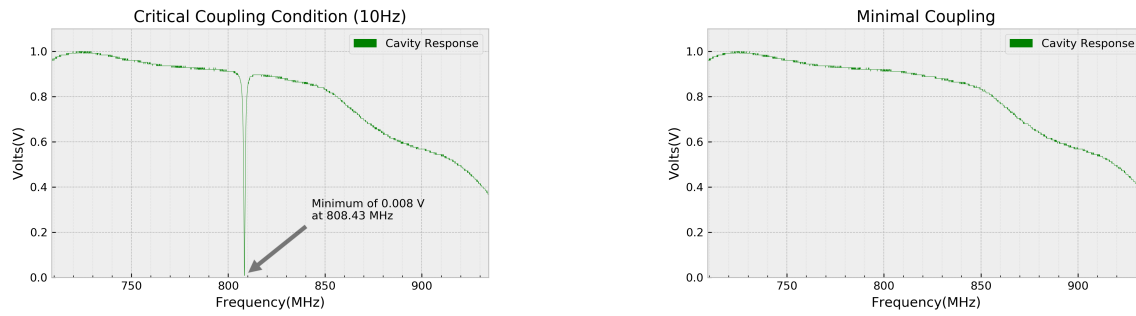


Figure 3: Raw output of a cavity reflection, with an interpolated voltage scale shown

Once we have coupled the input loop into the cavity we can change the orientation to create different coupling conditions. Here we present a critical coupling condition in Figure 4a and then we present a minimal coupling condition 4b. We note that these occur at 90° away from one another when turning the input coupling loop.

We can see the resonance frequency which is labelled in Figure 4a, scanning at 10Hz however scanning at a slower rate of 2Hz and decreasing our time bounds to more closely view the resonance drop we can get a more accurate value for the resonance frequency. This is shown in Figure 5, and gives us a frequency value of 809.89 MHz with a voltage of 0.120 V. We note that the non-zero value of the peak implies that there are some imperfections which are involved with the reading, and the lower value found in Figure 4a implies that the hand set value is not exactly a critical coupling condition. However this will not alter the location of the peak, so it should not cause any further issues.



(a) Critical coupling condition for reflection of the cavity at a voltage ramp of 10Hz

(b) Minimal coupling condition for reflection of the cavity at a voltage ramp of 10Hz

Figure 4: Different coupling conditions leading to different reflections.

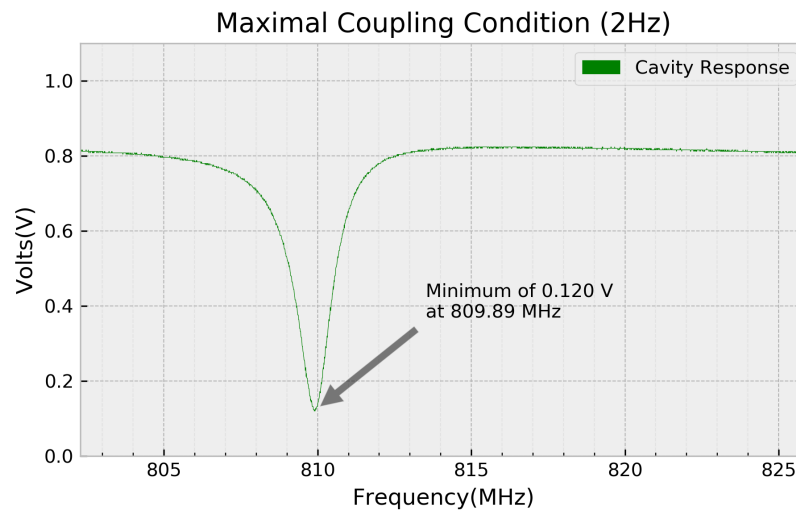
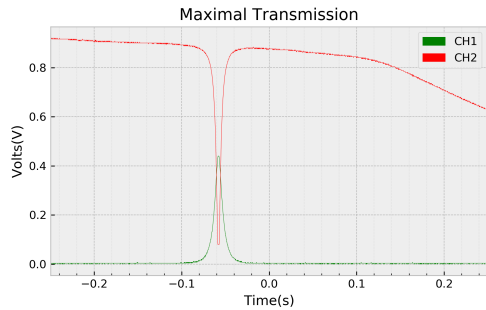


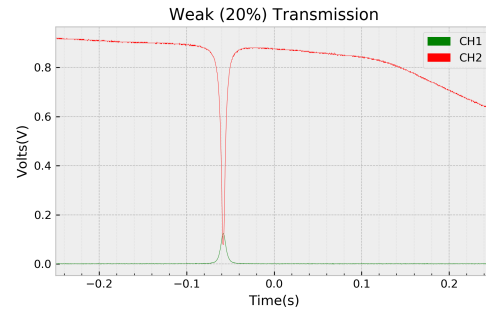
Figure 5: Critical coupling condition for reflection of the cavity at a voltage ramp of 2Hz

B.2 Cavity Transmission

We can then change the setup to allow for transmission. This is done according to the lab manual [2]. Therefore we then change the output loop to scan over the transmission conditions. This is shown in Figure 6. We can see that the received transmitted data is altered by the changes in the coupling loop's orientation, but nothing else is significantly altered.



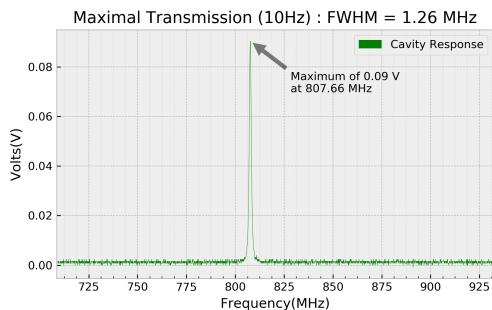
(a) Raw data of maximal transmission scan.



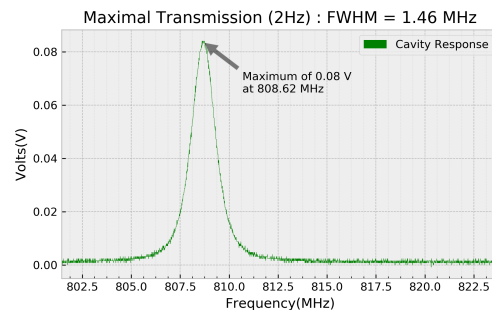
(b) Raw data of weak coupling transmission scan.

Figure 6: Transmission scopes, showing cavity transmission at different coupling conditions

We then take a value which is 20% of the maximal transmitted signal to create a weak coupling condition. This is completed as the impedance of the oscilloscope can affect resonance. This weak coupling condition is roughly found to occur at 15° away from minimal transmission settings. We will set the coupling loops in this configuration for the remainder of the lab.



(a) Transmission signal at 10Hz input



(b) Transmission signal at 2Hz input

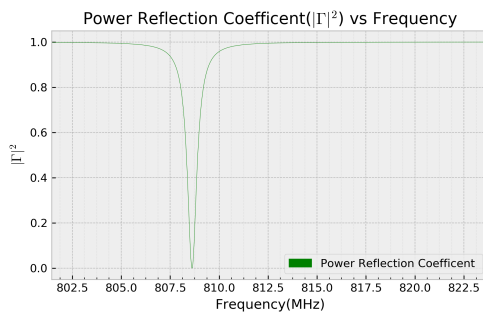
Figure 7: Transmission signals at different repeat frequencies

We determine that the resonant frequency of the cavity is at 808.62 MHz from our more accurate graph at 2Hz, but we then see that the less accurate graph gives us a different value of resonance at 807.66 MHz, it is expected that we will have some difference as the cavity must take some period of time to achieve equilibrium, but the 10Hz value is suggested in the lab manual does not appear to be appropriate. We will state that the resonant frequency is 808 ± 2 MHz due to this difference in found resonant frequencies. We also find that the Full Width at Half Maximum (FWHM) of the signal is $1.5 \pm .2$ MHz, which was found through a custom script [1].

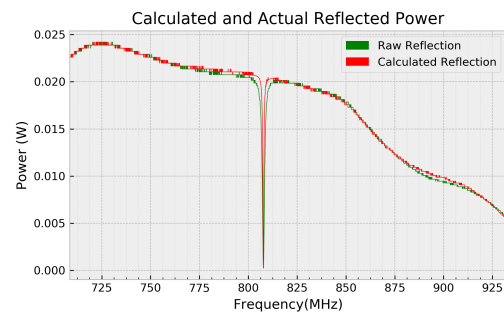
From this we are able to calculate a Quality factor of 560 ± 40 which is calculated based off of the formula of $Q_\ell = \frac{\text{Resonant Frequency}}{FWHM}$. We note that the error in this calculation is quite large as the FWHM from our two scans do not agree, hence we state a large error. From this measure we are able to calculate a graph of $|\Gamma|^2$ vs frequency, which we show in Figure 8. This is calculated through the formula, where ω_o is the angular resonance frequency and $\Delta\omega$ is the angular distance from resonance.

$$|\Gamma|^2 = 1 - \frac{1}{1 - 4Q_\ell^2(\Delta\omega/\omega_o)^2} \quad (9)$$

From this $|\Gamma|^2$ we are able to calculate an expected or theoretical graph of the reflected power, first we note that the power relationship must be computed which is based off of the experimentally found voltage to power relation given to us in the lab manual [2]. After translating voltage we use the P_{input} as the inputted signal without reflection, which can be seen in the minimal coupling condition (Figure 4b). We show this calculated reflection power in Figure 8b, we note that the found resonance frequency of this graph is used in calculations rather than the 2Hz signal to better compare the graphs.



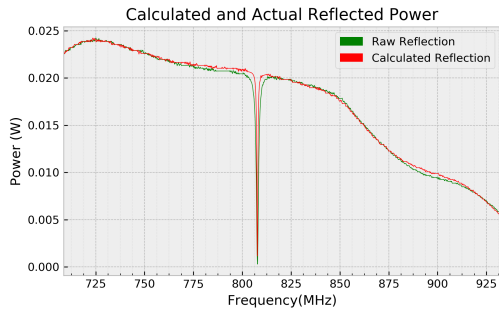
(a) Calculated $|\Gamma|^2$ values vs frequency



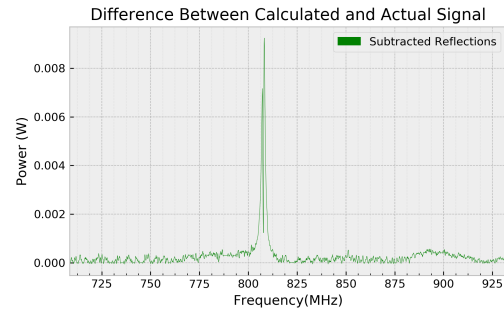
(b) Calculated vs Actual resonance conditions to compare the quality of the fit, original data also shown in Figure 4a

Figure 8: Gamma calculation and comparison graphs

We also note that in Figure 8b we see a large amount of noise, to reduce this and see the underlying signal we use a Savitzky-Golay filter. The results of this are shown in Figure 9a. We see that generally the shape of the signal is preserved, with a couple of deviations, but the largest effect is seen at resonance. This deviation at resonance shows that the resonance peak is broader than the calculated resonance. This effect is emphasized in Figure 9b where we subtract the calculated signal from the raw signal and filter it to show the difference in signal. We see from this that it is 10MHz broader than expected and the difference accounts for approximately 30% of the signal at peak. The broadness of the signal is guessed to be a lack of uniformity of the signal. Non-idealities of the cavity will shift the resonance slightly, which will cause a smearing of the inputted frequencies.



(a) Filtered calculated and raw signals



(b) Calculated and raw signal subtracted to see difference

Figure 9: Analysis graphs for the reflected signal

As a point of interest we also note that this resonance frequency found of 808 ± 2 MHz is simply the first resonant frequency. We will note that the reason that the cavity is called a quarter wavelength is that the resonance occurs when the condition $L = \lambda/4$ is satisfied. However we note that from theory we can derive this to be $L/n = \lambda/4$ where n is any positive integer, meaning our first frequency of $n=1 \implies f(n=1) = f_1 = 4/L \implies f(n=2) = 4n/L = 2 \times f_1$. Therefore these higher resonance frequencies will then be n multiples of our found resonant frequency, ex. $n=2 \implies 1616$ MHz. This also makes sense physically, or in relation to waves in any other medium as any n multiple of a resonant frequency will simply create $n-1$ nodes in the medium and create resonance at the higher frequency.

C Cavity Reflection Coefficient Phase

In this portion of the lab we measured the sin of the phase variations of the reflected signal from the cavity. This was done by multiplying the sampled signal before it hits the cavity.

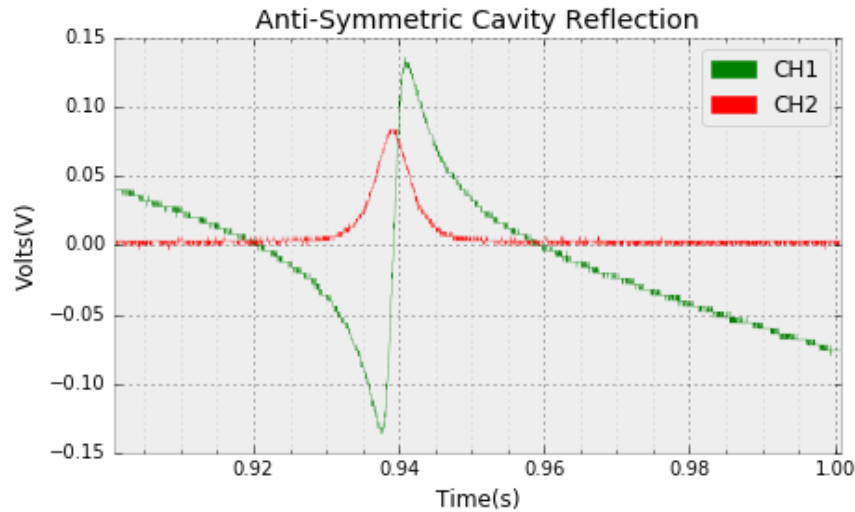


Figure 10: Raw output of an anti-symmetric cavity reflection, with an interpolated voltage scale shown

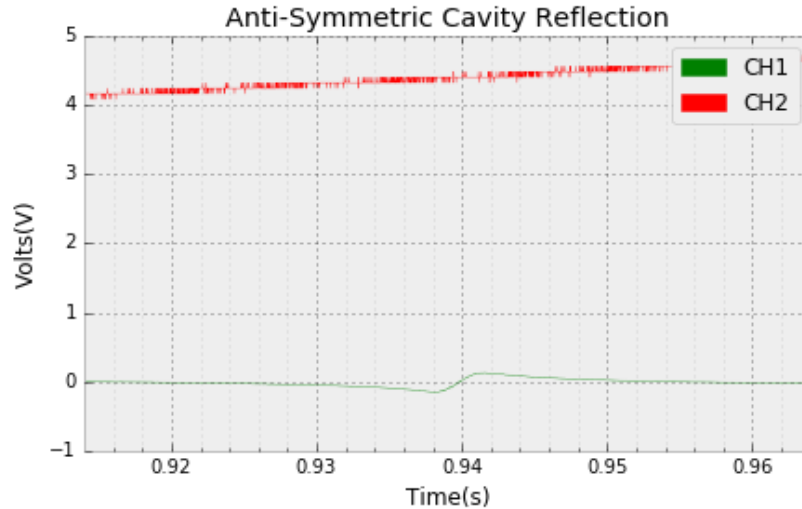


Figure 11: Raw output of an anti-symmetric cavity reflection, with an interpolated voltage scale shown

Above we have the outputted graphs we get when scanning the VCO frequency over the resonance. The dispersion-like variation in the output of the mixer is visible.

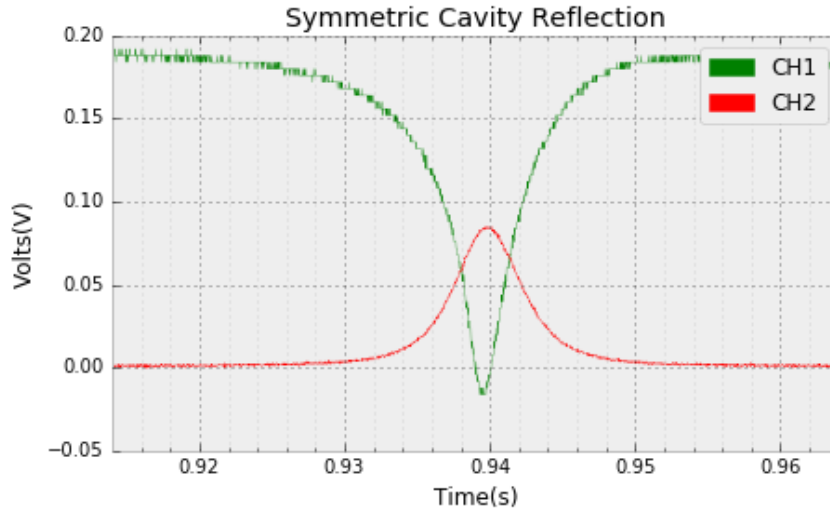


Figure 12: Raw output of a symmetric cavity reflection, with an interpolated voltage scale shown

This is the graph of the symmetric cavity reflection. The difference between this length and the length of the trombone when looking at the Anti-Symmetric reflection is that this length is shorter than that one.

C.1 Additional Questions

1.

$$|\Gamma|^2 = 1 - \frac{1}{1 + 4Q_\ell^2(\Delta\omega/\omega_o)^2} \quad (10)$$

$$\frac{1}{Z} = \frac{1}{R} + jwC + \frac{1}{jwL} \quad (11)$$

$$\Gamma = \frac{Z - Z_0}{Z + Z_0} \quad (12)$$

$$|\Gamma|^2 = \left| \frac{Z - Z_0}{Z + Z_0} \right|^2 \quad (13)$$

$$\left| \frac{Z - Z_0}{Z + Z_0} \right|^2 = 1 - \frac{1}{1 + 4Q_\ell^2(\Delta\omega/\omega_o)^2} \quad (14)$$

$$\frac{1}{1 + 4Q_\ell^2(\Delta\omega/\omega_o)^2} = 1 - \frac{(Z - Z_0)^2}{(Z + Z_0)^2} \quad (15)$$

$$\frac{1}{1 + 4Q_\ell^2(\Delta\omega/\omega_o)^2} = 1 - \frac{(Z - Z_0)^2}{(Z + Z_0)^2} \quad (16)$$

$$\frac{1}{1 + 4Q_\ell^2(\Delta\omega/\omega_o)^2} = \frac{4Z_0Z}{(Z + Z_0)^2} \quad (17)$$

$$1 + 4Q_\ell^2(\Delta\omega/\omega_o)^2 = \frac{(Z + Z_0)^2}{4Z_0Z} \quad (18)$$

$$4Q_\ell^2(\Delta\omega/\omega_o)^2 = \frac{(Z + Z_0)^2}{4Z_0Z} - 1 \quad (19)$$

$$4Q_\ell^2(\Delta\omega/\omega_o)^2 = \frac{(Z - Z_0)^2}{4Z_0Z} \quad (20)$$

$$Q_\ell^2 = \frac{1}{4} \frac{(Z - Z_0)^2}{4Z_0Z} (\omega_o/\Delta\omega)^2 \quad (21)$$

$$Q_\ell = \frac{1}{2} \frac{(Z - Z_0)}{2\sqrt{Z_0Z}} (\omega_o/\Delta\omega) \quad (22)$$

Where $\Delta\omega = \frac{R}{L}$

$$Q_\ell = \frac{w_0 L}{R} \frac{1}{4} \frac{(Z - Z_0)}{\sqrt{Z_0Z}} \quad (23)$$

Here we have that $Z = \frac{1}{\frac{1}{R} + j\omega C + \frac{1}{j\omega L}}$

2. From above we see that

$$\Gamma = \frac{Z - Z_0}{Z + Z_0} \quad (24)$$

With $Z = \frac{1}{\frac{1}{R} + j\omega C + \frac{1}{j\omega L}}$ from this we will then be able to compare the mixer output and the imaginary part of the Gamma function.

D Stabilization of the VCO to the Cavity

This part is originally for the measurement of the linear thermal expansion coefficients of copper and aluminum. To measure thermal expansion coefficient, the temperature of copper cavity or aluminum rod should be varied, and corresponding resonance frequency need to be measured. Therefore, V_{tune} and corresponding VCO frequency must be locked to the resonance frequency of copper cavity or aluminum rod, so the frequency can be measured with only varying the temperature of the cavity or rod. To lock the V_{tune} , integrator feedback control is inserted between the mixer and V_{tune} port of the VCO. By turning off the lock, modifying gain with the cream coloured small driver, and turning on the lock again, the immediate jump to the resonant frequency from the modified frequency has observed. Since the temperature of the cavity could not be modified because of the equipment problem, the only frequency lock observed is lock to the resonance frequency used in previous sections, with the same room temperature.

References

- [1] 460_e10_rf source code. https://github.com/tokiyoshi/460_E10_RF, 2017.
- [2] James D. D. Martin and Instructors. Radio-frequency electronics. *PHYS 360/460 Lab Manuals*, pages 1–4, 2014/05/13. URL: http://science.uwaterloo.ca/~jddmarti/teaching/phys360_460/rf_exp/rf_exp.html.