Next-generation all-solid-state battery (#ASSB)

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Mathematical modelling for the next-generation All-solid-state batteries: Nucleation (SE|SSE)^(*)-interface

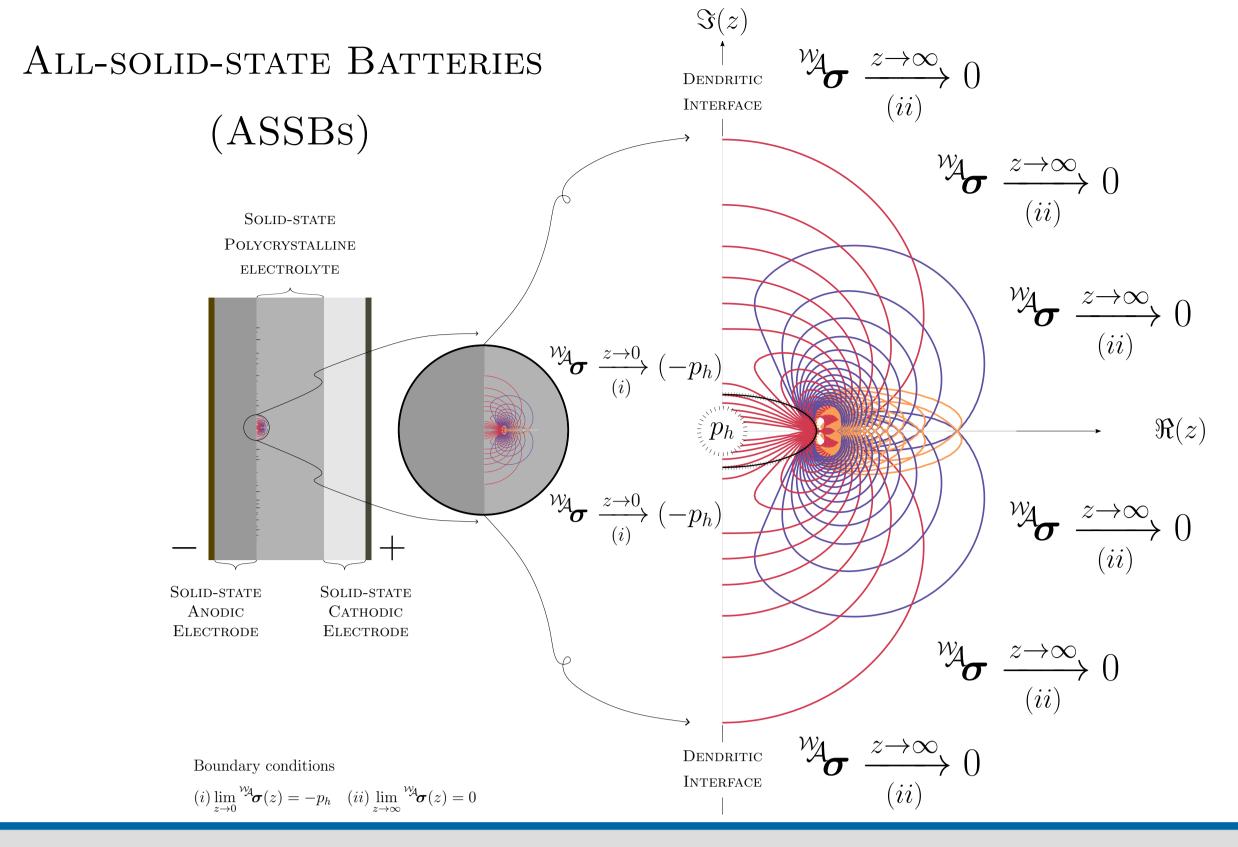
Rechargeable Lithium-ion battery (LIB) is at the heart of every electric vehicle (EV), portable electronic device, and energy storage system [1]. Nowadays, LIBs enable human life more efficient and help to solve global environment issues thanks to EVs' zero emission. However, conventional LIB (c-LIB) is sensible to temperature and pressure, hence, flammable and explosive, which is undesirable. This bottleneck is mainly due to liquid-based electrolyte found in c-LIBs.

All-solid-state battery (ASSB) is one of promising candidates to overcome bottlenecks of c-LIBs. Thanks to solid-state electrolyte (SSE), ASSB is highly stable towards temperature and pressure. Nevertheless, Limetal dendrite triggered at (SE|SSE)-interface is the main drawback of ASSB since these dendritic threads extrapolate into SSE grain boundary network, causing crevice, degradation of ionic conductivity, and the probability of short-circuit, which is unfavorable.

Next-generation All-solid-state battery (ng-ASSB) with a consideration of nucleation criterion defined by

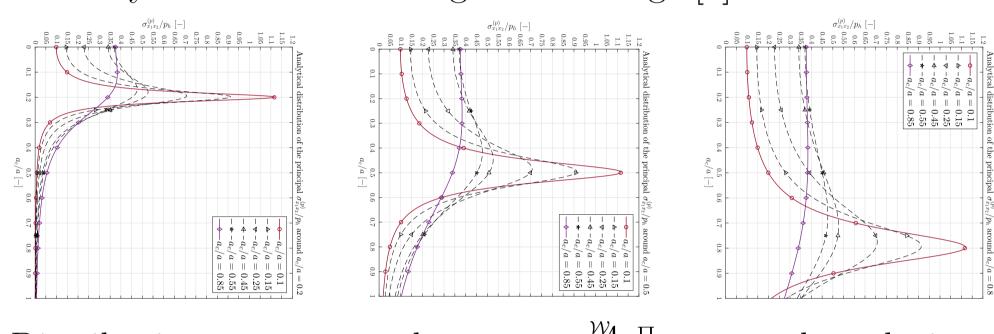
$$a_{ ext{Griffith}} := a^* = \arg\min_{a \in \mathbb{R}} \left. \iint_{\Omega} f(a, \boldsymbol{u}, \theta; \lambda, \mu, \boldsymbol{d} \otimes \boldsymbol{d}) \, d\Omega - \iint_{\Gamma} f(a; \gamma) \, d\Gamma \right|_{\boldsymbol{u}}$$

where \boldsymbol{u} displacement field, θ temperature field, a crevice length, λ, μ Lamé constants, $\boldsymbol{d} \otimes \boldsymbol{d}$ embedded misorientation structural tensor, and γ cracking-surface energy density, can help to improve ASSB performance.



Interface Analysis

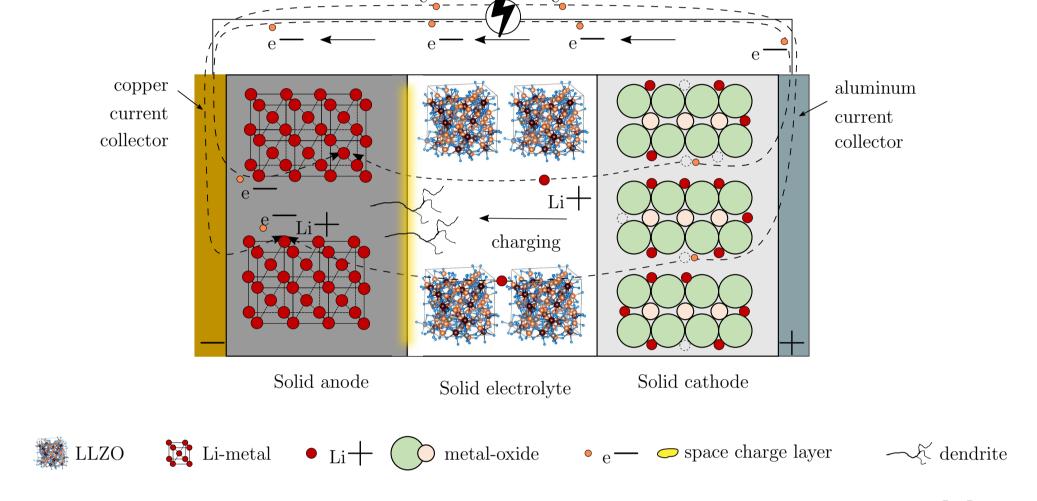
Interface between solid electrode and solid-state electrolyte (SE|SSE) taking place at space charge layer (SCL) [2] found in ASSBs critically exhibits mechanical and electrochemical instability [3]. This evidence points directly to the fact that the soft metallic li anode is erroneously prone to triggering dendrites, under cycles of electric charge & discharge [5].



<u>Distribution</u>: ana. max. shear stress ${}^{\mathcal{W}}\!\!\sigma_{x_1x_2}^{\Pi}$ around crack tip a_c .

Next-generation All-solid-state battery

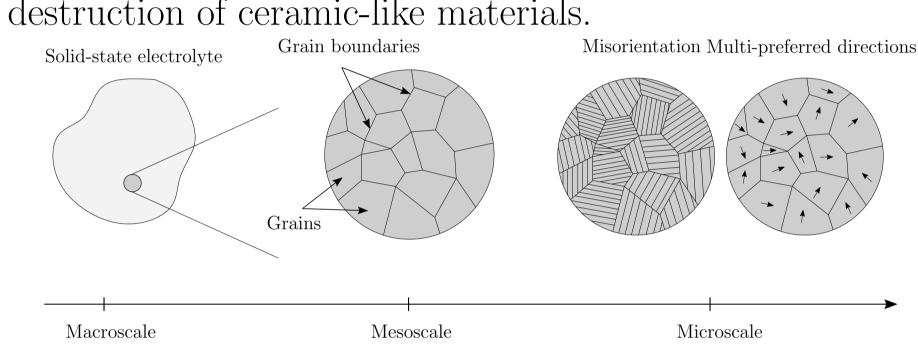
Nucleation criterion governs the instable (SE|SSE)-interface [3]



- **Thermodynamic consistency** is satisfied, followed by [2].
- ✓ Closure problem is fulfilled by 15 moments, followed by [4].

Embedded structural-tensor SSE

Polycrystalline garnet-typed SSE such as LLZO exhibit a network of grain boundaries, and grains with various sizes and shapes under microscopic observation. Therefore, this type of microstructure is potentially prone to nuance destruction of ceramic-like materials.



Consequentially, dendrites contribute to degradation of ionic conductivity and cracks via tracing along grain boundaries.

Nucleation interface: Taking place at the critical dendritic interface

Coupled fields: Displacement vector field and temperature scalar field

$$\boldsymbol{u}: \begin{cases} \Omega \times \mathbb{R}_{+} \to \mathbb{R}^{3}, \\ (\boldsymbol{x},t) \mapsto \boldsymbol{u}(\boldsymbol{x},t), \end{cases} \quad \theta: \begin{cases} \Omega \times \mathbb{R}_{+} \to \mathbb{R}, \\ (\boldsymbol{x},t) \mapsto \theta(\boldsymbol{x},t), \end{cases} \quad \theta: \begin{cases} \Omega \times \mathbb{R}_{+} \to \mathbb{R}, \\ (\boldsymbol{x},t) \mapsto \theta(\boldsymbol{x},t), \end{cases}$$

Governing conservation equations

$$\frac{d}{dt} \int_{\Omega}^{\Omega} (\cdot) \ d\Omega = \int_{\Omega}^{\Omega} (\cdot)^{\text{action}} \ d\Omega + \int_{\partial \Omega}^{\Omega} (\cdot)^{\text{action}} \ d\partial\Omega + \int_{\Omega}^{\Omega} (\cdot)^{\text{production/source/sink}} \ d\Omega$$

 $\rho(\boldsymbol{x},t)$ is mass density per unit volume (puv); $\boldsymbol{b}(\boldsymbol{x},t)$ body force puv; $\boldsymbol{v}(\boldsymbol{x},t)$ velocity; $e(\boldsymbol{x},t)$ internal energy puv; $\boldsymbol{q}(\boldsymbol{x},t)$ heat flux; $r(\boldsymbol{x},t)$ heat source puv; $\boldsymbol{\sigma}$ Cauchy stress and $\boldsymbol{\varepsilon}$ infinitesimal strain. Helmholtz energy functional

$$a_{ ext{Griffith}} := a^* = \arg\min_{a \in \mathbb{R}} \left. \iint_{\Omega} f(a, \boldsymbol{u}; \lambda, \mu, \boldsymbol{d} \otimes \boldsymbol{d}) \, d\Omega - \left. \iint_{\Gamma} f(a; \gamma) \, d\Gamma \right|_{\boldsymbol{u}^{(s)}}$$

Governing PDE

$$a_{\text{Griffith}} := a^* = \arg\min_{a \in \mathbb{R}} \left. \iint_{\Omega} f(a, \boldsymbol{u}; \lambda, \mu, \boldsymbol{d} \otimes \boldsymbol{d}) \, d\Omega - \left. \iint_{\Gamma} f(a; \gamma) \, d\Gamma \right|_{\boldsymbol{u}^{(s)}}$$

abc

Strain energy: Interface between solid electrode and solid-state electrolyte (SE|SSE) taking place at space charge

$$\iiint_{\Omega} f(a, \boldsymbol{u}; \lambda, \mu, \boldsymbol{d} \otimes \boldsymbol{d}) d\Omega$$

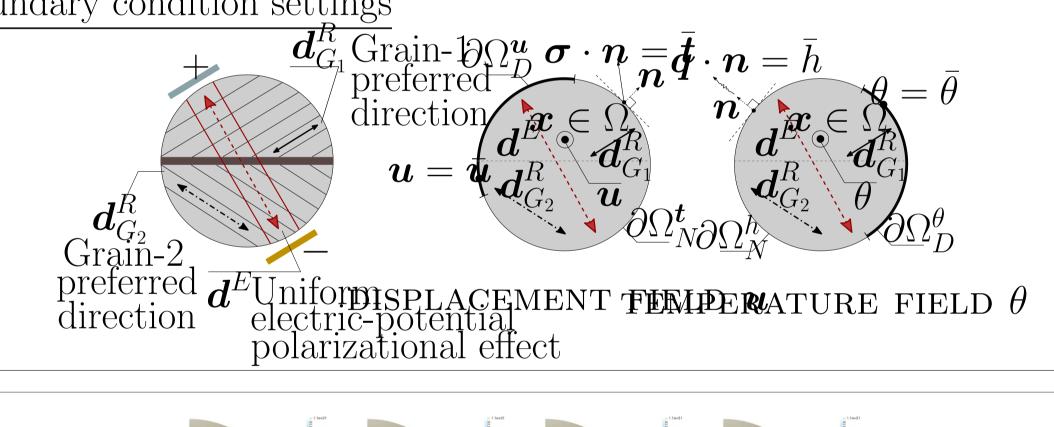
Surface energy: Interface between solid electrode and solid-state electrolyte (SE|SSE) taking place

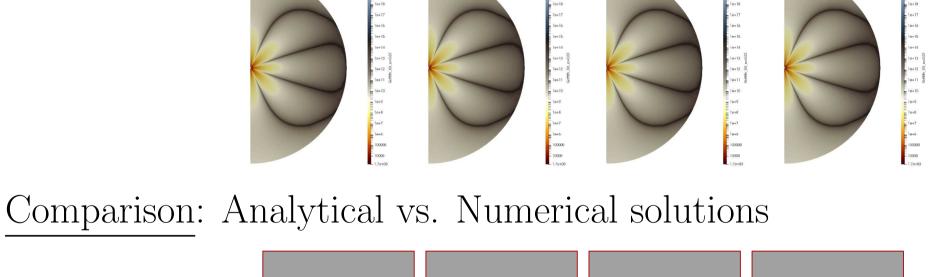
$$\iint_{\Gamma} f(a;\gamma) \, d\Gamma$$

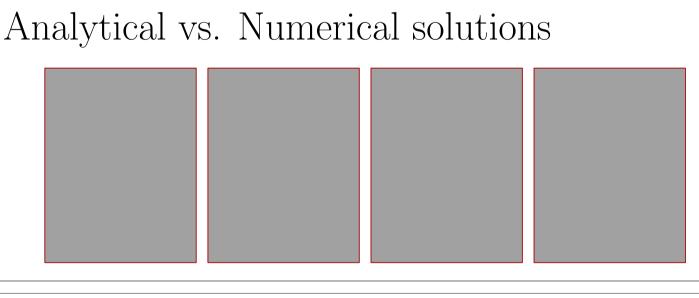
Therefore

$$\rho \, \partial_{t^2}^2 \boldsymbol{u}^{(s)} + \nabla \cdot \left(\overset{4}{\mathbb{C}} f_{(\lambda,\mu)}^{\mathbb{D}(\Omega)} : \nabla \boldsymbol{u}^{(s)} \right) + \rho \nabla V_e = \boldsymbol{0},$$
s.t. $a_{\text{Griffith}} := a^* = \arg\min_{a \in \mathbb{R}} \left. \iint_{\Omega} f(a, \boldsymbol{u}; \lambda, \mu, \boldsymbol{d} \otimes \boldsymbol{d}) \, d\Omega - \int_{\Gamma} f(a; \gamma) \, d\Gamma \right|_{\boldsymbol{u}^{(s)}}$
abc

Boundary condition settings







FEM: Strain energy density Partial differential equation $\nabla \cdot \left(\mathbb{C}^{f^{GL}}(y) \nabla_s \boldsymbol{u} \right) + \rho \boldsymbol{b} = 0$ Displacement solution $\varepsilon_{ij} = \frac{1}{2} \left(u_{i,j} + u_{j,i} \right)$ $oldsymbol{\sigma}_{ij} = \mathbb{C}_{ijkl}^{f^{GL}}\!(y) \; arepsilon_{kl}$ Strain energy density $\mathcal{E}_{ ext{strain}} := rac{1}{2} \sigma_{ij} arepsilon_{ij}$ abc

Airy-Westergaard function used for max. shear stress analysis

$$\mathcal{V}_{\mathcal{A}}: \mathbb{C} \to \mathbb{C}, z \mapsto \mathcal{V}_{\mathcal{A}}(z) := \Re(\iint_{\Gamma} \mathcal{K}^{(\star)} dz) + x_2 \Im(\oint_{\Gamma} \mathcal{K}^{(\star)} dz), \mathcal{K}^c(z) := -p_h + p_h/\sqrt{1 - a^2/z^2},$$

where $\{p_h, a\} \in \mathbb{R}_+$ is the.

FEM implementation: element matrix \mathbf{K}^e approx. by Gauss quadrature; indices imply 4+2=6 for-loop:

$$K_{ik}^{e^{\alpha\beta}} = \int_{\Omega^{\xi}} \left(\mathcal{L}_{1}^{\alpha} \, \mathbb{C}_{i1k1}^{fGL}(y) \, \mathcal{R}_{1}^{\beta} + \mathcal{L}_{1}^{\alpha} \, \mathbb{C}_{i1k2}^{fGL}(y) \, \mathcal{R}_{2}^{\beta} + \mathcal{L}_{2}^{\alpha} \, \mathbb{C}_{i2k1}^{fGL}(y) \, \mathcal{R}_{1}^{\beta} + \mathcal{L}_{2}^{\alpha} \, \mathbb{C}_{i2k2}^{fGL}(y) \, \mathcal{R}_{2}^{\beta} \right) \det(\boldsymbol{J}) \, d\Omega^{\xi}$$

where \mathcal{L}_i^{α} and \mathcal{R}_l^{β} are gradients of basis functions at node α^{th} and β^{th} , respectively.

Contact

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