

# SOME SET PROPERTIES UNDERLYING GEOMETRY AND PHYSICS

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ABSTRACT. Volume and distance equations are proved to be instances of ordered sets of combinations (n-tuples). The combinatorial properties can limit volume and distance to a set of 3 dimensions. More dimensions have different types (are members of other sets), with ratios of a distance unit to units of time, mass, and charge. The proofs and ratios are used to: derive well-known gravity, charge, electromagnetic equations, special and general relativity equations, and quantum physics equations; derive the gravity, charge, vacuum permittivity, vacuum permeability, Planck, and fine structure constants; add quantum extensions to gravity and charge equations. All the proofs are verified in Rocq.

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## 1. INTRODUCTION

The Riemann integral, Lebesgue integral, and Lebesgue measure define Euclidean volume as the product of interval sizes  $\subset \mathbb{R}^n$  [8] [15]. And Euclidean distance, vector magnitude, and the many “spaces”: vector, inner product, metric, Hausdorff, Cauchy, etc. are definitions, in analysis [18] [8] [15].

Justification of the definitions uses finger-pointing to geometry because definitions can only *describe* aspects of volume and distance. Deriving volume and distance equations from an abstract set and limit-based foundation exposes the principles that *explain*, cause, the volume and distance equations, which provides new insights and tools for analysis and physics.

All the proofs in this article have been verified using the Rocq proof verification system [14]. The formal proofs are in the Rocq files, “euclidrelations.v” and “threed.v,” at: <https://github.com/treeck/RASRGeometry>.

Using calculus and  $\sigma$ -algebras (for example, the Lebesgue measure) to prove the volume and distance theorems in this article would result in circular logic. Therefore, a “ruler” measure of intervals,  $[a, b] \subset \mathbb{R}$ , will be used to prove the theorems.

Where  $|x_i|$  is the cardinal of (number of elements in) a countable set,  $x_i$ , the countable number of ordered combinations (n-tuples) is  $v_c$ . The ruler measure will be used to prove the Euclidean volume relation:

$$\begin{aligned} \forall x_i \in \{x_1, \dots, x_n\} = X, \quad \bigcap_{x_i \in X} x_i = \emptyset : \quad v_c = \prod_{i=1}^n |x_i| \\ \Leftrightarrow \quad v = \prod_{i=1}^n s_i, \quad s_i = b_i - a_i, \quad [a_i, b_i] \subset \mathbb{R}. \end{aligned} \quad (1.1)$$

For all  $n > 1$ , there are an infinite number of combinations of domain values,  $s_i \cdots s_n$ , that sum to the same range value,  $v$ . The volume function,  $v = \prod_{i=1}^n s_i$  is only bijective where,  $\exists d : v = f(d)$  and  $d = f^{-1}(v) = f^{-1}(\prod_{i=1}^n s_i)$ .

For the bijective case of countable n-volumes:

$$\exists d_c, v_c, |x_i| \in \{0, \mathbb{N}\} : v_c = \prod_{i=1}^n |x_i| = \prod_{i=1}^n d_c = d_c^n, \quad (1.2)$$

And  $d = f^{-1}(v)$  and  $v = \sum_{i=1}^m v_i \Rightarrow d = f^{-1}(\sum_{i=1}^m v_i)$ . The ruler measure will be used to prove that:

$$d_c^n = \sum_{i=1}^m v_{c_i} = \sum_{i=1}^m (\prod_{j=1}^n |x_{i,j}|) \Leftrightarrow d^n = \sum_{i=1}^m v_i = \sum_{i=1}^m (\prod_{j=1}^n s_{i,j}). \quad (1.3)$$

The  $n = 2$  case is the basis of the inner product. Where each  $v_{c_i}$  is also a bijective function,  $v_{c_i} = d_{c_i}^n$ , the ruler measure will be used to prove that:

$$d_c^n = \sum_{i=1}^m d_{c_i}^n \Leftrightarrow d^n = \sum_{i=1}^m d_i^n. \quad (1.4)$$

$|d|$  is the  $p$ -norm (Minkowski distance) [12], which will be proved to imply the metric space properties [15]. The  $n = 2$  case is, obviously, the Euclidean distance.

Volume and distance are derived from sets of ordered combinations (n-tuples). Volume and distance have another combinatorial (permutation) property.

Calculating volume requires multiplying a sequentially ordered set of domain values. And calculating distance requires summing a sequentially ordered set of values. The commutative properties of multiplication and addition allows sequencing an ordered set in all  $n!$  permutations.

You cannot repeatably sequence the domain values in the same order, unless you place the values in a sequential order. Further, the *only* sequential order, where you can start with any value and sequence in a repeatable order, is a cyclic order.

The second sequenced member must be either the *immediate* cyclic successor or *immediate* cyclic predecessor. And repeatable, sequencing of a cyclic set in all  $n!$  permutations, is a symmetry, where every set member is either an *immediate* cyclic successor or an *immediate* cyclic predecessor to every other set member, herein, referred to as an “immediate symmetric” cyclic set. An immediate symmetric cyclic set will be proved to have  $n \leq 3$  members.

An immediate cyclic set of 3 “distance” dimensions requires more dimensions to have non-distance types (be members of other sets). Let  $\tau = \{t \text{ (time)}, m \text{ (mass)}, q \text{ (charge)}\}$  be an immediate cyclic set of type “non-distance” dimensions, where for each subinterval (unit) length,  $r_c$ , of distance interval length,  $r$ , there are unit

lengths:  $t_c$  of time interval length,  $t$ ;  $m_c$  of mass interval length,  $m$ ; and  $q_c$  of charge interval length,  $q$ , such that:  $r = (r_c/t_c)t = (r_c/m_c)m = (r_c/q_c)q$ .

The proofs and the 3 direct proportion ratios are used to provide simple derivations of the Newton, Gauss, and Poisson gravity equations, Coulomb charge, Gauss, Lorentz, and Faraday electromagnetic equations [and the constants: gravity ( $G$ ), charge ( $k_e$ ), vacuum permittivity ( $\varepsilon_0$ ), and vacuum permeability ( $\mu_0$ )]. They are also used to derive all the special relativity equations, the Schwarzschild time dilation and black hole metric equations pointing to a simplified method of finding solutions to Einstein's general relativity equations.

Next, algebraic manipulation of the 3 direct proportion ratios yields 3 inverse proportion ratios,  $r = t_c r_c / t = m_c r_c / m = q_c r_c / q$ . The Planck units and fine structure constant,  $\alpha$ , are derived from ratios of subtypes:  $\alpha = q_e^2 / q_p^2$  where  $q_e$  is the elementary (electron) charge and  $q_p$ :  $q_c^2 / q_p^2 = 2\pi$ , is the Planck charge unit.

The combination of the direct and inverse proportion ratios are used to derive the Planck relation, the Planck constant,  $h = (m_c r_c)(r_c / t_c)$ , the Compton, position-space Schrödinger, and Dirac wave equations. The inverse proportion ratios are also used to add quantum extensions to some general relativity and classical physics equations. Finally, the derivations expose some incorrect assumptions commonly used in physics today.

## 2. RULER MEASURE AND CONVERGENCE

A ruler (measuring stick) measures the size of each interval *approximately* as the sum of the nearest integer number,  $p$ , of size  $\kappa$  subintervals. The ruler is both an inner and outer measure of an interval.

**Definition 2.1.** Ruler measure,  $M = \sum_{i=1}^p \kappa = p\kappa$ , where  $\forall [a, b] \subset \mathbb{R}$ ,  $s = b - a \wedge 0 < \kappa \leq 1 \wedge (p = \text{floor}(s/\kappa) \vee p = \text{ceiling}(s/\kappa))$ .

**Theorem 2.2.** *Ruler convergence:*  $M = \lim_{\kappa \rightarrow 0} p\kappa = s$ .

The formal proof, “limit\_c\_0\_M.eq\_exact\_size,” is in the file, euclidrelations.v.

*Proof.* (epsilon-delta proof)

By definition of the floor function,  $\text{floor}(x) = \max(\{y : y \leq x, y \in \mathbb{Z}, x \in \mathbb{R}\})$ :

$$\forall 0 < \kappa \leq 1, p = \text{floor}(s/\kappa) \wedge 0 \leq |\text{floor}(s/\kappa) - s/\kappa| < 1 \Rightarrow |p - s/\kappa| < 1. \quad (2.1)$$

Multiply both sides of inequality 2.1 by  $\kappa$ :

$$\forall 0 < \kappa \leq 1, |p - s/\kappa| < 1 \Rightarrow |p\kappa - s| < |\kappa| = |\kappa - 0|. \quad (2.2)$$

$$\begin{aligned} \forall \epsilon = \delta \quad \wedge \quad |p\kappa - s| < |\kappa - 0| < \delta \\ \Rightarrow \quad |\kappa - 0| < \delta \quad \wedge \quad |p\kappa - s| < \delta = \epsilon \quad := \quad M = \lim_{\kappa \rightarrow 0} p\kappa = s. \quad \square \end{aligned} \quad (2.3)$$

The following is an example of ruler convergence for the interval,  $[0, \pi]$ :  $s = \pi - 0$ , and  $p = \text{floor}(s/\kappa) \Rightarrow p \cdot \kappa = 3.1_{\kappa=10^{-1}}, 3.14_{\kappa=10^{-2}}, 3.141_{\kappa=10^{-3}}, \dots, \pi_{\lim_{\kappa \rightarrow 0}}$ .

**Lemma 2.3.**  $\forall n \geq 1, 0 < \kappa \leq 1 \Rightarrow \lim_{\kappa \rightarrow 0} \kappa^n = \lim_{\kappa \rightarrow 0} \kappa$ .

*Proof.* The formal proof, “lim\_c\_to\_n.eq\_lim\_c,” is in the Rocq file, euclidrelations.v.

$$n \geq 1 \quad \wedge \quad 0 < \kappa \leq 1 \quad \Rightarrow \quad 0 < \kappa^n < \kappa \quad \Rightarrow \quad |\kappa - \kappa^n| < |\kappa| = |\kappa - 0|. \quad (2.4)$$

$$\begin{aligned} \forall \epsilon = \delta \quad \wedge \quad |\kappa - \kappa^n| < |\kappa - 0| < \delta \\ \Rightarrow \quad |\kappa - 0| < \delta \quad \wedge \quad |\kappa - \kappa^n| < \delta = \epsilon \quad := \quad \lim_{\kappa \rightarrow 0} \kappa^n = 0. \end{aligned} \quad (2.5)$$

$$\lim_{\kappa \rightarrow 0} \kappa^n = 0 \quad \wedge \quad \lim_{\kappa \rightarrow 0} \kappa = 0 \quad \Rightarrow \quad \lim_{\kappa \rightarrow 0} \kappa^n = \lim_{\kappa \rightarrow 0} \kappa. \quad (2.6)$$

□

### 3. VOLUME

**Definition 3.1.** A countable n-volume is the number of ordered combinations (n-tuples),  $v_c$ , of the members of  $n$  number of disjoint, countable domain sets,  $x_i$ :

$$x_i \in \{x_1, \dots, x_n\} = X, |x_i| \in \{0, \mathbb{N}\} : \bigcap_{x_i \in X} x_i = \emptyset \quad \wedge \quad v_c = \prod_{i=1}^n |x_i|. \quad (3.1)$$

**Theorem 3.2.** *Euclidean volume,*

$$\begin{aligned} \forall [a_i, b_i] \in \{[a_1, b_1], \dots [a_n, b_n]\}, [v_a, v_b] \subset \mathbb{R}, s_i = b_i - a_i, v = v_b - v_a : \\ v_c = \prod_{i=1}^n |x_i| \quad \Leftrightarrow \quad v = \prod_{i=1}^n s_i. \end{aligned} \quad (3.2)$$

The formal proof, “Euclidean\_volume,” is in the Rocq file, euclidrelations.v.

*Proof.*

$$v_c = \prod_{i=1}^n |x_i| \Leftrightarrow v_c \kappa = (\prod_{i=1}^n |x_i|) \kappa \Leftrightarrow \lim_{\kappa \rightarrow 0} v_c \kappa = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i|) \kappa. \quad (3.3)$$

Apply the ruler (2.1) and ruler convergence (2.2) to equation 3.3:

$$\begin{aligned} \exists v, \kappa \in \mathbb{R} : v_c = \text{floor}(v/\kappa) \quad \Rightarrow \quad v = \lim_{\kappa \rightarrow 0} v_c \kappa \quad \wedge \\ \lim_{\kappa \rightarrow 0} v_c \kappa = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i|) \kappa \quad \Rightarrow \quad v = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i|) \kappa. \end{aligned} \quad (3.4)$$

Apply lemma 2.3 to equation 3.4:

$$\begin{aligned} v = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i|) \kappa \quad \wedge \quad \lim_{\kappa \rightarrow 0} \kappa^n = \lim_{\kappa \rightarrow 0} \kappa \\ \Rightarrow \quad v = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i|) \kappa^n = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i| \kappa). \end{aligned} \quad (3.5)$$

Apply the ruler (2.1) and ruler convergence (2.2) to  $s_i$ :

$$\exists s_i, \kappa \in \mathbb{R} : \text{floor}(s_i/\kappa) = |x_i| \quad \Rightarrow \quad \lim_{\kappa \rightarrow 0} (|x_i| \kappa) = s_i. \quad (3.6)$$

$$\begin{aligned} v = \lim_{\kappa \rightarrow 0} (\prod_{i=1}^n |x_i| \kappa) \quad \wedge \quad \lim_{\kappa \rightarrow 0} (|x_i| \kappa) = s_i \\ \Leftrightarrow \quad v = \lim_{\kappa \rightarrow 0} (|x_i| \kappa) = \prod_{i=1}^n s_i \quad \square \end{aligned} \quad (3.7)$$

### 4. DISTANCE

**Definition 4.1.** Countable distance,

$$\begin{aligned} \exists n \in \mathbb{N}, v_c, d_c \in \{0, \mathbb{N}\}, x_i \in \{x_1, \dots, x_n\} = X : \bigcap_{x_i \in X} x_i = \emptyset \quad \wedge \\ d_c = |x_1| = \dots = |x_n| \quad \wedge \quad v_c = \prod_{i=1}^n |x_i| = \prod_{i=1}^n d_c = d_c^n. \end{aligned} \quad (4.1)$$

**Lemma 4.2.** *A volume is the sum of volumes,*

$$v_c = d_c^n = \sum_{i=1}^m v_{c_i} \quad \Leftrightarrow \quad v = \sum_{i=1}^m v_i, \quad v, v_i \in \mathbb{R}.$$

The formal proof, “sum\_of\_volumes,” is in the Rocq file, euclidrelations.v.

*Proof.* From the condition of this theorem:

$$v_c = \sum_{i=1}^m v_{c_i} \Leftrightarrow \lim_{\kappa \rightarrow 0} v_c \kappa = \lim_{\kappa \rightarrow 0} \sum_{j=1}^m (v_{c_i} \kappa). \quad (4.2)$$

Apply lemma 2.3 to equation 4.2:

$$\begin{aligned} \lim_{\kappa \rightarrow 0} v_c \kappa &= \lim_{\kappa \rightarrow 0} (\sum_{j=1}^m v_{c_i}) \kappa \quad \wedge \quad \lim_{\kappa \rightarrow 0} \kappa^n = \lim_{\kappa \rightarrow 0} \kappa \\ \Leftrightarrow \lim_{\kappa \rightarrow 0} v_c \kappa &= \lim_{\kappa \rightarrow 0} \sum_{j=1}^m (v_{c_i}) \kappa^n \Leftrightarrow \lim_{\kappa \rightarrow 0} v_c \kappa = \lim_{\kappa \rightarrow 0} \sum_{j=1}^m (v_{c_i} \kappa). \end{aligned} \quad (4.3)$$

Apply the ruler (2.1) and ruler convergence theorem (2.2) to equation 4.3:

$$\begin{aligned} \exists v, v_i : v &= \text{floor}(d/\kappa), v = \lim_{\kappa \rightarrow 0} v_c \kappa \\ \wedge v_{c_i} &= \text{floor}(v_i/\kappa), v_i = \lim_{\kappa \rightarrow 0} v_{c_i} \kappa \quad \wedge \quad \lim_{\kappa \rightarrow 0} (d_c \kappa)^n = \lim_{\kappa \rightarrow 0} \sum_{j=1}^m (v_{c_i} \kappa) \\ \Leftrightarrow v &= \lim_{\kappa \rightarrow 0} (d_c \kappa)^n = \lim_{\kappa \rightarrow 0} \sum_{j=1}^m (v_{c_i} \kappa) = \sum_{j=1}^m v_i^n. \quad \square \end{aligned} \quad (4.4)$$

#### 4.1. Sum of volumes distance.

**Theorem 4.3.** *Sum of volumes distance:*

$$v_c = d_c^n = \sum_{i=1}^m v_{c_i} \Leftrightarrow d^n = \sum_{i=1}^m (\prod_{j=1}^n s_{ij}).$$

The formal proof, “sum\_of\_volumes\_distance,” is in the Rocq file, euclidrelations.v.

*Proof.* From lemma 4.2 and the Euclidean volume theorem 3.2:

$$\begin{aligned} v_c = d_c^n = \sum_{i=1}^m v_{c_i} &\Leftrightarrow d^n = \sum_{i=1}^m (\prod_{j=1}^n v_i) \quad \wedge \quad v_i = \prod_{j=1}^n s_{ij} \\ v_c = d_c^n = \sum_{i=1}^m v_{c_i} &\Leftrightarrow d^n = \sum_{i=1}^m (\prod_{j=1}^n s_{ij}). \quad \square \end{aligned} \quad (4.5)$$

#### 4.2. Minkowski distance ( $p$ -norm).

**Theorem 4.4.** *Minkowski distance ( $p$ -norm):*

$$v_c = d_c^n = \sum_{i=1}^m v_{c_i} = \sum_{i=1}^m d_{c_i}^n \Leftrightarrow d^n = \sum_{i=1}^m d_i^n.$$

The formal proof, “Minkowski\_distance,” is in the Rocq file, euclidrelations.v.

*Proof.* From lemma 4.2 and the Euclidean volume theorem 3.2:

$$\begin{aligned} v_c = d_c^n = \sum_{i=1}^m v_{c_i} &\Leftrightarrow d^n = \sum_{i=1}^m v_i \quad \wedge \quad v_i = \prod_{j=1}^n d_i = d_i^n \\ v_c = d_c^n = \sum_{i=1}^m v_{c_i} &\Leftrightarrow d^n = \sum_{i=1}^m d_i^n. \quad \square \end{aligned} \quad (4.6)$$

**4.3. Distance inequality.** The formal proof, distance.inequality, is in the Rocq file, euclidrelations.v.

**Theorem 4.5.** *Distance inequality*

$$\forall n \in \mathbb{N}, v_a, v_b \geq 0 : (v_a + v_b)^{1/n} \leq v_a^{1/n} + v_b^{1/n}.$$

*Proof.* Expand  $(v_a^{1/n} + v_b^{1/n})^n$  using the binomial expansion:

$$\begin{aligned} \forall v_a, v_b \geq 0 : v_a + v_b &\leq v_a + v_b + \\ \sum_{i=1}^n \binom{n}{i} (v_a^{1/n})^{n-k} (v_b^{1/n})^k &+ \sum_{i=1}^n \binom{n}{i} (v_a^{1/n})^k (v_b^{1/n})^{n-k} = (v_a^{1/n} + v_b^{1/n})^n. \end{aligned} \quad (4.7)$$

Take the  $n^{\text{th}}$  root of both sides of the inequality 4.7:

$$\forall v_a, v_b \geq 0, n \in \mathbb{N} : v_a + v_b \leq (v_a^{1/n} + v_b^{1/n})^n \Rightarrow (v_a + v_b)^{1/n} \leq v_a^{1/n} + v_b^{1/n}. \quad (4.8)$$

□

**4.4. Distance sum inequality.** The formal proof, `distance_sum_inequality`, is in the Rocq file, `euclidrelations.v`.

**Theorem 4.6.** *Distance sum inequality*

$$\forall m, n \in \mathbb{N}, a_i, b_i \geq 0 : (\sum_{i=1}^m (a_i^n + b_i^n))^{1/n} \leq (\sum_{i=1}^m a_i^n)^{1/n} + (\sum_{i=1}^m b_i^n)^{1/n}.$$

*Proof.* Apply the distance inequality (4.5):

$$\begin{aligned} \forall m, n \in \mathbb{N}, v_a, v_b \geq 0 : \quad v_a &= \sum_{i=1}^m a_i^n \quad \wedge \quad v_b = \sum_{i=1}^m b_i^n \quad \wedge \\ (v_a + v_b)^{1/n} &\leq v_a^{1/n} + v_b^{1/n} \quad \Rightarrow \quad ((\sum_{i=1}^m a_i^n) + (\sum_{i=1}^m b_i^n))^{1/n} = \\ &(\sum_{i=1}^m (a_i^n + b_i^n))^{1/n} \leq (\sum_{i=1}^m a_i^n)^{1/n} + (\sum_{i=1}^m b_i^n)^{1/n}. \quad \square \end{aligned} \quad (4.9)$$

**4.5. Metric Space.** All Minkowski distances ( $p$ -norms) imply the metric space properties. The formal proofs: `triangle_inequality`, `symmetry`, `non_negativity`, and `identity_of_indiscernibles` are in the Rocq file, `euclidrelations.v`.

**Theorem 4.7.** *Triangle Inequality:*

$$d(s_1, s_2) = (\sum_{i=1}^2 s_i^p)^{1/p} \Rightarrow d(u, w) \leq d(u, v) + d(v, w).$$

*Proof.*  $\forall p \geq 1, k > 1, u = s_1, w = s_2, v = w/k$ :

$$(u^p + w^p)^{1/p} \leq ((u^p + w^p) + 2v^p)^{1/p} = ((u^p + v^p) + (v^p + w^p))^{1/p}. \quad (4.10)$$

Apply the distance inequality (4.5) to the inequality 4.10:

$$\begin{aligned} (u^p + w^p)^{1/p} &\leq ((u^p + v^p) + (v^p + w^p))^{1/p} \quad \wedge \quad (v_a + v_b)^{1/n} \leq v_a^{1/n} + v_b^{1/n} \\ &\quad \wedge \quad v_a = u^p + v^p \quad \wedge \quad v_b = v^p + w^p \\ \Rightarrow \quad (u^p + w^p)^{1/p} &\leq ((u^p + v^p) + (v^p + w^p))^{1/p} \leq (u^p + v^p)^{1/p} + (v^p + w^p)^{1/p} \\ &\Rightarrow \quad d(u, w) = (u^p + w^p)^{1/p} \leq \\ &\quad (u^p + v^p)^{1/p} + (v^p + w^p)^{1/p} = d(u, v) + d(v, w). \quad \square \end{aligned} \quad (4.11)$$

**Theorem 4.8.** *Symmetry:*  $d(s_1, s_2) = (\sum_{i=1}^2 s_i^p)^{1/p} \Rightarrow d(u, v) = d(v, u)$ .

*Proof.* By the commutative law of addition:

$$\begin{aligned} \forall p : p \geq 1, \quad d(s_1, s_2) &= (\sum_{i=1}^2 s_i^p)^{1/p} = (s_1^p + s_2^p)^{1/p} \\ &\Rightarrow \quad d(u, v) = (u^p + v^p)^{1/p} = (v^p + u^p)^{1/p} = d(v, u). \quad \square \end{aligned} \quad (4.12)$$

**Theorem 4.9.** *Non-negativity:*  $d(s_1, s_2) = (\sum_{i=1}^2 s_i^p)^{1/p} \Rightarrow d(u, w) \geq 0$ .

*Proof.* By definition, the length of an interval is always  $\geq 0$ :

$$\forall [a_1, b_1], [a_2, b_2], \quad u = b_1 - a_1, v = b_2 - a_2, \quad \Rightarrow \quad u \geq 0, v \geq 0. \quad (4.13)$$

$$p \geq 1, u, v \geq 0 \quad \Rightarrow \quad d(u, v) = (u^p + v^p)^{1/p} \geq 0. \quad (4.14)$$

□

**Theorem 4.10.** *Identity of Indiscernibles:*  $d(u, u) = 0$ .

*Proof.* From the non-negativity property (4.9):

$$\begin{aligned} d(u, w) \geq 0 \quad \wedge \quad d(u, v) \geq 0 \quad \wedge \quad d(v, w) \geq 0 \\ \Rightarrow \quad \exists d(u, w) = d(u, v) = d(v, w) = 0. \end{aligned} \quad (4.15)$$

$$d(u, w) = d(v, w) = 0 \quad \Rightarrow \quad u = v. \quad (4.16)$$

$$d(u, v) = 0 \quad \wedge \quad u = v \quad \Rightarrow \quad d(u, u) = 0. \quad (4.17)$$

□

**4.6. Set properties limiting a set to at most 3 members.** The following definitions and proof use first order logic. A Horn clause-like expression is used, here, to make the proof easier to read. By convention, the proof goal is on the left side and supporting facts are on the right side of the implication sign ( $\leftarrow$ ). The formal proofs in the Rocq file `threed.v` are:

**Lemmas:** `adj111`, `adj122`, `adj212`, `adj123`, `adj133`, `adj233`, `adj213`, `adj313`, `adj323`, and `not_all_mutually_adjacent_gt_3`.

**Definition 4.11.** Immediate Cyclic Successor of  $m$  is  $n$ :

$$\begin{aligned} \forall x_m, x_n \in \{x_1, \dots, x_{\text{setsize}}\} : \\ \text{Successor}(m, n, \text{setsize}) \leftarrow (m = \text{setsize} \wedge n = 1) \vee (n = m + 1 \leq \text{setsize}). \end{aligned} \quad (4.18)$$

**Definition 4.12.** Immediate Cyclic Predecessor of  $m$  is  $n$ :

$$\begin{aligned} \forall x_m, x_n \in \{x_1, \dots, x_{\text{setsize}}\} : \\ \text{Predecessor}(m, n, \text{setsize}) \leftarrow (m = 1 \wedge n = \text{setsize}) \vee (n = m - 1 \geq 1). \end{aligned} \quad (4.19)$$

**Definition 4.13.** Adjacent: Member  $m$  is sequentially adjacent to member  $n$  if the immediate cyclic successor of  $m$  is  $n$  or the immediate cyclic predecessor of  $m$  is  $n$ . Notionally:

$$\begin{aligned} \forall x_m, x_n \in \{x_1, \dots, x_{\text{setsize}}\} : \\ \text{Adjacent}(m, n, \text{setsize}) \leftarrow \text{Successor}(m, n, \text{setsize}) \vee \text{Predecessor}(m, n, \text{setsize}). \end{aligned} \quad (4.20)$$

**Definition 4.14.** Immediate Symmetric (every set member is sequentially adjacent to every other member):

$$\forall x_m, x_n \in \{x_1, \dots, x_{\text{setsize}}\} : \quad \text{Adjacent}(m, n, \text{setsize}). \quad (4.21)$$

**Theorem 4.15.** *An immediate symmetric cyclic set is limited to at most 3 members.*

*Proof.*

Every member is adjacent to every other member, where  $\text{setsize} \in \{1, 2, 3\}$ :

$$\text{Adjacent}(1, 1, 1) \leftarrow \text{Successor}(1, 1, 1) \leftarrow (m = \text{setsize} \wedge n = 1). \quad (4.22)$$

$$\text{Adjacent}(1, 2, 2) \leftarrow \text{Successor}(1, 2, 2) \leftarrow (n = m + 1 \leq \text{setsize}). \quad (4.23)$$

$$\text{Adjacent}(2, 1, 2) \leftarrow \text{Successor}(2, 1, 2) \leftarrow (n = \text{setsize} \wedge m = 1). \quad (4.24)$$

$$\text{Adjacent}(1, 2, 3) \leftarrow \text{Successor}(1, 2, 3) \leftarrow (n = m + 1 \leq \text{setsize}). \quad (4.25)$$

$$\text{Adjacent}(2, 1, 3) \leftarrow \text{Predecessor}(2, 1, 3) \leftarrow (n = m - 1 \geq 1). \quad (4.26)$$

$$\text{Adjacent}(3, 1, 3) \leftarrow \text{Successor}(3, 1, 3) \leftarrow (n = \text{setsize} \wedge m = 1). \quad (4.27)$$

$$\text{Adjacent}(1, 3, 3) \leftarrow \text{Predecessor}(1, 3, 3) \leftarrow (m = 1 \wedge n = \text{setsize}). \quad (4.28)$$

$$\text{Adjacent}(2, 3, 3) \leftarrow \text{Successor}(2, 3, 3) \leftarrow (n = m + 1 \leq \text{setsize}). \quad (4.29)$$

$$\text{Adjacent}(3, 2, 3) \leftarrow \text{Predecessor}(3, 2, 3) \leftarrow (n = m - 1 \geq 1). \quad (4.30)$$

Member 2 is the only immediate successor of member 1 for all  $setsize \geq 3$ , which implies member 3 is not ( $\neg$ ) an immediate successor of member 1 for all  $setsize \geq 3$ :

$$\neg Successor(1, 3, setsize \geq 3) \\ \leftarrow Successor(1, 2, setsize \geq 3) \leftarrow (n = m + 1 \leq setsize). \quad (4.31)$$

Member  $n = setsize > 3$  is the only immediate predecessor of member 1, which implies member 3 is not ( $\neg$ ) an immediate predecessor of member 1 for all  $setsize > 3$ :

$$\neg Predecessor(1, 3, setsize \geq 3) \\ \leftarrow Predecessor(1, setsize, setsize > 3) \leftarrow (m = 1 \wedge n = setsize > 3). \quad (4.32)$$

For all  $setsize > 3$ , some elements are not ( $\neg$ ) sequentially adjacent to every other element (not immediate symmetric):

$$\neg Adjacent(1, 3, setsize > 3) \\ \leftarrow \neg Successor(1, 3, setsize > 3) \wedge \neg Predecessor(1, 3, setsize > 3). \quad \square \quad (4.33)$$

The Symmetric goal matches Adjacent goals 4.22 and fails for all “setsize” greater than three.

## 5. APPLICATIONS TO PHYSICS

Where distance is an immediate cyclic set of dimensions, the 3D proof requires more dimensions to have non-distance types, where for each unit length,  $r_c$ , of distance interval length,  $r$ , there are unit lengths:  $t_c$  of time interval length,  $t$ ;  $m_c$  of mass interval length,  $m$ ; and  $q_c$  of charge interval length,  $q$ , such that:

$$r = (r_c/t_c)t = (r_c/m_c)m = (r_c/q_c)q. \quad (5.1)$$

**5.1. Derivation of the constant,  $G$ , and the gravity laws of Newton, Gauss, and Poisson.** From equation 5.1:

$$r = (r_c/m_c)m \quad \wedge \quad r = (r_c/t_c)t = ct \quad \Rightarrow \quad r/(ct)^2 = (r_c/m_c)m/r^2 \\ \Rightarrow \quad r/t^2 = ((r_c/m_c)c^2)m/r^2 = Gm/r^2, \quad (5.2)$$

where the constant,  $G = (r_c/m_c)c^2$ , conforms to the SI units:  $m^3 \cdot kg^{-1} \cdot s^{-2}$  [13].

Newton’s law follows from multiplying both sides of equation 5.2 by  $m$ :

$$r/t^2 = Gm/r^2 \Leftrightarrow F := mr/t^2 = Gm^2/r^2. \quad (5.3)$$

$$F = Gm^2/r^2 \wedge \forall m \in \mathbb{R} : \exists m_1, m_2 \in \mathbb{R} : m_1 m_2 = m^2 \Rightarrow F = Gm_1 m_2 / r^2. \quad (5.4)$$

From equation 5.2, Gauss’s gravity field,  $\mathbf{g}$  and Poisson’s gravity field,  $\nabla \Phi(r, t)$ :

$$\mathbf{g} = -\nabla \Phi(\vec{r}, t) := r/t^2 = Gm/r^2 \\ \Rightarrow \quad \nabla \cdot \mathbf{g} = \nabla^2 \Phi(\vec{r}, t) = -2Gm/r^3 = (-2Gm/r^3)(2\pi/2\pi) \quad \wedge \quad \rho = m/2\pi r^3 \\ \Rightarrow \quad \nabla \cdot \mathbf{g} = \nabla^2 \Phi(\vec{r}, t) = -4\pi G\rho. \quad (5.5)$$



## 5.2. Derivation of Coulomb's charge constant, $k_e$ and charge force.

$$\begin{aligned} \forall q \in \mathbb{R} : \exists q_1, q_2 \in \mathbb{R} : q_1 q_2 = q^2 \quad \wedge \quad r = (r_c/q_c)q \\ \Rightarrow \quad \exists q_1, q_2 \in \mathbb{R} : q_1 q_2 = q^2 = ((q_c/r_c)r)^2 \quad \Rightarrow \quad (r_c/q_c)^2 q_1 q_2 / r^2 = 1. \end{aligned} \quad (5.6)$$

$$\begin{aligned} r = (r_c/t_c)t = ct \quad \wedge \quad r = (r_c/m_c)m = ct \\ \Rightarrow \quad mr = (m_c/r_c)rct = (m_c/r_c)(ct)^2 \quad \Rightarrow \quad ((r_c/m_c)/c^2)mr/t^2 = 1. \end{aligned} \quad (5.7)$$

$$\begin{aligned} ((r_c/m_c)/c^2)mr/t^2 = 1 \quad \wedge \quad (r_c/q_c)^2 q_1 q_2 / r^2 = 1 \\ \Rightarrow \quad F := mr/t^2 = ((m_c/r_c)c^2)(r_c/q_c)^2 q_1 q_2 / r^2 = k_e q_1 q_2 / r^2. \end{aligned} \quad (5.8)$$

where Coulomb's constant,  $k_e = ((m_c/r_c)c^2)(r_c/q_c)^2$ , conforms to the base SI units:  $kg \cdot m^3 \cdot s^{-2} \cdot C^{-2}$ , which is equivalent to the charge SI units:  $N \cdot m^2 \cdot C^{-2}$  [7].

## 5.3. Vacuum permittivity, $\varepsilon_0$ , and Gauss's law for electric fields. From Coulomb's charge force equation 5.8:

$$\begin{aligned} \exists q, E \in \mathbb{R} : F = k_e q_1 q_2 / r^2 = k_e q^2 / r^2 := qE \\ \Rightarrow \quad E = k_e q / r^2 \quad \Leftrightarrow \quad \mathbf{E} = k_e q / \mathbf{r}^2, \end{aligned} \quad (5.9)$$

where  $E$  has the SI units  $N \cdot C^{-1}$ .

$$\mathbf{E} = k_e q / \mathbf{r}^2 \quad \Rightarrow \quad \nabla \cdot \mathbf{E} = -2k_e q / \mathbf{r}^3. \quad (5.10)$$

$$\begin{aligned} \nabla \cdot \mathbf{E} = -(2k_e q / \mathbf{r}^3)(2\pi/2\pi) \quad \wedge \quad \rho = q/2\pi \mathbf{r}^3 \quad \wedge \quad \varepsilon_0 := 1/4\pi k_e \\ \Rightarrow \quad \nabla \cdot \mathbf{E} = -4\pi k_e \rho = -\rho/\varepsilon_0, \end{aligned} \quad (5.11)$$

which is Gauss's electric field law [7].

**5.4. Magnetic field.** An elementary particle possessing both mass and charge has the mass-charge wavelength,  $\lambda = \sqrt{(k_m/m)^2 + (k_q/q)^2}$ , where  $\vec{\lambda}$ , is the direction (orientation) of an elementary particle-wave.

Consider electron, A, held in a stationary position and wave orientation. As electron B moves past electron A in a straight line, a torque force acts on the mass-charge wave orientation of electron B, which changes with respect to time and distance, which will be resisted due to the intrinsic angular momentum of the mass-charge wave.

From Coulomb's charge force equation 5.8:

$$\begin{aligned} \exists q \in \mathbb{R} : F = k_e q_1 q_2 / r^2 = k_e q^2 / r^2 \quad \wedge \quad r = (r_c/t_c)t = ct \\ \Rightarrow \quad Fct = r(k_e q^2 / r^2) \quad \Leftrightarrow \quad F = (r/ct)(k_e q^2 / r^2). \end{aligned} \quad (5.12)$$

For a unit radius around the charged particle-wave, the torque force,  $F_T = F \sin(\theta)$ :

$$\begin{aligned} F = (r/ct)(k_e q^2 / r^2) \quad \wedge \quad F_T = F \sin(\theta) \\ \Rightarrow \quad F_T = (r/ct)(k_e q^2 / r^2) \sin(\theta) = (r \sin(\theta)/t)((k_e/c)q^2 / r^2). \end{aligned} \quad (5.13)$$

$$\begin{aligned} F_T = (r \sin(\theta)/t)((k_e/c)q^2 / r^2) \quad \wedge \quad v = (r \sin(\theta)/t) \\ \Rightarrow \quad F_T = v(k_e/c)q^2 / r^2 := qvB \quad \Leftrightarrow \quad \mathbf{F}_T = q\vec{v} \times \mathbf{B}, \end{aligned} \quad (5.14)$$

where  $v$  is the speed of electron B,  $c$  is the speed of light, and  $B = (k_e/c)q/r^2$ , conforms to the base SI units:  $kg \cdot s^{-1} \cdot C^{-1} = kg \cdot s^{-2} \cdot A^{-1} = T$ .

**5.5. Lorentz's force law.** Summing Gauss's electric force equation 5.9 plus the magnetic force equation 5.14 yields the Lorentz law:

$$\mathbf{F}_E = q\mathbf{E} \quad \wedge \quad \mathbf{F}_B = q\vec{v} \times \mathbf{B} \quad \Rightarrow \quad \mathbf{F}_{total} = \mathbf{F}_E + \mathbf{F}_B = q(\mathbf{E} + \vec{v} \times \mathbf{B}). \quad (5.15)$$

**5.6. Vacuum permeability,  $\mu_0$ , and Faraday's law.** From the magnetic field equation 5.14:

$$\mathbf{B} = (k_e/c)q/r^2 \quad \wedge \quad r = ct \quad \Rightarrow \quad \mathbf{B} = (k_e/c^3)q/t^2. \quad (5.16)$$

$$\mathbf{B} = (k_e/c^3)q/t^2 \quad \Rightarrow \quad \partial\mathbf{B}/\partial t = -(2k_e/c^3)q/t^3. \quad (5.17)$$

$$\partial\mathbf{B}/\partial t = -(2k_e/c^3)q/t^3 \quad \wedge \quad r = ct \quad \Rightarrow \quad \partial\mathbf{B}/\partial t = -2k_eq/r^3. \quad (5.18)$$

From equation 5.9:

$$\mathbf{E} = k_eq/r^2 \quad \Rightarrow \quad \nabla \times \mathbf{E} = 2k_eq/r^3. \quad (5.19)$$

Combining equations 5.19 and 5.18 yields Faraday's law [7]:

$$\nabla \times \mathbf{E} = 2k_eq/r^3 \quad \wedge \quad \partial\mathbf{B}/\partial t = -2k_eq/r^3 \quad \Rightarrow \quad \nabla \times \mathbf{E} = -\partial\mathbf{B}/\partial t. \quad (5.20)$$

$$\begin{aligned} \partial\mathbf{B}/\partial t = -(2k_eq/r^3)(2\pi/2\pi) \quad \wedge \quad \rho = q/2\pi r^3 \quad \wedge \quad \mu_0 := 4\pi k_e/c^2 \\ \Rightarrow \quad \partial\mathbf{B}/\partial t = -4\pi k_e\rho = -\mu_0\rho. \end{aligned} \quad (5.21)$$

**5.7. Space-time-mass-charge.** Let  $r$  be an Euclidean distance. Then by the Minkowski distance theorem (4.4),  $r^2 = \sum_{i=1}^m r_i^2$ . Let,  $r' = r_1$  and  $r_v = (\sum_{i=1}^{m-1} r_i^2)$ . From equation 5.1, there are ratios  $\mu$  and  $\nu$  such that:

$$\begin{aligned} \forall \tau \in \{t, m, q\}, r^2 = r'^2 + r_v^2, \exists \mu, \nu : \quad \wedge \quad r = \mu\tau \quad \wedge \quad r_v = \nu\tau \\ \Rightarrow \quad (\mu\tau)^2 = r'^2 + (\nu\tau)^2 \quad \Rightarrow \quad r' = \sqrt{(\mu\tau)^2 - (\nu\tau)^2} = \mu\tau\sqrt{1 - (\nu/\mu)^2}. \end{aligned} \quad (5.22)$$

Rest frame distance,  $r'$ , contracts relative to stationary frame distance,  $r$ , as  $\nu \rightarrow \mu$ :

$$r' = \mu\tau\sqrt{1 - (\nu/\mu)^2} \quad \wedge \quad \mu\tau = r \quad \Rightarrow \quad r' = r\sqrt{1 - (\nu/\mu)^2}. \quad (5.23)$$

Stationary frame type,  $\tau$ , dilates relative to the rest frame type,  $\tau'$ , as  $\nu \rightarrow \mu$ :

$$\mu\tau = r'/\sqrt{1 - (\nu/\mu)^2} \quad \wedge \quad r' = \mu\tau' \quad \Rightarrow \quad \tau = \tau'/\sqrt{1 - (\nu/\mu)^2}. \quad (5.24)$$

Where  $\tau$  is type, time, the space-like flat Minkowski spacetime event interval is:

$$\begin{aligned} dr^2 = dr'^2 + dr_v^2 \quad \wedge \quad dr_v^2 = dr_1^2 + dr_2^2 + dr_3^2 \quad \wedge \quad d(\mu\tau) = dr \\ \Rightarrow \quad dr'^2 = d(\mu\tau)^2 - dr_1^2 - dr_2^2 - dr_3^2. \end{aligned} \quad (5.25)$$

**5.8. Derivation of Schwarzschild's gravitational time dilation and black hole metric.** [17] [1] From equations 5.23 and 5.1:

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - (v^2/c^2)(r/r)} \quad \wedge \quad r = (r_c/m_c)m \\ \Rightarrow \quad \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - ((r_c/m_c)m)v^2/rc^2}. \end{aligned} \quad (5.26)$$

Where  $v_{escape}$  is the escape velocity:

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - ((r_c/m_c)m)v^2/rc^2} \quad \wedge \quad KE = mv^2/2 = mv_{escape}^2 \\ \Rightarrow \quad \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - 2(r_c/m_c)mv_{escape}^2/rc^2}. \end{aligned} \quad (5.27)$$

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \lim_{v_{escape} \rightarrow c} \sqrt{1 - 2(r_c/m_c)mv_{escape}^2/rc^2} \\ &= \sqrt{1 - 2(r_c/m_c)mc^2/rc^2}. \end{aligned} \quad (5.28)$$

Combining equation 5.28 with the derivation of  $G$  (5.4):

$$\begin{aligned} (r_c/m_c)c^2 &= G \quad \wedge \quad \sqrt{1 - (v^2/c^2)} = \sqrt{1 - 2(r_c/m_c)mc^2/rc^2} \\ \Rightarrow \quad \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - 2Gm/rc^2}. \end{aligned} \quad (5.29)$$

Combining equation 5.29 with equation 5.24 yields Schwarzschild's gravitational time dilation:

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - 2Gm/rc^2} \quad \wedge \quad t' = t\sqrt{1 - (v^2/c^2)} \\ \Rightarrow \quad t' &= t\sqrt{1 - 2Gm/rc^2}. \end{aligned} \quad (5.30)$$

Schwarzschild defined the black hole event horizon radius,  $r_s := 2Gm/c^2$ . From equations 5.23 and 5.31:

$$\begin{aligned} r' &= r\sqrt{1 - (v/c)^2} \quad \wedge \quad \sqrt{1 - (v/c)^2} = \sqrt{1 - 2Gm/rc^2} \\ \Rightarrow \quad r' &= r\sqrt{1 - 2Gm/rc^2} = r\sqrt{1 - r_s/r}. \end{aligned} \quad (5.31)$$

Using the time-like spacetime interval, where  $ds^2$  is negative:

$$\begin{aligned} r' &= r\sqrt{1 - r_s/r} \quad \wedge \quad ds^2 = dr'^2 - dr^2 \\ \Rightarrow \quad ds^2 &= (\sqrt{1 - r_s/r}dr')^2 - (dr/\sqrt{1 - r_s/r})^2 = (1 - r_s/r)dr'^2 - (1 - r_s/r)^{-1}dr^2. \end{aligned} \quad (5.32)$$

$$\begin{aligned} ds^2 &= (1 - r_s/r)dr'^2 - (1 - r_s/r)^{-1}dr^2 \quad \wedge \quad dr' = d(ct) \quad \wedge \quad c = 1 \\ \Rightarrow \quad ds^2 &= (1 - r_s/r)dt^2 - (1 - r_s/r)^{-1}dr^2. \end{aligned} \quad (5.33)$$

Translating from 2D to 4D yields Schwarzschild's black hole metric:

$$\begin{aligned} ds^2 &= (1 - r_s/r)dt^2 - (1 - r_s/r)^{-1}dr^2 = f(r, t) \\ \Rightarrow \quad ds^2 &= (1 - r_s/r)dt^2 - (1 - r_s/r)^{-1}dr^2 - r^2(d\theta^2 + \sin^2\theta d\phi^2) = f(r, t, \theta, \phi) \\ \Rightarrow \quad g_{\mu,\nu} &= \text{diag}[1 - r_s/r, (1 - r_s/r)^{-1}, r^2(d\theta^2), r^2(\sin^2\theta d\phi^2)]. \end{aligned} \quad (5.34)$$

**5.9. Simplifying Einstein's general relativity (field) equation.** Step 1) Use the ratios to define functions returning scalar values for each component of the metric,  $g_{\nu,\mu}$ , in Einstein's field equations [6] [18]: All functions derived from the ratios, where the units on each side of the equation balance, are valid metrics, for example, the previous Schwarzschild black hole metric derivation using the ratios (5.8).

Step 2) Express the EFE as 2D tensors: As shown in equation 5.34, the Schwarzschild metric was first derived as a 2D metric and then expanded to a 4D metric. Further, the 4D flat spacetime interval equation (5.25) is an instance of the 2D equation,  $dr'^2 = d(ct)^2 - dr_v^2$ , where  $dr_v^2$  is the magnitude of a 3-dimensional vector.

The 2D metric tensor allows using the much simpler 2D Ricci curvature and scalar curvature. And the 2D tensors reduce the number of independent equations to solve.

Step 3) One simple method to translate from 2D to 4D is to use spherical coordinates, where  $r$  and  $t$  remain unchanged and two added dimensions are the angles,  $\phi$ , and  $\theta$ . For example, the 2D Schwarzschild metric was translated to 4D using this method in equation 5.34.

**5.10. 3 fundamental direct proportion ratios.**  $c_t$ ,  $c_m$ , and  $c_q$ :

$$c_t = r_c/t_c \approx 2.99792458 \cdot 10^8 m \ s^{-1}. \quad (5.35)$$

$$G = (r_c/m_c)c_t^2 = c_m c_t^2 \Rightarrow c_m = r_c/m_c \approx 7.4261602691 \cdot 10^{-28} m \ kg^{-1}. \quad (5.36)$$

$$k_e = (c_t^2/c_m)(r_c/q_c)^2 \Rightarrow c_q = r_c/q_c \approx 8.6175172023 \cdot 10^{-18} m \ C^{-1}. \quad (5.37)$$

**5.11. 3 fundamental inverse proportion ratios.**  $k_t$ ,  $k_m$ , and  $k_q$ :

$$\begin{aligned} r/t = r_c/t_c, \quad r/m = r_c/m_c &\Rightarrow (r/t)/(r/m) = (r_c/t_c)/(r_c/m_c) \Rightarrow \\ (mr)/(tr) = (m_c r_c)/(t_c r_c) &\Rightarrow mr = m_c r_c = k_m, \quad tr = t_c r_c = k_t. \end{aligned} \quad (5.38)$$

$$\begin{aligned} r/t = r_c/t_c, \quad r/q = r_c/q_c &\Rightarrow (r/t)/(r/q) = (r_c/t_c)/(r_c/q_c) \Rightarrow \\ (qr)/(tr) = (q_c r_c)/(t_c r_c) &\Rightarrow qr = q_c r_c = k_q, \quad tr = t_c r_c = k_t. \end{aligned} \quad (5.39)$$

**5.12. Planck relation and constant,  $h$ .** [9] Applying both the direct proportion ratio (5.35), and inverse proportion ratio (5.38):

$$r = ct \quad \wedge \quad m = k_m/r \Rightarrow m(ct)^2 = (k_m/r)r^2 = k_m r. \quad (5.40)$$

$$\begin{aligned} m(ct)^2 = k_m r \quad \wedge \quad r/t = r_c/t_c = c \\ \Rightarrow E := mc^2 = k_m r/t^2 = (k_m(r/t)) (1/t) = (k_m c)(1/t) = hf = \hbar\omega, \end{aligned} \quad (5.41)$$

where the Planck constant,  $h = k_m c$ , the frequency,  $f = 1/t$ , the reduced Planck constant,  $\hbar = h/2\pi$ , and the angular frequency,  $\omega = 2\pi f = 2\pi/t$ .

$$k_m = m_c r_c = h/c \approx 2.2102190943 \cdot 10^{-42} kg \ m. \quad (5.42)$$

$$k_t = t_c r_c = k_m c_m / c_t \approx 5.4749346710 \cdot 10^{-78} s \ m. \quad (5.43)$$

$$k_q = q_c r_c = k_t c_t / c_q \approx 1.9046601056 \cdot 10^{-52} C \ m. \quad (5.44)$$

**5.13. Compton wavelength.** [9] From equations 5.38 and 5.41:

$$mr = k_m \quad \wedge \quad h = k_m c \Rightarrow r = k_m/m = (k_m/m)(c/c) = h/mc. \quad (5.45)$$

5.14. **4 quantum units.** Distance ( $r_c$ ), time ( $t_c$ ), mass ( $m_c$ ), and charge ( $q_c$ ):

$$r_c = \sqrt{r_c^2} = \sqrt{c_t k_t} = \sqrt{c_m k_m} = \sqrt{c_q k_q} \approx 4.0513505432 \cdot 10^{-35} \text{ m}. \quad (5.46)$$

$$t_c = r_c/c_t \approx 1.3513850782 \cdot 10^{-43} \text{ s}. \quad (5.47)$$

$$m_c = r_c/c_m \approx 5.4555118613 \cdot 10^{-8} \text{ kg}. \quad (5.48)$$

$$q_c = r_c/c_q \approx 4.7012967286 \cdot 10^{-18} \text{ C}. \quad (5.49)$$

5.15. **Subtype ratios.**  $\frac{F_1}{F_2} = \frac{K\tau_1^2/r^2}{K\tau_2^2/r^2} = \frac{\tau_1^2}{\tau_2^2}$ .

$$\text{Planck length, } r_p : r_c^2/r_p^2 = 2\pi \Rightarrow r_p = r_c/\sqrt{2\pi} \approx 1.6162550244 \cdot 10^{-35} \text{ m}. \quad (5.50)$$

$$\text{Planck time, } t_p : t_c^2/t_p^2 = 2\pi \Rightarrow t_p = t_c/\sqrt{2\pi} \approx 5.3912464472 \cdot 10^{-44} \text{ s}. \quad (5.51)$$

$$\text{Planck mass, } m_p : m_c^2/m_p^2 = 2\pi \Rightarrow m_p = m_c/\sqrt{2\pi} \approx 2.176434343 \cdot 10^{-8} \text{ kg}. \quad (5.52)$$

$$\text{Planck charge, } q_p : q_c^2/q_p^2 = 2\pi \Rightarrow q_p = q_c/\sqrt{2\pi} \approx 1.875546038 \cdot 10^{-18} \text{ C}. \quad (5.53)$$

Where  $q_e$  is the elementary (electron) charge ( $1.60217663 \cdot 10^{-19} \text{ C}$ ), the fine structure constant,  $\alpha$  is:

$$\alpha q_c^2/q_e^2 = 2\pi \Rightarrow \alpha = 2\pi q_e^2/q_c^2 = q_e^2/(q_c/\sqrt{2\pi})^2 = q_e^2/q_p^2 \approx 0.0072973526. \quad (5.54)$$

5.16. **Schrödinger's equation.** Start with the previously derived Planck relation 5.41 and multiply the kinetic energy component by  $mc/mc$ :

$$h/t = mc^2 \Rightarrow \exists V(r, t) : h/t = h/2t + V(r, t) \Rightarrow h/t = hmc/2mct + V(r, t). \quad (5.55)$$

And from the distance-to-time (speed of light) ratio (5.35):

$$h/t = hmc/2mct + V(r, t) \wedge r = ct \Rightarrow h/t = hmc^2/2mcr + V(r, t). \quad (5.56)$$

$$h/t = hmc^2/2mcr + V(r, t) \wedge h/t = mc^2 \Rightarrow h/t = h^2/2mcrt + V(r, t). \quad (5.57)$$

$$h/t = h^2/2mcrt + V(r, t) \wedge r = ct \Rightarrow h/t = h^2/2mr^2 + V(r, t). \quad (5.58)$$

Replace the Planck constant in equation 5.58 with the reduced Planck constant:

$$h/t = h^2/2mr^2 + V(r, t) \wedge \hbar = h/2\pi \Rightarrow 2\pi\hbar/t = (2\pi)^2\hbar^2/2mr^2 + V(r, t). \quad (5.59)$$

Multiply both sides of equation 5.59 by a function,  $\Psi(r, t)$ .

$$\begin{aligned} 2\pi\hbar/t &= (2\pi)^2\hbar^2/2mr^2 + V(r, t) \\ \Rightarrow (2\pi\hbar/t)\Psi(r, t) &= ((2\pi)^2\hbar^2/2mr^2)\Psi(r, t) + V(r, t)\Psi(r, t). \end{aligned} \quad (5.60)$$

$$\begin{aligned} (2\pi\hbar/t)\Psi(r, t) &= ((2\pi)^2\hbar^2/2mr^2)\Psi(r, t) + V(r, t)\Psi(r, t) \wedge \\ \forall \Psi(r, t) : \partial^2\Psi(r, t)/\partial r^2 &= (-(2\pi)^2/r^2)\Psi(r, t) \wedge \partial\Psi(r, t)/\partial t = (i 2\pi/t)\Psi(r, t) \\ \Rightarrow i\hbar\partial\Psi(r, t)/\partial t &= -(\hbar^2/2m)\partial^2\Psi(r, t)/\partial r^2 + V(r, t)\Psi(r, t), \end{aligned} \quad (5.61)$$

which is the position-space Schrödinger's equation in one dimension of space [16] [9].

$$\begin{aligned} i\hbar\partial\Psi(r, t)/\partial t &= -(\hbar^2/2m)\partial^2\Psi(r, t)/\partial r^2 + V(r, t)\Psi(r, t) \wedge ||\vec{r}'|| = r \\ \Rightarrow \exists \vec{r} : i\hbar\partial\Psi(\vec{r}, t)/\partial t &= -(\hbar^2/2m)\partial^2\Psi(\vec{r}, t)/\partial \vec{r}^2 + V(\vec{r}, t)\Psi(\vec{r}, t), \end{aligned} \quad (5.62)$$

which is the position-space Schrödinger's equation in 3 dimensions of space [16] [9].

**5.17. Dirac's wave equation.** Using the derived Planck relation 5.41:

$$\begin{aligned} mc^2 = h/t \quad \Rightarrow \quad \exists V(r, t) : mc^2/2 + V(r, t) = h/t \\ \Rightarrow \quad 2h/t - 2V(r, t) = mc^2. \end{aligned} \quad (5.63)$$

$$\begin{aligned} \forall V(r, t) : V(r, t) = ih/t \quad \wedge \quad r = ct \quad \wedge \quad 2h/t - 2V(r, t) = mc^2 \\ \Rightarrow \quad 2h/t - i2hc/r = mc^2. \end{aligned} \quad (5.64)$$

Use the charge ratio,  $c_q$ , and time ratio,  $c_t = c$  to multiply each term on the left side of equation 5.64 by 1:

$$\begin{aligned} qc_q/r = qc_q/ct = 1 \quad \wedge \quad 2h/t - i2hc/r = mc^2 \\ \Rightarrow \quad 2h(-qc_q/c)/t^2 - i2h((-qc_q/c)/r^2)c = mc^2. \end{aligned} \quad (5.65)$$

where a negative sign is added to  $q$  to indicate an attractive force between an electron and a nucleus.

Applying a quantum amplitude equation in complex form to equation 5.66, where  $\tau$  is a unit time and  $\rho$  is a unit wavelength:

$$\begin{aligned} A_0 = (c_q/c)((1/t)\tau - i(1/r)\rho) \wedge 2h(-qc_q/c)/t^2 - i2h((-qc_q/c)/r^2)c = mc^2 \\ \Rightarrow \quad 2h\partial(-qA_0)/\partial t - i2h(\partial(-qA_0)/\partial r)c = mc^2. \end{aligned} \quad (5.66)$$

Translating equation 5.66 to moving coordinates via the Lorentz factor,  $\gamma_0 = 1/\sqrt{1 - (v/c)^2}$ :

$$\begin{aligned} 2h\partial(-qA_0)/\partial t - i2h(\partial(-qA_0)/\partial r)c = mc^2 \\ \Rightarrow \quad \gamma_0 2h\partial(-qA_0)/\partial t - \gamma_0 i2h(\partial(-qA_0)/\partial r)c = mc^2. \end{aligned} \quad (5.67)$$

Multiplying both sides of equation 5.67 by  $\Psi(r, t)$ :

$$\begin{aligned} \gamma_0 2h\partial(-qA_0)/\partial t - \gamma_0 i2h(\partial(-qA_0)/\partial r)c = mc^2 \\ \Rightarrow \quad \gamma_0 2h(\partial(-qA_0)/\partial t)\Psi(r, t) - \gamma_0 i2h(\partial(-qA_0)/\partial r)c\Psi(r, t) = mc^2\Psi(r, t). \end{aligned} \quad (5.68)$$

Applying the vectors to equation 5.68:

$$\begin{aligned} \gamma_0 2h(\partial(-qA_0)/\partial t)\Psi(r, t) - \gamma_0 i2h(\partial(-qA_0)/\partial r)c\Psi(r, t) = mc^2\Psi(r, t) \wedge \\ ||\vec{r}|| = r \quad \wedge \quad ||\vec{A}|| = A_0 \quad \wedge \quad ||\vec{\gamma}|| = \gamma_0 \quad \wedge \quad \Leftrightarrow \quad \exists \vec{r}, \vec{A}, \vec{\gamma} : \\ \gamma_0 2h(\partial(-qA_0)/\partial t)\Psi(r, t) - \vec{\gamma} \cdot i2h(\partial(-q\vec{A})/\partial r)c\Psi(\vec{r}, t) = mc^2\Psi(\vec{r}, t). \end{aligned} \quad (5.69)$$

Adding a  $\frac{1}{2}$  angular rotation (spin- $\frac{1}{2}$ ) of  $\pi$  to equation 5.66 allows substituting the reduced Planck constant,  $\hbar = h/2\pi$ , into equation 5.69, which yields Dirac's wave equation [5] [9]:

$$\begin{aligned} \gamma_0 2h(\partial(-qA_0)/\partial t)\Psi(r, t) - \vec{\gamma} \cdot i2h(\partial(-q\vec{A})/\partial r)c\Psi(\vec{r}, t) = mc^2\Psi(\vec{r}, t) \\ \wedge A_0 = \pi(c_q/c)((1/t)\tau - i(1/r)\rho) \\ \Rightarrow \quad \gamma_0 \hbar(\partial(-qA_0)/\partial t)\Psi(r, t) - \vec{\gamma} \cdot i\hbar(\partial(-q\vec{A})/\partial r)c\Psi(\vec{r}, t) = mc^2\Psi(\vec{r}, t). \end{aligned} \quad (5.70)$$

**5.18. Total mass.** The total mass of a particle is  $m = \sqrt{m_0^2 + m_{ke}^2}$ , where  $m_0$  is the rest mass and  $m_{ke}$  is the kinetic energy-equivalent mass. Applying both the direct (5.35) and inverse proportion ratios (5.38):

$$\begin{aligned} m_0 &= r/(r_c/m_c) = r/c_m \quad \wedge \quad m_{ke} = (m_c r_c)/r = k_m/r \quad \wedge \\ m &= \sqrt{m_0^2 + m_{ke}^2} \quad \Rightarrow \quad m = \sqrt{(r/c_m)^2 + (k_m/r)^2}. \end{aligned} \quad (5.71)$$

**5.19. Quantum extension to general relativity.** The simplest way to demonstrate how to add quantum physics to general relativity is by extending Schwarzschild's time dilation equation and black hole metric (5.8). Start by changing equation 5.26 in the Schwarzschild derivation:

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - (v^2/c^2)(r/r)} \quad \wedge \quad r = \sqrt{(c_m m)^2 + (k_m/m)^2} = Q_m \\ &\Rightarrow \quad \sqrt{1 - (v^2/c^2)} = \sqrt{1 - Q_m v^2/rc^2}. \end{aligned} \quad (5.72)$$

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - Q_m v^2/rc^2} \quad \wedge \quad KE = mv^2/2 = mv_{escape}^2 \\ &\Rightarrow \quad \sqrt{1 - (v^2/c^2)} = \sqrt{1 - 2Q_m v_{escape}^2/rc^2}. \end{aligned} \quad (5.73)$$

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \lim_{v_{escape} \rightarrow c} \sqrt{1 - 2Q_m v_{escape}^2/rc^2} \\ &\Rightarrow \quad \sqrt{1 - (v^2/c^2)} = \sqrt{1 - 2Q_m c^2/rc^2} = \sqrt{1 - 2Q_m/r}. \end{aligned} \quad (5.74)$$

Combining equation 5.74 with equation 5.24 yields Schwarzschild's gravitational time dilation with a quantum mass effect:

$$\begin{aligned} \sqrt{1 - (v^2/c^2)} &= \sqrt{1 - 2Q_m/r} \quad \wedge \quad t' = t\sqrt{1 - (v^2/c^2)} \\ &\Rightarrow \quad t' = t\sqrt{1 - 2Q_m/r}. \end{aligned} \quad (5.75)$$

Schwarzschild defined the black hole event horizon radius,  $r_s := 2Gm/c^2$ . The radius with the quantum extension is  $r_s := 2Q_m$ . At this point the exact same equations 5.31 through 5.34 yield what looks like the same Schwarzschild black hole metric.

**5.20. Quantum extension to Newton's gravity force.** The quantum mass effect is easier to understand in the context Newton's gravity equation than in general relativity, because the metric equations and solutions in the EFEs are much more complex. From equations 5.76 and 5.1:

$$\begin{aligned} m/\sqrt{(r/c_m)^2 + (k_m/r)^2} &= 1 \quad \wedge \quad r^2/(ct)^2 = 1 \\ &\Rightarrow \quad r^2/(ct)^2 = m/\sqrt{(r/c_m)^2 + (k_m/r)^2} \\ &\Rightarrow \quad r^2/t^2 = c^2 m/\sqrt{(r/c_m)^2 + (k_m/r)^2}. \end{aligned} \quad (5.76)$$

$$\begin{aligned} r^2/t^2 &= c^2 m/\sqrt{(r/c_m)^2 + (k_m/r)^2} \\ &\Rightarrow \quad (m/r)(r^2/t^2) = (m/r)(c^2 m/\sqrt{(r/c_m)^2 + (k_m/r)^2}) \\ &\Rightarrow \quad F := mr/t^2 = c^2 m^2/(r\sqrt{(r/c_m)^2 + (k_m/r)^2}) = c^2 m^2/\sqrt{(r^4/c_m^2) + k_m^2}. \end{aligned} \quad (5.77)$$

$$F = c^2 m^2 / \sqrt{(r^4/c_m^2) + k_m^2} \quad \wedge \quad \forall m \in \mathbb{R}, \exists m_1, m_2 \in \mathbb{R} : m_1 m_2 = m^2 \\ \Rightarrow F = c^2 m_1 m_2 / \sqrt{(r^4/c_m^2) + k_m^2}. \quad (5.78)$$

### 5.21. Quantum extension to Coulomb's force.

$$q^2 / ((r/c_q)^2 + (k_q/r)^2) = 1 \quad \wedge \quad r^2 / (ct)^2 = 1 \\ \Rightarrow r^2 / (ct)^2 = q^2 / ((r/c_q)^2 + (k_q/r)^2) \\ \Rightarrow r^2 / t^2 = c^2 q^2 / ((r/c_q)^2 + (k_q/r)^2). \quad (5.79)$$

$$\forall q \in \mathbb{R} : \exists q_1, q_2 \in \mathbb{R} : q_1 q_2 = q^2 \quad \wedge \quad r^2 / t^2 = c^2 q^2 / ((r/c_q)^2 + (k_q/r)^2) \\ \Rightarrow \exists q_1, q_2 \in \mathbb{R} : r^2 / t^2 = c^2 q_1 q_2 / ((r/c_q)^2 + (k_q/r)^2) \\ \Rightarrow r / t^2 = c^2 q_1 q_2 / (r((r/c_q)^2 + (k_q/r)^2)). \quad (5.80)$$

$$r / t^2 = c^2 q_1 q_2 / (r((r/c_q)^2 + (k_q/r)^2)) \quad \wedge \quad m = \sqrt{(r/c_m)^2 + (k_m/r)^2} \\ \Rightarrow F := mr / t^2 = c^2 q_1 q_2 \sqrt{(r/c_m)^2 + (k_m/r)^2} / (r((r/c_q)^2 + (k_q/r)^2)) \\ = c^2 q_1 q_2 \sqrt{(r^4/c_m^2) + k_m^2} / ((r^4/c_q^2) + k_q^2). \quad (5.81)$$

## 6. INSIGHTS AND IMPLICATIONS

- (1) The ruler measure (2.1) and convergence theorem (2.2) were shown to be useful tools for proving that a real-valued equation is the only instance of an abstract, countable set relation and that set relation is the only instance of that same real-valued equation.
- (2) Combinatorics, the ordered set of combinations of countable, disjoint sets (n-tuples),  $v_c = \prod_{i=1}^n |x_i|$ , was proven to imply the Euclidean volume equation (3.2), the sum of volumes equation (4.3) (which includes the inner product), and the Minkowski distance equation (4.4) (which includes the Manhattan and Euclidean distance equations), without relying on the geometric primitives and relations in Euclidean geometry [10], axiomatic geometry [11], and vector analysis [18].
- (3) Where the total n-volume is both the sum and subtraction of n-volumes, the  $n = 2$  case is the vector inner product. The distributive and associate laws of multiplication and addition allow the  $\pm$  signed volumes to be represented as each domain interval length multiplied by a  $\pm$ -signed unit values,  $\alpha_i, \beta_i \in \{-1, 1\}$ :

$$d^2 = \sum_{i=1}^m (a_i \alpha_i)(b_i \beta_i) := \mathbf{a} \cdot \mathbf{b}. \quad (6.1)$$

- (4) Defining all Euclidean and non-Euclidean distance measures as:

$$d : f(d, n) = v = \sum_{i=1}^m v_i, \quad d = f^{-1}(v, n) = f(\sum_{i=1}^m v_i)^{-1}. \quad (6.2)$$

- (a) shows the intimate relation between distance and volume that definitions, like inner product space and metric space, ignore [18] [8] [15];
- (b) is a more simple and concise definition of a distance measure that includes all the properties used in the definitions of inner product space and metric space [18] [8] [15];



- (5) Euclid's proof that Euclidean distance is the smallest distance between two distinct points equate Euclidean distance to a straight line, where it is assumed that the straight line length is the smallest distance [10]. And analytic proofs sum infinitesimal distances,  $ds = \sqrt{dx^2 + dy^2}$ , where the Euler-Lagrange equation is used find the minimum solution, which is the straight line equation,  $y = mx + b$  [2].

Without using the notion of a straight line: All distance measures in an Euclidean volume have corresponding Minkowski distances (4.4). For all 2-volumes, all Minkowski distances are limited to  $n \in \{1, 2\}$ :  $n = 1$  is the larger (Manhattan) distance case,  $d = \sum_{i=1}^m s_i$ .  $n = 2$  is the smaller (Euclidean) distance case,  $d = (\sum_{i=1}^m s_i^2)^{1/2}$ . That is, by the Cauchy-Schwarz inequality:  $\sum_{i=1}^m s_i \geq (\sum_{i=1}^m s_i^2)^{1/2}$ . And where  $1 \leq n \leq 2$ ,  $d$  decreases monotonically as  $n \rightarrow 2$ .

- (6) The left side of the distance sum inequality (4.6),

$$(\sum_{i=1}^m (a_i^n + b_i^n))^{1/n} \leq (\sum_{i=1}^m a_i^n)^{1/n} + (\sum_{i=1}^m b_i^n)^{1/n}, \quad (6.3)$$

differs from the left side of Minkowski's sum inequality [12]:

$$(\sum_{i=1}^m (a_i^n + b_i^n)^{\mathbf{n}})^{1/n} \leq (\sum_{i=1}^m a_i^n)^{1/n} + (\sum_{i=1}^m b_i^n)^{1/n}. \quad (6.4)$$

- (a) The two inequalities are only the same where  $n = 1$ .
  - (b) The distance sum inequality (4.6) is a more fundamental inequality because the proof does not require the convexity and Hölder's inequality assumptions required to prove the Minkowski sum inequality [12].
  - (c) The distance sum inequality term,  $\forall n > 1$ :  $d = v^{1/n} = (a_i^n + b_i^n)^{1/n}$ , is the Minkowski distance spanning the  $n$ -volume,  $v = a_i^n + b_i^n$ , which makes it directly related to geometry (for example, the metric space triangle inequality was derived from the  $m = 1$  case for all  $n \geq 1$  (4.7)). But the Minkowski sum inequality term,  $\forall n > 1$ :  $d = v^{1/n} \neq v = (a_i^n + b_i^n) = ((a_i^n + b_i^n)^{\mathbf{n}})^{1/n}$ .
- (7) Combinatorics, repeatable sequencing through an ordered set to yield all  $n!$  permutations of its members (without jumping around) was proved to be a cyclic set having  $n \leq 3$  members (4.15). Higher dimensions must have different types (members of different sets).
- (a) For example, the vector inner product space can only be extended beyond 3 dimensions if and only if the higher dimensions have non-distance types, for example, time.
  - (b) Order and symmetry probably limit the number of non-distance types  $\subset \mathbb{R}$  to 3, for example: time, mass, and charge.
  - (c) Each of 3 immediate symmetric cyclic dimensions of space can have at most 3 immediate symmetric cyclic state values, for example, an immediate symmetric cyclic set of 3 vector orientations,  $\{-1, 1\}$ .
  - (d) If the states are not ordered (a bag of states), then a state value is undetermined until observed (like Schrödinger's poisoned cat being both alive and dead until the box is opened [16]). For a bag of states, there is **no** "axiom of choice" [4], an axiom often used in math proofs that allows selecting a particular set element (in this case, selecting a particular state).
  - (e) A discrete value has measure 0 (no size). The ratio of a time or distance interval length to zero is undefined, which is the reason quantum

entangled particles change discrete state values together independent of time and distance.

- (8) For each unit,  $r_c$ , of a 3-dimensional distance interval having a length,  $r$ , there are units of other types of intervals forming unit ratios (5.10):  $c_t = r_c/t_c$ ,  $c_m = r_c/m_c$ ,  $c_q = r_c/q_c \Leftrightarrow$  the inverse proportion ratios (5.11):  $k_t = r_c t_c$ ,  $k_m = r_c m_c$ ,  $k_t = r_c q_c$ .

(a) **The Cartesian principle:** The direct proportion ratios state that for each unit of one type there is a corresponding unit of another type.

(b) **The conservation principle:** The indirect proportion ratios state that a change in size of one type requires an inverse change in another type.

- (9) The derivations empirical and hypothesized laws from the ratios expose some **incorrect assumptions** currently used in physics:

(a) Empirical and hypothesized laws use an **opaque** constant,  $K$ , that is defined to make an equation, where the units balance,  $y = Kf(x_1, \dots, x_n)$ . The opacity has led to the **incorrect** assumptions of those constants being fundamental constants.

The derivations, in this article, show that the gravity,  $G = c_m c_t^2$  (5.4), charge,  $k_e = (c_q^2/c_m) c_t^2$  (5.8), and Planck  $h = k_m c_t$  (5.41) constants were all derived directly from (composed of) the ratios. Further,  $\varepsilon_0 := 1/4\pi k_e$  (5.11) and  $\mu_0 := 1/c_t^2 \varepsilon_0 = 4\pi k_e/c_t^2$  (5.21).

And it will be shown, later in this section, that where the quantum effects become measurable, that the constants  $G$ ,  $k_e$ ,  $\varepsilon_0$ , and  $\mu_0$  are no longer valid.

Therefore,  $G$ ,  $k_e$ ,  $\varepsilon_0$ ,  $\mu_0$ , and  $h$  are **not** fundamental constants.

(b) In quantum physics, the reduced Planck constant,  $\hbar = h/2\pi$ , is commonly used instead of the full Planck constant because it simplifies finding solutions to some wave equations. But the derivations of the Schrödinger (5.16) and Dirac wave equations (5.17), in this article, shows that using the reduced Planck constant is only valid if the equations contain compensating  $\pi$ -based terms. In Schrödinger's equation, the compensating  $2\pi$  terms:  $\partial^2 \Psi(r, t)/\partial r^2 = -(2\pi)^2/r^2 \Psi(r, t)$  and  $\partial \Psi(r, t)/\partial t = (i 2\pi/t) \Psi(r, t)$ . And Dirac's equation, requires the compensating  $\pi$  term:  $A_0 = \pi(c_q/c)((1/t)\tau + i(1/r)\rho)$ .

(c) The notions of point-like particle and fields might be incorrect. Note that the Gauss's gravity and electric field equations were derived without using Gauss's theorem, notions of pressure on a surface, and field. Action at a distance can be explained in terms of mass and charge bending spacetime and quantum entanglement.

From a mathematical (transfinite) perspective, either a point-like particle is a point, which has measure zero and will pass through any slit or is infinitesimal, has a size greater than zero, but still smaller than any constant finite-sized slit. Empirical evidence suggests slits can block particles (for example photons), which implies a non-point-like size. If particles have a wavelength greater than zero, then they probably have a particle size greater than zero, for example, a mass radius,  $r = k_m/2\pi m$ , and circumference/wavelength,  $\lambda = 2\pi r$ . This has implications with respect to: torque applied to a particle, whether

particle can rotate or spin with angular momentum, and have orientations.

- (10) Using the ratios instead of empirical constants in equations would show the shared principles underlying the different laws of physics. For example, the speed of light ratio,  $c_t$ , is a component of the constants:  $G = c_m c_t^2$ ,  $k_e = (c_q^2/c_m)c_t^2$ ,  $\varepsilon_0 = 1/(4\pi(c_q^2/c_m)c_t^2)$ ,  $h = k_m c_t$ .
- (11) Using the ratios instead of the opaque constants can sometimes simplify equations. For example: the Compton wavelength equation,  $r = h/mc$ , simplified to  $r = k_m/m$ ; and the fine structure constant,  $\alpha = q_e^2/2\varepsilon_0 hc$  simplified to  $\alpha = q_e^2/q_p^2$ .
- (12) Empirical laws, like Newton's gravity law (5.3) and Coulomb's charge law (5.7), and the Planck relation (5.41) are *descriptions* of relations. Deriving the equations from the ratios *explains* the relations. Further, all the other derivations, in this article, were much shorter and simpler than other derivations, which shows that the ratios are an important new tool for physicists and engineers.
- (13) The derivations of:  $\nabla \cdot \mathbf{g} = -4\pi G\rho$  from  $\nabla \cdot \mathbf{g} = -2Gm/r^3$  (5.5),  $\nabla \cdot \mathbf{E} = -\rho/\varepsilon_0$  (5.11) from  $\nabla \cdot \mathbf{E} = -2k_e q/r^3$  (5.10), and  $\partial \mathbf{B}/\partial t = -\mu_0 \rho$  from  $\partial \mathbf{B}/\partial t = -2k_e q/r^3$  (5.21), show that the use of mass and charge density,  $\rho$ , and the definitions of  $\varepsilon_0$  and  $\mu_0$  are unnecessary complications that obfuscate the pattern,  $\partial f(x, y, r)/\partial r = -2k_{x,y}y/r^3$ , and the inverse square pattern,  $f(x, y, r) = k_{x,y}y/r^2$ . Likewise, the  $4\pi G$  in  $\kappa = 2(4\pi G)/c^4$  and the energy density in the stress-energy tensor,  $T_{\mu,\nu}$ , in Einstein's field equations [18] also obfuscates the inverse square assumption.
- (14) The derivation of the Lorentz (5.15) and Faraday laws (5.20) shows that:  $F_E = qE$  and  $F_B = qvB = q(v/c)E \Rightarrow c = E/B$ .
- (15) The magnetic field,  $B$ , was derived from the torque of a moving charged particle on another charged particle's wave orientation (5.14). The following are consequences:
  - (a) Magnetism might be a pseudo (fictitious) force, where a moving charge causes a torque on the wave orientation of other charged particles.
  - (b) The intrinsic angular momentum,  $\pm\hbar/2$ , is the angular momentum of the mass-charge wave. When the wave orientation rotates, it will tend to be  $\pm\pi$  radians, which would correspond to  $\pm\frac{1}{2}$ -spin.  
The fractional wave orientation rotations of charged particles in the nucleus of an atom would be due to the geometric arrangement of the charged particles.
- (16) The derivation of Schwarzschild's time dilation and black hole metric (5.8) [17] [1] using ratios:
  - (a) was much shorter and simpler than Schwarzschild's derivation;
  - (b) did not require the complexity of manipulating Christoffel symbols, calculating determinants, partial derivatives, etc.;
  - (c) points to a way of simplifying the finding of solutions to Einstein's field equations: 1) the ratios can be used to derive the components of the metric,  $g_{\mu,\nu}$ , independent of Einstein's field equations, and 2) the field equations can be first solved as 2-dimensional tensors and then generalized to 4-dimensional tensors (5.9).

- (17) Using the quantum units,  $r_c$  and  $t_c$ :  $r_c/t_c^2 \approx 2.2184088232 \cdot 10^{51} \text{ m s}^{-2}$ , which suggests a maximum acceleration for masses.
- (18) The simplification of  $\mu_0$  into the quantum units shows two interesting relationships:

$$\begin{aligned} \mu_0 &:= \frac{1}{c_t^2 \varepsilon_0} = \frac{4\pi k_e}{c_t^2} = 4\pi \frac{c_q^2}{c_m} = 4\pi \frac{(r_c/q_c)^2}{r_c/m_c} = 4\pi \frac{m_c r_c}{q_c^2} = 4\pi \frac{k_m}{q_c^2} \\ &\approx 4\pi \frac{2.2102190930 \cdot 10^{-42}}{2.2102190930 \cdot 10^{-35}} = 4\pi \cdot 10^{-7} \text{ kg m C}^{-2} = 4\pi \cdot 10^{-7} \text{ H m}^{-1}. \end{aligned} \quad (6.5)$$

- (a) The first time  $k_m = m_c r_c$  appears is in the derivation of the Planck relation and Planck constant,  $h = k_m c$  (5.12), the second time in the Compton wavelength,  $r = k_m/m$  (5.13). And now,  $k_m$  appears as a components of  $k_e$  and  $\mu_0$ .
- (b) It is an open question why  $\frac{c_q^2}{c_m} = \frac{(r_c/q_c)^2}{r_c/m_c} = \frac{k_m}{q_c^2} = 1.0 \cdot 10^{-7}$  exactly.
- (19) Two subtypes are related via the ratios of the same super-type (5.15).
- (a) For example, the quantum charge and reduced Planck charge units are related via the ratio:  $q_c^2/q_p^2 = 2\pi \Rightarrow q_p = q_c/\sqrt{2\pi}$ .
- (b) The CODATA electron coupling version of the fine structure constant,  $\alpha$  is defined as:  $\alpha = q_e^2/4\pi\varepsilon_0\hbar c = q_e^2/2\varepsilon_0\hbar c$  [3].
- (i) The derivation of  $\alpha$ , in this article (5.15), is much simpler because it is the ratio of two subtypes of charge: elementary (electron) charge,  $q_e^2$  and Planck charge,  $q_p^2$ :  $\alpha = q_e^2/q_p^2 \approx 0.0072973526$ , which is the empirical CODATA value [3].
- (ii) The following steps show that the CODATA definition reduces to the ratio-derived equation:

$$\begin{aligned} \varepsilon_0 &:= 1/4\pi k_e = 1/(4\pi(c_q^2/c_m)c_t^2) \quad \wedge \quad h = k_m c_t \\ &\Rightarrow \quad \varepsilon_0 \hbar c = k_m c_t^2/(4\pi(c_q^2/c_m)c_t^2) = k_m/(4\pi(c_q^2/c_m)) \\ &\quad = m_c r_c/(4\pi((r_c/q_c)^2/(r_c/m_c))) = q_c^2/4\pi = q_p^2/2. \end{aligned} \quad (6.6)$$

$$\alpha = q_e^2/2\varepsilon_0\hbar c \quad \wedge \quad \varepsilon_0\hbar c = q_p^2/2 \quad \Rightarrow \quad \alpha = q_e^2/2(q_p^2/2) = q_e^2/q_p^2. \quad (6.7)$$

- (iii) As shown above, CODATA defines the fine structure constant in terms of a relationship to the Planck constant, hence, the ratio containing the reduced Planck unit,  $q_p$ :  $\alpha = q_e^2/q_p^2$ .

The quantum unit,  $q_c$ , appears naturally in the derivation of  $k_e$ , where  $\varepsilon_0 := 1/4\pi k_e$ . Therefore, a better definition to describe particle interaction with a charge (electromagnetic) wave is:  $\alpha = q_e^2/q_c^2$ , where the current CODATA value would be divided by  $2\pi$ .

- (iv) Other fine structure constants can also be expressed more simply as the ratios of two subtypes of fields, for example, an electron gravity coupling constant can be expressed as the ratio of a stationary electron mass to a quantum mass unit:  $\alpha_m = m_e^2/m_p^2$  or  $\alpha_m = m_e^2/m_c^2$ .

- (20) Special and general relativity assume covariance, which states that the laws of physics are invariant in every coordinate frame of reference [6]. Where

the space near each coordinate point is Euclidean-like (for example, Riemann and pseudo-Riemann surfaces), the ratios are constant near every coordinate frame of reference. And because the laws of physics are derived from the ratios, the laws of physics are the same near every coordinate frame of reference.

- (a) The ratio-based derivations of the spacetime equations, in this article (5.7), do not rely on the Lorentz transformations or Einstein's postulates [6]. The derivations do not even require the notion of light.
- (b) The ratio-based derivations are also valid for spacemass and space-charge.
- (c) The special relativity time dilation equation 5.24 was derived from the distance-to-time ratio,  $r = (r_c/t_c)t$ , and combined with the distance-to-mass ratio,  $r = (r_c/m_c)m$ , (5.10) yielded Schwarzschild's gravitational time dilation and black hole metric equations (5.31).
- (21) The derivation of Schrödinger (5.16) and Dirac wave equations (5.17), in this article, differs from other derivations:
  - (a) The derivations here are more rigorous because the energy-momentum term,  $\hbar^2/2m$ , was derived, in this article, from the Planck relation (5.58), where the Planck relation was also rigorously derived from the ratios (5.12). Other derivations assume (define) the energy-momentum relation as:  $(\mathbf{p} \cdot \mathbf{p})/2m = \hbar^2/2m$  [16] [7].
- (22) The quantum extensions to: Schwarzschild's time dilation (5.74) black hole metric (5.34), Newton's gravity force (5.78), and Coulomb's charge force (5.81) make quantifiable predictions:
  - (a) The gravitation and charge forces peak at finite amounts as  $r \rightarrow 0$ :  $\lim_{r \rightarrow 0} F = c^2 m_1 m_2 / k_m$  and  $\lim_{r \rightarrow 0} F = c^2 q_1 q_2 k_m / k_q^2$ . Finite maximum gravity and charge forces: 1) allows radioactivity, finite sloped energy well walls, and possibly black hole evaporation; 2) eliminates the problem of forces going to infinity as  $r \rightarrow 0$ .
  - (b) The Schwarzschild metric, gravity, and charge equations reduce to the classic equations, where the distance between masses and charges is sufficiently large or the masses and charges sufficiently large that the quantum effect is not measurable. **Note** that  $G$ ,  $k_e$ ,  $\varepsilon_0$ ,  $\mu_0$ , and  $\kappa$  (Einstein's constant, which contains  $G$ ) are not valid, where the quantum effects becomes measurable.
  - (c) And the covariant tensor components, in Einstein's field equations, that had the units  $1/\text{distance}^2$ , will now have the more complex units,  $1/\sqrt{(\text{distance}^4/c_m^2) + k_m^2}$ .
  - (d)  $1/\sqrt{(\text{distance}^4/c_m^2) + k_m^2}$  implies that as distance  $\rightarrow 0$ , spacetime curvature peaks at a finite amount, which predicts that black holes probably have sizes  $> 0$  (are probably not singularities). If there was a "big bang", then it might not have originated from a singularity.

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