MA30044/MA40044/MA50181, Mathematical Methods I, 2020

Problem Sheet 2: Advanced Sturm-Liouville Equations

SOLUTIONS

Q1.

(i)
$$\phi_0 = 1$$
: $\phi'_0 = 0, \lambda_0 \phi_0 = 0 \implies \lambda_0 = 0$;
 $\phi_1 = x$: $((1 - x^2)\phi'_1(x))' + \lambda_1 \phi_1(x) = (1 - x^2)' + \lambda_1 x = -2x + \lambda_1 x \implies \lambda_1 = 2$;
 $\phi_2 = \frac{1}{2}(3x^2 - 1)$: $((1 - x^2)\phi'_2(x))' + \lambda_2 \phi_2(x) = 3((1 - x^2)x)' + \lambda_2 \frac{3x^2 - 1}{2}$
 $= 3(1 - 3x^2) + \lambda_2 \frac{3x^2 - 1}{2} - 6\frac{3x^2 - 1}{2} + \lambda_2 \frac{3x^2 - 1}{2} \implies \lambda_2 = 6$.

(ii)
$$\int_{-1}^{1} \phi_0(x)\phi_1(x)dx = \int_{-1}^{1} x dx = \frac{x^2}{2} \Big|_{-1}^{1} = 0;$$

$$\int_{-1}^{1} \phi_0(x)\phi_2(x)dx = \int_{-1}^{1} \frac{3x^2 - 1}{2} dx = \frac{x^3 - x}{2} \Big|_{-1}^{1} = 0;$$

$$\int_{-1}^{1} \phi_1(x)\phi_2(x)dx = \int_{-1}^{1} \frac{3x^2 - 1}{2} x dx = \int_{-1}^{1} \frac{3x^3 - x}{2} x dx = \left[\frac{3}{8}x^4 - \frac{1}{4}x^2\right]_{-1}^{1} = 0.$$

Hence, the functions ϕ_j , j=1,2,3 are orthogonal to each other, and so the system $\{\phi_0,\phi_1,\phi_2\}$ is orthogonal.

(iii)
$$\phi_3(x) = x^3 + ax \implies \phi_3'(x) = 3x^2 + a,$$

and so

$$((1-x^2)\phi_3'(x))' = ((1-x^2)(3x^2+a))' = (-3x^4+(3-a)x^2+a)' = -12x^3+2(3-a)x,$$

and

$$((1-x^2)\phi_3'(x))' + \lambda_3\phi_3(x) = -12x^3 + 2(3-a)x + \lambda_3(x^3+ax) = 0.$$

Comparing the coefficients on x^3 , we obtain $\lambda_3 = 12$, and comparing the coefficients on x, we obtain 2(3-a) + 12a = 0, from which a = -3/5. Summarising,

$$\phi_3(x) = \frac{1}{5}(5x^3 - 3x), \qquad \lambda_3 = 12.$$

$$\phi_4(x) = x^4 + bx^2 + d \implies \phi'_4(x) = 4x^3 + 2bx,$$

and so

$$((1-x^2)\phi_4'(x))' = ((1-x^2)(4x^3+2bx))' = (-4x^5+(4-2b)x^3+2bx)' = -20x^4+6(2-b)x^2+2b,$$

and

$$((1-x^2)\phi_4'(x))' + \lambda_4\phi_4(x) = -20x^4 + 6(2-b)x^2 + 2b + \lambda_4(x^4 + bx^2 + d) = 0.$$
 (1)

Comparing the coefficients on x^4 , we obtain $\lambda_4 = 20$, and (1) takes the form

$$6(2-b)x^2 + 2b + 20bx^2 + 20d = 0. (2)$$

Comparing the coefficients on x^2 in (2), we obtain 6(2-b) + 20b = 0, from which b = -6/7, and finally comparing the coefficients on x^0 in (2), we obtain 2b + 20d = 0, from which d = 3/35. Summarising, we have

 $\phi_4(x) = \frac{1}{35}(35x^4 - 30x^2 + 3), \qquad \lambda_4 = 20.$

(iv) The only zero of $\phi_1(x)$ is x = 0.

The zeros of $\phi_2 = 3x^2 - 1$ are $x = \pm 1/\sqrt{3} = \pm 0.577350$, and the zeros of $\phi_3(x) = 5x^3 - 3x$ are at x = 0 and those x on which $5x^2 - 3$ vanishes, so $x = \pm \sqrt{3/5} = \pm 0.774597$.

Finally, to determine the zeros of ϕ_4 set $z=x^2$ in $35x^4-30x^2+3=0$. This results in the equation

$$z^2 - \frac{6}{7}z + \frac{3}{35} = 0,$$

SO

$$z = \frac{3}{7} \pm \sqrt{\frac{9}{49} - \frac{3}{35}} = \frac{1}{7} \left(3 \pm \sqrt{\frac{24}{5}} \right).$$

This gives four zeros for ϕ_4 :

$$x = \pm \sqrt{\frac{1}{7} \left(3 \pm \sqrt{\frac{24}{5}}\right)} = \begin{cases} \pm 0.861136, \\ \pm 0.339981. \end{cases}$$

 $\mathbf{Q2}.$

$$-(1-x^2)u''(x) + xu'(x) = \lambda u(x), \quad x \in (-1,1).$$
(3)

(i) Here, in terms of the notation used in the lectures, we have

$$\alpha(x)=1-x^2, \quad \beta(x)=-x, \quad \gamma(x)=0, \quad \delta(x)=1.$$

Hence,

$$p(x) = \exp\left(-\int_0^x \frac{ydy}{1 - y^2}\right) = \exp\left(\frac{1}{2}\int_0^x \left(-\frac{ydy}{1 - y^2}\right)\right) = \exp\left(\frac{1}{2}\log(1 - x^2)\right) = \sqrt{1 - x^2},$$

$$q(x) = 0, \qquad r(x) = \frac{\sqrt{1 - x^2}}{1 - x^2} = \frac{1}{\sqrt{1 - x^2}}.$$

Therefore, the SL form of the equation (3) is

$$-(\sqrt{1-x^2}u'(x))' = \frac{\lambda}{\sqrt{1-x^2}}u(x).$$

(ii) Consider the function $u(x) = \cos(\sqrt{\lambda}\cos^{-1}(x))$. Applying the chain rule and using the fact that $(\cos^{-1}(x))' = -1/\sqrt{1-x^2}$, we obtain

$$u'(x) = \sqrt{\lambda} \frac{1}{\sqrt{1 - x^2}} \sin(\sqrt{\lambda} \cos^{-1}(x)), \tag{4}$$

and hence

$$\sqrt{1-x^2}u'(x) = \sqrt{\lambda}\sin(\sqrt{\lambda}\cos^{-1}(x)).$$

Taking derivative once again, we have

$$-(\sqrt{1-x^2}u'(x))' = \frac{\lambda}{\sqrt{1-x^2}}u(x), \qquad x \in (-1,1).$$

The calculation for $u(x) = \sin(\sqrt{\lambda}\sin^{-1}(x))$ is similar.

The linear independence of

$$\cos(\sqrt{\lambda}\cos^{-1}(x)), \qquad \sin(\sqrt{\lambda}\sin^{-1}(x))$$

can be shown, for example, by considering the system

$$\begin{cases} \alpha \cos(\sqrt{\lambda}\cos^{-1}(x_1)) + \beta \sin(\sqrt{\lambda}\sin^{-1}(x_1)) = 0, \\ \alpha \cos(\sqrt{\lambda}\cos^{-1}(x_2)) + \beta \sin(\sqrt{\lambda}\sin^{-1}(x_2)) = 0, \end{cases}$$

with $x_1, x_2 \in (-1, 1)$. As there are pairs of points x_1, x_2 for which the determinant of this system does not vanish, one necessarily has $\alpha = \beta = 0$, from which the claim follows.

(iii) Part of the difficulty of this question is to sort out boundedness. Consider solutions to (3) of the form

$$u(x) = a\cos(\sqrt{\lambda}\cos^{-1}(x)) \tag{5}$$

with some constant $a \neq 0$. The derivative of (5), given by (4), is bounded near x = 1. Indeed, setting $y = \cos^{-1}(x)$, we have

$$\sqrt{\lambda} \frac{\sin(\sqrt{\lambda}y)}{\sqrt{1 - \cos^2(y)}} = \sqrt{\lambda} \frac{\sin(\sqrt{\lambda}y)}{\sin(y)},\tag{6}$$

which is a bounded function near $y = \cos^{-1}(1) = 0$ [can be verified by L'Hopital's rule or similar].

It remains to analyse the behaviour of (4) near x=-1. Since $x\to -1$ means $y\to \pi$, then this suggests we should change the variables according to $y=\pi-\cos^{-1}(x)$. Now the critical point is y=0.

This results in

$$\sqrt{\lambda} \frac{\sin(\sqrt{\lambda}(\pi - y))}{\sqrt{1 - \cos^2(\pi - y)}} = \sqrt{\lambda} \frac{\sin(\sqrt{\lambda}(\pi - y))}{\sin(y)}.$$
 (7)

rather than the expression (6). The expression (7) is bounded near $y = \pi - \cos^{-1}(-1) = 0$ if and only if the numerator vanishes at y = 0, which yields the condition $\sqrt{\lambda}\pi = n\pi$, $n \in \mathbb{Z}$, or equivalently $\lambda = n^2$, $n = 0, 1, 2 \dots$

(iv) Recall that $T_n(x) = \cos(n\cos^{-1}(x))$, $x \in [-1, 1]$. Taking the inner product of T_n and T_m with weight $1/\sqrt{1-x^2}$, we have

$$I_{mn} = \int_{-1}^{1} \frac{T_n(x)T_m(x)}{\sqrt{1-x^2}} dx = \int_{-1}^{1} \cos(n\cos^{-1}(x))\cos(m\cos^{-1}(x)) \frac{dx}{\sqrt{1-x^2}}.$$

Making the change of variables $y = \cos^{-1}(x)$, so $dy = -dx/\sqrt{1-x^2}$, yields

$$I_{mn} = \int_0^{\pi} \cos(ny) \cos(my) dy = 0, \qquad n \neq m,$$

as required.

(v)
$$T_0(x) = \cos(0 \cdot \cos^{-1}(x)) = \cos(0) = 1 \implies T_0(x) = 1;$$
$$T_1(x) = \cos(\cos^{-1}(x)) = x \implies T_1(x) = x.$$

Next, using the notation $\theta = \cos^{-1}(x)$, we have

$$T_2(x) = \cos(2\cos^{-1}(x)) = \cos(2\theta) = 2\cos^2(\theta) - 1 = 2x^2 - 1 \implies T_2(x) = 2x^2 - 1.$$

Similarly, using the formulae $\cos(3\theta) = 4\cos^3(\theta) - 3\cos(\theta)$ and $\cos(4\theta) = 8\cos^4(\theta) - 8\cos^2(\theta) + 1$, we obtain

$$T_3(x) = \cos(3\cos^{-1}(x)) = \cos(3\theta) = 4x^3 - 3x,$$

 $T_4(x) = \cos(4\cos^{-1}(x)) = \cos(4\theta) = 8x^4 - 8x^2 + 1.$

Q3.

(i) Here, in terms of the notation used in the lectures, we have

$$\alpha(x) = 1, \quad \beta(x) = -\frac{x}{2}, \quad \gamma(x) = 0, \quad \delta(x) = 1.$$

Hence,

$$p(x) = \exp\left(-\frac{1}{2}\int_0^x y dy\right) = \exp\left(-\frac{x^2}{4}\right), \qquad q(x) = 0, \qquad r(x) = p(x) = \exp\left(-\frac{x^2}{4}\right).$$

Therefore, the SL form of the equation (3) is

$$-\left(\exp\left(-\frac{x^2}{4}\right)u'(x)\right)' = \lambda \exp\left(-\frac{x^2}{4}\right)u(x).$$

(ii) $H_n(x) = x^n + p(x), p \in \mathbb{P}_{n-1}$ (the space of polynomials of degree not exceeding n-1). Notice first that

$$xH'_n(x) = nx^n + p_1(x),$$
 $p_1 \in \mathbb{P}_{n-1},$
 $H''_n(x) \in \mathbb{P}_{n-2}.$

Therefore,

$$H_n''(x) - \frac{1}{2}xH_n'(x) + \lambda_n H_n(x) = -\frac{1}{2}nx^n + \lambda_n x^n + \left(H_n''(x) - \frac{1}{2}p_1(x) + \lambda_n p(x)\right),\tag{8}$$

where

$$H_n''(x) - \frac{1}{2}p_1(x) + \lambda_n p(x) \in \mathbb{P}_{n-1}.$$

In order for the expression (8) to vanish, one must have

$$-\frac{1}{2}n + \lambda_n = 0,$$

hence $\lambda_n = n/2$.

Setting the values n = 0, 1, 2, 3, we obtain

$$\lambda_{0} = 0, \qquad H_{0}(x) = 1,$$

$$\lambda_{1} = \frac{1}{2}, \qquad H_{1}(x) = x + a \implies H''_{1}(x) - \frac{1}{2}xH'_{1}(x) + \frac{1}{2}H_{1}(x) = -\frac{1}{2}x + \frac{1}{2}x + \frac{1}{2}a$$

$$\implies a = 0 \implies H_{1}(x) = x,$$

$$\lambda_{2} = 1, \qquad H_{2}(x) = x^{2} + ax + b \implies H''_{2}(x) - \frac{1}{2}xH'_{2}(x) + H_{2}(x)$$

$$= 2 - \frac{1}{2}(2x^{2} + ax) + x^{2} + ax + b = 2 - \frac{1}{2}ax + ax + b = 0$$

$$\implies a = 0, b = -2 \implies H_{1}(x) = x^{2} - 2,$$

$$\lambda_{3} = \frac{3}{2}, \qquad H_{3}(x) = x^{3} + ax^{2} + bx + c \implies H''_{3}(x) - \frac{1}{2}xH'_{3}(x) + \frac{3}{2}H_{3}(x)$$

$$= 6x + 2a - \frac{1}{2}(3x^{3} + 2ax^{2} + bx) + \frac{3}{2}(x^{3} + ax^{2} + bx + c)$$

$$= \left(-a + \frac{3}{2}a\right)x^{2} + \left(6 - \frac{b}{2} + \frac{3}{2}b\right)x + 2a + \frac{3}{2}c = 0$$

$$\implies a = c = 0, b = -6 \implies H_{1}(x) = x^{3} - 6x,$$

(iii)
$$\langle H_0, H_1 \rangle_r = \int_{-\infty}^{\infty} x \exp\left(-\frac{x^2}{4}\right) dx = 0, \tag{9}$$

as the integrand is an odd function, and the integration interval is symmetric with respect to zero.

$$\langle H_1, H_2 \rangle_r = \int_{-\infty}^{\infty} x(x^2 - 2) \exp\left(-\frac{x^2}{4}\right) dx = 0,$$

in the same way as (9).

Finally, using integration by parts, we obtain

$$\langle H_0, H_2 \rangle_r = \int_{-\infty}^{\infty} (x^2 - 2) \exp\left(-\frac{x^2}{4}\right) dx = I_1 + I_2.$$

Let us study both integrals separately. We have,

$$I_{1} = \int_{-\infty}^{\infty} x^{2} \exp\left(-\frac{x^{2}}{4}\right) dx$$

$$= -2 \int_{-\infty}^{\infty} x \frac{d}{dx} \left\{ \exp\left(-\frac{x^{2}}{4}\right) \right\} dx$$

$$= -2 \left[xe^{-x^{2}/4} \right]_{-\infty}^{\infty} - \int_{-\infty}^{\infty} e^{-x^{2}/4} dx \right] = 2 \int_{-\infty}^{\infty} e^{-x^{2}/4} dx = -I_{2},$$

since the first term is equal to zero. Therefore $I_1 + I_2 = -I_2 + I_2$.

(iv) $H_1(x)$ has one zero at $x=0, H_2(x)$ has zeros at $x=\pm\sqrt{2}, H_3(x)$ has zeros at $x=\pm\sqrt{6}$.

Suppose the system $\{\phi_n\}$ is complete. If there was $f \neq 0$ such that

$$\langle f, \phi_n \rangle_r = 0$$
 for all $n = 0, 1, 2, \dots$

then one would have

$$\min E_N = ||f||_r^2 - \sum_{n=0}^N \langle f, \phi_n \rangle_r^2 = ||f||_r^2 \quad \forall N,$$

and hence $\min E_N \to ||f||_r^2 > 0$ as $N \to \infty$, which contradicts the completeness of $\{\phi_n\}$.