The very first idea is to make this document a learning tool for git, latex and, of course, of Electronics. I will be very proud of ourselves if we reach to maintain English as its primary language.

Introduction to Electronics
Formulas and utilities

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1 Semiconductors

Semiconductors are material with an electrical conductivity value falling between that of a conductor and an insulator. An important property of semiconductors is that, during a temperature increment, their resistance decreases. Common semiconductors material are:

- Silicon, Si. Z (i.e. atomic number) = 14. IV (14th) group.

- Germanium, Ge. Atomic number = 32. IV (14th) group

- Gallium Arsenide, GaAs, a compound of Gallium and Arsenic.

1.1 Movement analysis of electric charge in a semiconductor bar

Given a bar of length ℓ, and an Electric field E, every electron in valence band will be affected by a force $F = q \cdot E$, and thus its acceleration a will be, until the electron is subjected to the Electric field, $a(t) = \frac{F}{m} = \frac{q \cdot E}{m}$; being this quantity a constant respect to time, the speed will be a straight line with angular coefficient $m = a = \frac{E}{q \cdot m}$:

$$v(t) = \frac{q \cdot E}{m} \cdot t = a \cdot t \tag{1}$$

However, scientists have observed that during its path to freedom (i.e. the end of the bar), the electrons hit some nucleus, this resetting their speed (and, by consequence, kinetic energy) to zero.

look up the axis measures and ticks to make it realistic.
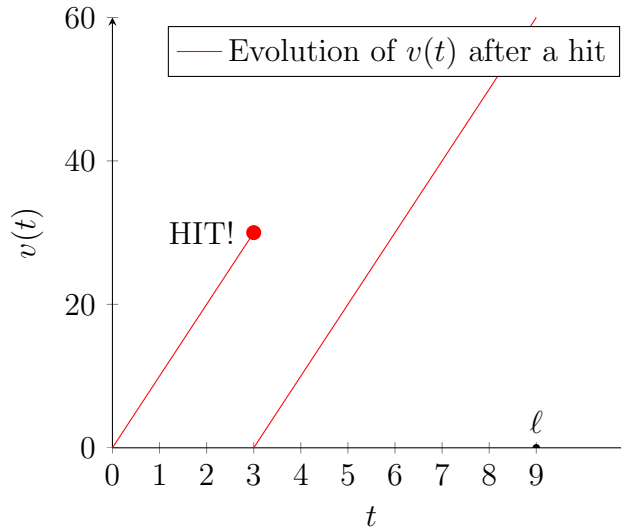
In a difficult way, the average speed of the electrons can be expressed by the area underlying the red line divided by the total amount of time: $\bar{v}$ of function $v(t)$ is $\frac{1}{t_1} \int_0^{t_1} v(t) \cdot dt$.

In an easier way we can observe that between the $i-1^{th}$ and the ith hit $\bar{v}_i = \frac{h}{2} = \frac{a \cdot \Delta t}{2}$, and so $\bar{v} = \frac{\bar{h}}{n} = \frac{\sum a \cdot \Delta t}{n} = a \cdot \bar{t}$, where n is the number of hits and $\bar{t}$ is the *average flying time*. By consequence

$$\bar{v} = a \cdot \bar{t} = \frac{\bar{t} \cdot q}{m} \cdot E = \mu \cdot E \tag{2}$$

Where we introduce the Electron mobility μ, that depends on $\bar{t}$, which is a material constant.

$$\mu = [\frac{cm^2}{V \cdot s}] \tag{3}$$

Mobility is a constant even for **holes** and usually it is related to electron mobility my the law

$$\mu_p = 3\mu_n \tag{4}$$

where μ_p is the *electron mobility* and μ_n is the *holes mobility* . This doesn't surprise ourself because holes "move" in *valence band*, while electrons move in *conductance band*, in which the movement is less obstructed and attracted by nuclei.

1.2 Semiconductors: a nuclear approach

In 1913 Niehls Bohr proposed a model of atom in which *quantization of the angular momentum* explained some interesting results regarding the light emitted by Hydrogen atoms heated. This model is very useful to explain the internal structure of atoms and can tell us something even about Silicon. A neutral Silicon atom possesses 14 protons and many electrons.

- the first 10 of them lies in the first two energy level: they fit perfectly around the nucleus and are generally not involved in chemical reactions (because of the very high amount of energy necessary to "interact" with them).

- the remaining 4 electrons are called *valence electrons* and are at higher energy levels. They are involved in the *covalent bond.*

When temperature is near at the absolute zero (-273^oC) the Silicon atoms assume a crystal form in which every *valence electron* is shared with the nearest 4 atoms of Silicon with a covalent bond: in this circumstance Silicon act as an **insulator**. At higher temperatures, (i. e. furnishing energy to the atom) covalent bounds can break.

1.2.1 Electrons and Holes

When the break happens, an electron *jumps* from valence band to conduction band and a hole is generated. The conduction-band electron is now a **charge carrier** and the atom it left is a charged positively (cation).

1.2.2 Valence and conduction bands

Without the necessary energy, electrons remain under the *Fermi level*, in a zone called *Valence Band.* when the *jump* occurs, electrons step over the *Fermi level* and stay in the *conduction band.* In the conduction band charge carrier have less constraints and can easily move in the crystalline lattice, depending on the intensity of the electro-magnetic field.

1.2.3 Silicon properties

- The amount of charge carriers in a cm^3 (i.e. electrons in the conduction band) at ambient temperature $n = 10^{10cm^{-3}}$. It is reasonably a little number compared to the global amount of atoms in a Silicon cm^3 (5 $10^2$2). [Think that if atoms were people on earth, only 7 people would result a cation with the loss of their electron].

- It is compatible with SiO_2 and Al

1.3 Doping

The doping process consists in significantly change the number of *electrons* or *holes* in a semiconductor. This result is achieved inserting atoms of the doping element in the crystalline arrangement of the doped.

1.3.1 Utilities and numbers

- $N_d \cong 10^{12} : 10^{20} cm^{-3}$ number of donors

- $p[\dfrac{electrons}{cm^3}]$ number of holes

- $n[\dfrac{electrons}{cm^3}]$ number of electrons

TO ADD: the part regarding mobility and resistance.

2 Diodes
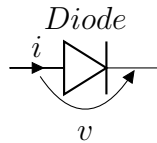
Diodes are defined as a p-n junction. Is is made by two parts: the left part (according to the figure) is a p-doped semiconductor in which there are holes left. The other part, on the contrary, is n-doped in which there are electrons left.

3   FORMULAS

3.1   Electric field and semiconductors

Electrostatic force between two charges $q_1 q_2$ at a distance r:

$$F = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \tag{5}$$

The Electric field (as the force which a charge q is subjected to in a point of the space divided by the charge q)

$$E = \frac{F}{q} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \tag{6}$$

Note that the Electric field is a **CONSERVATIVE FIELD**
A conservative field is a field in which the *Work* done for placing an element from a point a to a point b depends only from the different location of a and b and not by the route done.

The *Work* in a conservative field is defined as follows:

$$W_{a\to b} = \int_a^b \overrightarrow{F}(x)dx = q(V_b - V_a)d \tag{7}$$

The *Potential energy* of a charge subjected to an Electric field is

$$U = qV[eV] \tag{8}$$

and the measure unit is the electronVolt eV, where $1eV = 1,6\ 10^{-19}J$, the energy necessary to move an electron from a point a to a point b with $V_a - V_b = 1V$.

The *Electric Potential* at a point r in an Electric Field E is given by

$$V_E = -\int_C E \cdot \mathrm{d}\ell \tag{9}$$

Where C is an arbitrary path from the point r to the point with zero potential.

The *Electric Current* is defined as i (from the french *intensité du courant*):

$$I = \frac{Q}{t} = \frac{V}{R} \quad [Amp] \tag{10}$$

4 Engineer's corner

Question 1: How much is $1C$?

At first attempt we can say that 1 *Coulomb* is *very* approximately 10^{19} times the fundamental charge unit (the electron). However it is a bit difficult rendering the idea of 10^{19} electrons. Thinking about batteries, a common AAA battery has a capacity of 540 mAh, which are equivalent to 1944 C.

A common AA battery has a charge of 9 kC.

In everyday terms, a Coulomb is $\frac{1}{2000}$ times the charge occurring to empty a normal AAA battery.

Problem 6.   a) Suppose an entire function f is bounded by M along $|z| = R$. Show that the coefficients C_k in its power series expansion about 0 satisfy

$$|C_k| \le \frac{M}{R^k}.$$

b) Suppose a polynomial is bounded by 1 in the unit disc. Show that all its coefficients are bounded by 1.

Solution Part a): Since f is an entire function it can be expressed as an infinite power series, i.e.

$$f(z) = \sum_{k=0}^{\infty} \frac{f^{(k)}(0)}{k!} z^k = \sum_{k=0}^{\infty} C_k z^k.$$

If we recall Cauchy's Integral we have

$$f(z) = \frac{1}{2\pi i} \int_\gamma \frac{f(w)}{w - z} \, dw,$$

carefully notice that $\frac{1}{w-z} = \frac{1}{w} \cdot \frac{1}{1 - \frac{z}{w}}$ can be written as a geometric series. We have

$$\frac{1}{2\pi i} \int_\gamma \frac{f(w)}{w - z} \, dw = \frac{1}{2\pi i} \int_\gamma \left\{ \frac{f(w)}{w} \cdot \left(\frac{1}{1 - \frac{z}{w}} \right) \right\} dw$$

$$= \frac{1}{2\pi i} \int_\gamma \left\{ \frac{f(w)}{w} \cdot \left(1 + \frac{z}{w} + \frac{z^2}{w^2} + \frac{z^3}{w^3} + \cdots \right) \right\} dw$$

$$= \left(\frac{1}{2\pi i} \int_\gamma \frac{f(w)}{w} \, dw \right) z^0 + \left(\frac{1}{2\pi i} \int_\gamma \frac{f(w)}{w^2} \, dw \right) z^1 + \left(\frac{1}{2\pi i} \int_\gamma \frac{f(w)}{w^3} \, dw \right) z^2 \cdots$$

Now take the modulus of C_k to get

$$|C_k| = \left| \frac{1}{2\pi i} \int_\gamma \frac{f(w)}{w^{k+1}} \, dw \right| \le \frac{1}{2\pi} \int_\gamma \frac{|f(w)|}{|w^{k+1}|} \, |dw| \le \frac{M}{2\pi} \int_\gamma \frac{|dw|}{|w^{k+1}|}$$

Then integrate along $\gamma(\theta) = Re^{i\theta}$ for $\theta \in [0, 2\pi]$ to get

$$|C_k| \le \frac{M}{2\pi} \int_0^{2\pi} \frac{|iRe^{i\theta} \, d\theta|}{|R^{k+1} e^{ik\theta}|} = \frac{M}{2\pi \cdot R^k} \int_0^{2\pi} d\theta = \frac{M}{R^k}.$$

Hence, $|C_k| \le \frac{M}{R^k}$.

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