# UNIVERSITY OF OSLO

# Faculty of Mathematics and Natural Sciences

Mid-term exam for AST3220 — Cosmology I

Day of exam: Friday 27th of March 2015

Time for exam: 12.15 - 15.15

This problem set consists of 4 pages.

Attachments: None

Allowed aids: All non-communicative aids.

Make sure that the problem set is complete before you start answering the questions.

Throughout the equation set we will work in units of c=1 We will also assume that the value of the scale factor today is unity:

$$a_0 = 1$$

#### Problem 1

#### Redshift dependent equations

Recall that the redshift z is defined by the equation

$$\frac{1}{a} = 1 + z \tag{1}$$

This definition provides a one-to-one correspondence between redshift and scale factor so long as a > 0 and z > -1. We can use this fact to redefine equations in terms of redshift.

a) Show that the adiabatic expansion equation in terms of redshift becomes:

$$\dot{\rho} = \frac{3\dot{z}}{(1+z)} \left(\rho + p\right) \tag{2}$$

- b) Use the version of the adiabatic expansion equation above to find an expression for  $\rho(z)$  for dust and cosmological constant, without first calculating  $\rho(a)$ .
- c) Use your results from b) and the definitions of the density parameters  $\Omega_{m0}$  and  $\Omega_{\Lambda0}$  to write the first Friedmann equation for flat Universe with dust and cosmological constant ( $\Lambda$ CDM) using only  $\Omega_{m0}$ , z,  $H_0$  and H.
- d) Use the definition of the comoving coordinate and make a variable change from t to a to z to prove that the comoving coordinate is given by:

$$r = \frac{c}{H_0} \int_0^z \frac{dz'}{E(z')} \tag{3}$$

where  $E(z) \equiv \frac{H(z)}{H_0}$ . Find E(z) for the Universe described in c).

#### Problem 2

## $R_h = t$ and power law cosmologies

The  $R_h = t$  and power law cosmologies, are alternative models to  $\Lambda$ CDM.  $R_h = t$  cosmology is defined as a universe with a linear increase in Hubble radius

$$R_h = t, (4)$$

where the Hubble radius  $R_h = \frac{1}{H}$ . Power law cosmology is defined by assuming a power law function of a with respect to t, i.e.

$$a = \left(\frac{t}{t_0}\right)^n \tag{5}$$

They are both proposed to describe the Universe at late time, so at low redshift.

- a) Find H(t) for  $R_h = t$  and power law cosmology.
- b)  $R_h = t$  is equivalent to a power law cosmology for a special value of n, which value?
- c) Find the function  $E(z) \equiv \frac{H(z)}{H_0}$  for power law cosmologies.
- d) Show that the comoving distance r(z) in power law cosmology is given by:

$$r(z) = \frac{1}{H_0} \times \begin{cases} \frac{(1+z)^{1-1/n}-1}{1-1/n}, & n \neq 1, \\ \ln(1+z), & n = 1. \end{cases}$$

#### Problem 3

### Supernova observations

We want to use supernovae observations to compare the  $\Lambda$ CDM model to the alternative models described above. When we observe supernovae, we measure their magnitude m on the sky, and their redshift. The observed magnitude is given by:

$$m = M + 5\log_{10}(H_0 d_L(z)) \tag{6}$$

where M is the absolute magnitude, which is assumed to be constant for supernovae type Ia, and  $d_L$  is the luminosity distance

- a) Explain why the differences in observed magnitude in universes governed by different models, can be found by specifying just r(z).
- b) Assume that z >> 1, and find approximate behaviour of r(z) as a function of z for the  $\Lambda$ CDM universe in this limit.
- c) Find r(z) in the large z limit for the power law cosmologies.
- d) The power law cosmologies are only used to describe cosmology in the late time of the Universe history. Do your answers in b) and c) help you understand why?