



\hat{n} - polarization axis

for transverse waves here
 $\hat{n} \cdot \hat{z} = 0$

$$\hat{n} = \cos\theta \hat{x} + \sin\theta \hat{y}$$

In ch. 7 we derived

(no ρ , no J , free space)

$$\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\nabla^2 \vec{B} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}} = 3 \times 10^8 \frac{m}{s}$$

Monochromatic Plane waves

↳ single wavelength
(single color)

↳ only dependence
on one coordinate
will choose z

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

so we get
$$\frac{\partial^2 \vec{E}}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\frac{\partial^2 \vec{B}}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2}$$

⇒ solutions to above equations is $\vec{E} = \vec{E}_0 \cos(kz - \omega t + \delta)$

$$(3) \quad \vec{E} = \text{Re} \left[\left(\vec{E}_0 e^{i(kz - \omega t)} \right) \right] = \text{Re} \left[\tilde{\vec{E}} \right]$$

$$(4) \quad \vec{B} = \text{Re} \left[\left(\vec{B}_0 e^{i(kz - \omega t)} \right) \right] = \text{Re} \left[\tilde{\vec{B}} \right]$$

complex amplitudes

Goal: find additional information about 3 and 4

sub into maxwell's equations

$$\left. \begin{aligned} \vec{\nabla} \cdot \vec{E} &= 0 \quad (\text{b/c } \rho=0) \\ \vec{\nabla} \cdot \vec{B} &= 0 \quad (\text{always}) \end{aligned} \right\}$$

$$\vec{\nabla} \cdot \tilde{\vec{E}} = \frac{\partial}{\partial z} \tilde{E}_z = \tilde{E}_{0,z} (ik) e^{i(kz - \omega t)} = 0$$

$$\text{Similarly } \vec{\nabla} \cdot \tilde{\vec{B}}(z,t) = \tilde{B}_{0,z} (ik) e^{i(kz - \omega t)} = 0$$

$$\text{so } \tilde{E}_{0,z} = \tilde{B}_{0,z} = 0$$

so EM waves are not longitudinal, they are transverse

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (\text{Faraday's law})$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$-\frac{\partial \vec{B}}{\partial t} = -i\omega \vec{B}_0 e^{i(kz - \omega t)}$$

$$\vec{\nabla} \times \vec{E} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & \frac{\partial}{\partial z} \\ \tilde{E}_x & \tilde{E}_y & 0 \end{vmatrix} = \hat{x} \left(-\frac{\partial}{\partial z} \tilde{E}_y \right) - \hat{y} \left(-\frac{\partial}{\partial z} \tilde{E}_x \right) \\ = \hat{y} \frac{\partial}{\partial z} \tilde{E}_x - \hat{x} \frac{\partial}{\partial z} \tilde{E}_y$$

$$\text{so } \frac{\partial}{\partial z} \tilde{E}_x \hat{y} - \frac{\partial}{\partial z} \tilde{E}_y \hat{x} = -i\omega \vec{B}_0 e^{i(kz - \omega t)} \\ = -i\omega \vec{B}$$

$$\text{so } -i\omega \tilde{B}_y = \frac{\partial}{\partial z} \tilde{E}_x \quad \text{and} \quad -i\omega \tilde{B}_x = \frac{\partial}{\partial z} \tilde{E}_y$$

$$\Rightarrow -\omega \tilde{B}_{0,y} = k \tilde{E}_{0,x}$$

$$-\omega \tilde{B}_{0,x} = -k \tilde{E}_{0,y}$$

$$\rightarrow \vec{B}_0 = \frac{k}{\omega} (\hat{z} \times \vec{E}_0) \quad \text{so } \boxed{|\vec{B}_0| = \frac{1}{c} |\vec{E}_0|}$$

(no phase shift between E and B waves)

$$\delta_E = \delta_B \quad (\text{in vacuum})$$