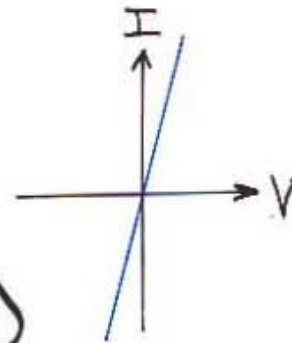
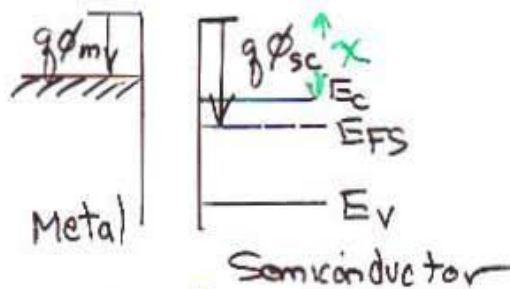


# Ohmic Contacts

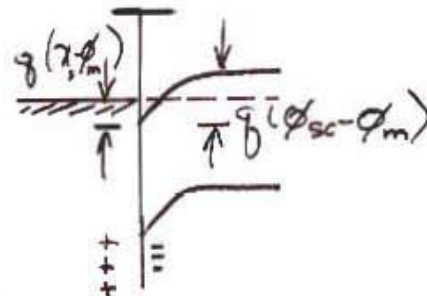
- low resistance
- pass current equally in both directions (no rectification)
- no depletion region



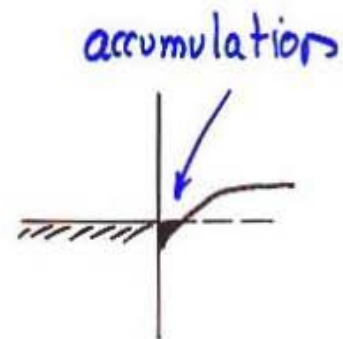
n-Type :  $\phi_m < \phi_{sc}$



Before



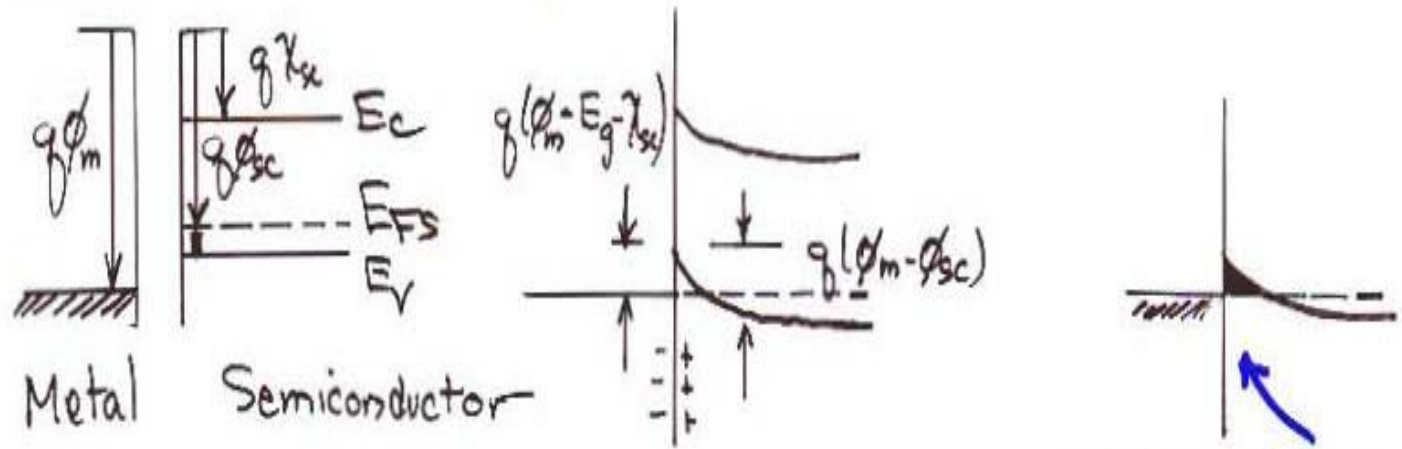
After



No barrier to  $e^-$   
low barrier to bulk

Accumulation of majority carriers.

P-Type:  $\phi_m > \phi_{sc}$



accumulation

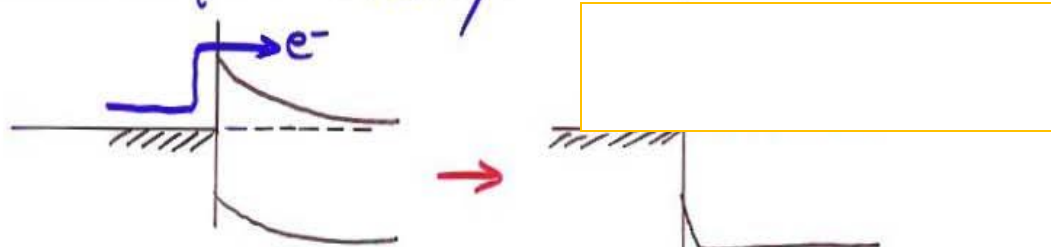
No barriers to holes  
Low barrier to bulk

of majority carriers

# Summary of Metal-Semiconductor Contacts

	Schottky	ohmic
n-Type	$\phi_m > \phi_{sc}$	$\phi_m < \phi_{sc}$
p-Type	$\phi_m < \phi_{sc}$	$\phi_m > \phi_{sc}$

How to "fool" Schottky?



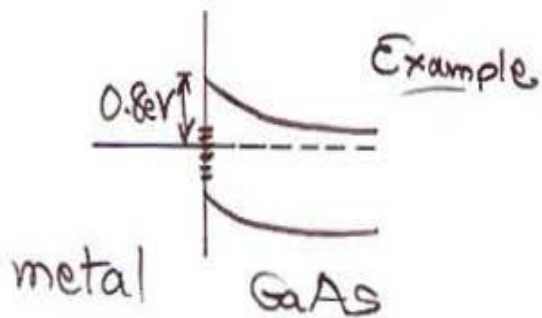
Charge  through very thin barriers.  
How?

Sometimes need such tricks to control "real" Schottky barriers.

Sometimes need such tricks to control "real" Schottky barriers.

$$\phi_{V_0} \neq \phi(\phi_m - \phi_{sc})$$

Due to Surface States!



$E_F$  "pinned" in gap  
Localized States due to:

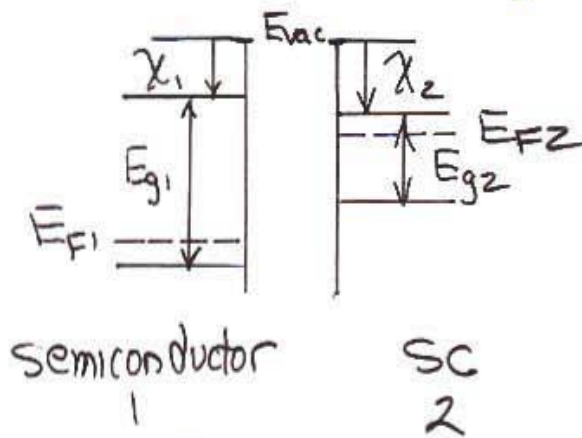
- dangling atomic bonds
- reactions
- diffusion

# Heterojunctions

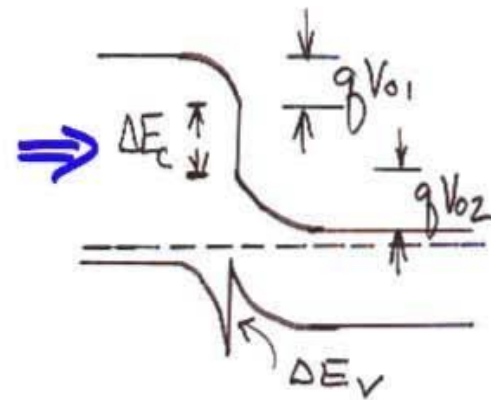
So far: single sc-sc junction: "homojunction"  
Metal-SC junction: Schottky junction

Consider now two<sup>sc</sup> junctions.

→ different  $E_g$ ,  $\chi$ ,  $\epsilon$ ,  $E_F$



Before



After



- Now get discontinuities in band edges  $\Delta E_c, \Delta E_v$
- Also, different  $qV_{01}$  and  $qV_{02}$

How to predict?

Important for designing  
next generation devices.

First,  $\Delta E_g = E_{g1} - E_{g2} = \Delta E_c + \Delta E_v$

and  $qV_0 = qV_{01} + qV_{02}$

and  $\frac{V_{01}}{V_{02}} = \frac{\epsilon_2 N_{d2}}{\epsilon_1 N_{d1}}$  ← a



From requirement that  $D_1 = D_2$

$$\epsilon_1 E_1 = \epsilon_2 E_2$$

$$\text{at boundary, } E_1 = E_{1 \max} = \frac{q}{\epsilon_1} N_{a1} X_1 \quad \text{S\&B Eq. 5-17}$$

$$E_2 = E_{2 \max} = \frac{q}{\epsilon_2} N_{d1} X_2$$

$$\text{so } N_{a1} X_1 = N_{d2} X_2 \quad (\text{same as with regular homojunctions})$$

$$\frac{dE_1}{dx} = \frac{1}{\epsilon_1} (-q N_a)$$

$$E_1 = \frac{1}{\epsilon_1} (-q N_a) x$$

$$V_{01} = \int_{x_1}^0 \frac{1}{\epsilon_1} (-q N_a) x dx = \frac{q N_a x_1^2}{2 \epsilon_1}$$

$$V_{02} = \frac{q N_d x_2^2}{2 \epsilon_2}$$

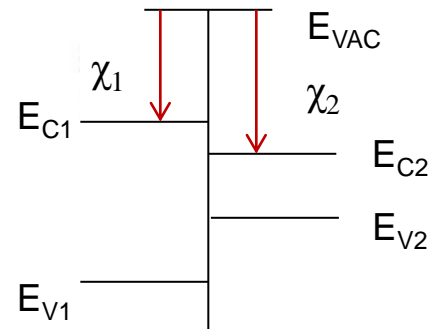
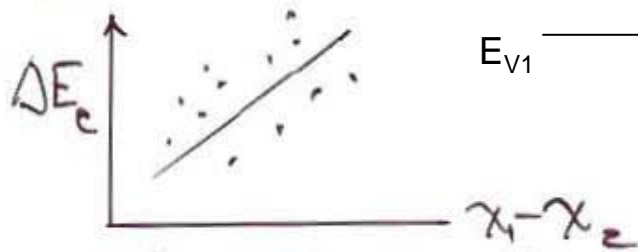
$$\frac{V_{01}}{V_{02}} = \frac{N_a x_1^2 / \epsilon_1}{N_d x_2^2 / \epsilon_2}$$

$$\frac{x_1}{x_2} = \frac{N_d}{N_a}$$

$$\frac{V_{01}}{V_{02}} = \frac{N_d \epsilon_2}{N_a \epsilon_1}$$



Ideal Case:  $\Delta E_c = g(\chi_2 - \chi_1)$   
 But not often seen  
 (Anderson Rule)



For III-V compounds (e.g., GaAs, AlAs, InAs)

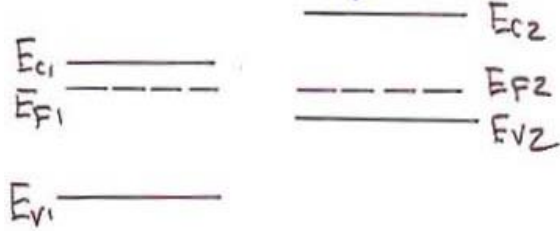
$\frac{2}{3}$  Rule:  $\Delta E_g : \frac{2}{3} \Delta E_c, \frac{1}{3} \Delta E_v$

In general, need experimental values.

Exact solution: Poisson's equation with doping, space charge,  $\Delta E_c$ ,  $\Delta E_v$ .

But can sketch approximate band diagrams without detailed calculation. (Art 101)

Step 1: Align  $E_F$  levels with bands flat



Leave room for transition region

Step 2: Locate  $x=0$  ( $N_1 N_1 = N_2 N_2$ )

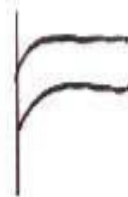


Electrons will flow from n to p and  
holes " " " p to n.

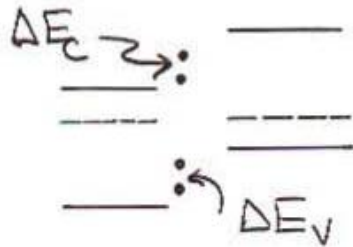
so n-Type side has



and p-Type side has



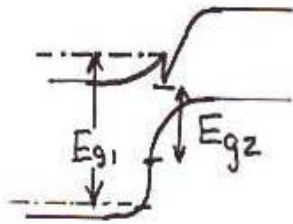
step 2b:  $P_{\text{utin}}$  known



(where the dots go determines the  $V_{01}$  and  $V_{02}$ . Only approximate calculation from Poisson's equation)

step 3: Connect up the valence bands and conduction bands with band bending

(only approximate, can calculate from Poisson's equation)



Keep  $E_{g1}$  constant at junction  
"  $E_{g2}$  " " "

$\Delta E_v$  and  $\Delta E_c$  make up the difference.

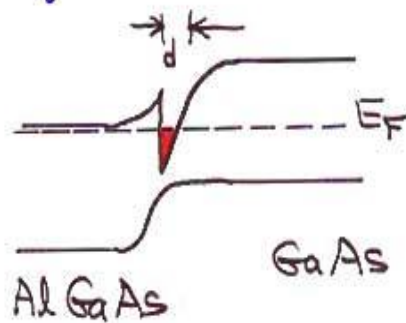
## Observations

1) In general,  $\Delta E_c \neq \Delta E_v$

so  $e^-$  flow versus  $h^+$  flow across junction can be different

In fact, we can control one over the other by choosing semiconductors appropriately.

2) Band bending at heterojunction can give quantum wells.



Electrons spill into well  
( $E_F$  above  $E_c$ )

$$d \lesssim 0.02 \mu\text{m} \quad (200 \text{\AA})$$

Very sharp levels with  $d$  is narrow. Very high mobility and transport



3) Can have n-n heterojunctions  
n-p  
p-n  
p-p

Potential barrier  
can now be  
controlled by  
material band gap  
and doping

Semiconductor  $E_g$  and doping are  
extra degrees of freedom.

Basis for

