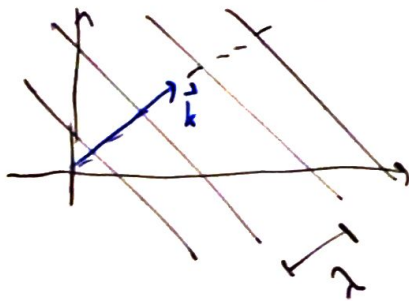


How to write a 2D plane wave



$$\psi(x, y, t)$$

when $|k| = \frac{2\pi}{\lambda}$

wavefronts defined by $\vec{k} \cdot \vec{r} = \text{const.}$

$$\psi(x, y, t) = A \exp(i(\vec{k} \cdot \vec{r} - \omega t))$$

write in components: $\vec{r} = x\hat{i} + y\hat{j}$ $\vec{k} = k_x\hat{i} + k_y\hat{j}$

$$\psi(x, y, t) = A \exp(i(k_x x + k_y y - \omega t))$$

Example: 2D wave travelling at 30° to x-axis

$$\lambda = 20\text{m} \quad v = 10\text{m/s}$$

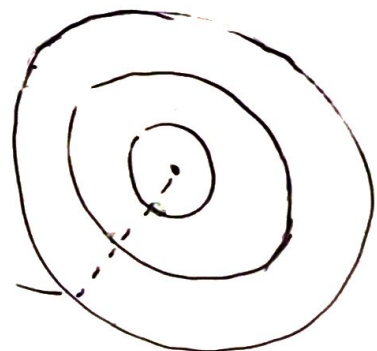
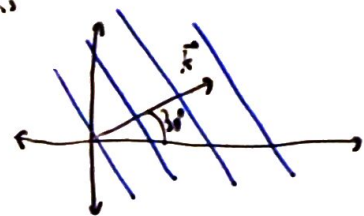
$$\psi(x, y, t) = k = \frac{2\pi}{\lambda} = \frac{\pi}{10} \text{m}^{-1}$$

$$\omega = 2\pi f = 2\pi \frac{v}{\lambda} = \pi \text{s}^{-1}$$

$$k_x = \frac{\pi}{10} \cos 30^\circ \quad k_y = \frac{\pi}{10} \sin 30^\circ = \frac{\pi}{20}$$

$$\psi(x, y, t) = A \exp(i(\frac{\sqrt{3}\pi}{20}x + \frac{\pi}{20}y - \pi t))$$

Spherical wave: $\psi(r, t) = \frac{A}{r} e^{i(kr - \omega t)}$



at large r , the wave begins to look plane

↓ getting smaller as it goes out

Electromagnetic waves



$$\vec{E} = \vec{E}_0 \sin(\vec{k} \cdot \vec{r} - \omega t)$$

$$\vec{B} = \vec{B}_0 \sin(\vec{k} \cdot \vec{r} - \omega t)$$

$$|\vec{E}_0| = c|\vec{B}_0|$$

there is an energy density to do with E and B fields

$$u_E = \frac{1}{2} \epsilon_0 E^2$$

$$u_B = \frac{1}{2} \frac{1}{\mu_0} B^2 \quad (\text{woah! } E, \mu)$$

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

Irradiance of the EM wave $I = \frac{1}{2} \epsilon_0 c E_0^2$ watts/m²

↑ This is what we observe

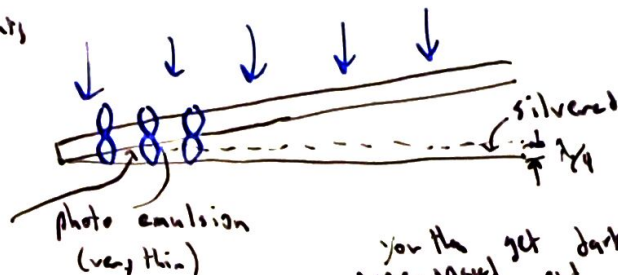
$$I \propto E_0^2$$

why so excited about the electric field?

Wiener experiments

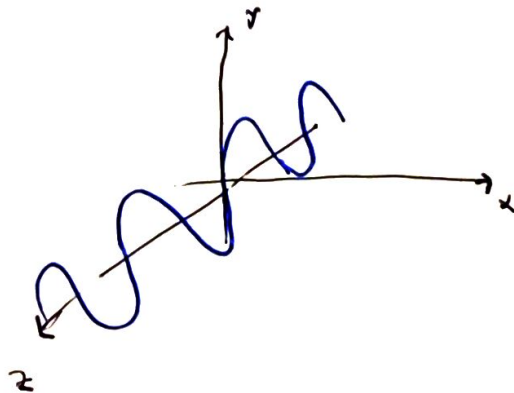
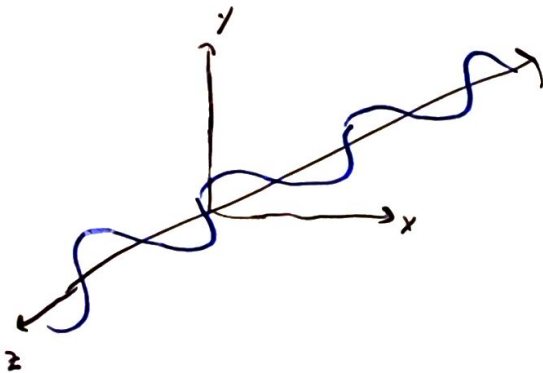
get standing waves

antinode on the photo emulsion

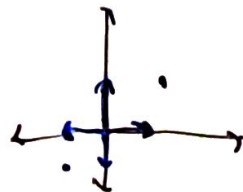
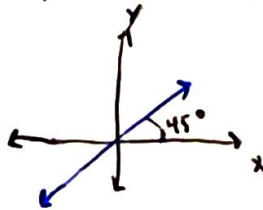


↳ the electric field does the exposure

polarization: $\vec{E} = E_0 \sin(kz - \omega t)$

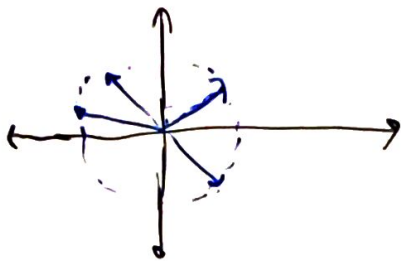


we define the direction of polarization by looking into the beam
e.g. 45° linear polarization



we would write this as $\vec{E} = E_0 \sin(kz - \omega t) \hat{x} + E_0 \sin(kz - \omega t) \hat{y}$

Circular polarization - eg. left circularly polarized

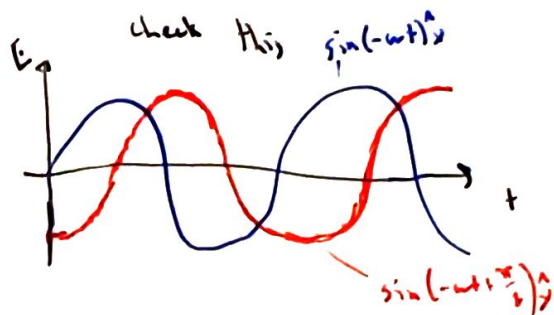


moves around a full circle in 1 wavelength

$$\vec{E} = E_0 \sin(kz - \omega t) \hat{x} + E_0 \sin(kz - \omega t + \frac{\pi}{2}) \hat{y}$$

need about dopler effect

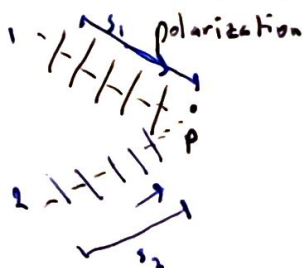
$$\frac{\lambda'}{\lambda} = 1 - \frac{v}{c}$$



Ch 7

Light interference

We pretty much treat the light as a scalar waves; we use just one



$$E_{1,p} = E_0 \cos(k s_1 - \omega t + \phi_1)$$

$$k = \frac{2\pi}{\lambda}$$

$$E_{2,p} = E_0 \cos(k s_2 - \omega t + \phi_2)$$

ϕ_1, ϕ_2 are the phase constants

$$E_p = E_{1,p} + E_{2,p} \text{ - instantaneous E-field at } p$$

we want the irradiance at p, power/area e.g. $\frac{\text{watts}}{\text{m}^2}$

Irradiance \propto amplitude²

$$\tilde{E}_{1,p} = E_1 \exp(i(k s_1 + \phi_1))$$

$$\tilde{E}_p = \tilde{E}_{1,p} + \tilde{E}_{2,p}$$

$$\tilde{E}_{2,p} = E_2 \exp(i(k s_2 + \phi_2))$$

$$\tilde{E}_p \tilde{E}_p^* = (\tilde{E}_{1,p} + \tilde{E}_{2,p})(\tilde{E}_{1,p} + \tilde{E}_{2,p})^*$$

$$\tilde{E}_p \tilde{E}_p^* = I$$

$$= I_1 + I_2 + E_1 E_2 \exp(i(k(s_1 - s_2) + (\phi_1 - \phi_2))) + E_1 E_2 \exp(i(k(s_2 - s_1) + (\phi_2 - \phi_1)))$$

$$E_1 = \sqrt{I_1} \quad E_2 = \sqrt{I_2}$$

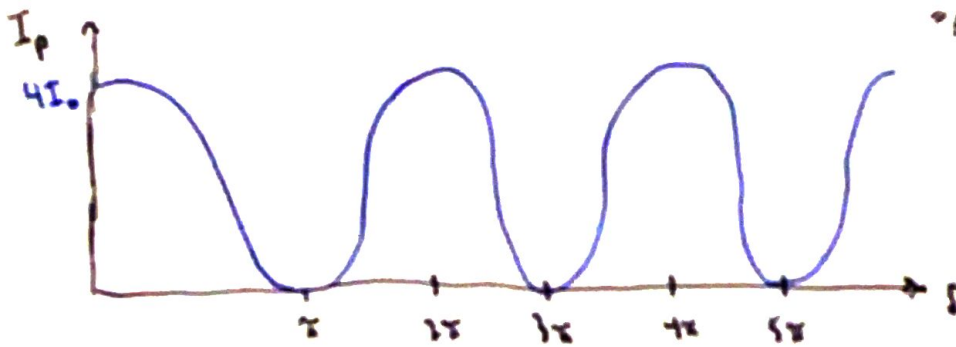
$$I = I_1 + I_2 + \sqrt{I_1 I_2} 2 \cos(\delta) \quad \text{when } \delta = k(s_2 - s_1) + \phi_2 - \phi_1$$

$$I_p = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\delta)$$

let $I_1 = I_2 = I_0$

then $I_p = 2I_0 + 2I_0 \cos(\delta)$

→ refer to this as
"fringes"



use double angle formula

$$I_p = 2I_0 + 2I_0 \cos^2(\delta/2)$$

$$= 4I_0 \cos^2(\delta/2)$$

thus 'cosine squared' fringes
- typical of 2 beam interference

If $I_1 \neq I_2$

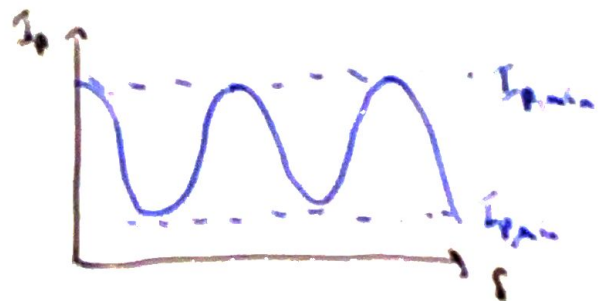
$$\hookrightarrow I_p = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\delta)$$

$$I_{p, \max} = I_1 + I_2 + 2\sqrt{I_1 I_2}$$

$$I_{p, \min} = I_1 + I_2 - 2\sqrt{I_1 I_2}$$

define visibility: $V = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$

$$0 < V < 1$$



So far have assumed that ϕ_1 & ϕ_2
are fixed ($\phi_2 - \phi_1$ constant)

But light from different sources (even different lasers)

will have ϕ_1 and ϕ_2 varying rapidly

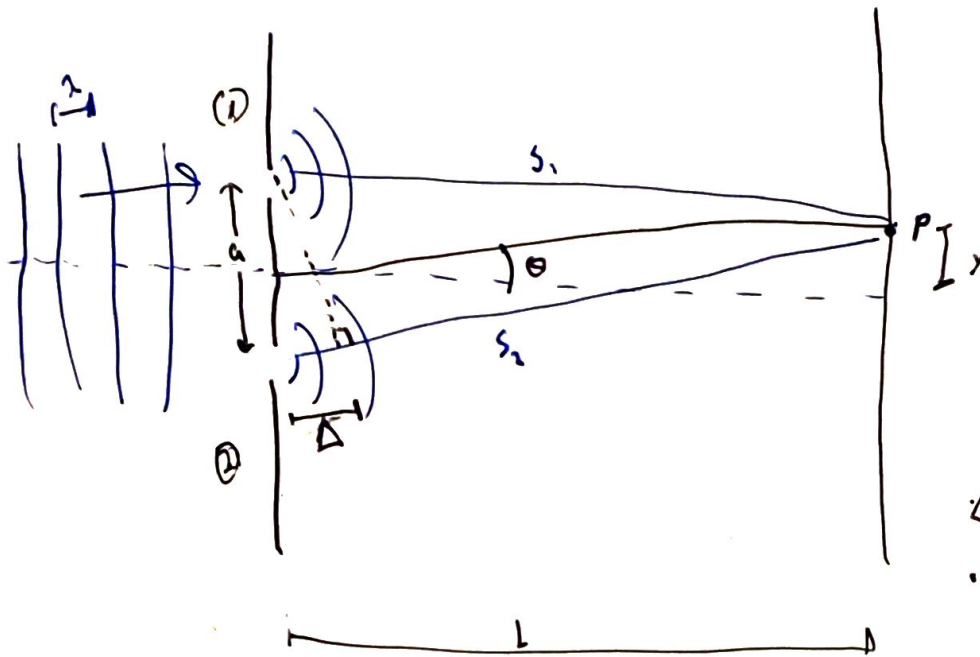
$$I_p = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\delta)$$

→ will average out to zero

$$I_p = I_1 + I_2 \quad \text{→ call this sources incoherent}$$

$$\text{average } I_p = I_1 + I_2$$

Young's 2 slits - an example of a wavefront division interferometer

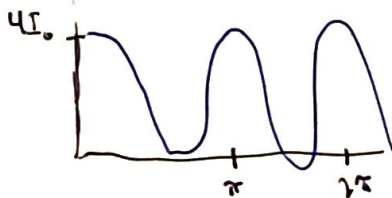


slits of equal size
 → give the same contribution
 $s_2 - s_1$ is not so big that
 the amplitudes at P do not
 differ significantly

$$\Delta = s_1 - s_2 \approx a \sin \theta$$

$$\text{let } \phi_1 = \phi_2$$

then $\delta = k(s_2 - s_1) = k\Delta$ so $I_P = 4I_0 \cos^2\left(\frac{\delta}{2}\right) = 4I_0 \cos^2\left(\frac{k}{2} a \sin \theta\right)$



so we get maxima when

$$\frac{k}{2} a \sin \theta = 0, \pi, 2\pi, \dots$$

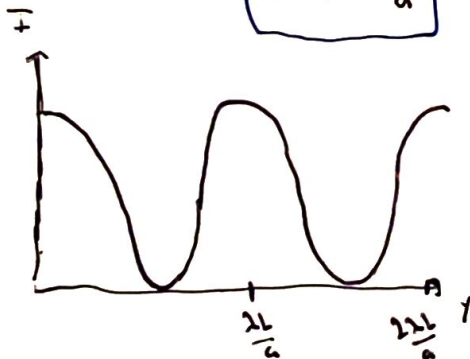
$$\text{so } \frac{\pi}{2} a \frac{y_{\max}}{L} = m\pi$$

$$\text{so } \Delta y_{\max} = \frac{\lambda}{a}$$

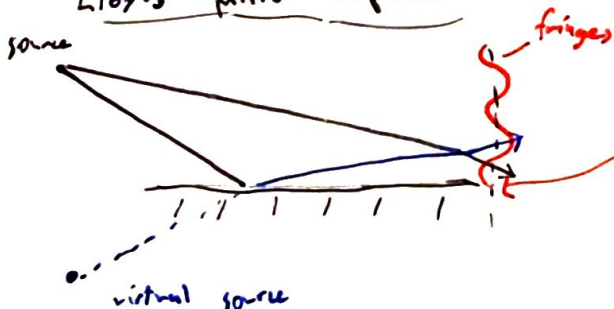
$$m=0, 1, 2, \dots \Rightarrow y_{\max} = \frac{mL\lambda}{a}$$

so larger $\lambda \rightarrow$ spots spread apart

decrease slit spacing \rightarrow spots spread apart



Lloyd's mirror experiment



destructive interference

- because reflected and transmitted light
 are 180° out of phase

→ phase change due to reflection