kinematic equations	
uniform circular motion	
kintec and potential energy	

$$\bar{v} = \frac{1}{2}(v_i + v_f)$$

$$\Delta v = a\Delta t + v_i$$

$$x(t) = \frac{1}{2}at^2 + v_i t + x_i$$

$$v_f^2 - v_i^2 = 2a\Delta x$$

$$a = \frac{v^2}{r}$$

$$E_{\rm trans} = \frac{1}{2}mv^2$$

$$E_{\rm rot} = \frac{1}{2}I\omega^2$$

$$E_{\rm gpe} = mgh$$

$$E_{\rm sho} = \frac{1}{2}kx^2$$

potential energy, force, work
gravitational force
angular momentum



$$\vec{F} = -Gm_1m_2r^{-2}\hat{r}_{12}$$

$$\vec{L} = \vec{r} \times \vec{p} = I \vec{\omega}$$

torque
constant angular velocity
spherical elements

$$\tau = \vec{r} \times \vec{F} = \frac{d\vec{L}}{dt}$$

$$\vec{F} = -m\Omega^2 r \hat{r}_{\text{centrifuge}} - 2m\vec{\Omega} \times \vec{v}$$

$$\begin{split} d\vec{l} &= dr\hat{r} + rd\theta\hat{\theta} + r\sin\theta d\phi\hat{\phi} \\ d\vec{S} &= r^2\sin\theta d\theta d\phi\hat{r} + r\sin\theta drd\phi\hat{\theta} + rdrd\theta\hat{\phi} \\ dV &= r^2\sin\theta drd\theta d\phi \end{split}$$

moment of inertia	
Lagrangian	
Hamiltonian	

$$I=\int s^2dm$$

$$I=I_{\rm CM}+Ms^2$$

$$I_z=I_x+I_y \mbox{ for } z \mbox{ perpendicular to the body lying in the } x-y \mbox{ plane}$$

$$L(\dot{q}_i, q_i, t) = T - U$$
$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} = \frac{\partial L}{\partial q_i}$$
$$p_i = \frac{\partial L}{\partial \dot{q}_i}$$

$$\begin{split} H(p,q) &= \sum_i p_i \dot{q}_i - L \\ H &= T + U \text{ if } U \text{ is independent of } \dot{q} \text{ and } t \\ \dot{p} &= -\frac{\partial H}{\partial q} \\ \dot{q} &= \frac{\partial H}{\partial p} \end{split}$$

effective orbital potential	
Kepler's laws	
damped oscillation	

$$E = T + V = \frac{1}{2}m\dot{r}^2 + \frac{l^2}{2mr^2} + U(r)$$

$$l = mr^2\dot{\phi}$$

$$\mu = \frac{m_1m_2}{m_1 + m_2}$$

$$E > 0 \implies \text{hyperbolic}$$

$$E = 0 \implies \text{parabolic}$$

$$E < 0 \implies \text{elliptical}$$

- Planets move in elliptical orbits with one focus at the sun
- Equal orbital area sweeps out equal orbital time
- $T = ka^{\frac{3}{2}}$ for all planets

$$0 = m\ddot{x} + b\dot{x} + kx$$

$$\beta = \frac{b}{2m}$$

$$\omega_0 = (\frac{k}{m})^{\frac{1}{2}}$$

- $\beta^2 > \omega_0^2$: overdamped decays exponentially
- $\beta^2 = \omega_0^2$: critically damed
- $\beta^2<\omega_0^2$: underdamped decays exponentially, modulated by a sinusoid at $\omega_1^2=\omega_0^2-\beta^2$

driven oscillation
oscillators
fluid dynamics

$$\begin{split} \omega_{\rm R} &= (\omega_0^2-2\beta^2)^{\frac{1}{2}} \\ D &\propto \left|\omega_0^2-\omega^2\right|^{-1} \text{ for undamped, driven oscillator} \end{split}$$

- pendulum $\omega = (\frac{g}{l})^{\frac{1}{2}}$
- physical pendulum $\omega = (\frac{mgR}{I})^{\frac{1}{2}}$

$$\frac{v^2}{2} + gz + \frac{p}{\rho} = \text{const}$$

p pressure, ρ fluid density, z height of a point along streamline $F = \rho V g$

Maxwell's equations
electrostatics
boundary conditions

$$\begin{split} \vec{\nabla} \cdot \vec{E} &= \frac{\rho}{\epsilon_0} \\ \vec{\nabla} \times \vec{E} &= \frac{-\partial \vec{B}}{\partial t} \\ \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{B} &= \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \\ \frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} &= 0 \\ \vec{E} &= -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t} \\ \vec{B} &= \vec{\nabla} \times \vec{A} \end{split}$$

$$\vec{E} = -\vec{\nabla}\phi \implies \phi = -\int_a^b \vec{E} \cdot d\vec{l}$$

$$\nabla^2 \phi = -\frac{\rho}{\epsilon_0} \implies \phi = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3 \vec{r}$$

$$\int_S \vec{E}(\vec{r}) \cdot d\vec{S} = \frac{Q_{\rm enc}}{\epsilon_0}$$

$$\int_C \vec{E}(\vec{r}) \cdot d\vec{l} = 0$$

Electrostatics

$$\begin{split} E_{\mathrm{out}}^{\parallel} - E_{\mathrm{in}}^{\parallel} &= 0 \\ E_{\mathrm{out}}^{\perp} - E_{\mathrm{in}}^{\perp} &= \frac{\sigma}{\epsilon_0} \end{split}$$

 ϕ is continuous

 $\partial \phi$ continuous when no surface charge is present

Magnetostatics

$$B_{\mathrm{out}}^{\parallel} - B_{\mathrm{in}}^{\parallel} = \mu_0 \vec{k} \times \hat{n} B_{\mathrm{out}}^{\perp} - B_{\mathrm{in}}^{\perp} \qquad \qquad = 0$$

 \vec{k} is the surface current density

conductors	
work and energy in (electro/magneto)statics	
capacitors	

 $\phi = {\rm const.}$ throughout conductor

 \implies E-field inside is zero

 $\implies Q_{\text{inside}} = 0$

 $\implies Q_{\rm net}$ confined to surface

 \implies E-field outside is perpendicular to surface

$$W = \frac{1}{2} \int \rho \phi d^3 r$$

$$U_E = \frac{\epsilon_0}{2} \int |\vec{E}|^2 d^3 r$$

$$U_B = \frac{1}{2\mu_0} \int |\vec{B}|^2 d^3 r$$

$$\begin{split} Q &= C\phi \\ C_{\text{parallel plate}} &= \frac{\epsilon_0 A}{d} \\ U &= \frac{1}{2} Q^2 C^{-1} = \frac{1}{2} C \phi^2 \end{split}$$

magnetostatics
magnetostatic fields
cyclotron

$$\int_{S} \vec{B} \cdot d\vec{s} = 0$$

$$\int_{C} \vec{B} \cdot d\vec{l} = \mu_{0} I_{\text{enc}}$$

$$\vec{F_{B}} = q\vec{v} \times \vec{B}$$

$$\vec{B}(\vec{r}) = \frac{\mu_{0} I}{4\pi} \int \frac{d\vec{l} \times \vec{r'}}{|\vec{r'}|^{3}}$$

$$\vec{B} = \mu_0 \frac{N}{L} I \hat{z} \quad \text{(solenoid)}$$

$$\vec{B} = \frac{\mu_0 N I}{2\pi r} \hat{\phi}$$

There is no field outside a toroid or solenoid.

$$\vec{F} = q\vec{v} \times \vec{B} = qvB\hat{x}$$

$$R = \frac{mv}{qB}$$

inductance	
inductance configurations	
multipoles	

$$\mathcal{E} = \int_{C} \vec{E} \cdot d\vec{l} = \frac{-\partial \Phi_{B}}{\partial t}$$

$$\Phi_{21} = M_{12}I_{1}$$

$$\Phi_{B} = LI$$

Induced curents oppose changes in magnetic flux

$$L = \frac{\mu_{0N^2A}}{l} \quad \text{solenoid}$$

$$\vec{p} = \int \vec{r} \rho(\vec{r}) d^3r$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \vec{r}}{|\vec{r}|^3}$$

$$\tau = \vec{p} \times \vec{E}$$

$$U = -\vec{p} \cdot \vec{E}$$

$$\vec{m} = I\vec{A}$$

$$\tau = \vec{m} \times \vec{B}$$

$$U = -\vec{m} \times \vec{B}$$

em in media	
em waves	
radiation	

$$\sigma_b = \vec{P} \cdot \hat{n}$$
$$\rho_b = -\vec{\nabla} \cdot \vec{P}$$
$$\epsilon = \kappa \epsilon_0$$

$$\begin{split} \vec{E}(\vec{r}) &= E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \hat{n} \\ \vec{B}(\vec{r}) &= B_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} (\hat{k} \times \hat{n}) \\ \vec{S} &= \frac{1}{\mu_0} (\vec{E} \times \vec{B}) = \frac{1}{2\mu_0} \mathrm{Re}[\vec{E} \times \vec{B}^*] \\ I &= \langle S \rangle = \frac{c\epsilon_0}{2} E_0^2 \end{split}$$

$$P = \frac{\mu_0 q^2 a^2}{6\pi c} \quad \text{accelerating point charge}$$

$$\langle S \rangle = \frac{\mu_0 p_0^2 \omega^4}{32\pi^2 c} \frac{\sin^2(\theta)}{r^2}$$

$$\langle P_E \rangle = \frac{\mu_0 p_0^2 \omega^4}{12\pi c}$$

$$\langle P_B \rangle = \frac{\mu_0 m_0^2 \omega^4}{12\pi c^3}$$

The bottom three equations are for a radiating dipole.

circui	its
wave equ	ation
wave config	curations

$$V_R = IR$$

$$V_C = \frac{Q}{C}$$

$$V_L = \frac{L}{\ddot{Q}}$$

$$R_{\text{series}} = R_i$$

$$\frac{1}{C_{\text{series}}} = \frac{1}{C_i}$$

$$L_{\text{series}} = L_i$$

$$\frac{1}{R_{\text{parallel}}} = \frac{1}{R_i}$$

$$C_{\text{parallel}} = C_i$$

$$\frac{1}{L} = \frac{1}{L}$$

$$\frac{L_{\text{parallel}}}{L_{\text{parallel}}} = \frac{L_{i}}{L_{i}}$$

$$\frac{1}{\overline{L_{\text{parallel}}}} = \frac{1}{L_i}$$

$$R = \frac{\rho l}{A}$$

$$\partial_t^2 \psi = v^2 \partial_x^2 \underline{\psi} IV$$

$$\psi(\vec{r}, t) = A \cos(\vec{k} \cdot \vec{l} \cdot \vec{r} \cdot \vec{c} \cdot \vec{v} \cdot t + \phi)$$

$$= \text{Re}[Ae^{i(\vec{k} \cdot \vec{r} - \omega t)}] v_{\text{phase}} = \frac{U}{N} L = \frac{1}{2}LI^2$$

$$v_{\text{group}} = \frac{\partial u}{\partial k} I = \frac{V}{N} (1 - e^{-\frac{t}{L/R}}) \quad \text{RL circuit}$$

$$\frac{\omega}{k} = \frac{c}{n} \lambda \rightarrow \text{Re} \frac{1}{N} \lim_{k \to \infty} \frac{1}{RC} \quad \text{RC circuit}$$

$$\lambda = \frac{2\pi}{Rk} C = (LC)^{-\frac{1}{2}}$$

- circuit elements=in-series see same current $\overset{2\pi}{\omega}$
- circuit elements=in2ntrallel see same voltage
- current in/out of a node is conserved
- voltage is conserved over elements around a closed loop

$$v=(rac{T}{\mu})^{rac{1}{2}}$$
 string w/ tension T and mass density μ $v=(rac{\kappa}{
ho})^{rac{1}{2}}$ sound w/ bulk modulus κ and density ho

polarization
interference
geometric optics

$$I = I_0 \cos^2(\theta)$$
$$\theta_B = \arctan(\frac{n_2}{n_1})$$

 $d\sin(\theta) = m\lambda$ max dub-slit

 $d\sin(\theta)=(m+\frac{1}{2})\lambda$ min dub-slit

 $a\sin(\theta)=m\lambda$ min single-slit

 $D\sin(\theta) = 1.22\lambda$ first circular diffraction minimum $d\sin(\theta) = \frac{n\lambda}{2}$ x-rays on crystal maxima

 $n_2 > n_1 \implies \pi$ phase shift

 $n_2 < n_1 \implies 0$ phase shift

$$\theta_i = \theta_r n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$

$$\frac{1}{s} + \frac{1}{s'} = \frac{1}{f}$$

$$f = \frac{R}{2} \text{ spherical mirror}$$

$$\frac{1}{f} = (n-1)(\frac{1}{R_1} - \frac{1}{R_2})$$

$$m = \frac{-s'}{s}$$

misc optics
basic stat mech
entropy

$$I \propto I_0 \lambda^{-4} a^6$$
$$f = (\frac{v + v_r}{v - v_s}) f_0$$

$$p_{i} = \frac{e^{-\beta E_{i}}}{Z}$$

$$Z = \sum_{j} e^{-\beta E_{j}}$$

$$\beta = (k_{B}T)^{-1}$$

$$\langle O \rangle = \sum_{i} p_{i}O_{i}$$

$$\langle E \rangle = -\partial_{\beta} \ln Z$$

$$Z_{N} = \frac{1}{N!h^{3N}} \int e^{-\beta H(\vec{p}_{1}, \dots \vec{p}_{n}; \vec{x}_{1}, \dots \vec{x}_{n})} d^{3}p_{1} \dots d^{3}p_{n} d^{3}x_{1} \dots d^{3}x_{n}$$

$$k_{B} = 1.380649e - 23JK^{-1}$$

$$\begin{split} S &= k_B \mathrm{ln} \Omega \\ S &= -k_B \sum_i p_i \mathrm{ln} p_i = \partial_T (k_B T \mathrm{ln} Z) \\ S &= N k_B \mathrm{ln} \frac{V T^{3/2}}{N} + \mathrm{const.~(monoatomic~ideal~gas)} \end{split}$$

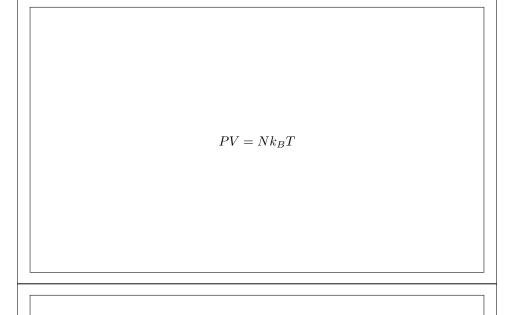
equipartition theorem	
misc combinatorics	
thermodynamic laws	

Each quadratic term (d.o.f.) in the Hamiltonian for a particle contributes $k_BT/2$ to the internal energy of a particle

$$(N, M) = \frac{N!}{(N - M)!M!}$$
$$\ln n! = n \ln n - n \ n \text{ large}$$

- There does not exist a process in which the sole effect is to transfer heat from a body at lower temperature to a body at higher temperature $\Delta S \geq \int \frac{\delta Q}{T}$
- Entropy is zero at absolute zero temperature. There is only one microstate at absolute zero temperature

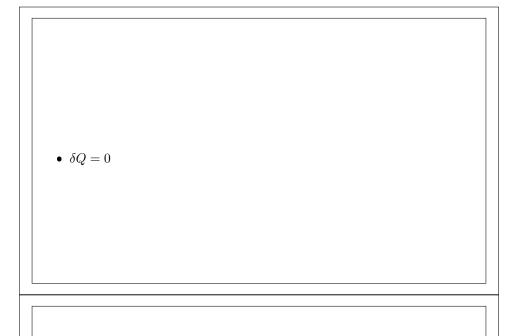
ideal gas law
reversible process
quasistatic



- $\bullet\,$ system is in equilibrium at every instant
- $\delta W = PdV$
- $\bullet \ \delta Q = TdS$

 \bullet in thermal equilibrium at every instant, but not necessarily reversible

adiabatic	
isentropic / reversible adiabatic	
IsoX	



- $\bullet\,$ reversible and adiabatic
- $\bullet \ \iff \Delta S = 0$
- $PV^{\gamma} = \text{const.}$; $\gamma = \frac{C_P}{C_V}$

• X held constant

free expansion	
thermodynamic identities	
heat capacity	

•
$$\Delta T = 0$$

- \bullet irreversible
- $\bullet \ PV = P'V'$

$$dU = TdS - PdV$$

$$T = \frac{\partial U}{\partial S}|_{V}$$

$$P = -\frac{\partial U}{\partial V}|_{S}$$

$$\frac{\partial P}{\partial S}|_{V} = -\frac{\partial T}{\partial V}|_{S}$$

$$C_V = \frac{\partial Q}{\partial T}|_V = \frac{\partial U}{\partial T}$$

$$C_P = \frac{\partial Q}{\partial T}|_P$$

$$C_P - C_V = Nk_B \text{ (ideal gas)}$$

$$Q = mc\Delta T$$

$$c_{\text{water}} = 4.18 \text{ JK}^{-1}\text{g}^{-1}$$

model systems	
monoatomic ideal gas	
quantum stat mech	

$$e = 1 - |\frac{Q_C}{Q_H}|$$

$$e = 1 - \frac{T_C}{T_H} \text{ theoretical max, carnot}$$

- use thermodynamic identities to calculate areas
- $\bullet\,$ clockwise paths in P V and T S planes do positive work

$$Z_N = \frac{V^N}{N!h^{3N}} (2\pi m k_B T)^{3^N/2}$$

$$U = \frac{3}{2} N k_B T$$

$$v_{\rm rms} = (\frac{3k_B T}{m})^{1/2}$$

$$K = \gamma P$$

$$F_{\text{FD}}(E_i) = \frac{1}{e^{(E_i - \mu)/k_B T} + 1}$$
$$F_{\text{BE}}(E_i) = \frac{1}{e^{(E_i - \mu)/k_B T} - 1}$$
$$\langle N \rangle = \sum_i g(E_i) F(E_i)$$