# Quantum computation as a new scheme for problem-solving

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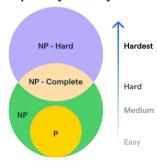






- P-class problems: decision problems in polynomial resources.
- NP-class problems: decision problems whose proof is verifiable in polynomial resources.
- NP-Complete: "Generalizations" of NP problems.
- NP-Hard: which are at least as hard as the hardest problems in NP

# Computational Complexity Theory





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Consider the integer factorization problem... It is suspected to be neither  ${\bf P}$  nor  ${\bf NP}$ -complete.



In: Proceedings, 35th Annual Symposium on Foundations of Computer Science, Santa Fe, NM, November 20–22, 1994, IEEE Computer Society Press, pp. 124–134.

#### Algorithms for Quantum Computation: Discrete Logarithms and Factoring

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### **Brief History of Quantum Mechanics**



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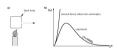


Figure: Classical Mechanics fails to describe blackbody radiation

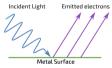


Figure: Electrons are only ejected when there is *enough* light hitting the surface, contrary to classical predictions.

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Figure: Max-Planck (1858-1947)



Figure: Einstein in 1905

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From then on, physicists expanded QM by exploring

#### **Motivation**



• What would happen if storing and computation devices became ever so small that quantum effects start to become noticeable?



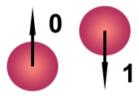


Figure: The classical bit



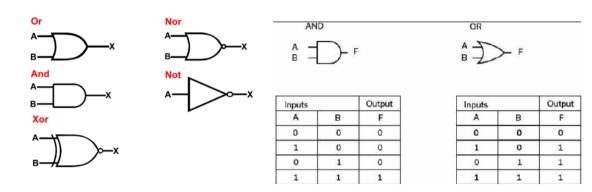


Figure: Caption



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$$\wedge:\{0,1\}^2\longrightarrow\{0,1\}$$
 with 
$$\land (0,0)=0$$
 
$$\land (0,1)=0$$
 
$$\land (1,0)=0$$
 
$$\land (1,1)=1$$

Figure: Caption



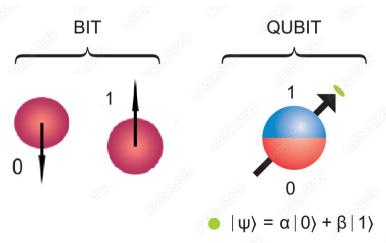
An full circuit can then be thought of as a general function

$$F: \{0,1\}^n \to \{0,1\}^m$$
.

#### The Quantum-Bit



Let us present the Quantum Bit or (QuBit)



## **The Qubit**



A question arises...

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• this leads to the measurement process.





Figure: System of qubits at initial time t = 0.



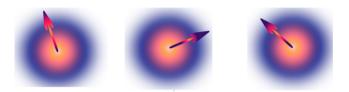


Figure: System of qubits at initial time t = 0.



Figure: System of qubits measured at some later time  $t>t_0$ .



If we repeat the experiment, we might get...

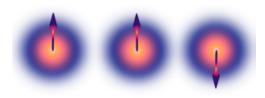


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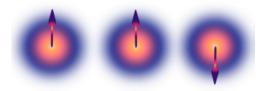


Figure: We might call this state  $|\Psi\rangle=|\uparrow\uparrow\downarrow\rangle$ 

We can build a probability distribution  $P_i(a = 0_i, 1_i) = |\langle a|\Psi\rangle|^2$ .

#### **Single Qubit Gates**



$$X = \sigma_x = ext{NOT} = egin{bmatrix} 0 & 1 \ 1 & 0 \end{bmatrix}, \ Y = \sigma_y = egin{bmatrix} 0 & -i \ i & 0 \end{bmatrix}, \ Z = \sigma_z = egin{bmatrix} 1 & 0 \ 0 & -1 \end{bmatrix}.$$

#### **Two-Qubit Gates**



Of special interest is the CNOT gate, given by

$$\text{CNOT} = \left(\begin{array}{cccc} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{array}\right)$$



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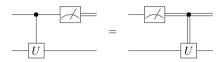


Figure: Deferred measurement



Now, we will study an application: Superdense coding

# Secure Quantum Comunications: Superdense **Coding**



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### **Grover's Search Algorithms**



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Consider a list of N items, with an item  $\omega$  of interest...

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- In the worst-case scenario, N (unstructured list).

With quantum computation, we need only  $\mathcal{O}(\sqrt{N})$  operations  $\Rightarrow$  Grover's search Algorithm

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1. Initialization.

Repeat until  $|\omega\rangle$  is found:

- 2. We apply a quantum operation called *oracle*.
- 3. We apply an additional reflection.

End(Repeat)

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There is no knowledge about the state of the system  $\rightarrow |s\rangle = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle$ .

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We initialize via the operation  $|S\rangle = H^{\otimes n}$ 

# **Grove: Oracle (first reflection)**



Let the oracle be:

$$U_{\omega} |x\rangle = \left\{ egin{array}{l} |x
angle & ext{if } x 
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Then, the new state is

$$|\psi_1\rangle = U_\omega |S\rangle = U_\omega H^\otimes |x\rangle$$
.

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Then we apply the reflection  $U_s=2\left|S\right>\left< S\right|-1$  onto the previous state,

$$|\psi_2\rangle = \mathcal{U}_{\mathrm{S}} \, |\psi_1\rangle = \mathcal{U}_{\mathrm{S}} \mathcal{U}_{\omega} \mathcal{H}^{\otimes} \, |\mathbf{x}\rangle \ .$$

# **Grover: Repetition of these steps**



Then, we repeat these steps s.t. the whole operation now reads

$$|\psi_i\rangle = (U_{\mathcal{S}}U_{\omega)} |\psi_{i-1}\rangle,$$

or, more succintly,

$$|\psi_{r(N)}\rangle = (U_{S}U_{\omega})^{r(N)}H^{\otimes}|x\rangle,$$

where  $r(N) \approx \frac{\pi}{4} \sqrt{N}$ , with N the total number of states.

### Example for N=4:



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• We want to study a protocol with  $\mathit{N}=4$  states, which we will label  $\{\ket{00},\ket{01},\ket{10},\ket{11}\}.$ 



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- Target state:  $|\omega\rangle = |11\rangle$ .
- For N = 4,  $r(N) \approx 1.57 \approx 2$ .



from qiskit import IBMQ, Aer, transpile, execute from qiskit import QuantumCircuit, ClassicalRegister, QuantumRegister from qiskit.providers.ibmq import least\_busy



```
n = 2
grover_circuit = QuantumCircuit(n)
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grover_circuit = QuantumCircuit(n)

def initialize_s(qc, qubits):
    for q in qubits:
        qc.h(q)
    return qc
```



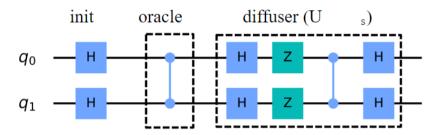
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```
grover_circuit = initalize_s(grover_circuit, [0,1])
grover_circuit.h([0,1])
grover_circuit.z([0,1])
grover_circuit.cz(0,1)
grover_circuit.h([0,1])
grover_circuit.draw()
```







```
grover_circuit.measure_all()
qasm_sim = Aer.get_backend('qasm_simulator')
result = qasm_sim.run(grover_circuit).result()
counts = result.get_counts()
plot_histogram(counts)
```



