## MA30044/MA40044/MA50181, Mathematical Methods I, 2020

## Problem Sheet 1: Introduction to Sturm-Liouville Equations

## SOLUTIONS

Q1.

(i) 
$$-u''(x) - 6u'(x) = \lambda u(x), \quad x \in (-\infty, \infty),$$

Here, in terms of the notation used in the lecture,

$$\alpha(x) = 1$$
,  $\beta(x) = 6$ ,  $\gamma(x) = 0$ ,  $\delta(x) = 1$ .

Hence, using the formulae obtained therein,

$$p(x) = e^{6x}, \quad q(x) = 0, \quad r(x) = e^{6x},$$

and therefore the SL form of the original equation is

$$-\left(e^{6x}u'(x)\right)' = \lambda e^{6x}u(x). \tag{1}$$

(ii) 
$$-x(1-x)u'' + 2xu'(x) = \lambda xu(x), \quad x \in [0,1].$$

Here, in terms of the notation used in the lecture,

$$\alpha(x) = x(1-x), \quad \beta(x) = -2x, \quad \gamma(x) = 0, \quad \delta(x) = x.$$

Hence, using the formulae obtained therein.

$$p(x) = \exp\left\{ \int_0^x \left( -\frac{2y}{y(1-y)} \right) dy \right\} = \exp\left\{ -2 \int_0^x \frac{dy}{1-y} \right\}.$$

Making the change of variables z = 1 - y, so dz = -dy, we proceed to conclude

$$p(x) = \exp\left(-2\int_{1-x}^{1} \frac{dz}{z}\right) = \exp\left(2\log(1-x)\right) = (1-x)^{2}.$$

Furthermore,

$$q(x) = 0$$
,  $r(x) = \frac{x(1-x)^2}{x(1-x)} = 1-x$ ,

and therefore the SL form of the original equation is

$$-((1-x)^2 u'(x))' = \lambda (1-x)u(x).$$
 (2)

(iii) 
$$-x^2u''(x) - xu'(x) = \lambda \frac{u(x)}{x}, \quad x \in (0, \infty).$$

Here, in terms of the notation used in the lecture,

$$\alpha(x) = x^2$$
,  $\beta(x) = x$ ,  $\gamma(x) = 0$ ,  $\delta(x) = \frac{1}{x}$ .

Hence, using the formulae obtained therein,

$$p(x) = \left(\int_{1}^{x} \frac{dy}{y}\right) = \exp(\log(x)) = x, \quad q(x) = 0, \quad r(x) = x\frac{1/x}{x^2} = \frac{1}{x^2},$$

and therefore the SL form of the original equation is

$$-(xu'(x))' = \lambda \frac{u(x)}{x^2}.$$
 (3)

**Q2**.

(i) 
$$-u''(x) = \lambda u(x), \quad x \in [0, \pi],$$
 (4)

$$u(0) = u(\pi) = 0, (5)$$

All solutions to (4) vanishing at zero (first condition in (5)) have the form  $u(x) = \sin(\sqrt{\lambda}x)$ . Using the second condition in (5), we obtain  $u(\pi) = 0$ , and hence  $\sqrt{\lambda}\pi = n\pi$ , n = 1, 2, ... Hence,

$$u(x) = \sin(\sqrt{\lambda}x), \qquad \lambda = n^2, \quad n \in \mathbb{N}.$$
 (6)

(ii) 
$$-u''(x) = \lambda u(x), \quad x \in [a, b], \tag{7}$$

$$u'(a) = u'(b) = 0,$$
 (8)

Note that we would typically write the general solution as

$$u(x) = A\cos(\sqrt{\lambda}x) + B\sin(\sqrt{\lambda}x),$$

but because of the boundary conditions at x = a, b, it is instead better to either write

$$u(x) = A\cos(\sqrt{\lambda}(x+B))$$
 or  $u(x) = A\sin(\sqrt{\lambda}(x+B))$ .

We do this because then B is easily selected. In this case, the derivative condition of u'(a) = 0 suggests that we choose a cosine expression and B = -a thus

$$u(x) = \cos(\sqrt{\lambda}(x-a)),$$

(setting the constant pre-factor to 1). Note that you did not necessarily have to do this but it makes the algebra a bit easier (since now the eigenfunction is a single cosine).

Applying the boundary condition then requires  $u'(x) = -\sqrt{\lambda} \sin(\sqrt{\lambda}(x-a))$ . Now the second condition in (8) yields  $\sqrt{\lambda} \sin(\sqrt{\lambda}(b-a)) = 0$ , and hence  $\sqrt{\lambda}(b-a) = n\pi$ , n = 1, 2, ..., or  $\lambda = 0$ . Summarising, the sought eigenfunction-eigenvalue pairs are given by

$$u(x) = \cos\left(n\pi \frac{x-a}{b-a}\right), \qquad \lambda = \frac{n^2\pi^2}{(b-a)^2}, \quad n \in \{0, 1, 2\dots\}$$
(9)

(iii) 
$$-u''(x) = \lambda u(x), \quad x \in \left[0, \frac{\pi}{2}\right], \tag{10}$$

$$u(0) = u'\left(\frac{\pi}{2}\right) = 0,\tag{11}$$

Using (10) and the first boundary condition in (11), we obtain  $u(x) = \sin(\sqrt{\lambda}x)$ , and hence  $u'(x) = -\sqrt{\lambda}\cos(\sqrt{\lambda}x)$ . Now from the second condition in (8) we have

$$u'\left(\frac{\pi}{2}\right)\sqrt{\lambda}\sin\left(\sqrt{\lambda}\frac{\pi}{2}\right) = 0,$$

and hence  $\sqrt{\lambda}(b-a)=n\pi, n=1,2,\ldots,$  or  $\lambda=0.$  Therefore,

$$u(x) = \sin((2n+1)x), \qquad \lambda = (2n+1)^2, \quad n \in \{0, 1, 2...\}$$
 (12)

(iv) 
$$-\left(e^{-x}u'(x)\right)' = \lambda e^{-x}u(x), \quad x \in \left[0, \frac{\pi}{2}\right], \tag{13}$$

$$u(0) = u(1) = 0, (14)$$

Differentiating and dividing by  $\exp(-x)$ , the equation (13) is written in the form

$$-u'' + u' = \lambda u. \tag{15}$$

This is an equation with constant coefficients, which generically have solutions of the form  $\exp(\alpha x)$ . Substituting this as an ansatz into (15) gives  $\alpha^2 - \alpha + \lambda = 0$ , from which

$$\alpha_{\pm} = \frac{1}{2} \pm \sqrt{\frac{1}{4} - \lambda}.$$

[Note: To make sense of the last expression, strictly speaking we need to specify the branch of the square root. For example, we can stipulate that  $\arg(\sqrt{z}) \in (-\pi/2, \pi/2]$ . However, the final answer will not depend on this choice, of course.]

The general solution to (13) is given by

$$u(x) = Ae^{\alpha + x} + Be^{\alpha - x} = e^{x/2} \left( A \exp\left(\sqrt{\frac{1}{4} - \lambda x}\right) + B \exp\left(\sqrt{\frac{1}{4} - \lambda x}\right) \right)$$

Using the first condition in (14) implies A + B = 0, so

$$u(x) = Ae^{x/2} \left( \exp\left(\sqrt{\frac{1}{4} - \lambda x}\right) - \exp\left(\sqrt{\frac{1}{4} - \lambda x}\right) \right)$$
 (16)

The second condition in (14) yields

$$\exp\left(\sqrt{\frac{1}{4}-\lambda}\right) - \exp\left(\sqrt{\frac{1}{4}-\lambda}\right) = 0,$$

or equivalently

$$\sin\left(\sqrt{\lambda - \frac{1}{4}}\right) = 0,$$

which implies

$$\lambda - \frac{1}{4} = n^2 \pi^2, \qquad n = 1, 2 \dots$$
 (17)

For the values  $\lambda$  given by (17), the expression (16) is of the form

$$u(x) = 2i\widetilde{A}e^{x/2}\sin(n\pi x).$$

[The relation between  $\widetilde{A}$  and A (in fact  $\widetilde{A}=A$  or  $\widetilde{A}=-A$ ) depends on which one of the two possibilities

$$\sqrt{\frac{1}{4} - \lambda} = i\sqrt{\lambda - \frac{1}{4}}, \qquad \sqrt{\frac{1}{4} - \lambda} = -i\sqrt{\lambda - \frac{1}{4}}$$

is realised under our choice of the branch of the square root.

Summarising, the eigenfunction-eigenvalue pairs are given by

$$u(x) = e^{x/2} \sin(n\pi x), \qquad \lambda = \frac{1}{4} + n^2 \pi^2, \quad n \in \{1, 2...\}$$
(18)

**Q3**.

$$-(pu'(x))' + qu(x) = \lambda ru(x), \quad x \in [a, b].$$

(i) Multiplying both sides by  $\overline{u(x)}$ , we write

$$-(pu')'\overline{u} + q|u|^2 = \lambda r|u|^2.$$
(19)

Let  $(\lambda, \phi)$  be an eigenvalue-eigenfunction pair. Set  $u = \phi$  in (19) and integrate over [a, b]:

$$-\int_{a}^{b} (p\phi')'\overline{\phi} + \int_{a}^{b} q|\phi|^{2} = \lambda \int_{a}^{b} r|\phi|^{2}.$$

Integrating by parts, we obtain

$$-\left[p\phi'\overline{\phi}\right]_a^b + \int_a^b p|\phi'|^2 + \int_a^b q|\phi|^2 = \lambda \int_a^b r|\phi|^2,$$

and therefore

$$\lambda \int_{a}^{b} r|\phi|^{2} = \int_{a}^{b} p|\phi'|^{2} + \int_{a}^{b} q|\phi|^{2} - \left[p\phi'\overline{\phi}\right]_{a}^{b}.$$
 (20)

(ii) In the case (a), we have  $\phi'(a) = \phi'(b) = 0$ , and therefore  $\left[p\phi'\overline{\phi}\right]_a^b = 0$ . In the case (b), we have  $\phi(a) = \phi(b) = 0$ , and therefore  $\left[p\phi'\overline{\phi}\right]_a^b = 0$  again.

Combining this with (20) yields

$$\lambda \int_{a}^{b} r|\phi|^{2} = \int_{a}^{b} p|\phi'|^{2} + \int_{a}^{b} q|\phi|^{2}.$$

Using the fact that

$$\int_{a}^{b} r|\phi|^{2} > 0, \qquad \int_{a}^{b} (p|\phi'|^{2} + q|\phi|^{2}) \ge 0,$$

we obtain  $\lambda \geq 0$ .

**Q4**.

$$-x^{2}u''(x) - xu'(x) = \lambda u(x), \qquad x > 0.$$
(21)

(i) Substituting the ansatz  $u(x) = x^{\alpha}$  into (21), we obtain

$$-x^{2}\alpha(\alpha-1)x^{\alpha-2} - x\alpha x^{\alpha-1} = \lambda x^{\alpha} \qquad \forall x > 0,$$

and hence

$$\alpha(\alpha - 1) + \alpha + \lambda = 0,$$

or equivalently

$$\alpha^2 + \lambda = 0$$
.

(ii) Using the notation from lectures, we note that for the equation (21) we have

$$\alpha(x) = x^2, \qquad \beta(x) = x, \qquad \gamma(x) = 0, \qquad \delta(x) = 1$$

Hence, by the theorem about writing equations in the SL form, we have

$$p(x) = \exp\left(\int_{1}^{x} \frac{dy}{y}\right) = \exp(\log(x)) = x, \qquad q(x) = 0, \qquad r(x) = \frac{x}{x^{2}} = \frac{1}{x}.$$

Therefore, the SL form of (21) is

$$-(xu'(x))' = \frac{\lambda}{x}u(x), \qquad x > 0.$$

(iii) Using Q3(ii), case (b), it follows from u(1) = u(e) = 0 that  $\lambda \ge 0$ . Suppose first that  $\lambda > 0$ . By the result of part (i), the general solution is

$$u(x) = c_{+}x^{i\sqrt{\lambda}} + c_{-}x^{-i\sqrt{\lambda}} \equiv c_{+}e^{i\sqrt{\lambda}\log(x)} + c_{-}e^{-i\sqrt{\lambda}\log(x)}.$$
 (22)

Setting first x = 1, then x = e, we obtain

$$\begin{cases} c_{+} + c_{-} = 0, \\ c_{+} e^{i\sqrt{\lambda}} + c_{-} e^{-i\sqrt{\lambda}} = 0. \end{cases}$$
 (23)

This system has a non-trivial solution pair  $(c_+, c_-)$  if and only if its determinant vanishes, which gives the condition  $\sin(\sqrt{\lambda}) = 0$ , and hence  $\lambda = n^2 \pi^2$ , n = 1, 2, ... If  $\lambda = 0$ , then the general solution is given by  $u(x) = c_+ + c_- \log(x)$ , which does not satisfy the boundary conditions for any  $c_+, c_-$ . So, the final answer is

$$\lambda = n^2 \pi^2, \qquad n = 1, 2, \dots$$
 (24)

(iv) The eigenfunctions corresponding to the eigenvalues (24) are found by combining the formula (22) for the general solution with the fact that for solutions of (23) we have  $c_{-} = -c_{+}$ :

$$u(x) = c_{+}e^{i\pi n \log(x)} - c_{+}e^{-i\pi n \log(x)} = c_{+}\sin(n\pi \log(x)).$$

So the (non-normalised) eigenfunctions are given by

$$\phi_n(x) = \sin(n\pi \log(x)), \qquad n = 1, 2, \dots$$

(v) 
$$\int_{1}^{e} \phi_{n}(x)\phi_{m}(x)\frac{dx}{x} = \int_{1}^{e} \sin(n\pi \log(x))\sin(m\pi \log(x))\frac{dx}{x}.$$

Changing the variable of integration:  $z = \log(x)$ , so dz = dx/x, we obtain

$$\int_0^1 \sin(n\pi z)\sin(m\pi z)dz = 0, \qquad n \neq m,$$

as was shown in the lectures.

**Q5**.

$$-Eu^{(4)} = \lambda u, \quad -L < x < L,$$
 (25)

$$u(-L) = u(L) = u''(-L) = u''(L) = 0.$$
(26)

(i) Setting  $u(x) = \exp(i\omega x)$  in (25), we obtain  $E\omega^4 + \lambda = 0$ , so the four possible values for  $\omega$  are

$$\omega_j = \left| \frac{\lambda}{E} \right|^{1/4} \alpha_j, \qquad j = 1, 2, 3, 4,$$

where  $\alpha_j$ , j = 1, 2, 3, 4, are the four roots of -1 of degree 4:

$$\alpha_{1,2} = \exp\left(\pm i\frac{\pi}{4}\right), \qquad \alpha_{3,4} = \exp\left(\pm i\frac{3\pi}{4}\right),$$

In other words,

$$\omega_{1,2} = \frac{R}{\sqrt{2}}(1 \pm i), \quad \omega_{3,4} = \frac{R}{\sqrt{2}}(-1 \pm i), \qquad R := \left|\frac{\lambda}{E}\right|^{1/4}.$$

(ii) The corresponding four exponential solutions of (25) are

$$\widetilde{\phi}_{1,2} = \exp\left(\frac{R}{\sqrt{2}}(1\pm i)x\right), \qquad \widetilde{\phi}_{3,4} = \exp\left(\frac{R}{\sqrt{2}}(-1\pm i)x\right).$$

These can be combined to obtain four linearly independent real-valued solutions:

$$\phi_1(x) = \frac{1}{2} (\widetilde{\phi}_1(x) + \widetilde{\phi}_2(x)) = \exp\left(\frac{R}{\sqrt{2}}x\right) \cos\left(\frac{R}{\sqrt{2}}x\right),$$

$$\phi_2(x) = \frac{1}{2} (\widetilde{\phi}_1(x) - \widetilde{\phi}_2(x)) = \exp\left(\frac{R}{\sqrt{2}}x\right) \sin\left(\frac{R}{\sqrt{2}}x\right),$$

$$\phi_3(x) = \frac{1}{2} (\widetilde{\phi}_3(x) + \widetilde{\phi}_4(x)) = \exp\left(-\frac{R}{\sqrt{2}}x\right) \cos\left(\frac{R}{\sqrt{2}}x\right),$$

$$\phi_4(x) = \frac{1}{2} (\widetilde{\phi}_3(x) - \widetilde{\phi}_4(x)) = \exp\left(-\frac{R}{\sqrt{2}}x\right) \sin\left(\frac{R}{\sqrt{2}}x\right).$$

(iii) Differentiating twice the functions  $\phi_j$ , j=1,2,3,4, we obtain

$$u''(x) = -\frac{R^2}{2} \left( a\phi_1(x) + b\phi_2(x) + c\phi_3(x) + d\phi_4(x) \right)$$

The boundary conditions (26) lead to:

$$u(L) = 0:$$

$$\left(a \exp\left[\frac{RL}{\sqrt{2}}\right] + c \exp\left[-\frac{RL}{\sqrt{2}}\right]\right) \cos\left[\frac{RL}{\sqrt{2}}\right] + \left(b \exp\left[\frac{RL}{\sqrt{2}}\right] + d \exp\left[-\frac{RL}{\sqrt{2}}\right]\right) \sin\left[\frac{RL}{\sqrt{2}}\right] = 0, \quad (27)$$

$$u(-L) = 0:$$

$$\left(a \exp\left[-\frac{RL}{\sqrt{2}}\right] + c \exp\left[\frac{RL}{\sqrt{2}}\right]\right) \cos\left[\frac{RL}{\sqrt{2}}\right] - \left(b \exp\left[-\frac{RL}{\sqrt{2}}\right] + d \exp\left[\frac{RL}{\sqrt{2}}\right]\right) \sin\left[\frac{RL}{\sqrt{2}}\right] = 0, \quad (28)$$

$$u''(L) = 0:$$

$$\left(a \exp\left[\frac{RL}{\sqrt{2}}\right] - c \exp\left[-\frac{RL}{\sqrt{2}}\right]\right) \cos\left[\frac{RL}{\sqrt{2}}\right] + \left(b \exp\left[\frac{RL}{\sqrt{2}}\right] + d \exp\left[-\frac{RL}{\sqrt{2}}\right]\right) \sin\left[\frac{RL}{\sqrt{2}}\right] = 0, \quad (29)$$

$$u(-L) = 0:$$

$$\left(a\exp\left[-\frac{RL}{\sqrt{2}}\right] - c\exp\left[\frac{RL}{\sqrt{2}}\right]\right)\cos\left[\frac{RL}{\sqrt{2}}\right] - \left(b\exp\left[-\frac{RL}{\sqrt{2}}\right] - d\exp\left[\frac{RL}{\sqrt{2}}\right]\right)\sin\left[\frac{RL}{\sqrt{2}}\right] = 0.$$
(30)

Together (27)–(30) form a system of four linear equations for a,b,c,d. The dependency on  $\lambda$  is implicit, since  $r=\left|\lambda/E\right|^{1/4}$ .