

NEXT-GENERATION ALL-SOLID-STATE BATTERY (#ASSB)

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Mathematical modelling for the next-generation All-solid-state batteries: Nucleation (SE|SSE)^(*)-interface

Rechargeable Lithium-ion battery (LIB) is at the heart of every electric vehicle (EV), portable electronic device, and energy storage system [1]. Nowadays, LIBs enable human life more efficient and help to solve global environment issues thanks to EVs' zero emission. However, conventional LIB (c-LIB) is sensible to temperature and pressure, hence, flammable and explosive, which is undesirable. This bottleneck is mainly due to **liquid-based electrolyte** found in c-LIBs.

All-solid-state battery (ASSB) is one of promising candidates to overcome bottlenecks of c-LIBs. Thanks to **solid-state electrolyte** (SSE), ASSB is highly stable towards temperature and pressure. Nevertheless, Li-metal dendrite triggered at (SE|SSE)-interface is the main drawback of ASSB since these dendritic threads extrapolate into SSE grain boundary network, causing crevice, degradation of ionic conductivity, and the probability of short-circuit, which is unfavorable.

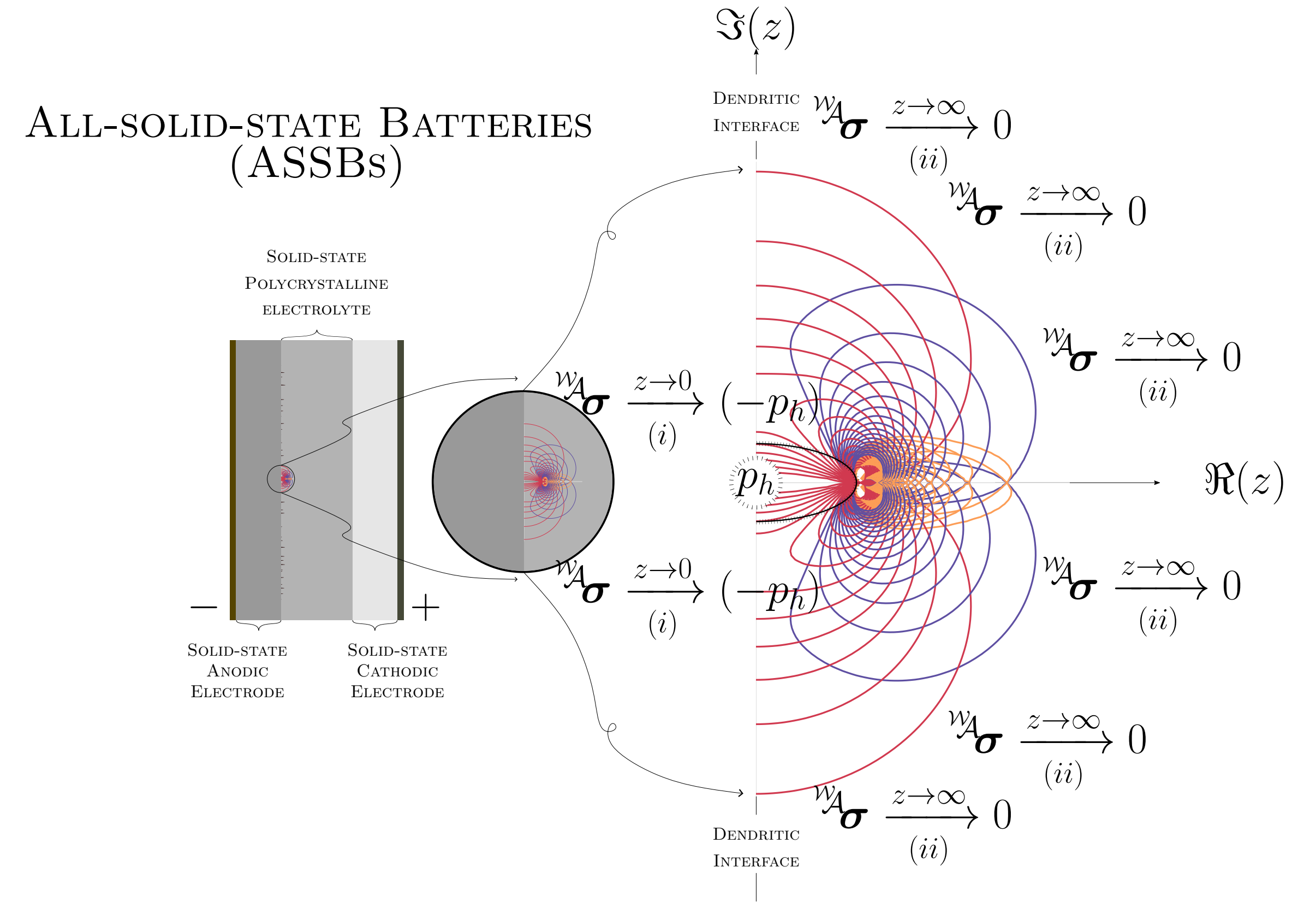
Next-generation All-solid-state battery (ng-ASSB) with a consideration of **nucleation criterion** defined by

$$a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}}$$

where, can help to improve ASSB performance.

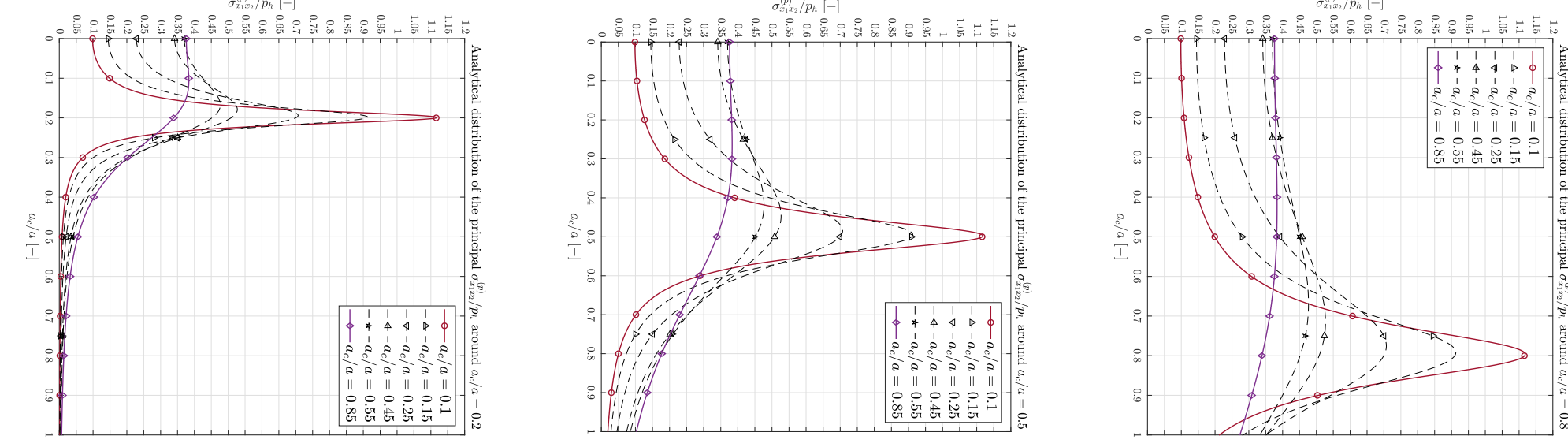
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Interface

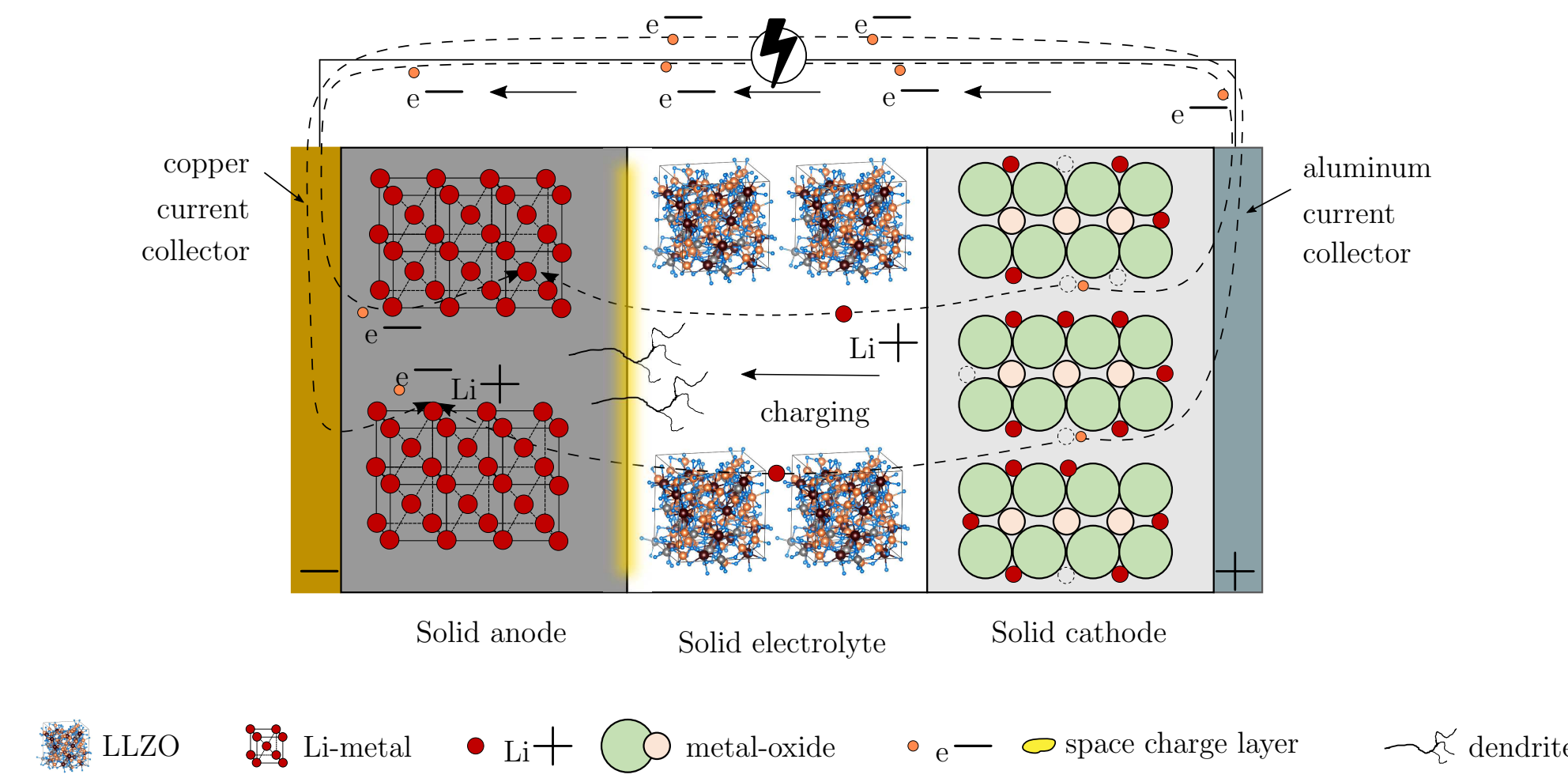
Interface between solid electrode and solid-state electrolyte (SE|SSE) taking place at space charge layer (SCL) [2] found in ASSBs critically exhibits mechanical and electrochemical instability [3]. This evidence points directly to the fact that the soft metallic li anode is erroneously prone to triggering dendrites, under cycles of electric charge & discharge [4].



Distribution: ana. max. shear stress γ_{σ}^{Π} around crack tip a_c .

Next-generation All-solid-state battery

Nucleation taking place at critical dendritic (SE|SSE)-interface

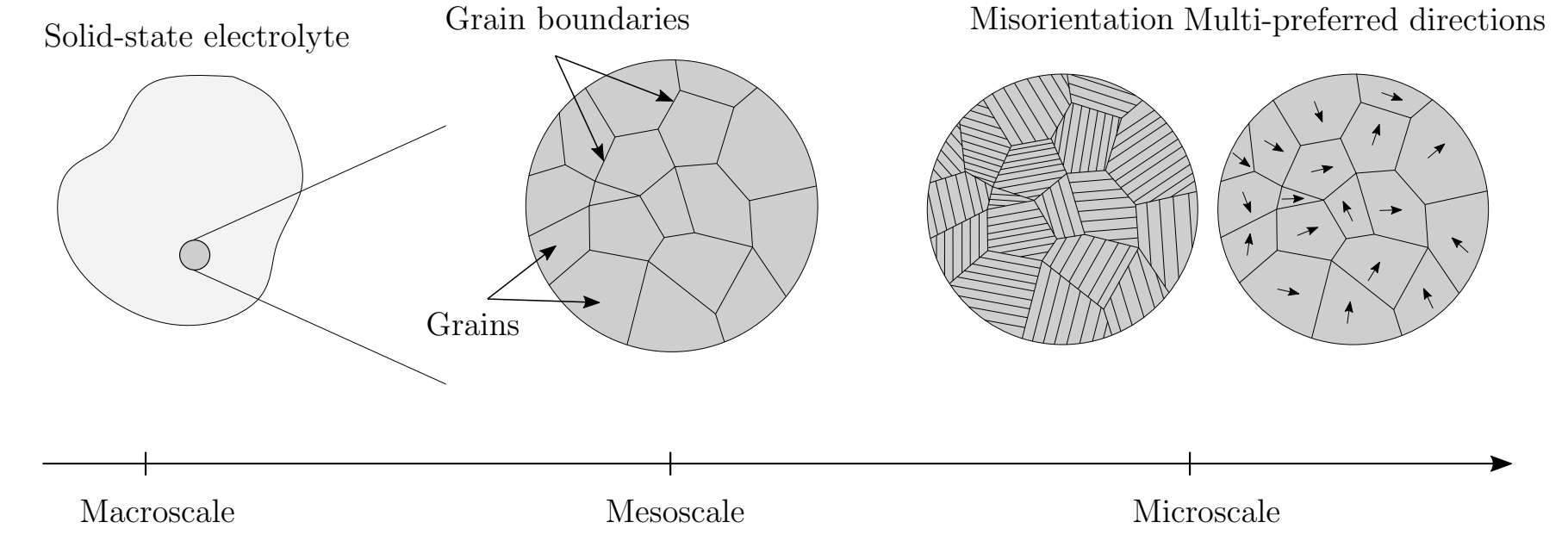


Nucleation taking place at critical dendritic (SE|SSE)-interface

Nucleation taking place at critical dendritic (SE|SSE)-interface

Embedded structural-tensor SSE

Polycrystalline garnet-typed SSE such as LLZO exhibit a network of grain boundaries, and grains with various sizes and shapes under microscopic observation. Therefore, this type of microstructure is potentially prone to nuance destruction of ceramic-like materials.



Consequently, dendrites contribute to degradation of ionic conductivity and cracks via tracing along grain boundaries.

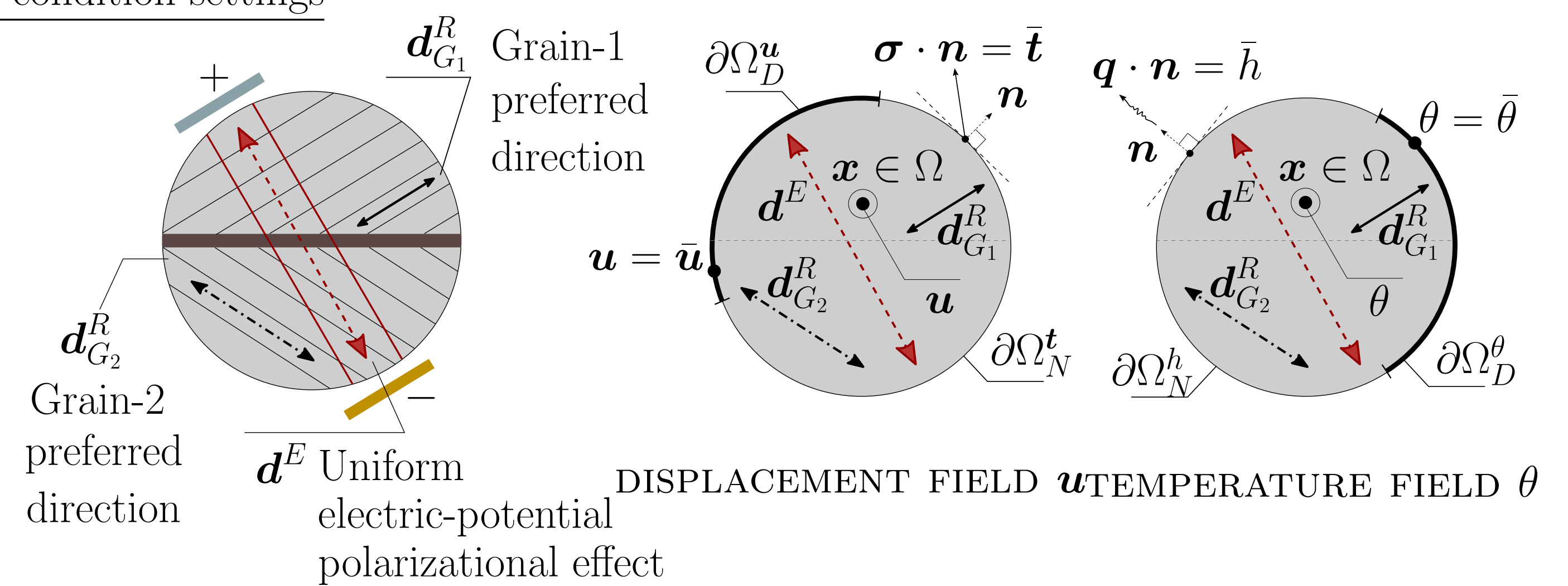
Nucleation interface: Taking place at the critical dendritic interface

Boundary condition settings

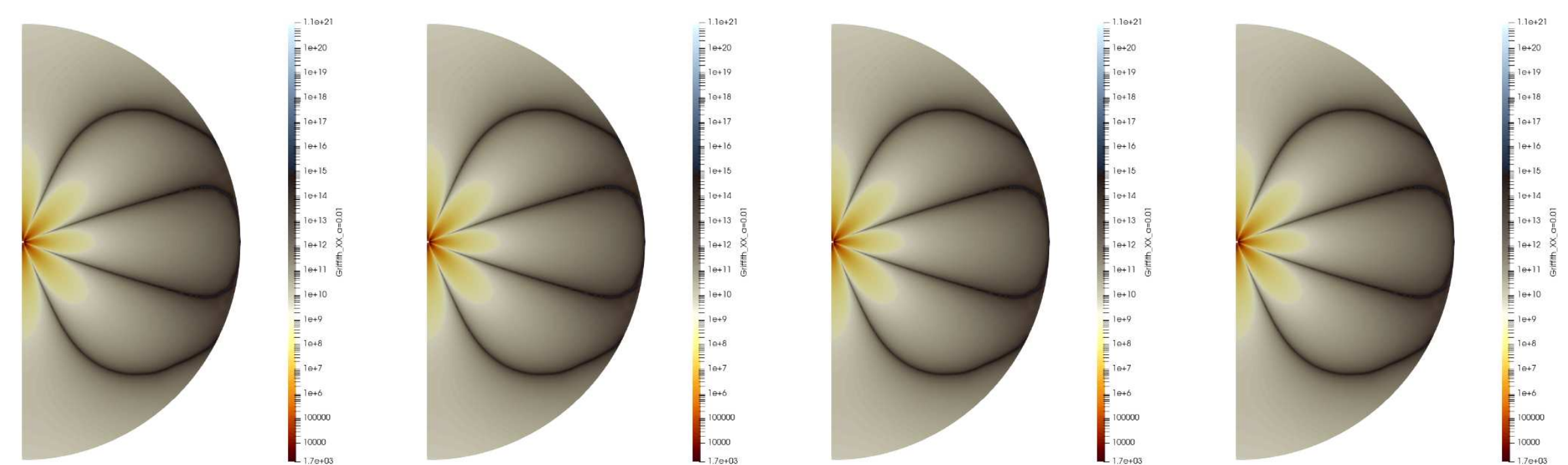
$$\begin{aligned} & \rho \partial_{t^2}^2 \mathbf{u}^{(s)} + \nabla \cdot \left(\mathbb{C} f_{(\lambda, \mu)}^{D(\Omega)} : \nabla \mathbf{u}^{(s)} \right) + \rho \nabla V_e = \mathbf{0}, \\ \text{s.t. } & a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}} \\ & \rho \partial_{t^2}^2 \mathbf{u}^{(s)} + \nabla \cdot \left(\mathbb{C} f_{(\lambda, \mu)}^{D(\Omega)} : \nabla \mathbf{u}^{(s)} \right) + \rho \nabla V_e = \mathbf{0}, \\ \text{s.t. } & a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}} \\ & \rho \partial_{t^2}^2 \mathbf{u}^{(s)} + \nabla \cdot \left(\mathbb{C} f_{(\lambda, \mu)}^{D(\Omega)} : \nabla \mathbf{u}^{(s)} \right) + \rho \nabla V_e = \mathbf{0}, \\ \text{s.t. } & a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}} \\ & \rho \partial_{t^2}^2 \mathbf{u}^{(s)} + \nabla \cdot \left(\mathbb{C} f_{(\lambda, \mu)}^{D(\Omega)} : \nabla \mathbf{u}^{(s)} \right) + \rho \nabla V_e = \mathbf{0}, \\ \text{s.t. } & a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}} \\ & \rho \partial_{t^2}^2 \mathbf{u}^{(s)} + \nabla \cdot \left(\mathbb{C} f_{(\lambda, \mu)}^{D(\Omega)} : \nabla \mathbf{u}^{(s)} \right) + \rho \nabla V_e = \mathbf{0}, \\ \text{s.t. } & a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}} \end{aligned}$$

Therefore

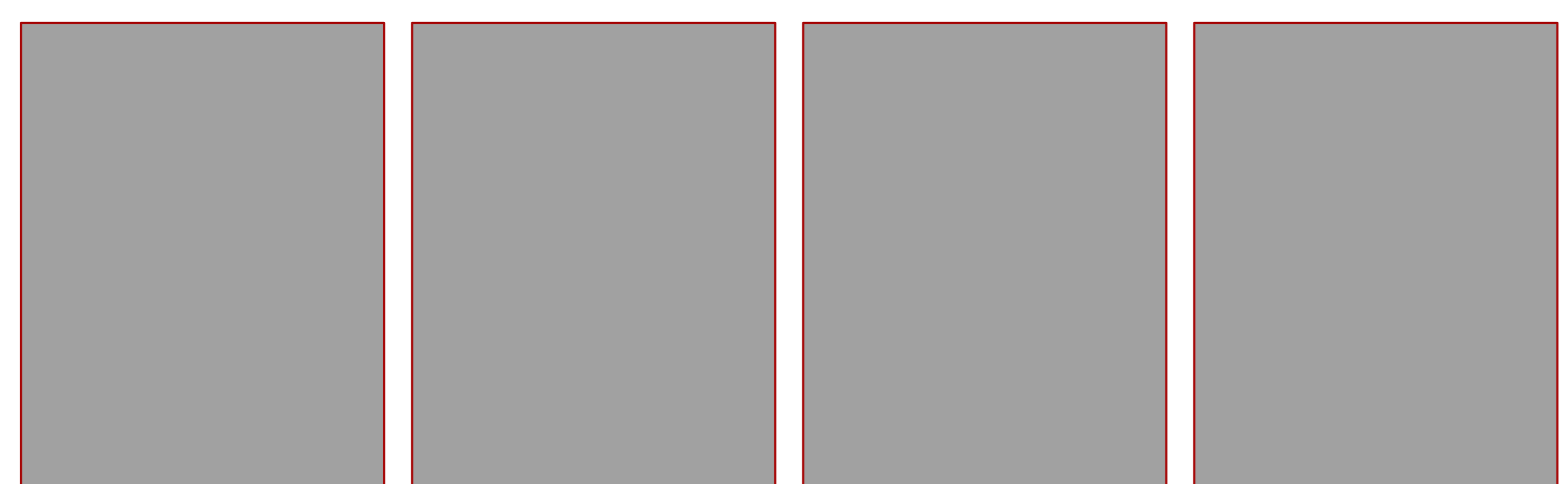
$$\therefore a_{\text{Griffith}} := a^* = \arg \min_{a \in \mathbb{R}} \iint_{\Omega} f(a, \mathbf{u}; \lambda, \mu, \mathbf{d} \otimes \mathbf{d}) d\Omega - \iint_{\Gamma} f(a; \gamma) d\Gamma \Big|_{\mathbf{u}^{(s)}}$$



The set of boundary conditions is likewise the path of the pressure-centric dendritic crack.

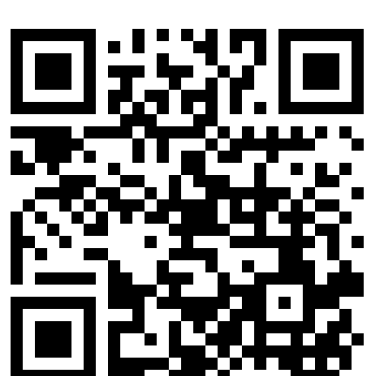


Comparison: Analytical vs. Numerical solutions



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References

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