## Extremal black holes in dynamical Chern-Simons gravity

Leo C. Stein (TAPIR, Caltech) with Robert McNees (Loyola University Chicago) and Nicolás Yunes (Montana State University)

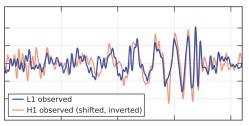


April APS Meeting 2016

[arXiv:1512.05453]

### Squad goals

- Before this year: precision tests of GR in weak field
- Now: first direct measurements of dynamical, strong field regime



- Future: precision tests of GR in the strong field (BBH)
  - Parametric vs. non-parametric. Know what we're looking for?
- Only have binary black hole mergers in GR! Some ideas.

### Modest goals

- Short term: Shapes of black holes in beyond-GR theories
- Necessary step before dynamics of black holes
- Our work: BHs in dynamical Chern-Simons gravity

# What is dynamical Chern-Simons gravity?

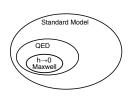
Chern-Simons = GR + pseudo-scalar + interaction

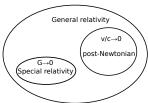
$$S = \int d^4x \sqrt{-g} \left[ \frac{m_{\rm pl}^2}{2} R - \frac{1}{2} (\partial \vartheta)^2 + \frac{m_{\rm pl}}{8} \ell^2 \vartheta *RR \right]$$

Anomaly cancellation, low-E string theory, LQG...(see Nico's review)

# Why study dynamical Chern-Simons gravity?

- Don't think dCS is fundamental theory (see Delsate+Hilditch+Witek)
- GR is a low-energy effective field theory

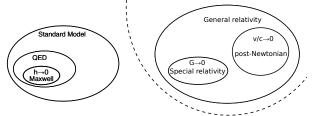




- Lowest-order EFT with parity-odd  $\vartheta$ , shift symmetry (long range)
- Phenomenology unique from other  $\mathbb{R}^2$  (e.g. Einstein-dilaton-Gauss-Bonnet)
- Only decoupling limit makes sense

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## Decoupling limit

Theory is GR  $+ \varepsilon \times$  deformation. Expand everything in  $\varepsilon$ 

- $\varepsilon^0$ : Vacuum GR
- $\varepsilon^1$ : Unfreeze new degrees of freedom

$$\square_{\mathrm{GR}} \vartheta^{(1)} = \mathrm{Src}[g_{\mathrm{GR}}]$$

•  $\varepsilon^2$ : Deformation to metric

$$G_{ab}^{(1)}[h^{(2)}] = \operatorname{Src}_{ab}[\vartheta^{(1)}, g_{GR}]$$

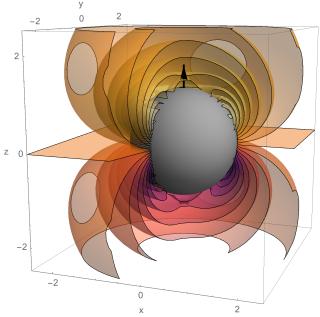
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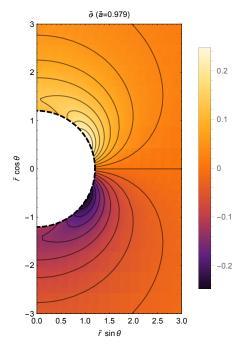
- Short term: Shapes of black holes in beyond-GR theories
- Necessary step before dynamics of black holes
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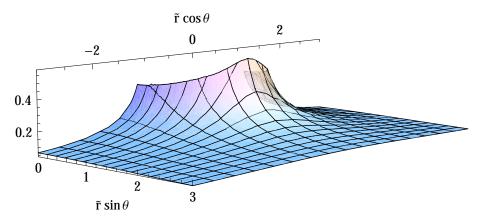
#### Black holes in dCS

- a=0 (Schwarzschild) is exact solution with  $\vartheta=0$
- · Analytically known solutions in decoupling limit
  - $a \ll M$  limit up to  $\mathcal{O}(a^2)$ , valid  $\forall r$  (see Yunes+Pretorius, Yagi+Yunes+Tanaka)
  - $r \gg M$  limit for l = 1, valid  $\forall a$  (see Yagi+Yunes+Tanaka)
- Numerical solutions for scalar sector  $\forall r, a < M$  Phys. Rev. D 90, 044061 (2014) [arXiv:1407.2350]
- Conjectured divergence at  $a \to M$
- Analytics for a = M?





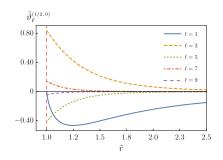




# What we did in [arXiv:1512.05453]

- Equations simplify for a=M
- Analytic scalar field solution for  $\vartheta_\ell$

$$\vartheta_\ell(r) = (\text{polyn., roots, logs, inv. trig})(r,\ell)$$

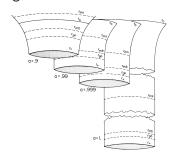


# What we did in [arXiv:1512.05453]

- Metric tensor deformation is hard
- Tensor deformation = trace + grad vect + transverse tracefree
- Formal solution for trace of metric deformation (gauge dependent)
- Monopole  $\ell = 0$  diverges on Kerr horizon!

#### Fresh results

- Near-horizon extremal Kerr (NHEK) (Sam's and Niels' talks)
- Zoom in on horizon region



• Acquire enhanced symmetry group  $SL(2,\mathbb{R}) \times U(1)$ 

#### Fresh results

• NHEK allows to find horizon solution analytically

$$\begin{split} \vartheta &= \, -\frac{2c(-7+2c^2+c^4)}{(1+c^2)^3} - 4\arctan c \\ \mathrm{tr} h^{\mathrm{Def}} &= \frac{4(301+482c^2+7526c^4+7120c^6+3221c^8+566c^{10})}{105(1+c^2)^6} \\ &- 24i\arctan(c)\arctan\left(\frac{-i+c}{i+c}\right) - \frac{3944}{105}\arctan(c^2) \\ &+ 6i\operatorname{Li}_2\left(-\frac{1+ic}{i+c}\right) - 6i\operatorname{Li}_2\left(\frac{1+ic}{i+c}\right) \end{split}$$

where  $c = \cos \theta$ 

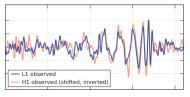
#### What does it mean?

- Confirmed conjectured divergence at Kerr horizon for  $a \to M$
- New conjecture: divergence hidden behind deformed horizon
- Need full metric to locate
- ullet Or: extremality condition may be shifted away from a=M
- Constraints from observations of rapidly spinning BHs

#### What's next?

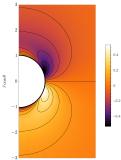


- Understand NHEK solution
- Find full metric tensor solutions analytically (extremal)
- Numerical metric tensor solutions (all a)
- Correction to BH thermo from dCS
- Observables: Accretion disk, black hole shadow
- Binary black hole merger in decoupling limit (with Maria Okounkova)





- Goal: Want to test GR in dynamical strong-field, BBH.
   Need to understand merger in deformations of theory.
- First: Must understand deformation to Kerr when theory deformed
- Previous work: Computed dCS scalar deformation to Kerr numerically
- This work: Computed scalar deformation to extremal Kerr analytically



 $\tilde{r} \sin \theta$ 

Details in [arXiv:1512.05453]

Bonus slides!

## Equations to solve

$$S = \int d^4x \sqrt{-g} \left[ \frac{m_{\rm pl}^2}{2} R - \frac{1}{2} (\partial \vartheta)^2 + \frac{m_{\rm pl}}{8} \ell^2 \vartheta RR \right]$$
$$\square \vartheta = -\frac{m_{\rm pl}}{8} \ell^2 RR$$
$$m_{\rm pl}^2 G_{ab} + m_{\rm pl} \ell^2 C_{ab} = T_{ab}^{(m)} + T_{ab}^{(\vartheta)}$$

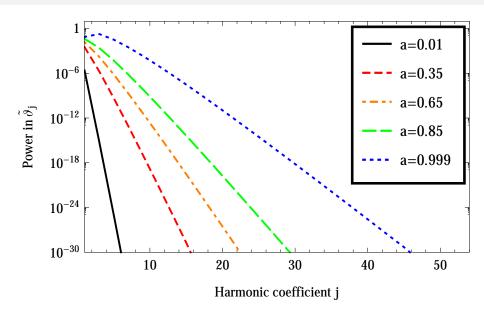
•  $g^{ab}C_{ab}=0$ 

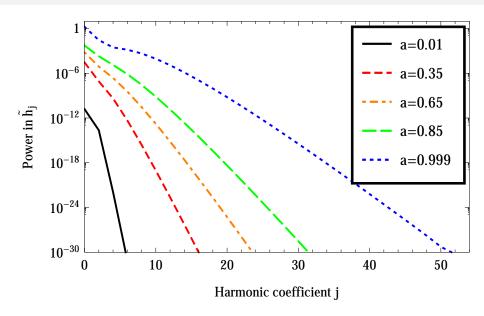
### Numerical approach

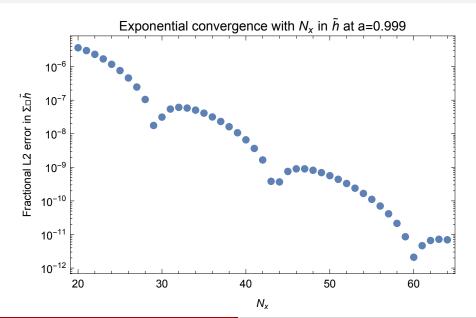
- Elliptic PDE. Could solve hyperbolic, parabolic, relaxation scheme
- Numerical separation of variables. Each j mode is an ODE.
- ullet Compactify r
- Pseudospectral collocation method
- Directly solve discrete ODE operator ("numerical Green's function")

## Numerical approach

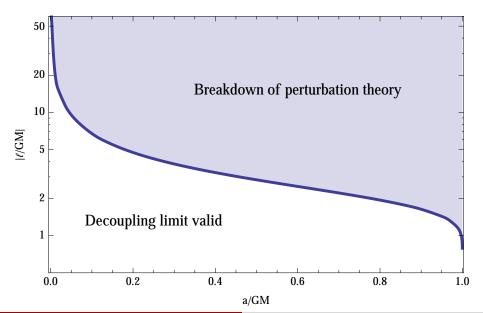
- For each a, find  $\vartheta(r,\theta;a)$ , compute  $(\partial\vartheta)^2$ , find  $h^{\mathrm{def}}(r,\theta;a)$
- ullet Evaluate  $\max |h^{
  m def}|$  and find regime of validity





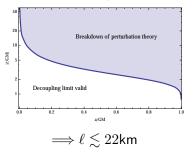


# Regime of validity



## Forecasting bounds

- Observation of BH indistinguishable from GR predictions
- Size of ℓ correction below breakdown (caveat: cancellation)
- GRO J1655–40:  $M = 6.30 \pm 0.27 M_{\odot}$ ,  $\tilde{a} \approx 0.65$ –0.75



ullet Better by  $10^7$  than Solar System bounds