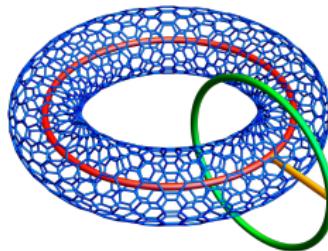


# Generalized string nets at finite temperature

**Julien Vidal**

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CNRS, Sorbonne Université, Paris



in collaboration with:

- Jean-Noël Fuchs and Anna Ritz-Zwilling (LPTMC, Paris)
- Steven H. Simon (Rudolf Peierls Centre for Theoretical Physics, Oxford)

J. Vidal, Phys. Rev. B **105**, L041110 (2022) / arXiv:2108.13425

A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, arXiv:2309.00343

## Topological quantum order in condensed matter in three dates

- 1989 : High- $T_c$  superconductors, FQHE ([X.-G. Wen, F. Wilczek, A. Zee](#))
- 1997 : Fault-tolerant quantum computation ([A. Kitaev, J. Preskill](#))
- 2005 : String-net condensation ([M. Levin, X.-G. Wen](#))

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## Concepts coming from 70's and 80's

- Lattice gauge theories ([Wegner, Wilson, Kogut](#))
- Conformal field theories ([Pasquier, Verlinde, Moore, Seiberg](#))
- Topological quantum field theories ([Witten](#))
- Knot theory ([Kauffman, Jones](#))

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## In this talk:

Two-dimensional gapped quantum systems with:

- Topology-dependent ground-state degeneracy
- Anyonic excitations

# Outline

- 1 The Levin-Wen model in a nutshell
- 2 Spectral degeneracies
- 3 Partition function

# Outline

1 The Levin-Wen model in a nutshell

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3 Partition function

# The Levin-Wen model in a nutshell

## The Levin-Wen (string-net) model

- Microscopic lattice model
- Trivalent graph (honeycomb lattice, two-leg ladder,...)

M. Levin and X.-G. Wen, Phys. Rev. B **71**, 045110 (2005)

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## Input: Unitary Fusion Category $\mathcal{C}$

- Set of objects (strings, labels, charges, superselection sectors, particles...)
- Fusion rules:  $a \times b = \sum_c N_{ab}^c c \rightarrow$  Quantum dimensions
- Associativity of the fusion rules  $\rightarrow F$ -symbols

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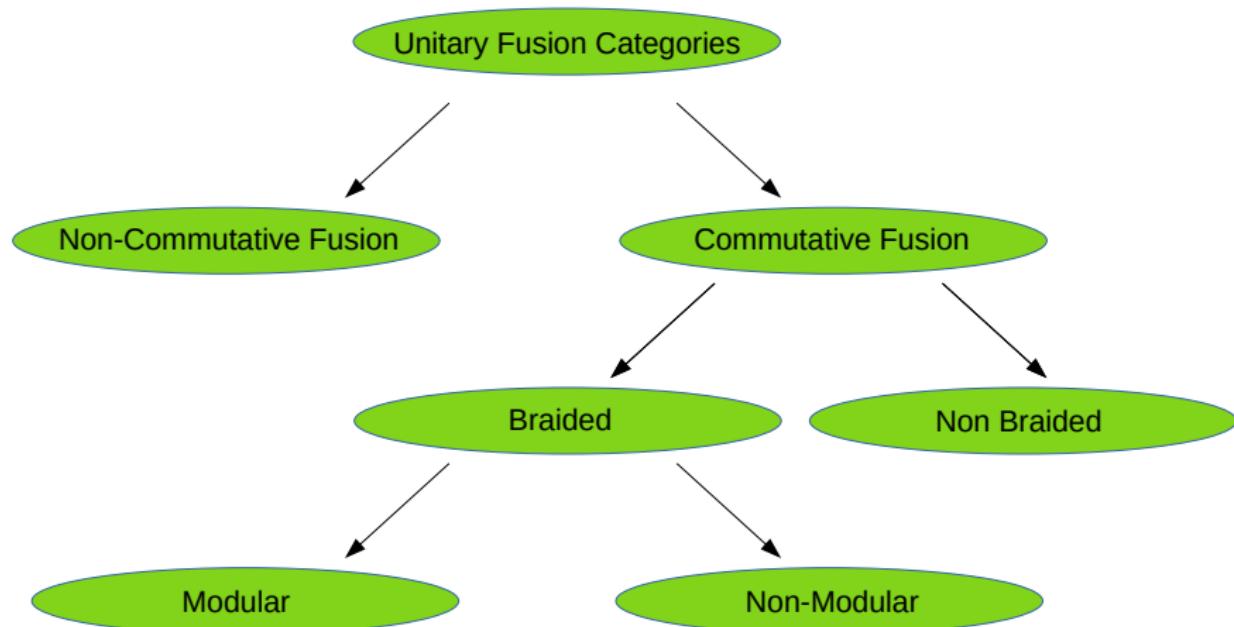
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## Output: Unitary Modular Tensor Category $\mathcal{Z}(\mathcal{C})$ (Drinfeld center of $\mathcal{C}$ )

- Anyon theory (other objects, other fusion rules, other...)
- Achiral topological phase (chiral central charge  $c \bmod 8 = 0$ )
- Different  $\mathcal{C}$  may have the same  $\mathcal{Z}(\mathcal{C})$  (Morita equivalence)

# The Levin-Wen model in a nutshell

A possible tree of UFC



# The Levin-Wen model in a nutshell

## Local constraints and Hilbert space

- Degrees of freedom are the objects of  $\mathcal{C}$
- Degrees of freedom are defined on the links of the trivalent graph
- Hilbert space  $\mathcal{H}$ : set of configurations respecting fusion rules at each vertex

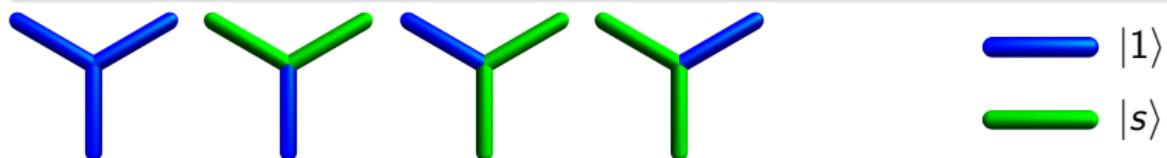
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## Ex: $\mathbb{Z}_2$ category

- Two labels:  $\{1, s\}$
- Fusion rules:  $1 \times 1 = 1$ ,  $1 \times s = s$ ,  $s \times s = 1$



Hilbert space dimension for any trivalent graph with  $N_v$  trivalent vertices

- $\text{Dim } \mathcal{H} = 2^{\frac{N_v}{2} + 1}$

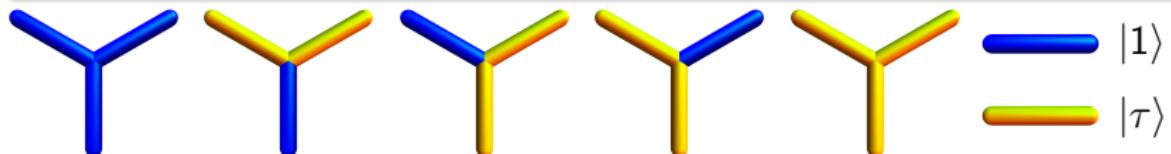
# The Levin-Wen model in a nutshell

## Local constraints and Hilbert space

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## Ex: Fibonacci category

- Two objects:  $\{1, \tau\}$
- Fusion rules:  $1 \times 1 = 1$ ,  $1 \times \tau = \tau$ ,  $\tau \times \tau = 1 + \tau$



## Hilbert space dimension for any trivalent graph with $N_v$ vertices

- $\dim \mathcal{H} = (1 + \varphi^2)^{\frac{N_v}{2}} + (1 + \varphi^{-2})^{\frac{N_v}{2}}$ ,  $\varphi = \frac{1+\sqrt{5}}{2}$  (golden ratio)

# The Levin-Wen model in a nutshell

## The Hamiltonian

$$H = - \sum_p P_p$$

- $H$  is a sum of local commuting projectors:  $[P_p, P_{p'}] = 0$
- $P_p$ : projector onto the vacuum of  $\mathcal{Z}(\mathcal{C})$  in the plaquette  $p$

$$P_p \quad \begin{array}{c} a \\ | \\ f-\zeta-\alpha-\beta-\gamma-\delta-\epsilon-e-d-c \\ | \\ b \end{array} = \sum_s \frac{d_s}{D^2} \sum' F_{s\alpha'\zeta'}^{a\zeta\alpha} F_{s\beta'\alpha'}^{b\alpha\beta} F_{s\gamma'\beta'}^{c\beta\gamma} F_{s\delta'\gamma'}^{d\gamma\delta} F_{s\epsilon'\delta'}^{e\delta\epsilon} F_{s\zeta'\epsilon'}^{f\epsilon\zeta}$$

$$\quad \begin{array}{c} a \\ | \\ f-\zeta'-\alpha'-\beta'-\gamma'-\delta'-\epsilon'-e'-d'-c \\ | \\ b \end{array}$$

- $d_s$ : quantum dimension of the string  $s \in \mathcal{C}$
- $D = (\sum_s d_s^2)^{1/2}$ : total quantum dimension of  $\mathcal{C}$

# Outline

1 The Levin-Wen model in a nutshell

2 Spectral degeneracies

3 Partition function

# Spectral degeneracies

On a genus- $g$  orientable compact surface (sphere, torus,...)

- Ground-state degeneracy:  $\mathcal{D}_0 = \sum_{i \in \mathcal{Z}(\mathcal{C})} \left( \frac{d_i}{D} \right)^\chi =$  Turaev-Viro invariant
- $D$ : total quantum dimension of  $\mathcal{Z}(\mathcal{C})$
- Euler-Poincaré characteristic:  $\chi = 2 - 2g$
- $g = 0$ :  $\mathcal{D}_0 = 1$
- $g = 1$ :  $\mathcal{D}_0 =$  Number of objects in  $\mathcal{Z}(\mathcal{C})$
- $g \geq 2$ :  $\mathcal{D}_0$  depends (non trivially) on  $\mathcal{Z}(\mathcal{C})$

Z. Kádár, A. Marzuoli, and M. Rasetti, Adv. Math. Phys. **2010**, 671039 (2010)

F. J. Burnell and S. H. Simon, Ann. Phys. **325**, 2550 (2010)

V. G. Turaev and O. Y. Viro, Topology **31**, 865 (1992)

G. Moore and N. Seiberg, Comm. Math. Phys. **123**, 177 (1989)

E. Verlinde, Nucl. Phys. B **300**, 360 (1988)

# Spectral degeneracies

Ex: Fibonacci category (commutative, braided, modular)

- Objects of  $\mathcal{C}$  :  $\{1, \tau\}$
- $d_1 = 1, d_\tau = \frac{1+\sqrt{5}}{2}$
- Objects of  $\mathcal{Z}(\mathcal{C})$  :  $\{(1,1), (1,\tau), (\tau,1), (\tau,\tau)\}$
- $d_{(i,j)} = d_i d_j$

M. D. Schulz, S. Dusuel, K. P. Schmidt, and J. Vidal, Phys. Rev. Lett. **110**, 147203 (2013)

Y. Hu, S. D. Stirling, and Y.-S. Wu, Phys. Rev. B **85**, 075107 (2012)

$g = 0$



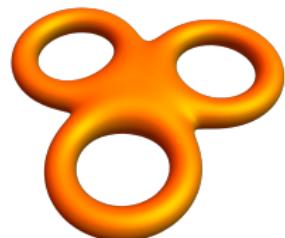
$g = 1$



$g = 2$



$g = 3$



$\mathcal{D}_0 = 1$

$\mathcal{D}_0 = 4$

$\mathcal{D}_0 = 25$

$\mathcal{D}_0 = 225$

# Spectral degeneracies

Ex:  $\text{Rep}(S_3)$  category (commutative, braided, non-modular)

- Objects of  $\mathcal{C}$  :  $\{1, 2, 3\}$
- $d_1 = 1, d_2 = 1, d_3 = 2$
- Objects of  $\mathcal{Z}(\mathcal{C})$ :  $\{A, B, C, D, E, F, G, H\}$
- $d_A = 1, d_B = 1, d_C = 2, d_D = 3, d_E = 3, d_F = 2, d_G = 2, d_H = 2$

A. Kitaev, Ann. Phys. 303, 2 (2003)

S. Beigi, P. W. Shor, D. Whalen, Commun. Math. Phys., 306, 663 (2011)

$g = 0$



$\mathcal{D}_0 = 1$

$g = 1$



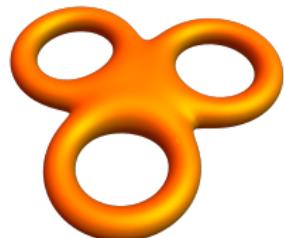
$\mathcal{D}_0 = 8$

$g = 2$



$\mathcal{D}_0 = 116$

$g = 3$



$\mathcal{D}_0 = 2948$

# Spectral degeneracies

Ex: Tambara-Yamagami category  $\text{TY}_3$  (commutative, non-braided)

- Objects of  $\mathcal{C} : \{1, 2, 3, \sigma\}$
- $d_1 = d_2 = d_3 = 1, d_\sigma = \sqrt{3}$
- Objects of  $\mathcal{Z}(\mathcal{C}) : \{\alpha_{i=1, \dots, 15}\}$
- $d_{i=1, \dots, 6} = 1, d_{i=7, \dots, 9} = 2, d_{i=10, \dots, 15} = \sqrt{3}$

D. Tambara and S. Yamagami, J. Algebra 209, 692 (1998)

S. Gelaki, D. Naidu, and D. Nikshych, Algebra Number Theory 3, 959 (2009)

$$g = 0$$



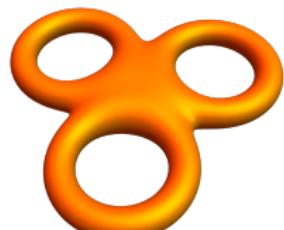
$$g = 1$$



$$g = 2$$



$$g = 3$$



$$\mathcal{D}_0 = 1$$

$$\mathcal{D}_0 = 15$$

$$\mathcal{D}_0 = 315$$

$$\mathcal{D}_0 = 8\,883$$

# Spectral degeneracies

Ex: Haagerup subfactor category  $\mathcal{H}_3$  (non-commutative)

- Objects of  $\mathcal{C}$  :  $\{1, \alpha, \alpha^*, \rho, \rho\alpha, \rho\alpha^*\}$
- $d_1 = d_\alpha = d_{\alpha^*} = 1, d_\rho = d_{\rho\alpha} = d_{\rho\alpha^*} = \frac{3+\sqrt{13}}{2}$
- Objects of  $\mathcal{Z}(\mathcal{C})$ :  $\{0, \mu^1, \mu^2, \mu^3, \mu^4, \mu^5, \mu^6, \pi^1, \pi^2, \sigma^1, \sigma^2, \sigma^3\}$
- $d_0 = 1, d_\mu = 3d_\rho, d_\pi = 3d_\rho + 1, d_\sigma = 3d_\rho + 2$

M. Asaeda and U. Haagerup, Comm. Math. Phys. **202**, 1 (1999)

S.-M. Hong, E. Rowell, and Z. Wang, Commun. Contemp. Math. **10**, 1049 (2008)

$$g = 0$$



$$\mathcal{D}_0 = 1$$

$$g = 1$$



$$\mathcal{D}_0 = 12$$

$$g = 2$$



$$\mathcal{D}_0 = 1401$$

$$g = 3$$



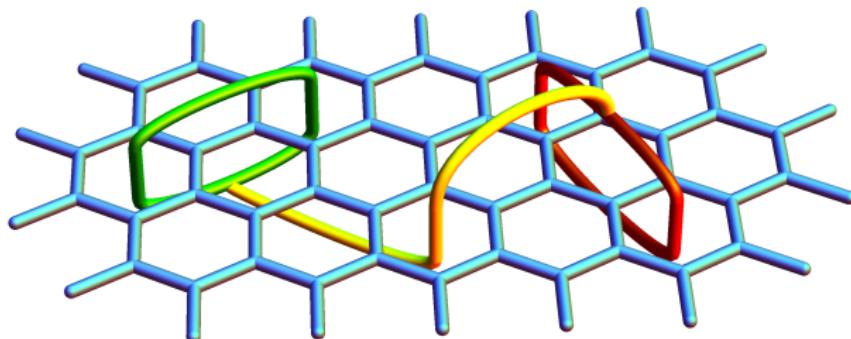
$$\mathcal{D}_0 = 1\,603\,329$$

# Spectral degeneracies

## Excitations of the Levin-Wen model (2005)

- Excitations are the nontrivial objects of  $\mathcal{Z}(\mathcal{C})$
- General construction of the Drinfeld center via the Ocneanu tube algebra
- Some objects only violate the plaquette constraints (fluxons)
- Some objects violate the vertex/plaquette constraints (not considered here !)

T. Lan and X.-G. Wen, Phys. Rev. B 90, 115119 (2014)



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## Old results for commutative $\mathcal{C}$

- Number of states with  $q$  fluxons on a genus- $g$  surface with  $N_p$  plaquettes:

$$\mathcal{D}_q = \binom{N_p}{q} \sum_{i \in \mathcal{Z}(\mathcal{C})} \left( \frac{d_i}{D} \right)^{\chi-q} \left( 1 - \frac{d_i}{D} \right)^q$$

G. Moore and N. Seiberg, Comm. Math. Phys. 123, 177 (1989)

J. Vidal, Phys. Rev. B 105, L041110 (2022)

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$$\mathcal{D}_q = \binom{N_p}{q} \sum_{i \in \mathcal{Z}(\mathcal{C})} \left( \frac{d_i}{D} \right)^{\chi-q} \left( n_i - \frac{d_i}{D} \right)^q$$

- $n_i$ : internal multiplicity of the particle  $i$  given by the tube algebra

C.-H. Lin, M. Levin, and F. J. Burnell, Phys. Rev. B **103**, 195155 (2021)  
A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, arXiv:2309.00343

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## Example of Morita equivalence: $\text{Rep}(S_3)$ and $\text{Vec}(S_3)$

- Same Drinfeld center  $\mathcal{Z}(S_3)$ :  $\{A, B, C, D, E, F, G, H\}$
- $\text{Rep}(S_3)$ :  $\mathbf{n} = (1, 0, 0, 1, 0, 1, 0, 0)$
- $\text{Vec}(S_3)$ :  $\mathbf{n} = (1, 1, 2, 0, 0, 0, 0, 0)$
- Different Hilbert space, different degeneracies,...

# Outline

- 1 The Levin-Wen model in a nutshell
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# Partition function

## Energy spectrum

- Levin-Wen Hamiltonian:  $H = - \sum_p P_p$
- Ground-state energy:  $E_0 = -N_p$
- $q$ -fluxon state energy:  $E_q = -N_p + q$

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## Exact finite-size and finite-temperature partition function

- Partition function:  $Z = \text{Tr}(e^{-\beta H}) = \sum_{q=0}^{N_p} D_q e^{-\beta E_q}$   $\left( \beta = \frac{1}{k_B T} \right)$

$$Z = \sum_{i \in \mathcal{Z}(C)} \left( \frac{d_i}{D} \right)^{\chi - N_p} \left[ n_i - \left( \frac{d_i}{D} \right) (1 - e^\beta) \right]^{N_p}$$

- $Z$  depends on the fusion rules:  $d_a \times d_b = \sum_c N_c^{ab} d_c$
- $Z$  depends on the surface topology:  $\chi = 2 - 2g$
- $Z$  depends on the number of plaquettes:  $N_p$

# Partition function

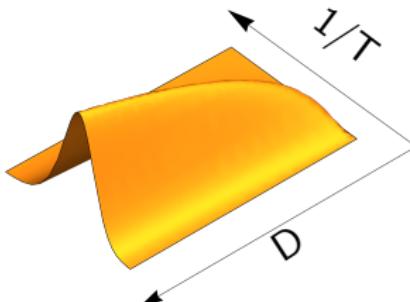
## Specific heat and absence of thermal phase transition

- Specific heat per plaquette :  $c = \frac{\beta^2}{N_p} \frac{\partial^2 \ln Z}{\partial \beta^2}$
- Thermodynamical limit ( $N_p \rightarrow \infty$ ):  $c = \frac{e^\beta \beta^2 (D - 1)}{(D - 1 + e^\beta)^2}$
- Only depends on the total quantum dimension  $D$

**No finite-temperature phase transition in the Levin-Wen model !**

J. Vidal, Phys. Rev. B 105, L041110 (2022)

A. Ritz-Zwilling, J.-N. Fuchs, S. H. Simon, and J. Vidal, in preparation



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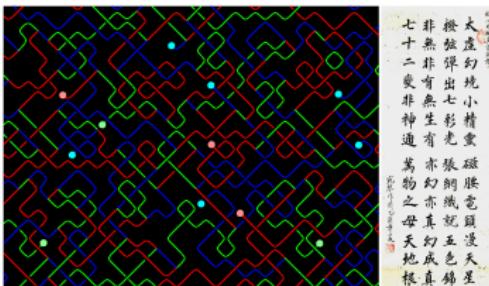
Absence of topological order at  $T > 0$  for any two-dimensional Hamiltonian which is a sum of local commuting projectors !

M. Hastings, Phys. Rev. Lett. 107, 210501 (2011)

# Outlook

## Summary and extensions

- Exact partition function of the Levin-Wen model
- More results available for:
  - any UFC with fusion multiplicities ( $N_{ab}^c \geq 1$ )
  - any topology with punctures (boundaries)
- Absence of finite-temperature phase transition
- Observables (e.g., Wegner-Wilson loops)
- Importance of charge excitations (violation of Gauss's law)



Courtesy of X.-G.Wen

## What's next ?

Forthcoming conference in Les Houches (french Alps): 1-12 April 2024

Topological order: Anyons and Fractons

Website: <https://topoanyons.sciencesconf.org/>