Appendix D

## D.1 Derivation of the Stiffness and Consistent mass matrices of a Euler Bernoulli beam.

The differential equation for an Euler-Bernoulli beam can be defined as follows,

$$\rho \frac{\delta^2 v}{\delta t^2} + \frac{\delta^2}{\delta x^2} \left( EI \frac{\delta^2 v}{\delta x^2} \right) = q(x, t)$$
 (D.1)

Where v(x,t) is the transverse displacement of the beam and q(x,t) is the externally applied pressure. The shear force V and bending moment M can be defined as follows,

$$V = -EI\left(\frac{\delta^3 v}{\delta x^3}\right) \tag{D.2}$$

$$M = EI\left(\frac{\delta^2 v}{\delta x^2}\right) \tag{D.3}$$

The finite element model of the beam uses a two noded bar element with two degrees of freedom per node, a transverse displacement and a rotation. As there are four degrees of freedom in the model it is assumed that a cubic polynomial can adequatly represent the displacement of the beam at a distance x from the left hand side. The displacement v is defined by,

$$v(x) = \alpha_0 + \alpha_1 x + \alpha_2 x^2 + \alpha_4 x^3$$
 (D.4)

The rotation is defined by,

$$\theta(x) = \frac{\delta v(x)}{\delta x} = \alpha_1 + 2\alpha_2 x + 3\alpha_3 x^2$$
 (D.5)

Defining a vector of displacements containing v and  $\theta$  as,

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$$\begin{bmatrix} v \\ \theta \end{bmatrix} = \begin{bmatrix} 1 & x & x^2 & x^3 \\ 0 & 1 & 2x & 3x \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{bmatrix}$$
 (D.6)

Applying the boundary conditions at nodes one and two, at node one x is equal to zero and at node two x is equal to L. The vector of elemental dispacements can be defined by,

$$\begin{bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & L & L^2 & L^3 \\ 0 & 1 & 2L & 3L^2 \end{bmatrix} \begin{bmatrix} \alpha_0 \\ \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{bmatrix}$$
 (D.7)

If the vector of elemental displacements is defined by d and the vector of coefficients by  $\alpha$ , equation (D.7) can be written as,

$$\{d\} = [A]\{\alpha\} \tag{D.8}$$

The vector of coefficients can then be solved by inverting the matrix A where,

$$[A]^{-1} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ -\frac{3}{L^2} & -\frac{2}{L} & \frac{3}{L^2} & -\frac{1}{L} \\ \frac{2}{L^3} & \frac{1}{L^2} & -\frac{2}{L^3} & \frac{1}{L^2} \end{bmatrix}$$
 (D.9)

$$\begin{bmatrix} \alpha_0 \\ \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ -\frac{3}{L^2} & -\frac{2}{L} & \frac{3}{L^2} & -\frac{1}{L} \\ \frac{2}{L^3} & \frac{1}{L^2} & -\frac{2}{L^3} & \frac{1}{L^2} \end{bmatrix} \begin{bmatrix} v_1 \\ \theta_1 \\ v_2 \\ \theta_2 \end{bmatrix}$$
(D.10)

Now substituting the coefficients from equation (D.10) into equation (D.4) gives,

$$v(x) = v_1 + \theta_1 x + \left( -\frac{3}{L^2} v_1 - \frac{2}{L} \theta_1 + \frac{3}{L^2} v_2 - \frac{1}{L} \theta_2 \right) x^2 + \left( \frac{2}{L^3} v_1 + \frac{1}{L^2} \theta_1 - \frac{2}{L^3} v_2 + \frac{1}{L^2} \theta_2 \right) x^3$$
(D.11)

Separating the varibales into  $v_1$ ,  $\theta_1$ ,  $v_2$  and  $\theta_2$  gives,

$$v(x) = v_1 \left( 1 - \frac{3x^2}{L^2} + \frac{2x^3}{L^3} \right) + \theta_1 \left( x - \frac{2x^2}{L} + \frac{x^3}{L^2} \right) + v_2 \left( \frac{3x^2}{L^2} - \frac{2x^3}{L^3} \right) + \theta_2 \left( \frac{-x^2}{L} + \frac{x^3}{L^2} \right)$$
(D.12)

Defining the displacement v by,

$$v(x) = N_1 v_1 + N_2 \theta_1 + N_3 v_2 + N_4 \theta_2$$
 (D.13)

Where N are the Hermite shape functions describing the behaviour of the beam,

$$N_1(x) = \left(1 - \frac{3x^2}{L^2} + \frac{2x^3}{L^3}\right)$$
 (D.14)

$$N_2(x) = \left(x - \frac{2x^2}{L} + \frac{x^3}{L^2}\right)$$
 (D.15)

$$N_3(x) = \left(\frac{3x^2}{L^2} - \frac{2x^3}{L^3}\right)$$
 (D.16)

$$N_4(x) = \left(\frac{-x^2}{L} + \frac{x^3}{L^2}\right)$$
 (D.17)

The strain at a particular point in the element can be defined by,

$$\varepsilon_{x} = \frac{\delta^{2} v}{\delta x^{2}} \tag{D.18}$$

The strain can be defined in terms of the shape functions and the vector of elemental displacements as,

$$\varepsilon_{x} = \frac{\delta^{2}}{\delta x^{2}} ([N]\{d\}) = \frac{\delta^{2}[N]}{\delta x^{2}} \{d\}$$
 (D.19)

Rewritting equation (D.19) as,

$$\varepsilon_{x} = [B]\{d\} \tag{D.20}$$

where [B] is the strain displacement matrix, relating nodal displacements to element strains. The strain in the element can now be written as,

$$\varepsilon_{x} = [B]\{d\} = \left\{ \frac{\delta^{2} N_{1}(x)}{\delta x^{2}} \quad \frac{\delta^{2} N_{2}(x)}{\delta x^{2}} \quad \frac{\delta^{2} N_{3}(x)}{\delta x^{2}} \quad \frac{\delta^{2} N_{4}(x)}{\delta x^{2}} \right\} \begin{cases} v_{1} \\ \theta_{1} \\ v_{2} \\ \theta_{2} \end{cases}$$
(D.21)

from the principle of minimum potential energy the elemental stiffnes matrix can be defined as,

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$$K^{e} = \int_{\sigma} [B]^{T} EI[B] dvol = EI \int_{0}^{L} [B]^{T} [B] dx$$
 (D.22)

$$K^{e} = EI \int_{0}^{L} \begin{cases} \frac{6}{L^{2}} - \frac{12x}{L^{3}} \\ \frac{4}{L} - \frac{6x}{L^{2}} \\ -\frac{6}{L^{2}} + \frac{12x}{L^{3}} \\ \frac{2}{L} - \frac{6x}{L^{2}} \end{cases} \begin{cases} \frac{6}{L^{2}} - \frac{12x}{L^{3}} & \frac{4}{L} - \frac{6x}{L^{2}} & \frac{-6}{L^{2}} + \frac{12x}{L^{3}} & \frac{2}{L} - \frac{6x}{L^{2}} \end{cases} dx$$
(D.23)

Check K(1,1),

$$K_{11}^e = EI \int_0^L \left(\frac{6}{L^2} - \frac{12x}{L^3}\right)^2 dx = \frac{12EI}{L^3}$$
 (D.24)

The stiffness matrix for the beam element can then be defined by,

$$K^{e} = \frac{EI}{L^{3}} \begin{bmatrix} 12 & 6L & -12 & 6L \\ 6L & 4L^{2} & -6L & 2L^{2} \\ -12 & -6L & 12 & -6l \\ 6L & 2L^{2} & -6L & 4L^{2} \end{bmatrix}$$
(D.25)

The consistent mass matrix can be defined by,

$$M^e = \int_0^L \rho A[N]^T [N] dx$$
 (D.26)

$$M^{e} = \rho A \int_{0}^{L} \left\{ 1 - \frac{3x^{2}}{L^{2}} + \frac{2x^{3}}{L^{3}} \\ x - \frac{2x^{2}}{L} + \frac{x^{3}}{L^{2}} \\ \frac{3x^{2}}{L^{2}} - \frac{2x^{3}}{L^{3}} \\ -\frac{x^{2}}{L} + \frac{x^{3}}{L^{2}} \right\} \left\{ \frac{3x^{2}}{L^{2}} - \frac{2x^{3}}{L^{3}} \\ \frac{-x^{2}}{L} + \frac{x^{3}}{L^{2}} \right\} dx$$
 (D.27)

Check M(1,1),

$$M_{44}^{e} = \rho A \int_{0}^{L} \left[ \frac{-x^{2}}{L} + \frac{x^{3}}{L^{2}} \right]^{2} dx = \frac{x^{5}}{5L^{2}} + \frac{x^{7}}{7L^{4}} - \frac{2x^{5}}{L^{3}} \Big|_{0}^{L} = \frac{\rho A L^{3}}{105} = \frac{4L^{3}\rho A}{420}$$
(D.28)

The consistent mass matrix for the beam element can now be defined by,

$$M_e = \rho A L / 420 \begin{pmatrix} 156 & 22L & 54 & -13L \\ 22L & 4L^2 & 13L & -3L^2 \\ 54 & 13L & 156 & -22L \\ -13L & -3L^2 & -22L & 4L^2 \end{pmatrix}$$
 (D.29)