



Computational Science:  
Introduction to Finite-Difference Time-Domain

# Formulation of Two-Dimensional FDTD Without PML

## Lecture Outline

- Code development sequence for 2D
- Maxwell's equations
- Finite-difference approximations
- Reduction to two dimensions
- Update equations
- Boundary Conditions
- Revised FDTD Algorithm for 2D Simulations

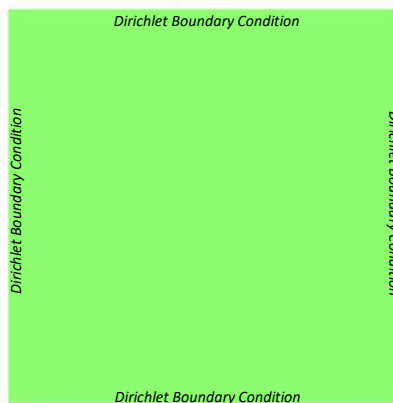
# Code Development Sequence

Slide 3

## Step 1 – Basic Update Equations

The basic update equations are implemented along with simple Dirichlet boundary conditions.

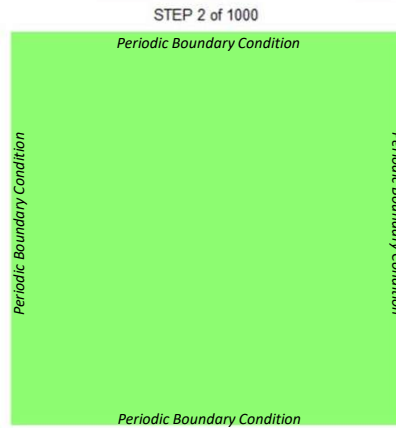
STEP 2 of 1000



Slide 4

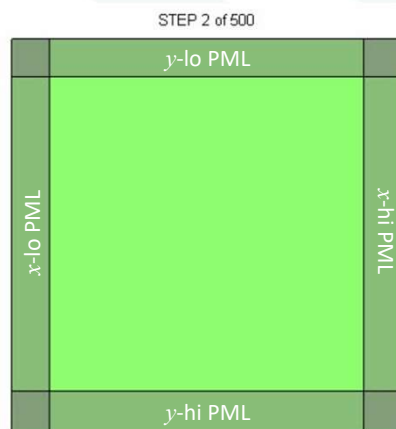
## Step 2 – Incorporate Periodic Boundaries

Periodic boundary conditions are incorporated so that a wave leaving the grid reenters the grid at the other side.



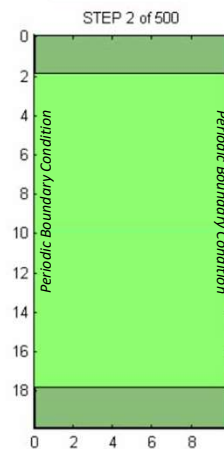
## Step 3 – Incorporate a PML

The perfectly matched layer (PML) absorbing boundary condition is incorporated to absorb outgoing waves.



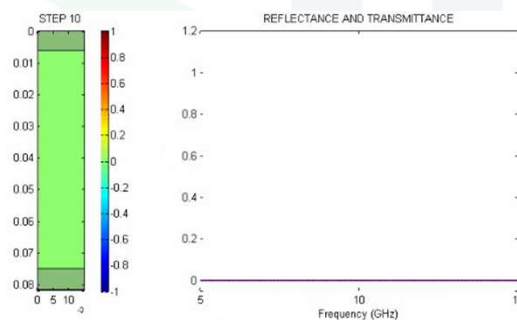
## Step 4 – Total-Field/Scattered-Field

Most periodic electromagnetic devices are modeled by using periodic boundaries for the horizontal axis and a PML for the vertical axis. We then implement TF/SF at the vertical center of the grid for testing.



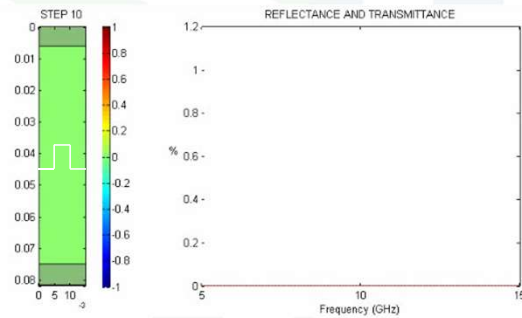
## Step 5 – Calculate TRN, REF, and CON

We move the TF/SF interface to a unit cell or two outside of the top PML. We include code to calculate Fourier transforms and to calculate transmittance, reflectance, and conservation of power.



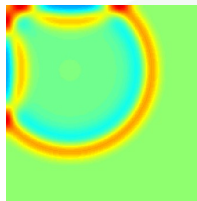
## Step 6 – Model a Device to Benchmark

We build a device on the grid that has a known solution. We run the simulation and duplicate the known results to benchmark our new code.

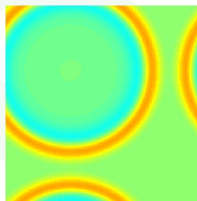


## Summary of Code Development Sequence

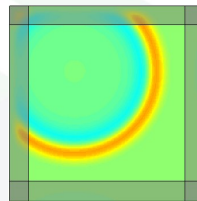
Step 1 – Basic Update  
+ Dirichlet



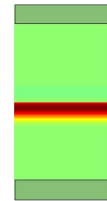
Step 2 – Basic Update  
+ Periodic BC



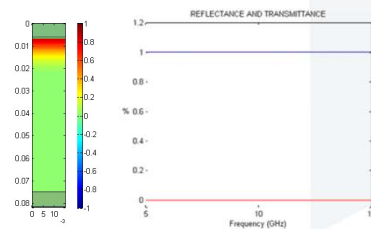
Step 3 – Add PML



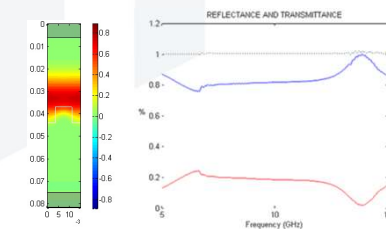
Step 4 – TF/SF



Step 5 – Calculate Response



Step 6 – Add a Device and Benchmark



# Maxwell's Equations

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## Maxwell's Equations

Start with Maxwell's equations in the following form:

$$\nabla \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \cdot \vec{D} = \rho_v$$

$$\vec{D}(t) = \epsilon_0 \epsilon_r \vec{E}(t)$$

$$\vec{B}(t) = \mu_0 \mu_r \vec{H}(t)$$

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## Normalize the Electric Fields

Now we will adopt the more conventional approach in FDTD and normalize the electric field according to:

$$\vec{\tilde{E}} = \sqrt{\frac{\epsilon_0}{\mu_0}} \vec{E} = \frac{1}{\eta_0} \vec{E}$$

To be consistent, other parameters related to the electric field must also be normalized.

$$\vec{\tilde{D}} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \vec{D} = c_0 \vec{D}$$

$$\vec{\tilde{P}} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \vec{P} = c_0 \vec{P}$$

$$\vec{\tilde{J}} = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \vec{J} = c_0 \vec{J}$$

$\vec{\tilde{J}}$  and  $\vec{\tilde{P}}$  won't be used in this course.

## Normalized Maxwell's Equations

Using the normalized fields, Maxwell's equations become

$$\nabla \times \vec{\tilde{E}} = -\frac{\mu_r}{c_0} \frac{\partial \vec{\tilde{H}}}{\partial t}$$

$$\nabla \times \vec{\tilde{H}} = \frac{1}{c_0} \frac{\partial \vec{\tilde{D}}}{\partial t}$$

These equations are independent of  $\epsilon_r$ .

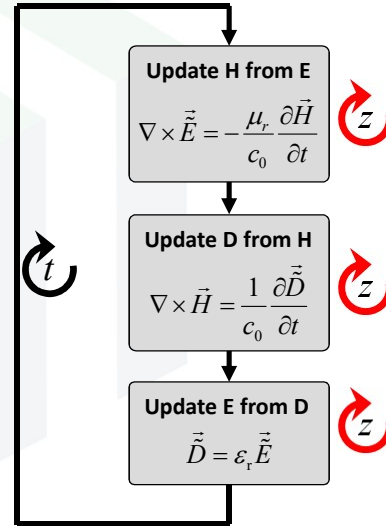
Update equations can be derived from these expressions without complicating them with what might need to be done to model  $\epsilon_r$ .

$$\vec{\tilde{D}} = \epsilon_r \vec{\tilde{E}}$$

This is a very simple equation that makes it much easier to model more sophisticated dielectrics that may be anisotropic, nonlinear, dispersive, all of the above, or something else altogether.

## HDE Algorithm

To make FDTD more modular, modify the steps to what is called the *HDE algorithm*.



## Expand Maxwell's Equations

$$\begin{aligned}
 \nabla \times \vec{E} &= -\frac{\mu_r}{c_0} \frac{\partial \vec{H}}{\partial t} \rightarrow \begin{aligned} \frac{\partial \tilde{E}_z}{\partial y} - \frac{\partial \tilde{E}_y}{\partial z} &= -\frac{1}{c_0} \left( \mu_{xx} \frac{\partial H_x}{\partial t} + \mu_{xy} \frac{\partial H_y}{\partial t} + \mu_{xz} \frac{\partial H_z}{\partial t} \right) \\ \frac{\partial \tilde{E}_x}{\partial z} - \frac{\partial \tilde{E}_z}{\partial x} &= -\frac{1}{c_0} \left( \mu_{yx} \frac{\partial H_x}{\partial t} + \mu_{yy} \frac{\partial H_y}{\partial t} + \mu_{yz} \frac{\partial H_z}{\partial t} \right) \\ \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} &= -\frac{1}{c_0} \left( \mu_{zx} \frac{\partial H_x}{\partial t} + \mu_{zy} \frac{\partial H_y}{\partial t} + \mu_{zz} \frac{\partial H_z}{\partial t} \right) \end{aligned} \\
 \nabla \times \vec{H} &= \frac{1}{c_0} \frac{\partial \vec{D}}{\partial t} \rightarrow \begin{aligned} \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} \\ \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} \\ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} \end{aligned} \\
 \vec{D} &= \epsilon_r \vec{E} \rightarrow \begin{aligned} \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x + \epsilon_{xy} \tilde{E}_y + \epsilon_{xz} \tilde{E}_z \\ \tilde{D}_y &= \epsilon_{yx} \tilde{E}_x + \epsilon_{yy} \tilde{E}_y + \epsilon_{yz} \tilde{E}_z \\ \tilde{D}_z &= \epsilon_{zx} \tilde{E}_x + \epsilon_{zy} \tilde{E}_y + \epsilon_{zz} \tilde{E}_z \end{aligned}
 \end{aligned}$$



## Assume Only Diagonal Tensors

$$\begin{aligned}
 \nabla \times \vec{E} &= -\frac{\mu_r}{c_0} \frac{\partial \vec{H}}{\partial t} \rightarrow \begin{aligned} \frac{\partial \tilde{E}_z}{\partial y} - \frac{\partial \tilde{E}_y}{\partial z} &= -\frac{1}{c_0} \left( \mu_{xx} \frac{\partial H_x}{\partial t} + \cancel{\mu_{xy} \frac{\partial H_y}{\partial t}} + \cancel{\mu_{xz} \frac{\partial H_z}{\partial t}} \right) \\ \frac{\partial \tilde{E}_x}{\partial z} - \frac{\partial \tilde{E}_z}{\partial x} &= -\frac{1}{c_0} \left( \cancel{\mu_{yx} \frac{\partial H_x}{\partial t}} + \mu_{yy} \frac{\partial H_y}{\partial t} + \cancel{\mu_{yz} \frac{\partial H_z}{\partial t}} \right) \\ \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} &= -\frac{1}{c_0} \left( \cancel{\mu_{zx} \frac{\partial H_x}{\partial t}} + \cancel{\mu_{zy} \frac{\partial H_y}{\partial t}} + \mu_{zz} \frac{\partial H_z}{\partial t} \right) \end{aligned} \\
 \nabla \times \vec{H} &= \frac{1}{c_0} \frac{\partial \vec{D}}{\partial t} \rightarrow \begin{aligned} \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} \\ \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} \\ \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} \end{aligned} \\
 \vec{D} &= \epsilon_r \vec{E} \rightarrow \begin{aligned} \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x + \cancel{\epsilon_{xy} \tilde{E}_y} + \cancel{\epsilon_{xz} \tilde{E}_z} \\ \tilde{D}_y &= \cancel{\epsilon_{yx} \tilde{E}_x} + \epsilon_{yy} \tilde{E}_y + \cancel{\epsilon_{yz} \tilde{E}_z} \\ \tilde{D}_z &= \cancel{\epsilon_{zx} \tilde{E}_x} + \cancel{\epsilon_{zy} \tilde{E}_y} + \epsilon_{zz} \tilde{E}_z \end{aligned}
 \end{aligned}$$

## Governing Equations

These are the primary governing equations from which the 2D and 3D FDTD methods will be formulated.

$$\begin{aligned}
 \frac{\partial \tilde{E}_z}{\partial y} - \frac{\partial \tilde{E}_y}{\partial z} &= -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t} & \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} \\
 \frac{\partial \tilde{E}_x}{\partial z} - \frac{\partial \tilde{E}_z}{\partial x} &= -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t} & \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} \\
 \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} &= -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} & \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t}
 \end{aligned}$$

$$\begin{aligned}
 \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x \\
 \tilde{D}_y &= \epsilon_{yy} \tilde{E}_y \\
 \tilde{D}_z &= \epsilon_{zz} \tilde{E}_z
 \end{aligned}$$

## Final Analytical Equations

It will be beneficial in the formulation and implementation to calculate the curl terms separately. Start by expressing Maxwell's equations in the following form.

$$\begin{aligned} C_x^E &= -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t} \\ C_y^E &= -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t} \\ C_z^E &= -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} \end{aligned}$$

$$\begin{aligned} C_x^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} \\ C_y^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} \\ C_z^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} \end{aligned}$$

$$\begin{aligned} \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x \\ \tilde{D}_y &= \epsilon_{yy} \tilde{E}_y \\ \tilde{D}_z &= \epsilon_{zz} \tilde{E}_z \end{aligned}$$

$$\begin{aligned} C_x^E &= \frac{\partial \tilde{E}_z}{\partial y} - \frac{\partial \tilde{E}_y}{\partial z} \\ C_y^E &= \frac{\partial \tilde{E}_x}{\partial z} - \frac{\partial \tilde{E}_z}{\partial x} \\ C_z^E &= \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} \end{aligned}$$

$$\begin{aligned} C_x^H &= \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \\ C_y^H &= \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} \\ C_z^H &= \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \end{aligned}$$

## Finite-Difference Approximations

## Finite-Difference Equations for $\vec{H}$

### Curl Terms

$$C_x^E = \frac{\partial \tilde{E}_z}{\partial y} - \frac{\partial \tilde{E}_y}{\partial z}$$

$$C_y^E = \frac{\partial \tilde{E}_x}{\partial z} - \frac{\partial \tilde{E}_z}{\partial x}$$

$$C_z^E = \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y}$$

$$C_x^E|_l^{i,j,k} = \frac{\tilde{E}_z|_l^{i,j,k+1} - \tilde{E}_z|_l^{i,j,k}}{\Delta y} - \frac{\tilde{E}_y|_l^{i,j,k+1} - \tilde{E}_y|_l^{i,j,k}}{\Delta z}$$

$$C_y^E|_l^{i,j,k} = \frac{\tilde{E}_x|_l^{i,j,k+1} - \tilde{E}_x|_l^{i,j,k}}{\Delta z} - \frac{\tilde{E}_z|_l^{i+1,j,k} - \tilde{E}_z|_l^{i,j,k}}{\Delta x}$$

$$C_z^E|_l^{i,j,k} = \frac{\tilde{E}_y|_l^{i+1,j,k} - \tilde{E}_y|_l^{i,j,k}}{\Delta x} - \frac{\tilde{E}_x|_l^{i,j,k+1} - \tilde{E}_x|_l^{i,j,k}}{\Delta y}$$

### Maxwell's Equations

$$C_x^E = -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t}$$

$$C_y^E = -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t}$$

$$C_z^E = -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t}$$

$$C_x^E|_l^{i,j,k} = -\frac{\mu_{xx}|_l^{i,j,k}}{c_0} \frac{H_x|_{l+\frac{\Delta x}{2}}^{i,j,k} - H_x|_{l-\frac{\Delta x}{2}}^{i,j,k}}{\Delta t}$$

$$C_y^E|_l^{i,j,k} = -\frac{\mu_{yy}|_l^{i,j,k}}{c_0} \frac{H_y|_{l+\frac{\Delta y}{2}}^{i,j,k} - H_y|_{l-\frac{\Delta y}{2}}^{i,j,k}}{\Delta t}$$

$$C_z^E|_l^{i,j,k} = -\frac{\mu_{zz}|_l^{i,j,k}}{c_0} \frac{H_z|_{l+\frac{\Delta z}{2}}^{i,j,k} - H_z|_{l-\frac{\Delta z}{2}}^{i,j,k}}{\Delta t}$$

## Finite-Difference Equations for $\vec{D}$

### Curl Terms

$$C_x^H = \frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z}$$

$$C_y^H = \frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x}$$

$$C_z^H = \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y}$$

$$C_x^H|_{l+\frac{\Delta x}{2}}^{i,j,k} = \frac{H_z|_{l+\frac{\Delta x}{2}}^{i,j,k} - H_z|_{l+\frac{\Delta x}{2}}^{i,j-1,k}}{\Delta y} - \frac{H_y|_{l+\frac{\Delta x}{2}}^{i,j,k} - H_y|_{l+\frac{\Delta x}{2}}^{i,j,k-1}}{\Delta z}$$

$$C_y^H|_{l+\frac{\Delta y}{2}}^{i,j,k} = \frac{H_x|_{l+\frac{\Delta y}{2}}^{i,j,k} - H_x|_{l+\frac{\Delta y}{2}}^{i,j,k-1}}{\Delta z} - \frac{H_z|_{l+\frac{\Delta y}{2}}^{i,j,k} - H_z|_{l+\frac{\Delta y}{2}}^{i-1,j,k}}{\Delta x}$$

$$C_z^H|_{l+\frac{\Delta z}{2}}^{i,j,k} = \frac{H_y|_{l+\frac{\Delta z}{2}}^{i,j,k} - H_y|_{l+\frac{\Delta z}{2}}^{i,j,k-1}}{\Delta x} - \frac{H_x|_{l+\frac{\Delta z}{2}}^{i,j,k} - H_x|_{l+\frac{\Delta z}{2}}^{i,j-1,k}}{\Delta y}$$

### Maxwell's Equations

$$C_x^H = \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t}$$

$$C_y^H = \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t}$$

$$C_z^H = \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t}$$

$$C_x^H|_{l+\frac{\Delta x}{2}}^{i,j,k} = \frac{1}{c_0} \frac{\tilde{D}_x|_{l+\Delta t}^{i,j,k} - \tilde{D}_x|_l^{i,j,k}}{\Delta t}$$

$$C_y^H|_{l+\frac{\Delta y}{2}}^{i,j,k} = \frac{1}{c_0} \frac{\tilde{D}_y|_{l+\Delta t}^{i,j,k} - \tilde{D}_y|_l^{i,j,k}}{\Delta t}$$

$$C_z^H|_{l+\frac{\Delta z}{2}}^{i,j,k} = \frac{1}{c_0} \frac{\tilde{D}_z|_{l+\Delta t}^{i,j,k} - \tilde{D}_z|_l^{i,j,k}}{\Delta t}$$

## Finite-Difference Equations for $\vec{\tilde{E}}$

$$\tilde{D}_x = \epsilon_{xx} \tilde{E}_x$$

$$\tilde{D}_y = \epsilon_{yy} \tilde{E}_y$$

$$\tilde{D}_z = \epsilon_{zz} \tilde{E}_z$$



$$\tilde{D}_x \Big|_t^{i,j,k} = \left( \epsilon_{xx} \Big|_t^{i,j,k} \right) \tilde{E}_x \Big|_t^{i,j,k}$$

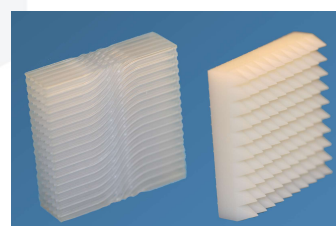
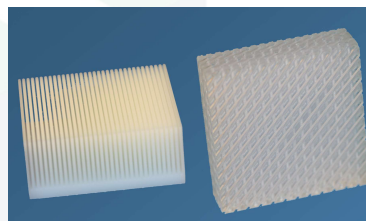
$$\tilde{D}_y \Big|_t^{i,j,k} = \left( \epsilon_{yy} \Big|_t^{i,j,k} \right) \tilde{E}_y \Big|_t^{i,j,k}$$

$$\tilde{D}_z \Big|_t^{i,j,k} = \left( \epsilon_{zz} \Big|_t^{i,j,k} \right) \tilde{E}_z \Big|_t^{i,j,k}$$

## Reduction to Two Dimensions

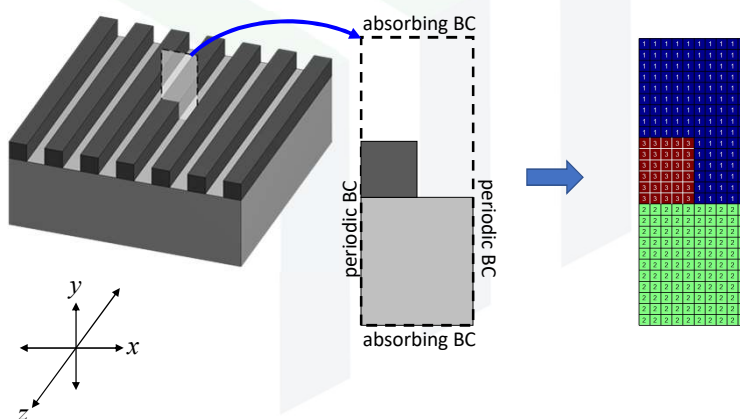
## Real Electromagnetic Devices

All physical devices are three-dimensional.



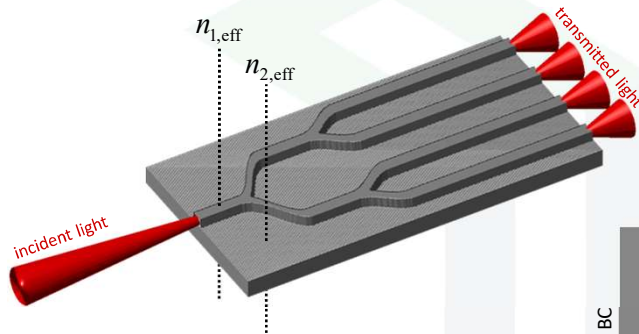
## 3D $\rightarrow$ 2D (Exact)

Sometimes it is possible to describe a physical device using just two dimensions. Doing so dramatically reduces the numerical complexity of the problem and is ALWAYS GOOD PRACTICE.



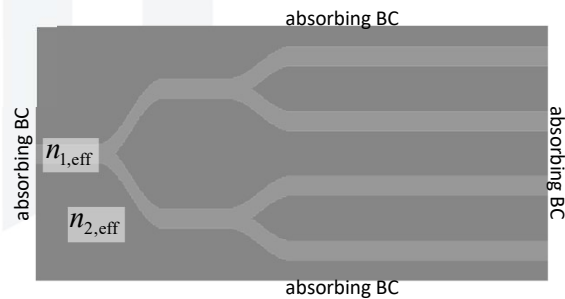
## 3D → 2D (Approximate)

Many times it is possible to approximate a 3D device in two dimensions. It is very good practice to at least perform the initial simulations in 2D and only moving to 3D to verify the final design.



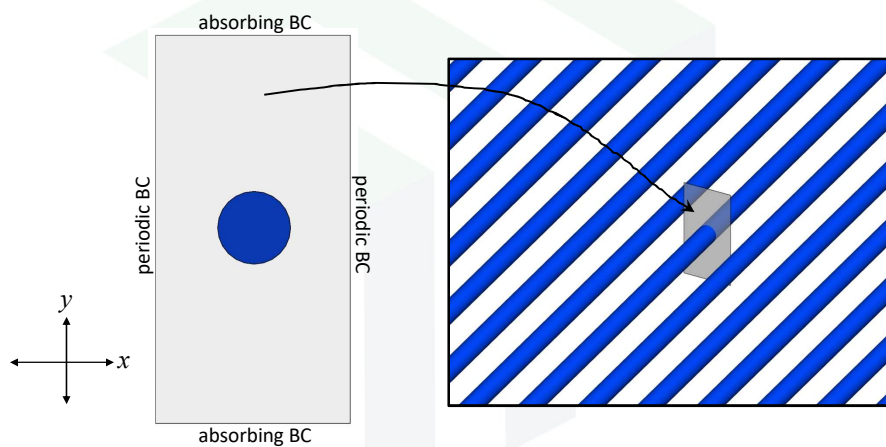
Effective indices are best computed by modeling the vertical cross section as a slab waveguide.

A simple average index can also produce good results.



## 2D Grids are Infinite in the 3<sup>rd</sup> Dimension

Anything represented on a 2D grid, is actually a device that is of infinite extent along the 3<sup>rd</sup> dimension.



Assuming the  $z$  direction is uniform and of infinite extent, then the field is also uniform and of infinite extent in the  $z$  direction.

$$\frac{\partial}{\partial z} = 0$$

## Assume Uniform in z Direction

For 2D devices,  $\partial/\partial z=0$  and Maxwell's equations reduce to

$$\begin{aligned}
 C_x^E &= \frac{\partial \tilde{E}_z}{\partial y} - \cancel{\frac{\partial \tilde{E}_y}{\partial z}} & C_x^E &= -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t} & C_x^E|_{l+\frac{\Delta x}{2}} &= \frac{\tilde{E}_z|_{l+\frac{\Delta x}{2}} - \tilde{E}_z|_{l-\frac{\Delta x}{2}}}{\Delta y} - \cancel{\frac{\tilde{E}_y|_{l+\frac{\Delta x}{2}} - \tilde{E}_y|_{l-\frac{\Delta x}{2}}}{\Delta z}} & C_x^E|_{l+\frac{\Delta x}{2}} &= -\frac{\mu_{xx}}{c_0} \frac{H_x|_{l+\frac{\Delta x}{2}} - H_x|_{l-\frac{\Delta x}{2}}}{\Delta t} \\
 C_y^E &= \cancel{\frac{\partial \tilde{E}_x}{\partial z}} - \frac{\partial \tilde{E}_z}{\partial x} & C_y^E &= -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t} & C_y^E|_{l+\frac{\Delta y}{2}} &= \cancel{\frac{\tilde{E}_x|_{l+\frac{\Delta y}{2}} - \tilde{E}_x|_{l-\frac{\Delta y}{2}}}{\Delta z}} - \frac{\tilde{E}_z|_{l+\frac{\Delta y}{2}} - \tilde{E}_z|_{l-\frac{\Delta y}{2}}}{\Delta x} & C_y^E|_{l+\frac{\Delta y}{2}} &= -\frac{\mu_{yy}}{c_0} \frac{H_y|_{l+\frac{\Delta y}{2}} - H_y|_{l-\frac{\Delta y}{2}}}{\Delta t} \\
 C_z^E &= \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} & C_z^E &= -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} & C_z^E|_{l+\frac{\Delta z}{2}} &= \frac{\tilde{E}_y|_{l+\frac{\Delta z}{2}} - \tilde{E}_y|_{l-\frac{\Delta z}{2}}}{\Delta x} - \frac{\tilde{E}_x|_{l+\frac{\Delta z}{2}} - \tilde{E}_x|_{l-\frac{\Delta z}{2}}}{\Delta y} & C_z^E|_{l+\frac{\Delta z}{2}} &= -\frac{\mu_{zz}}{c_0} \frac{H_z|_{l+\frac{\Delta z}{2}} - H_z|_{l-\frac{\Delta z}{2}}}{\Delta t} \\
 \\ 
 C_x^H &= \frac{\partial H_z}{\partial y} - \cancel{\frac{\partial H_y}{\partial z}} & C_x^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} & C_x^H|_{l+\frac{\Delta x}{2}} &= \frac{H_z|_{l+\frac{\Delta x}{2}} - H_z|_{l-\frac{\Delta x}{2}}}{\Delta y} - \cancel{\frac{H_y|_{l+\frac{\Delta x}{2}} - H_y|_{l-\frac{\Delta x}{2}}}{\Delta z}} & C_x^H|_{l+\frac{\Delta x}{2}} &= \frac{1}{c_0} \frac{\tilde{D}_x|_{l+\frac{\Delta x}{2}} - \tilde{D}_x|_{l-\frac{\Delta x}{2}}}{\Delta t} \\
 C_y^H &= \cancel{\frac{\partial H_x}{\partial z}} - \frac{\partial H_z}{\partial x} & C_y^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} & C_y^H|_{l+\frac{\Delta y}{2}} &= \cancel{\frac{H_x|_{l+\frac{\Delta y}{2}} - H_x|_{l-\frac{\Delta y}{2}}}{\Delta z}} - \frac{H_z|_{l+\frac{\Delta y}{2}} - H_z|_{l-\frac{\Delta y}{2}}}{\Delta x} & C_y^H|_{l+\frac{\Delta y}{2}} &= \frac{1}{c_0} \frac{\tilde{D}_y|_{l+\frac{\Delta y}{2}} - \tilde{D}_y|_{l-\frac{\Delta y}{2}}}{\Delta t} \\
 C_z^H &= \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} & C_z^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} & C_z^H|_{l+\frac{\Delta z}{2}} &= \frac{H_y|_{l+\frac{\Delta z}{2}} - H_y|_{l-\frac{\Delta z}{2}}}{\Delta x} - \frac{H_x|_{l+\frac{\Delta z}{2}} - H_x|_{l-\frac{\Delta z}{2}}}{\Delta y} & C_z^H|_{l+\frac{\Delta z}{2}} &= \frac{1}{c_0} \frac{\tilde{D}_z|_{l+\frac{\Delta z}{2}} - \tilde{D}_z|_{l-\frac{\Delta z}{2}}}{\Delta t} \\
 \\ 
 \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x & \tilde{D}_x|_{l+\frac{\Delta x}{2}} &= (\epsilon_{xx}|_{l+\frac{\Delta x}{2}}) \tilde{E}_x|_{l+\frac{\Delta x}{2}} \\
 \tilde{D}_y &= \epsilon_{yy} \tilde{E}_y & \tilde{D}_y|_{l+\frac{\Delta y}{2}} &= (\epsilon_{yy}|_{l+\frac{\Delta y}{2}}) \tilde{E}_y|_{l+\frac{\Delta y}{2}} \\
 \tilde{D}_z &= \epsilon_{zz} \tilde{E}_z & \tilde{D}_z|_{l+\frac{\Delta z}{2}} &= (\epsilon_{zz}|_{l+\frac{\Delta z}{2}}) \tilde{E}_z|_{l+\frac{\Delta z}{2}}
 \end{aligned}$$



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## 2D Finite-Difference Equations

We eliminate the z-derivative terms.

$$\begin{aligned}
 C_x^E &= \frac{\partial \tilde{E}_z}{\partial y} & C_x^E &= -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t} & C_x^E|_{l+\frac{\Delta x}{2}} &= \frac{\tilde{E}_z|_{l+\frac{\Delta x}{2}} - \tilde{E}_z|_{l-\frac{\Delta x}{2}}}{\Delta y} & C_x^E|_{l+\frac{\Delta x}{2}} &= -\frac{\mu_{xx}}{c_0} \frac{H_x|_{l+\frac{\Delta x}{2}} - H_x|_{l-\frac{\Delta x}{2}}}{\Delta t} \\
 C_y^E &= -\frac{\partial \tilde{E}_z}{\partial x} & C_y^E &= -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t} & C_y^E|_{l+\frac{\Delta y}{2}} &= -\frac{\tilde{E}_z|_{l+\frac{\Delta y}{2}} - \tilde{E}_z|_{l-\frac{\Delta y}{2}}}{\Delta x} & C_y^E|_{l+\frac{\Delta y}{2}} &= -\frac{\mu_{yy}}{c_0} \frac{H_y|_{l+\frac{\Delta y}{2}} - H_y|_{l-\frac{\Delta y}{2}}}{\Delta t} \\
 C_z^E &= \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} & C_z^E &= -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} & C_z^E|_{l+\frac{\Delta z}{2}} &= \frac{\tilde{E}_y|_{l+\frac{\Delta z}{2}} - \tilde{E}_y|_{l-\frac{\Delta z}{2}}}{\Delta x} - \frac{\tilde{E}_x|_{l+\frac{\Delta z}{2}} - \tilde{E}_x|_{l-\frac{\Delta z}{2}}}{\Delta y} & C_z^E|_{l+\frac{\Delta z}{2}} &= -\frac{\mu_{zz}}{c_0} \frac{H_z|_{l+\frac{\Delta z}{2}} - H_z|_{l-\frac{\Delta z}{2}}}{\Delta t} \\
 \\ 
 C_x^H &= \frac{\partial H_z}{\partial y} & C_x^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} & C_x^H|_{l+\frac{\Delta x}{2}} &= \frac{H_z|_{l+\frac{\Delta x}{2}} - H_z|_{l-\frac{\Delta x}{2}}}{\Delta y} & C_x^H|_{l+\frac{\Delta x}{2}} &= \frac{1}{c_0} \frac{\tilde{D}_x|_{l+\frac{\Delta x}{2}} - \tilde{D}_x|_{l-\frac{\Delta x}{2}}}{\Delta t} \\
 C_y^H &= -\frac{\partial H_z}{\partial x} & C_y^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} & C_y^H|_{l+\frac{\Delta y}{2}} &= -\frac{H_z|_{l+\frac{\Delta y}{2}} - H_z|_{l-\frac{\Delta y}{2}}}{\Delta x} & C_y^H|_{l+\frac{\Delta y}{2}} &= \frac{1}{c_0} \frac{\tilde{D}_y|_{l+\frac{\Delta y}{2}} - \tilde{D}_y|_{l-\frac{\Delta y}{2}}}{\Delta t} \\
 C_z^H &= \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} & C_z^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} & C_z^H|_{l+\frac{\Delta z}{2}} &= \frac{H_y|_{l+\frac{\Delta z}{2}} - H_y|_{l-\frac{\Delta z}{2}}}{\Delta x} - \frac{H_x|_{l+\frac{\Delta z}{2}} - H_x|_{l-\frac{\Delta z}{2}}}{\Delta y} & C_z^H|_{l+\frac{\Delta z}{2}} &= \frac{1}{c_0} \frac{\tilde{D}_z|_{l+\frac{\Delta z}{2}} - \tilde{D}_z|_{l-\frac{\Delta z}{2}}}{\Delta t} \\
 \\ 
 \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x & \tilde{D}_x|_{l+\frac{\Delta x}{2}} &= (\epsilon_{xx}|_{l+\frac{\Delta x}{2}}) \tilde{E}_x|_{l+\frac{\Delta x}{2}} \\
 \tilde{D}_y &= \epsilon_{yy} \tilde{E}_y & \tilde{D}_y|_{l+\frac{\Delta y}{2}} &= (\epsilon_{yy}|_{l+\frac{\Delta y}{2}}) \tilde{E}_y|_{l+\frac{\Delta y}{2}} \\
 \tilde{D}_z &= \epsilon_{zz} \tilde{E}_z & \tilde{D}_z|_{l+\frac{\Delta z}{2}} &= (\epsilon_{zz}|_{l+\frac{\Delta z}{2}}) \tilde{E}_z|_{l+\frac{\Delta z}{2}}
 \end{aligned}$$



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## We No Longer Need Grid Index $k$

For 2D devices, we only retain grid indices  $i$  and  $j$ .

$$\begin{array}{llll}
 C_x^E = \frac{\partial \tilde{E}_z}{\partial y} & C_x^E = -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t} & \rightarrow & C_x^{E|j} = \frac{\tilde{E}_z^{i,j+1} - \tilde{E}_z^{i,j}}{\Delta y} \\
 C_y^E = -\frac{\partial \tilde{E}_z}{\partial x} & C_y^E = -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t} & \rightarrow & C_y^{E|j} = -\frac{\tilde{E}_z^{i+1,j} - \tilde{E}_z^{i,j}}{\Delta x} \\
 C_z^E = \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} & C_z^E = -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} & \rightarrow & C_z^{E|j} = \frac{\tilde{E}_y^{i+1,j} - \tilde{E}_y^{i,j}}{\Delta x} - \frac{\tilde{E}_x^{i,j+1} - \tilde{E}_x^{i,j}}{\Delta y} \\
 \\ 
 C_x^H = \frac{\partial H_z}{\partial y} & C_x^H = \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} & \rightarrow & C_x^{H|j} = \frac{H_z^{i,j} - H_z^{i,j-1}}{\Delta y} \\
 C_y^H = -\frac{\partial H_z}{\partial x} & C_y^H = \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} & \rightarrow & C_y^{H|j} = -\frac{H_z^{i,j} - H_z^{i-1,j}}{\Delta x} \\
 C_z^H = \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} & C_z^H = \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} & \rightarrow & C_z^{H|j} = \frac{H_y^{i,j} - H_y^{i,j-1}}{\Delta x} - \frac{H_x^{i,j} - H_x^{i-1,j}}{\Delta y} \\
 \\ 
 \tilde{D}_x = \epsilon_{xx} \tilde{E}_x & \rightarrow & \tilde{D}_x^{i,j} = (\epsilon_{xx}^{i,j}) \tilde{E}_x^{i,j} \\
 \tilde{D}_y = \epsilon_{yy} \tilde{E}_y & \rightarrow & \tilde{D}_y^{i,j} = (\epsilon_{yy}^{i,j}) \tilde{E}_y^{i,j} \\
 \tilde{D}_z = \epsilon_{zz} \tilde{E}_z & \rightarrow & \tilde{D}_z^{i,j} = (\epsilon_{zz}^{i,j}) \tilde{E}_z^{i,j}
 \end{array}$$



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## Two Distinct Modes: $E_z$ and $H_z$

Maxwell's equations have decoupled into two distinct sets of equations.

$$\begin{array}{llll}
 C_x^E = \frac{\partial \tilde{E}_z}{\partial y} & C_x^E = -\frac{\mu_{xx}}{c_0} \frac{\partial H_x}{\partial t} & \rightarrow & C_x^{E|j} = \frac{\tilde{E}_z^{i,j+1} - \tilde{E}_z^{i,j}}{\Delta y} \\
 C_y^E = -\frac{\partial \tilde{E}_z}{\partial x} & C_y^E = -\frac{\mu_{yy}}{c_0} \frac{\partial H_y}{\partial t} & \rightarrow & C_y^{E|j} = -\frac{\tilde{E}_z^{i+1,j} - \tilde{E}_z^{i,j}}{\Delta x} \\
 C_z^E = \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} & C_z^E = -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} & \rightarrow & C_z^{E|j} = \frac{\tilde{E}_y^{i+1,j} - \tilde{E}_y^{i,j}}{\Delta x} - \frac{\tilde{E}_x^{i,j+1} - \tilde{E}_x^{i,j}}{\Delta y} \\
 \\ 
 C_x^H = \frac{\partial H_z}{\partial y} & C_x^H = \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} & \rightarrow & C_x^{H|j} = \frac{H_z^{i,j} - H_z^{i,j-1}}{\Delta y} \\
 C_y^H = -\frac{\partial H_z}{\partial x} & C_y^H = \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} & \rightarrow & C_y^{H|j} = -\frac{H_z^{i,j} - H_z^{i-1,j}}{\Delta x} \\
 C_z^H = \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} & C_z^H = \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} & \rightarrow & C_z^{H|j} = \frac{H_y^{i,j} - H_y^{i,j-1}}{\Delta x} - \frac{H_x^{i,j} - H_x^{i-1,j}}{\Delta y} \\
 \\ 
 \tilde{D}_x = \epsilon_{xx} \tilde{E}_x & \rightarrow & \tilde{D}_x^{i,j} = (\epsilon_{xx}^{i,j}) \tilde{E}_x^{i,j} \\
 \tilde{D}_y = \epsilon_{yy} \tilde{E}_y & \rightarrow & \tilde{D}_y^{i,j} = (\epsilon_{yy}^{i,j}) \tilde{E}_y^{i,j} \\
 \tilde{D}_z = \epsilon_{zz} \tilde{E}_z & \rightarrow & \tilde{D}_z^{i,j} = (\epsilon_{zz}^{i,j}) \tilde{E}_z^{i,j}
 \end{array}$$



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## E<sub>z</sub> Mode

$$\begin{aligned}
 C_x^E &= \frac{\partial \tilde{E}_z}{\partial y} & C_x^E &= -\frac{\mu_{xy}}{c_0} \frac{\partial H_x}{\partial t} & \rightarrow & C_x^{E|l,j} = \frac{\tilde{E}_z^{[j+1]} - \tilde{E}_z^{[j]}}{\Delta y} & C_x^{E|l,j} &= -\frac{\mu_{xy}}{c_0} \frac{H_x^{[j]}|_{l+\frac{\Delta x}{2}} - H_x^{[j]}|_{l-\frac{\Delta x}{2}}}{\Delta t} \\
 C_y^E &= -\frac{\partial \tilde{E}_z}{\partial x} & C_y^E &= -\frac{\mu_{yx}}{c_0} \frac{\partial H_y}{\partial t} & \rightarrow & C_y^{E|l,j} = -\frac{\tilde{E}_z^{[j+1]} - \tilde{E}_z^{[j]}}{\Delta x} & C_y^{E|l,j} &= -\frac{\mu_{yx}}{c_0} \frac{H_y^{[j]}|_{l+\frac{\Delta y}{2}} - H_y^{[j]}|_{l-\frac{\Delta y}{2}}}{\Delta t} \\
 C_z^H &= \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} & C_z^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_z}{\partial t} & \rightarrow & C_z^{H|l,j} = \frac{H_y^{[j]}|_{l+\frac{\Delta x}{2}} - H_y^{[j-1]}|_{l+\frac{\Delta x}{2}}}{\Delta x} - \frac{H_x^{[j]}|_{l+\frac{\Delta y}{2}} - H_x^{[j-1]}|_{l+\frac{\Delta y}{2}}}{\Delta y} & C_z^{H|l,j} &= \frac{1}{c_0} \frac{\tilde{D}_z^{[j]}|_{l+\frac{\Delta z}{2}} - \tilde{D}_z^{[j]}|_{l-\frac{\Delta z}{2}}}{\Delta t} \\
 & & \tilde{D}_z &= \epsilon_{zz} \tilde{E}_z & \rightarrow & \tilde{D}_z^{[j]}|_l = (\epsilon_{zz}^{[j]}) \tilde{E}_z^{[j]}|_l
 \end{aligned}$$

## H<sub>z</sub> Mode

$$\begin{aligned}
 C_z^E &= \frac{\partial \tilde{E}_y}{\partial x} - \frac{\partial \tilde{E}_x}{\partial y} & C_z^E &= -\frac{\mu_{zz}}{c_0} \frac{\partial H_z}{\partial t} & \rightarrow & C_z^{E|l,j} = \frac{\tilde{E}_y^{[j+1]} - \tilde{E}_y^{[j]}|_l - \tilde{E}_x^{[j+1]} - \tilde{E}_x^{[j]}|_l}{\Delta x} & C_z^{E|l,j} &= -\frac{\mu_{zz}}{c_0} \frac{H_z^{[j]}|_{l+\frac{\Delta z}{2}} - H_z^{[j]}|_{l-\frac{\Delta z}{2}}}{\Delta t} \\
 C_x^H &= \frac{\partial H_z}{\partial y} & C_x^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_x}{\partial t} & \rightarrow & C_x^{H|l,j} = \frac{H_z^{[j]}|_{l+\frac{\Delta y}{2}} - H_z^{[j-1]}|_{l+\frac{\Delta y}{2}}}{\Delta y} & C_x^{H|l,j} &= \frac{1}{c_0} \frac{\tilde{D}_x^{[j]}|_{l+\frac{\Delta x}{2}} - \tilde{D}_x^{[j]}|_{l-\frac{\Delta x}{2}}}{\Delta t} \\
 C_y^H &= -\frac{\partial H_z}{\partial x} & C_y^H &= \frac{1}{c_0} \frac{\partial \tilde{D}_y}{\partial t} & \rightarrow & C_y^{H|l,j} = -\frac{H_z^{[j]}|_{l+\frac{\Delta x}{2}} - H_z^{[j-1]}|_{l+\frac{\Delta x}{2}}}{\Delta x} & C_y^{H|l,j} &= \frac{1}{c_0} \frac{\tilde{D}_y^{[j]}|_{l+\frac{\Delta y}{2}} - \tilde{D}_y^{[j]}|_{l-\frac{\Delta y}{2}}}{\Delta t} \\
 & & \tilde{D}_x &= \epsilon_{xx} \tilde{E}_x & \rightarrow & \tilde{D}_x^{[j]}|_l = (\epsilon_{xx}^{[j]}) \tilde{E}_x^{[j]}|_l \\
 & & \tilde{D}_y &= \epsilon_{yy} \tilde{E}_y & \rightarrow & \tilde{D}_y^{[j]}|_l = (\epsilon_{yy}^{[j]}) \tilde{E}_y^{[j]}|_l
 \end{aligned}$$

# Update Equations

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## Update Equations for $E_z$ Mode

Solving the finite-difference equations for the future time values of the fields associated with the  $E_z$  mode leads to:

$$\begin{aligned}
 C_x^E \Big|_t^{i,j} &= -\frac{\mu_{xx} \Big|^{i,j}}{c_0} \frac{H_x \Big|_{t+\frac{\Delta t}{2}}^{i,j} - H_x \Big|_{t-\frac{\Delta t}{2}}^{i,j}}{\Delta t} & H_x \Big|_{t+\frac{\Delta t}{2}}^{i,j} &= H_x \Big|_{t-\frac{\Delta t}{2}}^{i,j} + \left( -\frac{c_0 \Delta t}{\mu_{xx} \Big|^{i,j}} \right) C_x^E \Big|_t^{i,j} \\
 C_y^E \Big|_t^{i,j} &= -\frac{\mu_{yy} \Big|^{i,j}}{c_0} \frac{H_y \Big|_{t+\frac{\Delta t}{2}}^{i,j} - H_y \Big|_{t-\frac{\Delta t}{2}}^{i,j}}{\Delta t} & H_y \Big|_{t+\frac{\Delta t}{2}}^{i,j} &= H_y \Big|_{t-\frac{\Delta t}{2}}^{i,j} + \left( -\frac{c_0 \Delta t}{\mu_{yy} \Big|^{i,j}} \right) C_y^E \Big|_t^{i,j} \\
 C_z^H \Big|_{t+\frac{\Delta t}{2}}^{i,j} &= \frac{1}{c_0} \frac{\tilde{D}_z \Big|_{t+\Delta t}^{i,j} - \tilde{D}_z \Big|_t^{i,j}}{\Delta t} & \tilde{D}_z \Big|_{t+\Delta t}^{i,j} &= \tilde{D}_z \Big|_t^{i,j} + (c_0 \Delta t) C_z^H \Big|_{t+\frac{\Delta t}{2}}^{i,j} \\
 \tilde{D}_z \Big|_t^{i,j} &= \left( \epsilon_{zz} \Big|^{i,j} \right) \tilde{E}_z \Big|_t^{i,j} & \tilde{E}_z \Big|_{t+\Delta t}^{i,j} &= \left( \frac{1}{\epsilon_{zz} \Big|^{i,j}} \right) \left( \tilde{D}_z \Big|_{t+\Delta t}^{i,j} \right)
 \end{aligned}$$

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## Update Equations for $H_z$ Mode

Solving the finite-difference equations for the future time values of the fields associated with the  $H_z$  mode leads to:

$$\begin{aligned}
 C_z^E \Big|_t^{i,j} &= -\frac{\mu_{zz} \Big|_t^{i,j}}{c_0} \frac{H_z \Big|_{t+\frac{\Delta t}{2}}^{i,j} - H_z \Big|_{t-\frac{\Delta t}{2}}^{i,j}}{\Delta t} \\
 C_x^H \Big|_{t+\frac{\Delta t}{2}}^{i,j} &= \frac{1}{c_0} \frac{\tilde{D}_x \Big|_{t+\Delta t}^{i,j} - \tilde{D}_x \Big|_t^{i,j}}{\Delta t} \\
 C_y^H \Big|_{t+\frac{\Delta t}{2}}^{i,j} &= \frac{1}{c_0} \frac{\tilde{D}_y \Big|_{t+\Delta t}^{i,j} - \tilde{D}_y \Big|_t^{i,j}}{\Delta t} \\
 \tilde{D}_x \Big|_t^{i,j} &= \left( \varepsilon_{xx} \Big|_t^{i,j} \right) \tilde{E}_x \Big|_t^{i,j} \\
 \tilde{D}_y \Big|_t^{i,j} &= \left( \varepsilon_{yy} \Big|_t^{i,j} \right) \tilde{E}_y \Big|_t^{i,j}
 \end{aligned}
 \quad \rightarrow \quad
 \begin{aligned}
 H_z \Big|_{t+\frac{\Delta t}{2}}^{i,j} &= H_z \Big|_{t-\frac{\Delta t}{2}}^{i,j} + \left( -\frac{c_0 \Delta t}{\mu_{zz} \Big|_t^{i,j}} \right) C_z^E \Big|_t^{i,j} \\
 \tilde{D}_x \Big|_{t+\Delta t}^{i,j} &= \tilde{D}_x \Big|_t^{i,j} + (c_0 \Delta t) C_x^H \Big|_{t+\frac{\Delta t}{2}}^{i,j} \\
 \tilde{D}_y \Big|_{t+\Delta t}^{i,j} &= \tilde{D}_y \Big|_t^{i,j} + (c_0 \Delta t) C_y^H \Big|_{t+\frac{\Delta t}{2}}^{i,j} \\
 \tilde{E}_x \Big|_{t+\Delta t}^{i,j} &= \left( \frac{1}{\varepsilon_{xx} \Big|_t^{i,j}} \right) \left( \tilde{D}_x \Big|_{t+\Delta t}^{i,j} \right) \\
 \tilde{E}_y \Big|_{t+\Delta t}^{i,j} &= \left( \frac{1}{\varepsilon_{yy} \Big|_t^{i,j}} \right) \left( \tilde{D}_y \Big|_{t+\Delta t}^{i,j} \right)
 \end{aligned}$$

## Boundary Conditions

## Where Are the Boundary Conditions?

All of the spatial derivatives appear only the curl calculations.

$$C_x^E \Big|_l^{i,j} = \frac{\tilde{E}_z \Big|_l^{i,j+1} - \tilde{E}_z \Big|_l^{i,j}}{\Delta y} \quad (y\text{-hi})$$

$$C_y^E \Big|_l^{i,j} = -\frac{\tilde{E}_z \Big|_l^{i+1,j} - \tilde{E}_z \Big|_l^{i,j}}{\Delta x} \quad (x\text{-hi})$$

$$C_z^E \Big|_l^{i,j} = \frac{\tilde{E}_y \Big|_l^{i+1,j} - \tilde{E}_y \Big|_l^{i,j}}{\Delta x} - \frac{\tilde{E}_x \Big|_l^{i,j+1} - \tilde{E}_x \Big|_l^{i,j}}{\Delta y} \quad (x\text{-hi and } y\text{-hi})$$

$$C_x^H \Big|_{l+\frac{\Delta x}{2}}^{i,j} = \frac{H_z \Big|_{l+\frac{\Delta x}{2}}^{i,j} - H_z \Big|_{l+\frac{\Delta x}{2}}^{i,j-1}}{\Delta y} \quad (y\text{-lo})$$

$$C_y^H \Big|_{l+\frac{\Delta x}{2}}^{i,j} = -\frac{H_z \Big|_{l+\frac{\Delta x}{2}}^{i,j} - H_z \Big|_{l+\frac{\Delta x}{2}}^{i-1,j}}{\Delta x} \quad (x\text{-lo})$$

$$C_z^H \Big|_{l+\frac{\Delta x}{2}}^{i,j} = \frac{H_y \Big|_{l+\frac{\Delta x}{2}}^{i,j} - H_y \Big|_{l+\frac{\Delta x}{2}}^{i-1,j}}{\Delta x} - \frac{H_x \Big|_{l+\frac{\Delta x}{2}}^{i,j} - H_x \Big|_{l+\frac{\Delta x}{2}}^{i,j-1}}{\Delta y} \quad (x\text{-lo and } y\text{-lo})$$

Boundary conditions are handled in the curl computations.

We have “modularized” the boundary conditions by isolating the curl calculations.

## Hint About Implementing Boundary Conditions

# DO NOT USE IF STATEMENTS!!!

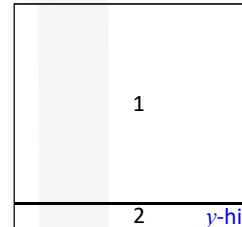
Your code will not be much shorter and it will run slower if you use `if` statements.

## Dirichlet Boundary Conditions for $CE_x$

For  $CE_x$ , the problem occurs at the  $y$ -hi side of the grid.

It is fixed explicitly.

$$C_x^E|_t^{i,j} = \begin{cases} \frac{\tilde{E}_z|_t^{i,j+1} - \tilde{E}_z|_t^{i,j}}{\Delta y} & 1. j < N_y \\ 0 & 2. j = N_y \end{cases}$$



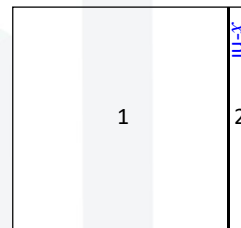
```
% Compute CEx
for nx = 1 : Nx
    for ny = 1 : Ny-1
        CEx(nx,ny) = (Ez (nx,ny+1) - Ez (nx,ny)) / dy;
    end
    CEx(nx,Ny) = (0 - Ez (nx, Ny)) / dy;
end
```

## Dirichlet Boundary Conditions for $CE_y$

For  $CE_y$ , the problem occurs at the  $x$ -hi side of the grid.

It is fixed explicitly.

$$C_y^E|_t^{i,j} = \begin{cases} -\frac{\tilde{E}_z|_t^{i+1,j} - \tilde{E}_z|_t^{i,j}}{\Delta x} & 1. i < N_x \\ 0 & 2. i = N_x \end{cases}$$



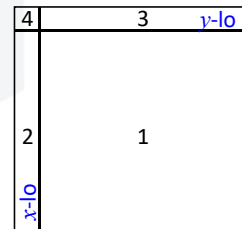
```
% Compute CEy
for ny = 1 : Ny
    for nx = 1 : Nx-1
        CEy(nx,ny) = - (Ez (nx+1,ny) - Ez (nx,ny)) / dx;
    end
    CEy(Nx,ny) = - (0 - Ez (Nx,ny)) / dx;
end
```

## Dirichlet Boundary Conditions for CHz

For CHz, the problem occurs at both the  $x$ -lo side of the grid and  $y$ -lo.

Both of these are fixed explicitly.

$$C_z^{H|^{i,j}}|_{t+\Delta t/2} = \begin{cases} \frac{H_y|_{t+\frac{\Delta t}{2}}^{i,j} - H_y|_{t+\frac{\Delta t}{2}}^{i-1,j}}{\Delta x} - \frac{H_x|_{t+\frac{\Delta t}{2}}^{i,j} - H_x|_{t+\frac{\Delta t}{2}}^{i,j-1}}{\Delta y} & 1. \text{ for } i > 1 \text{ and } j > 1 \\ \frac{H_y|_{t+\frac{\Delta t}{2}}^{1,j} - 0}{\Delta x} - \frac{H_x|_{t+\frac{\Delta t}{2}}^{1,j} - H_x|_{t+\frac{\Delta t}{2}}^{1,j-1}}{\Delta y} & 2. \text{ for } i = 1 \text{ and } j > 1 \\ \frac{H_y|_{t+\frac{\Delta t}{2}}^{i,1} - H_y|_{t+\frac{\Delta t}{2}}^{i-1,1}}{\Delta x} - \frac{H_x|_{t+\frac{\Delta t}{2}}^{i,1} - 0}{\Delta y} & 3. \text{ for } i > 1 \text{ and } j = 1 \\ \frac{H_y|_{t+\frac{\Delta t}{2}}^{1,1} - 0}{\Delta x} - \frac{H_x|_{t+\frac{\Delta t}{2}}^{1,1} - 0}{\Delta y} & 4. \text{ for } i = 1 \text{ and } j = 1 \end{cases}$$



## MATLAB Code for CHz (1 of 3)

First, just blindly implement the curl calculation...

```
% Compute CHz
for ny = 1 : Ny
    for nx = 1 : Nx
        CHz(nx,ny) = (Hy(nx,ny) - Hy(nx-1,ny))/dx ...
                    - (Hx(nx,ny) - Hx(nx,ny-1))/dy;
    end
end
```

As expected, this will produce an error when trying to access  $H_y(0, ny)$  or  $H_x(nx, 0)$ .

## MATLAB Code for CHz (2 of 3)

Second, the problem at  $n_x = 1$  is handled explicitly by copying the code inside the  $n_x$  loop, pasting it above, and handling the problem.

```
% Compute CHz
for ny = 1 : Ny
    CHz(1,ny) = (Hy(1,ny) - 0)/dx ...
               - (Hx(1,ny) - Hx(1,ny-1))/dy;
    for nx = 2 : Nx
        CHz(nx,ny) = (Hy(nx,ny) - Hy(nx-1,ny))/dx ...
                    - (Hx(nx,ny) - Hx(nx,ny-1))/dy;
    end
end
```

There is still an error at  $H_x(1, 0)$  and  $H_x(nx, 0)$

## MATLAB Code for CHz (3 of 3)

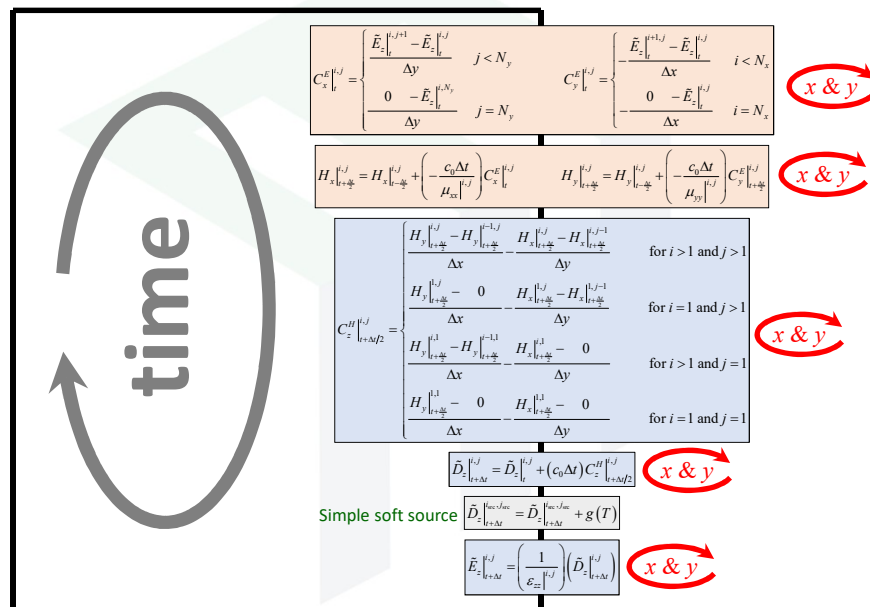
Third, the problem at  $n_y = 1$  is handled explicitly by copying the code inside the  $n_y$  loop, pasting it above, and handling the problem.

```
% Compute CHz
CHz(1,1) = (Hy(1,1) - 0)/dx ...
           - (Hx(1,1) - 0)/dy;
for nx = 2 : Nx
    CHz(nx,1) = (Hy(nx,1) - Hy(nx-1,1))/dx ...
               - (Hx(nx,1) - 0)/dy;
end
for ny = 2 : Ny
    CHz(1,ny) = (Hy(1,ny) - 0)/dx ...
               - (Hx(1,ny) - Hx(1,ny-1))/dy;
    for nx = 2 : Nx
        CHz(nx,ny) = (Hy(nx,ny) - Hy(nx-1,ny))/dx ...
                    - (Hx(nx,ny) - Hx(nx,ny-1))/dy;
    end
end
```

# Revised FDTD Algorithm for 2D Simulations

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## FDTD Algorithm for $E_z$ Mode



Slide 48



# Animation of Basic 2D FDTD Algorithm

