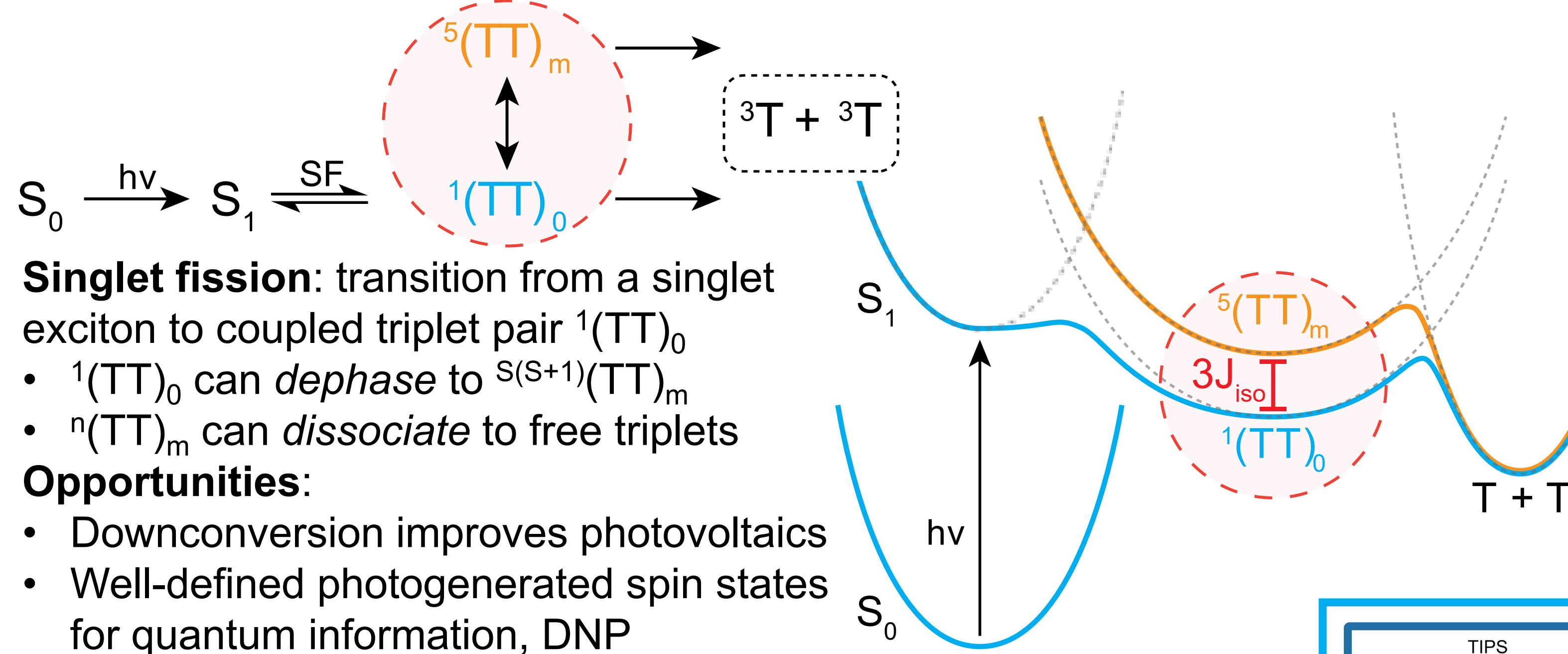


# Spin Dynamics in Singlet Fission

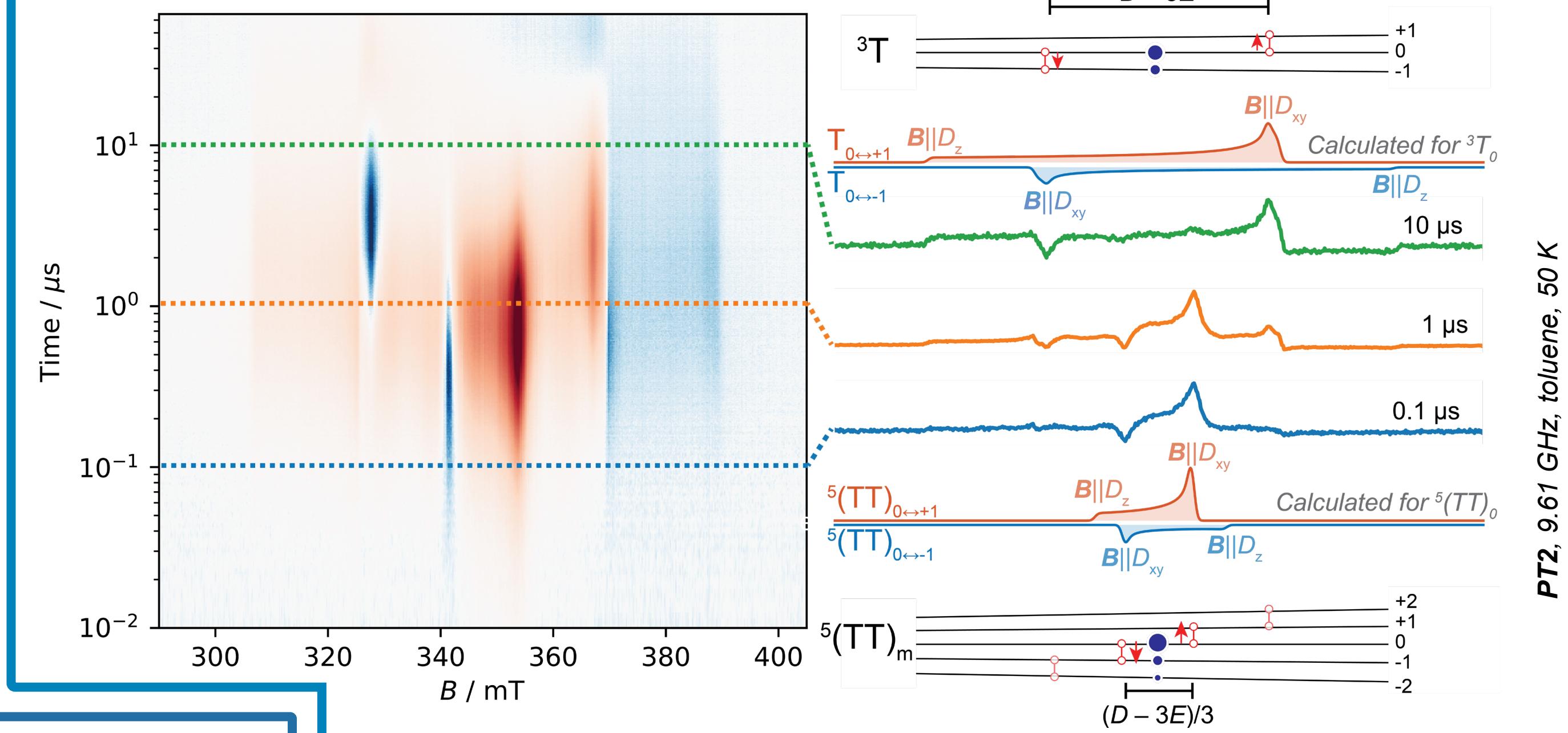
Thomas S. C. MacDonald<sup>a</sup>, Miles I. Collins<sup>a</sup>, Luis M. Campos<sup>b</sup>, Elango Kumarasamy<sup>b</sup>, Samuel N. Sander<sup>b</sup>, Matthew Y. Sfeir<sup>c</sup>, Murad Tayebjee<sup>a</sup>, Dane R. McCamey<sup>a</sup>

<sup>a</sup>ARC Centre of Excellence in Exciton Science, School of Physics, UNSW Sydney; <sup>b</sup>Department of Chemistry, Columbia University, New York, New York 10027, USA; <sup>c</sup>Center for Functional Nanomaterials, Brookhaven National Laboratory, Upton, New York 11973, USA

## Singlet fission in molecular dimers

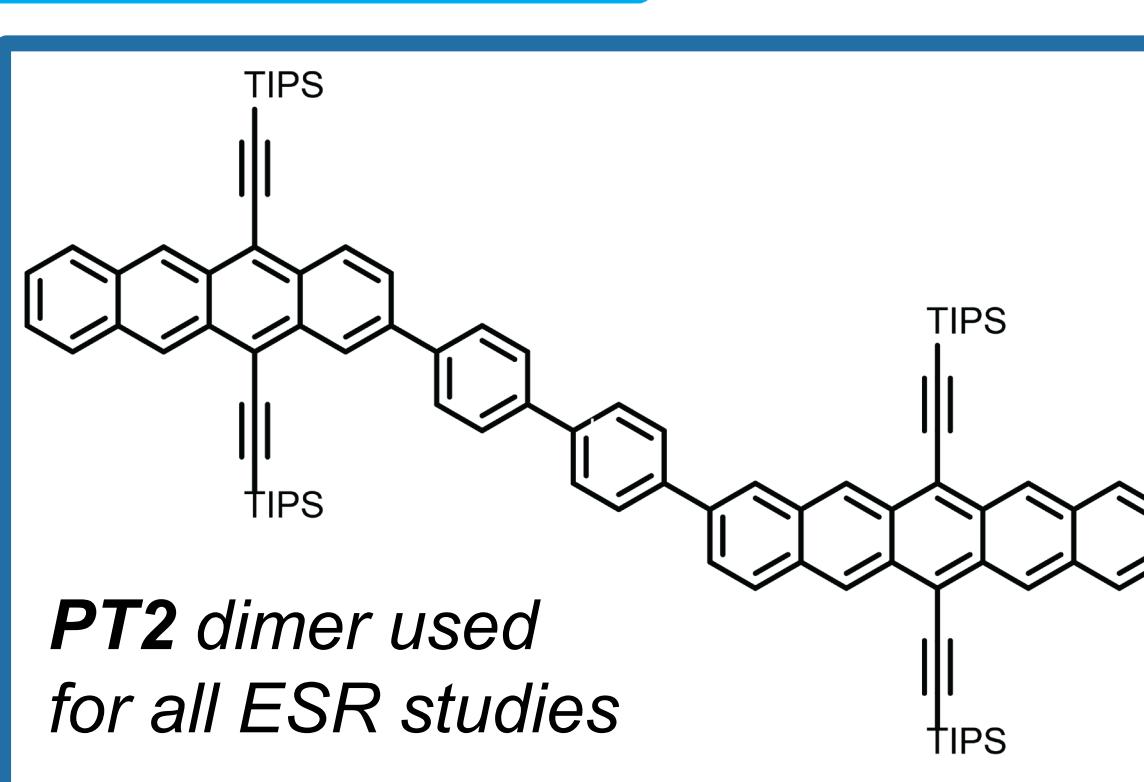


## Transient cw-ESR of singlet fission

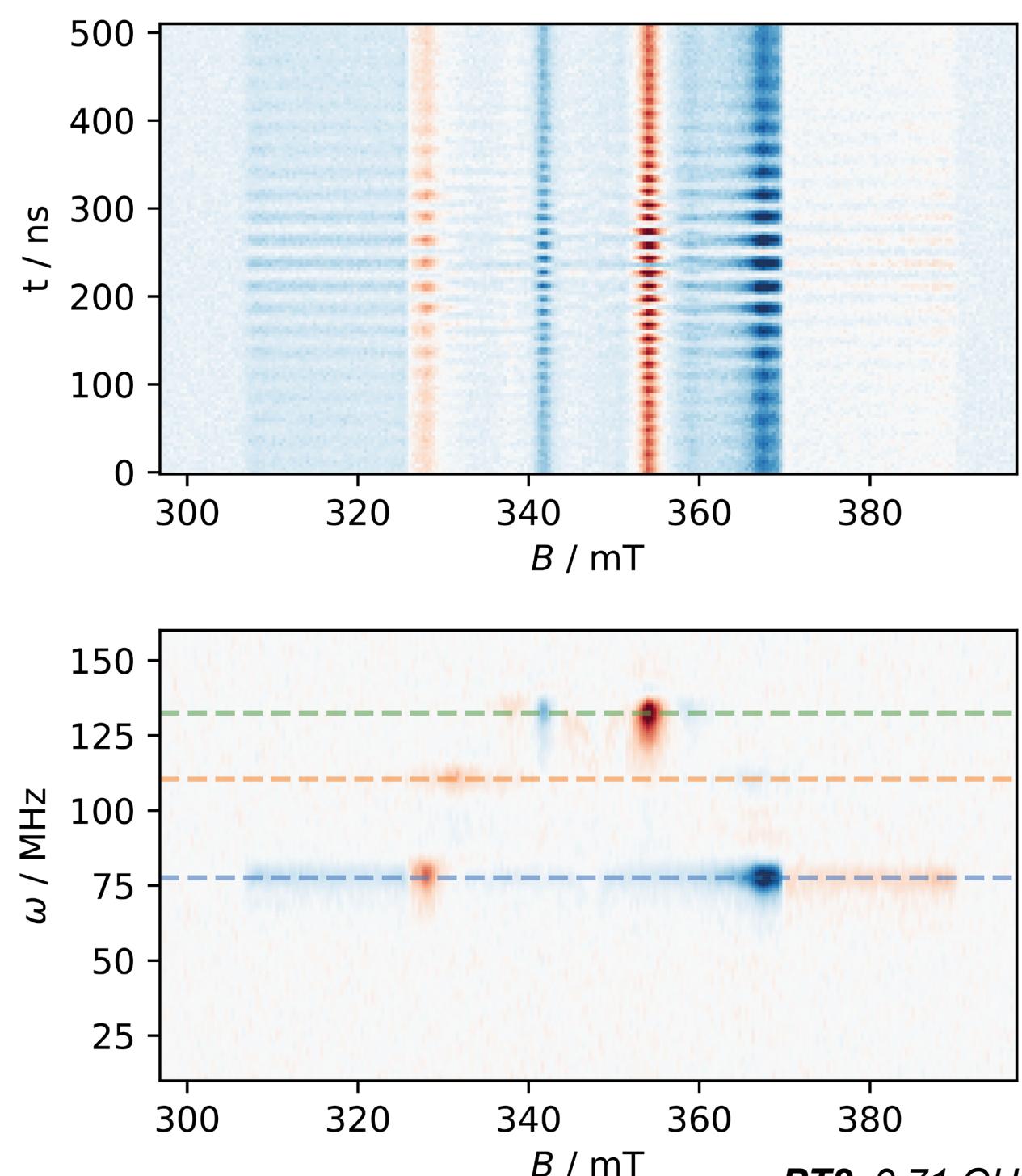


## Resolving transitions with nutation p-ESR

SF cw-ESR spectra are **complex and overlapping**. Nutation p-ESR lets us **identify** (1D) or **resolve** (2D) transitions by field or time.



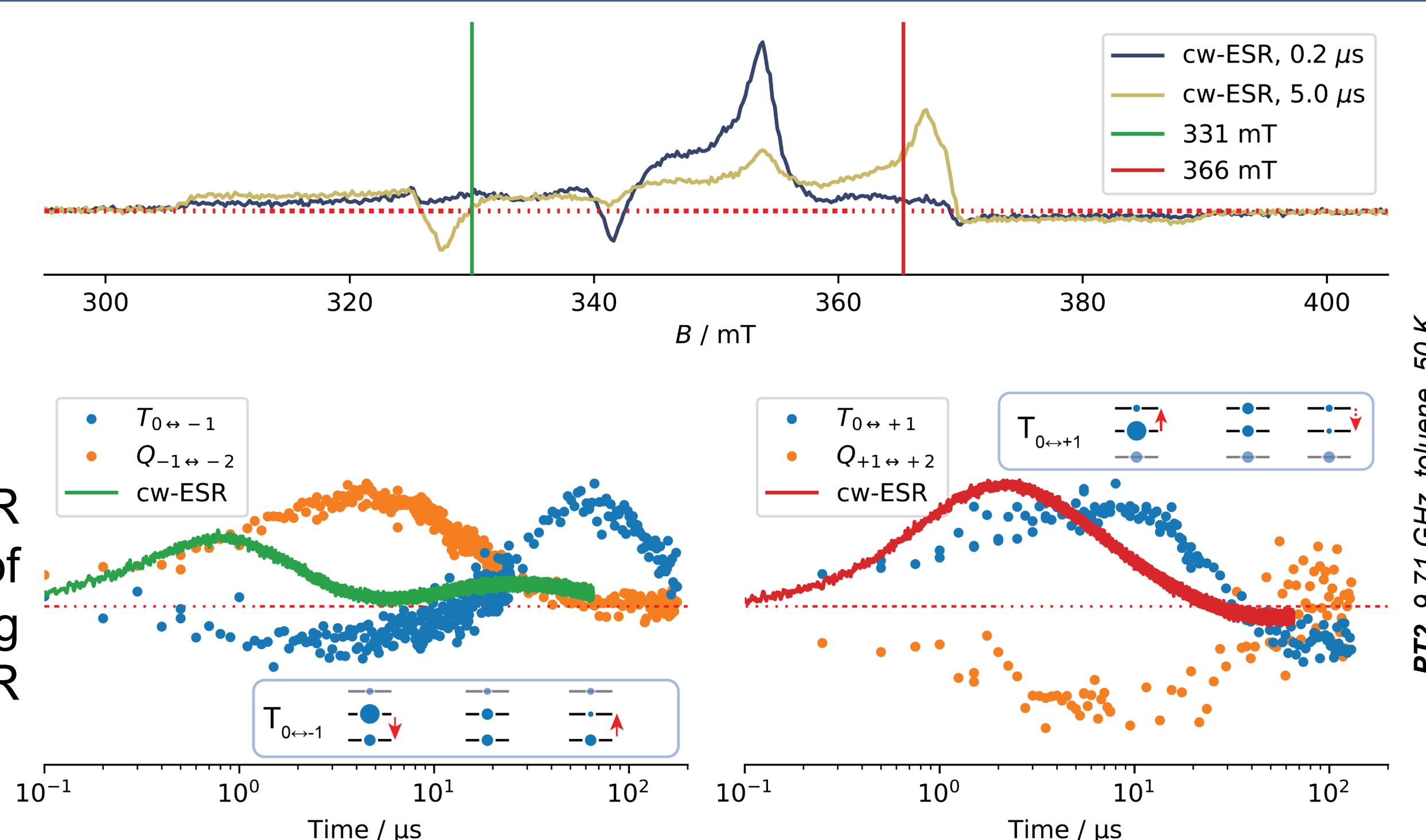
- Transient cw-ESR reveals evolving spin populations**
- Narrow  $^5(\text{TT})_{0\leftrightarrow\pm 1}$  transitions within <100 ns
  - Broader  $^3T_{0\leftrightarrow\pm 1}$  transitions 1 – 10s of  $\mu\text{s}$
  - Mainly  $m_s=0$ , net absorptive: more  $m_s<0$  than  $m_s>0$
  - Assigned by fitting, analogy: cw-ESR information poor
  - Overlapping signals complicate analysis



— cw-ESR at 11  $\mu\text{s}$   
—  $^3T_{0\leftrightarrow\pm 1}$  ( $\omega = 78 \text{ MHz}$ )  
—  $^5Q_{\pm 1\leftrightarrow\pm 2}$  ( $\omega = 111 \text{ MHz}$ )  
—  $^5Q_{0\leftrightarrow\pm 1}$  ( $\omega = 132 \text{ MHz}$ )

**Left:** 2D field-swept PEANUT p-ESR resolves signals by S and  $|m_s|$ , revealing surprise  $m_s = \pm 2$  quintets. P-ESR quintets seem less absorptive p-ESR triplets (or cw-ESR)

**Right:** 2D time-swept p-ESR nutation resolves dynamics of overlapping transitions, finding slower kinetics than cw-ESR

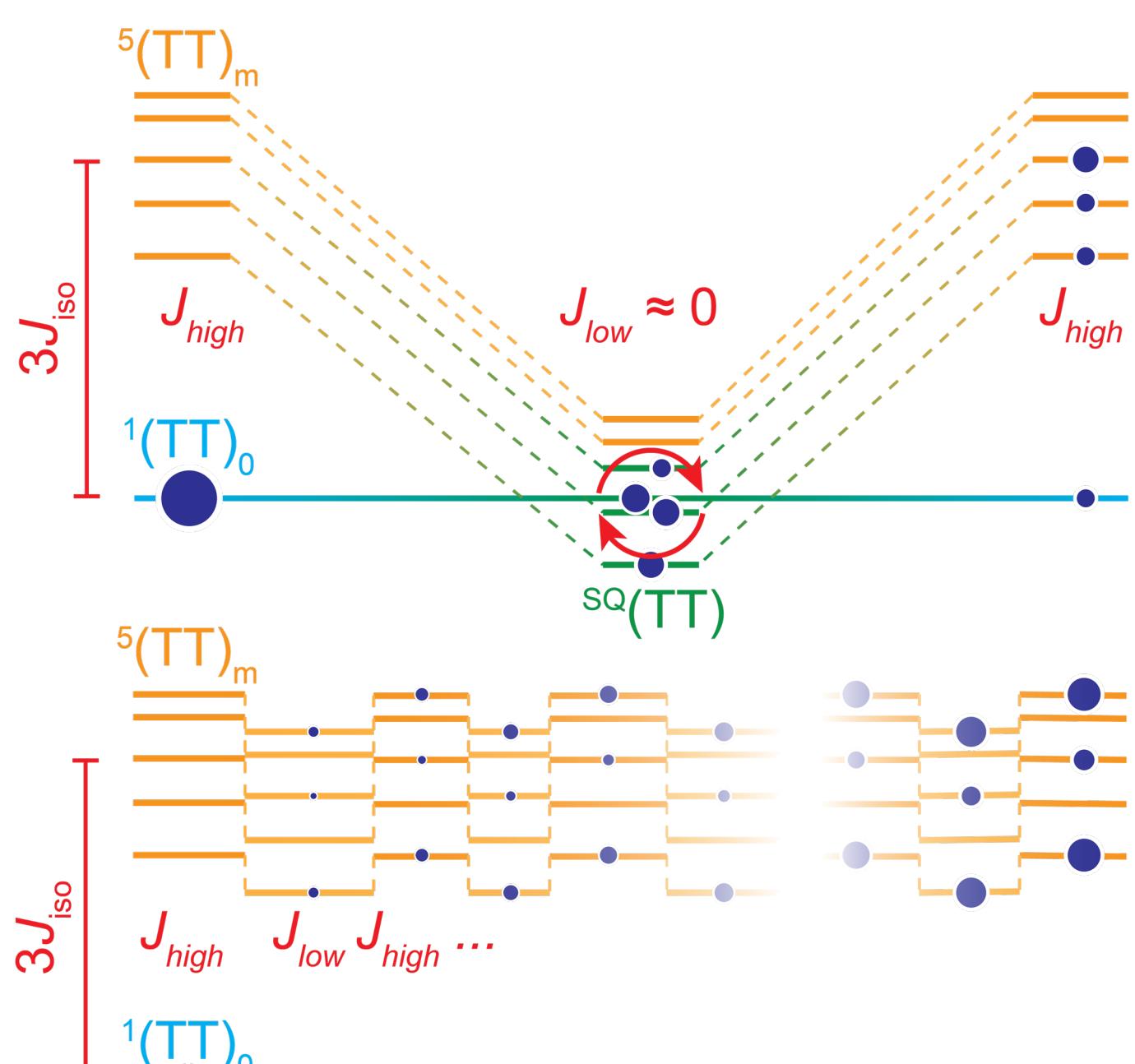


## Quintets form with dynamic $J_{iso}$

Narrow  $(D-3E)/3$  quintet spectra are only possible when inter-triplet  $J_{iso}$  is large, but large  $J_{iso}$  prevents  $^1(\text{TT})_0 \leftrightarrow ^5(\text{TT})_m$  mixing: suggests a **time-dependent  $J_{iso}$** . In a two-triplet basis:

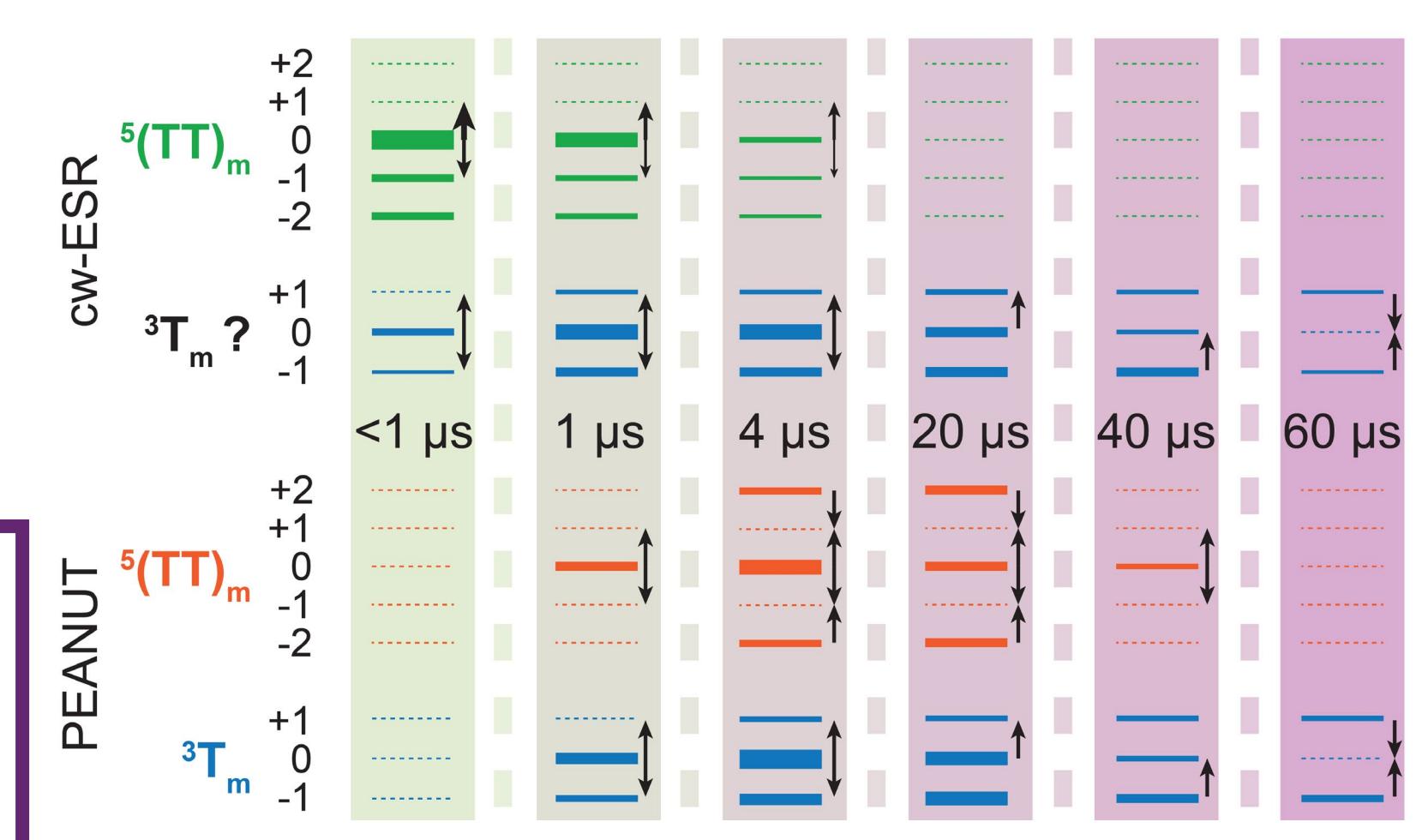
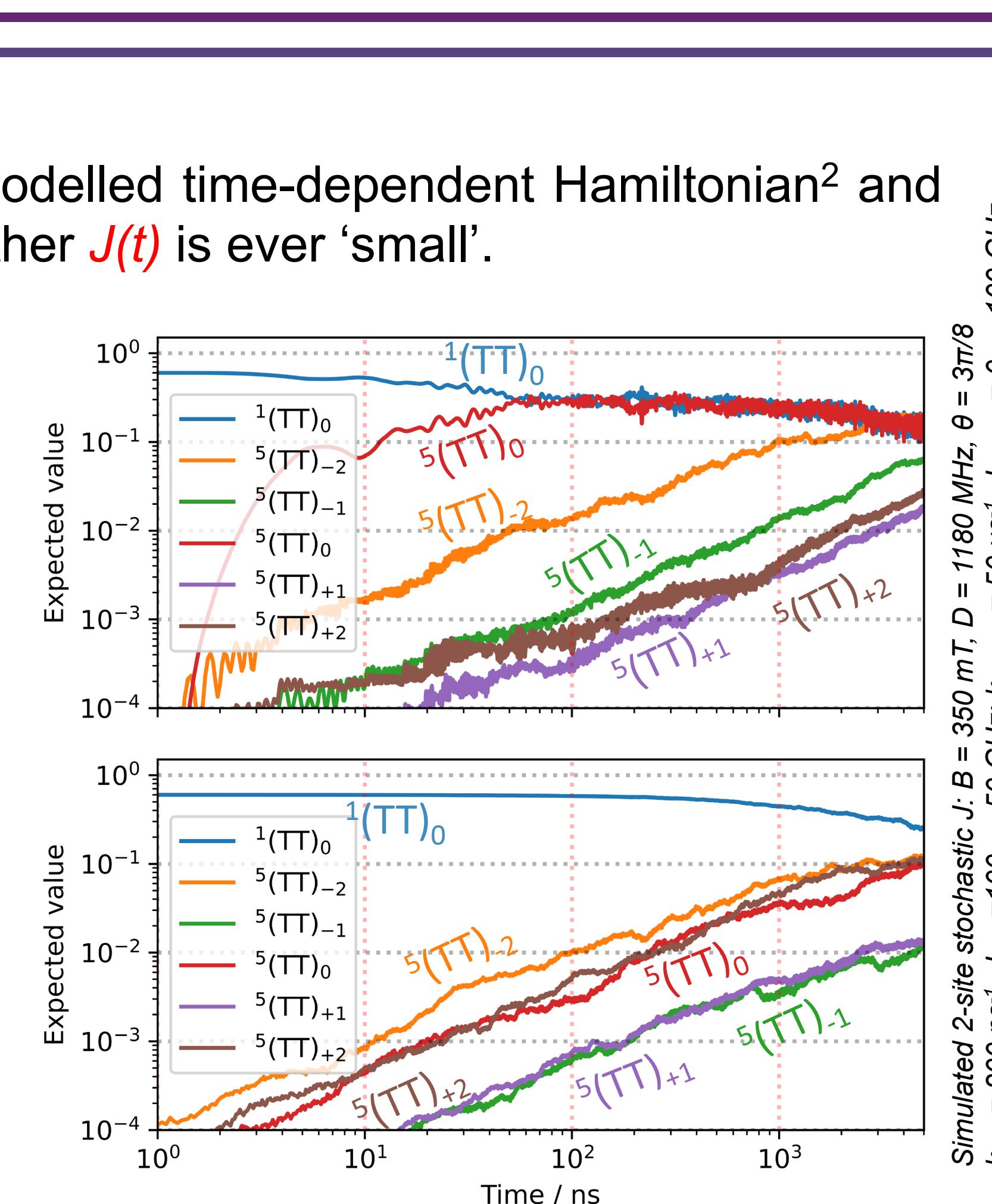
$$\hat{H}_{\text{spin}} = \hat{H}_{ee}(t) + \hat{H}_{zee} + \hat{H}_{zfs} \\ = J_{iso}(t) \hat{S}_1 \cdot \hat{S}_2 + \sum_{i=1,2} \mu_B g B_0 \cdot \hat{S}_i + \hat{S}_i \cdot D_i \cdot \hat{S}_i$$

We simulate  $^5(\text{TT})_m$  formation by solving the TD Schrödinger for a modelled time-dependent Hamiltonian<sup>2</sup> and find **two distinct modes** of  $^1(\text{TT})_0 \leftrightarrow ^5(\text{TT})_m$  mixing, depending on whether  $J(t)$  is ever ‘small’.



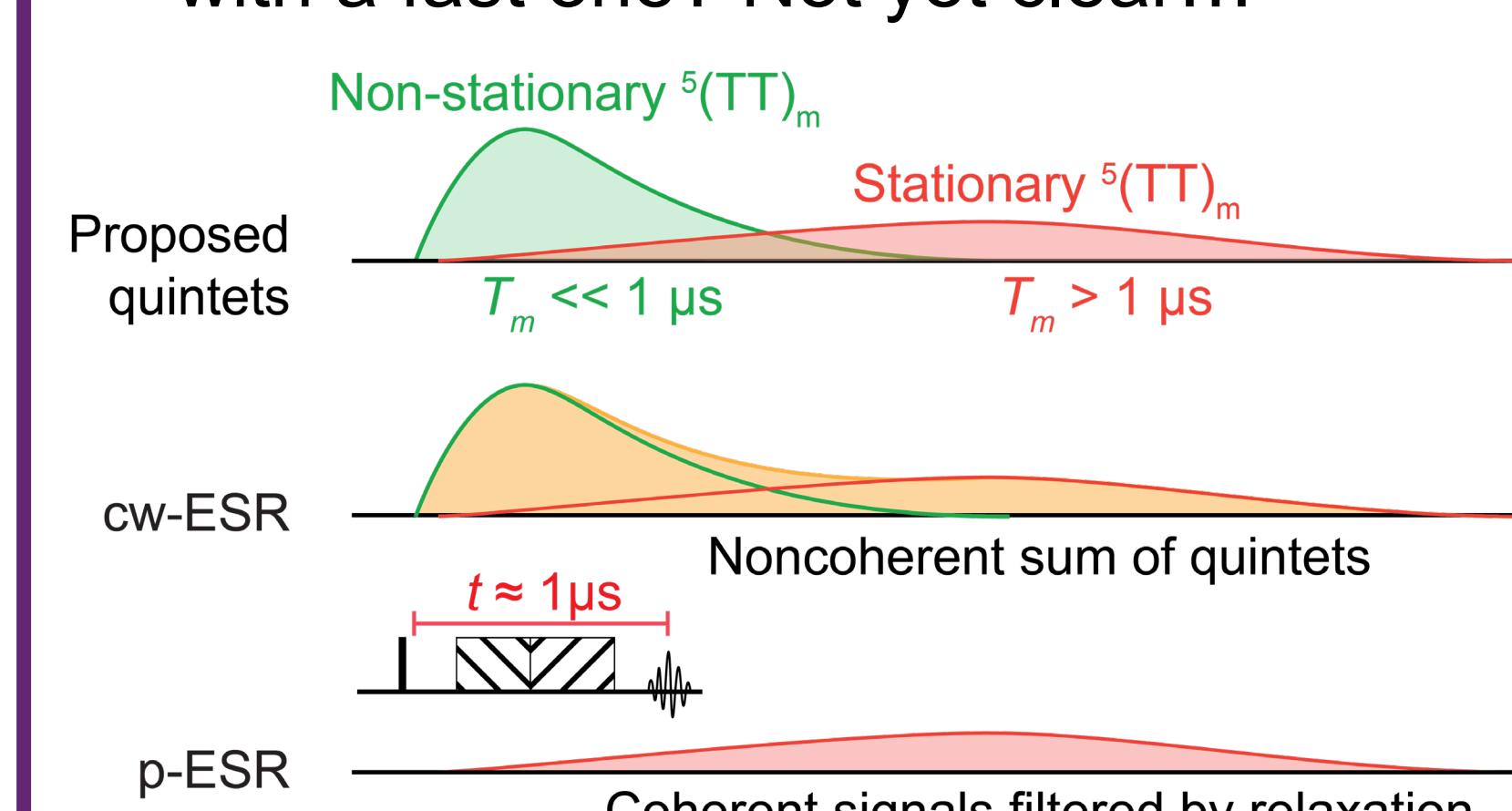
## Conclusions: how to reconcile cw-ESR & p-ESR?

- Cw-ESR signals are prompt, net-absorptive; consistent with ‘non-stationary’ formation
- P-ESR signals are delayed, comparable  $\pm m_s$  character; consistent with ‘stationary’ formation



## Populations, dynamics consistent with coexisting quintet pathways

- Non-stationary prompt quintets relax quickly → filtered out of p-ESR
- How does a slow process compete with a fast one? Not yet clear...



## References

- <sup>1</sup>
- Tayebjee, M. J. Y. et al.
- Nature Physics*
- 13**
- , 182–188 (2017).
- 
- <sup>2</sup>
- Collins, M. I., et al.
- J. Chem. Phys.*
- 151**
- , 164104 (2019).

## Acknowledgements

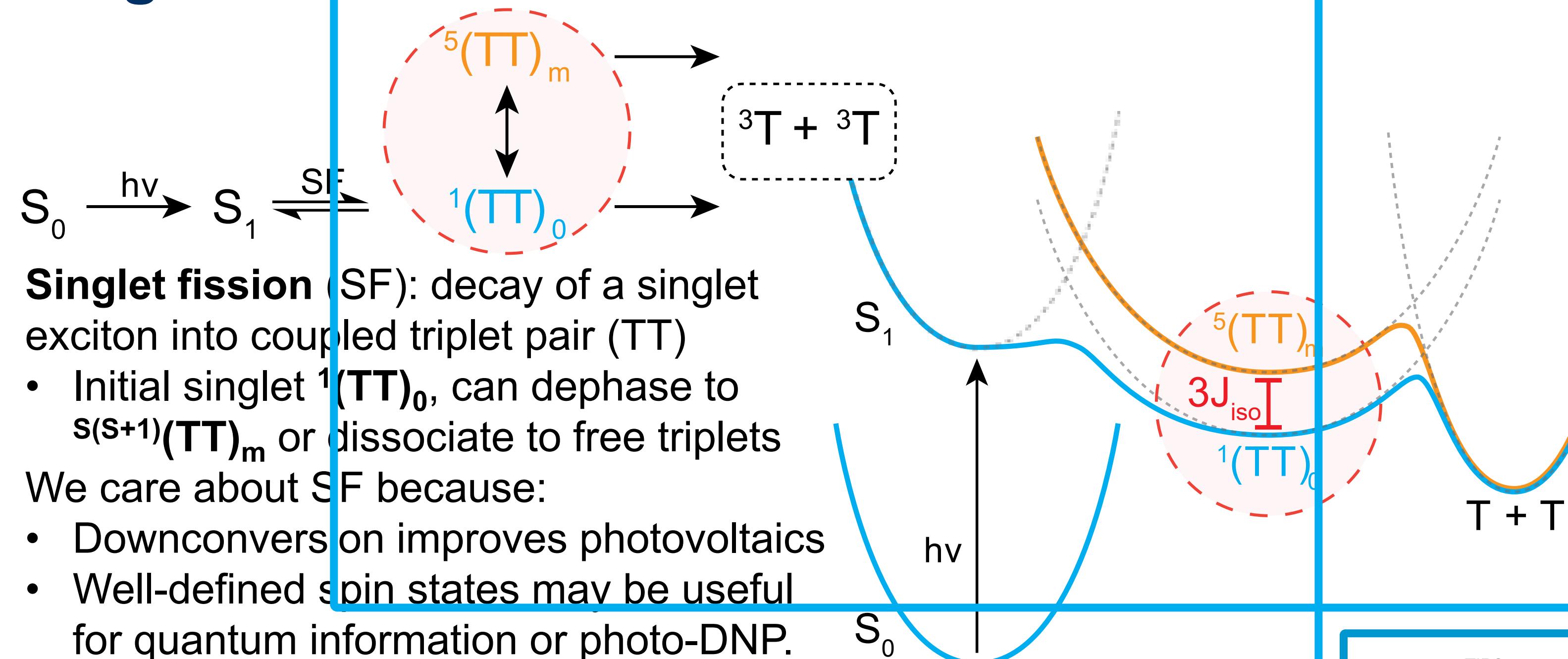
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# Spin Dynamics in Singlet Fission

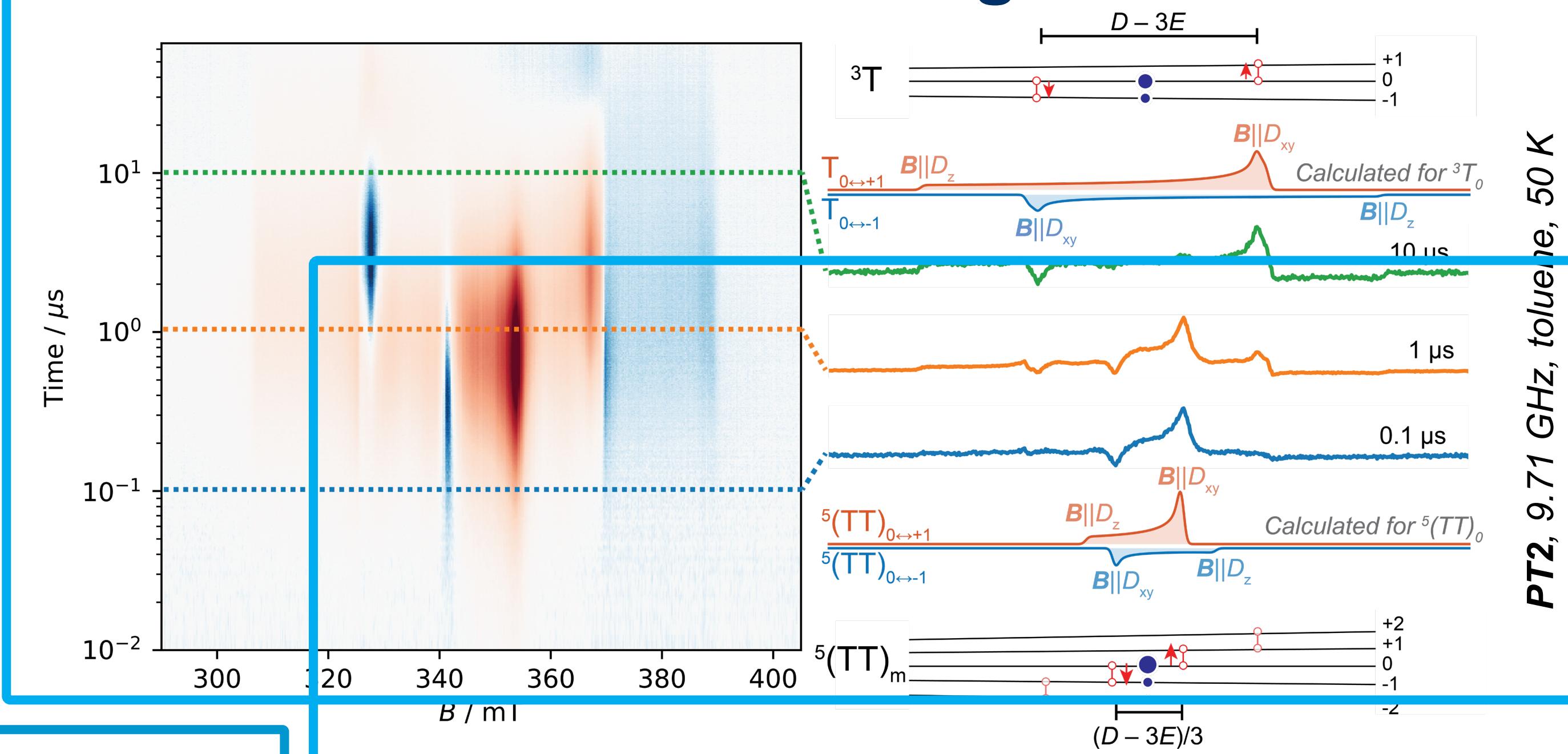
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## Singlet fission in molecular dimers

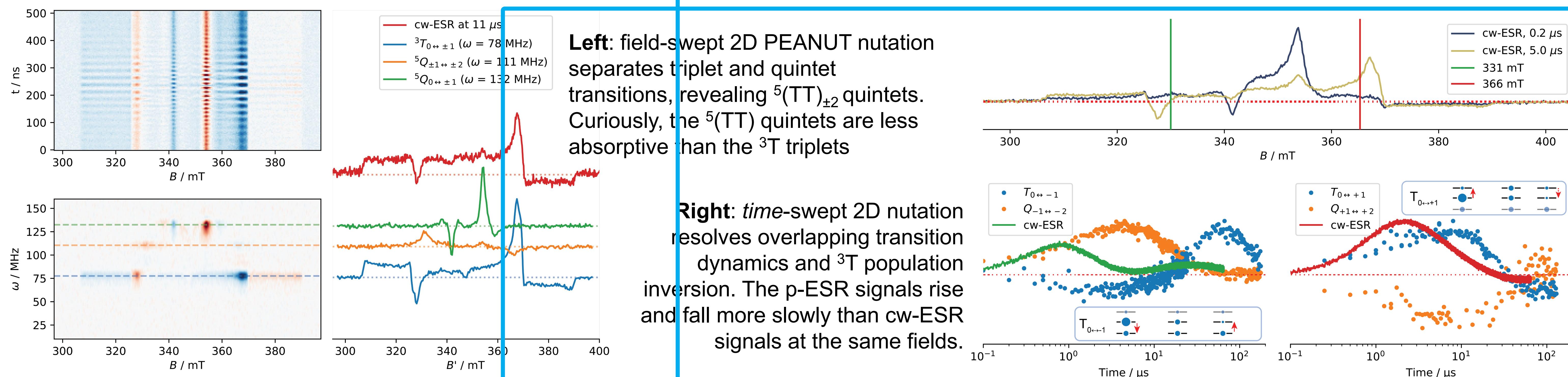
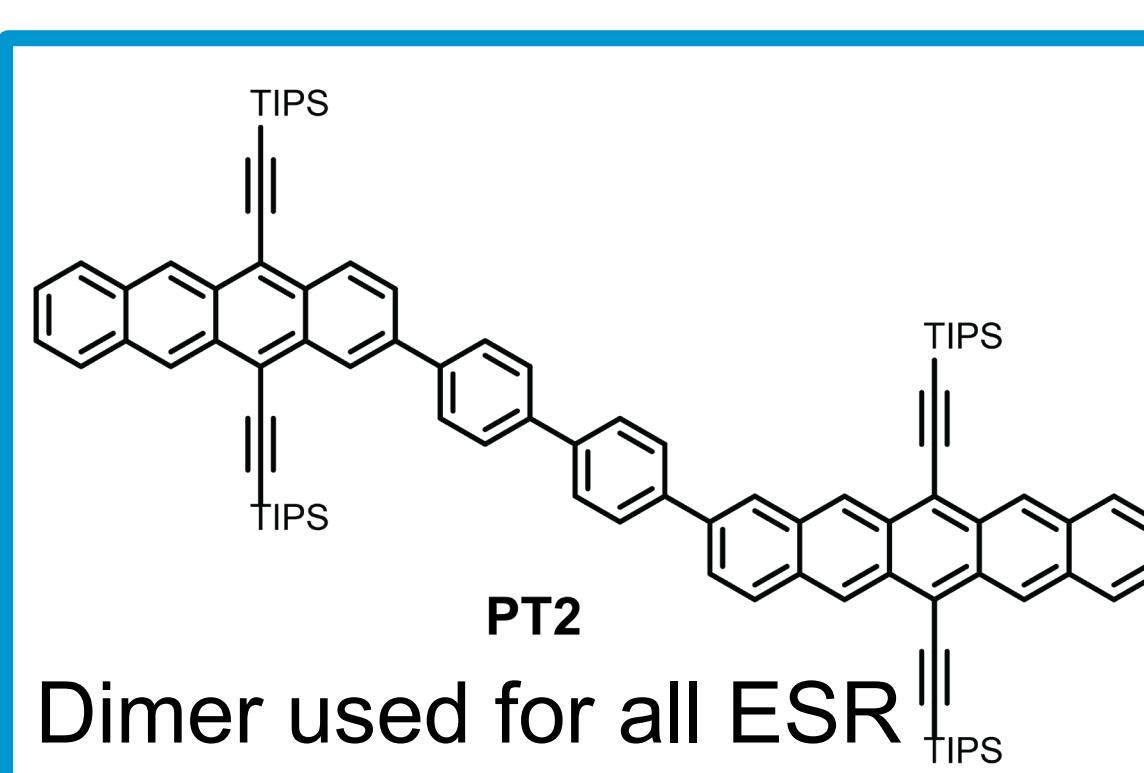


## Transient cw-ESR of singlet fission



## Resolving transitions with nutation p-ESR

SF forms a complex mixture of overlapping ESR signals, and transient cw-ESR can't identify the transitions involved. We use nutation p-ESR to resolve transitions over  $B$  field or time.

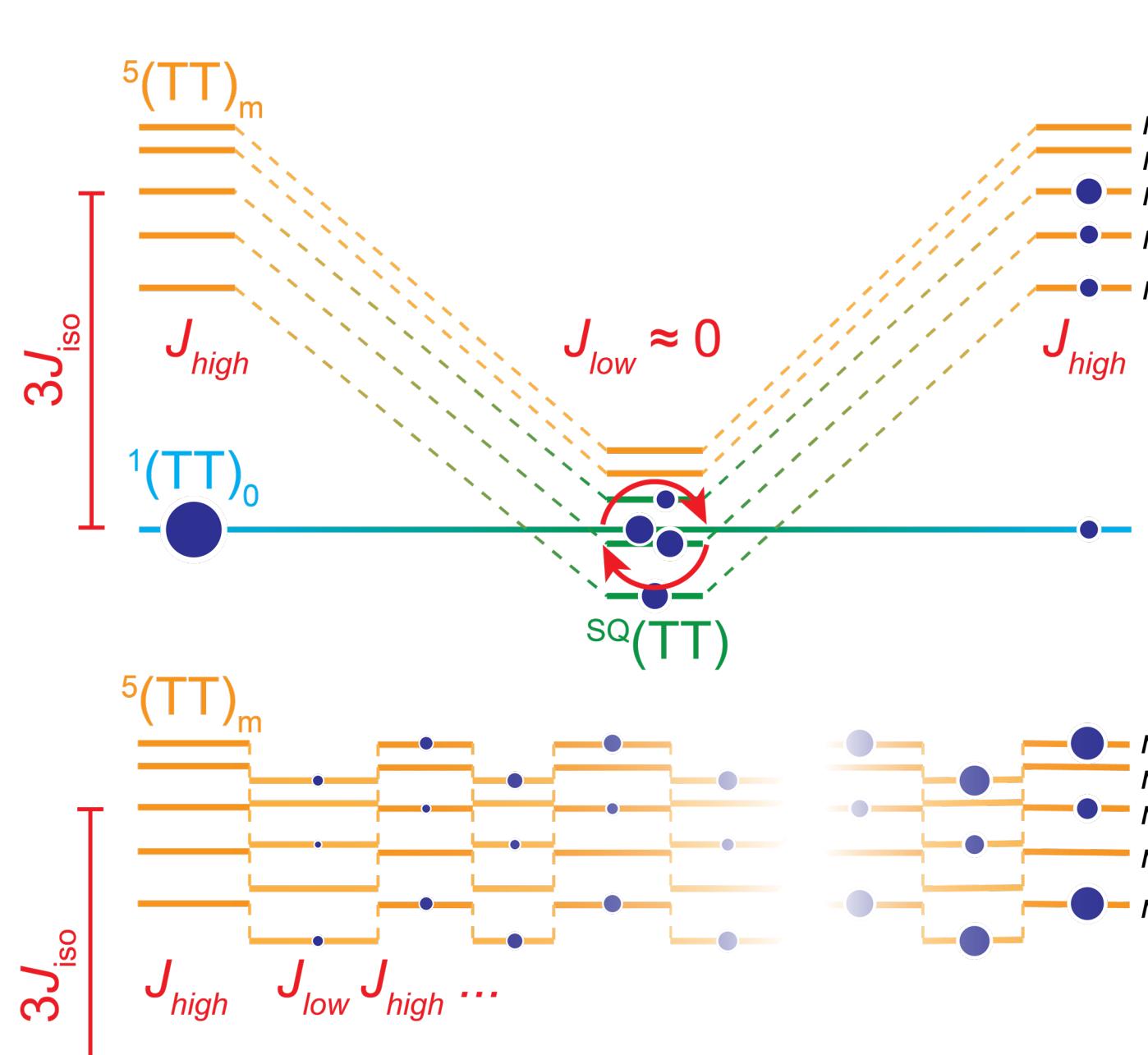


## Quintet formation requires dynamic $J_{iso}$

- $\hat{H}_{ZFS}$  only mixes  $1(TT)_0$  with  $5(TT)_m$  when spin-spin term  $\hat{H}_{ee}$  is small – but narrow quintet ESR spectra require large  $\hat{H}_{ee}$ !
- Solve with a time-dependent  $\hat{H}_{ee}$ . In the two-triplet basis:

$$\hat{H}_{spin} = \hat{H}_{ee} + \hat{H}_{zee} + \hat{H}_{zfs} = J_{iso}(t) \hat{S}_1 \cdot \hat{S}_2 + \sum_{i=1,2} \mu_B g \mathbf{B}_0 \cdot \hat{S}_i + \hat{S}_i \cdot \mathbf{D}_i \cdot \hat{S}_i$$

- Simulate quintet formation by solving the time-dependent hamiltonian<sup>2</sup> with a model  $J(t)$  term
  - Model? Site hopping, molecular dynamics, open quantum systems... Stay tuned!
- We see two distinct pathways to quintet formation, depending on whether  $J(t)$  is ever ‘small’

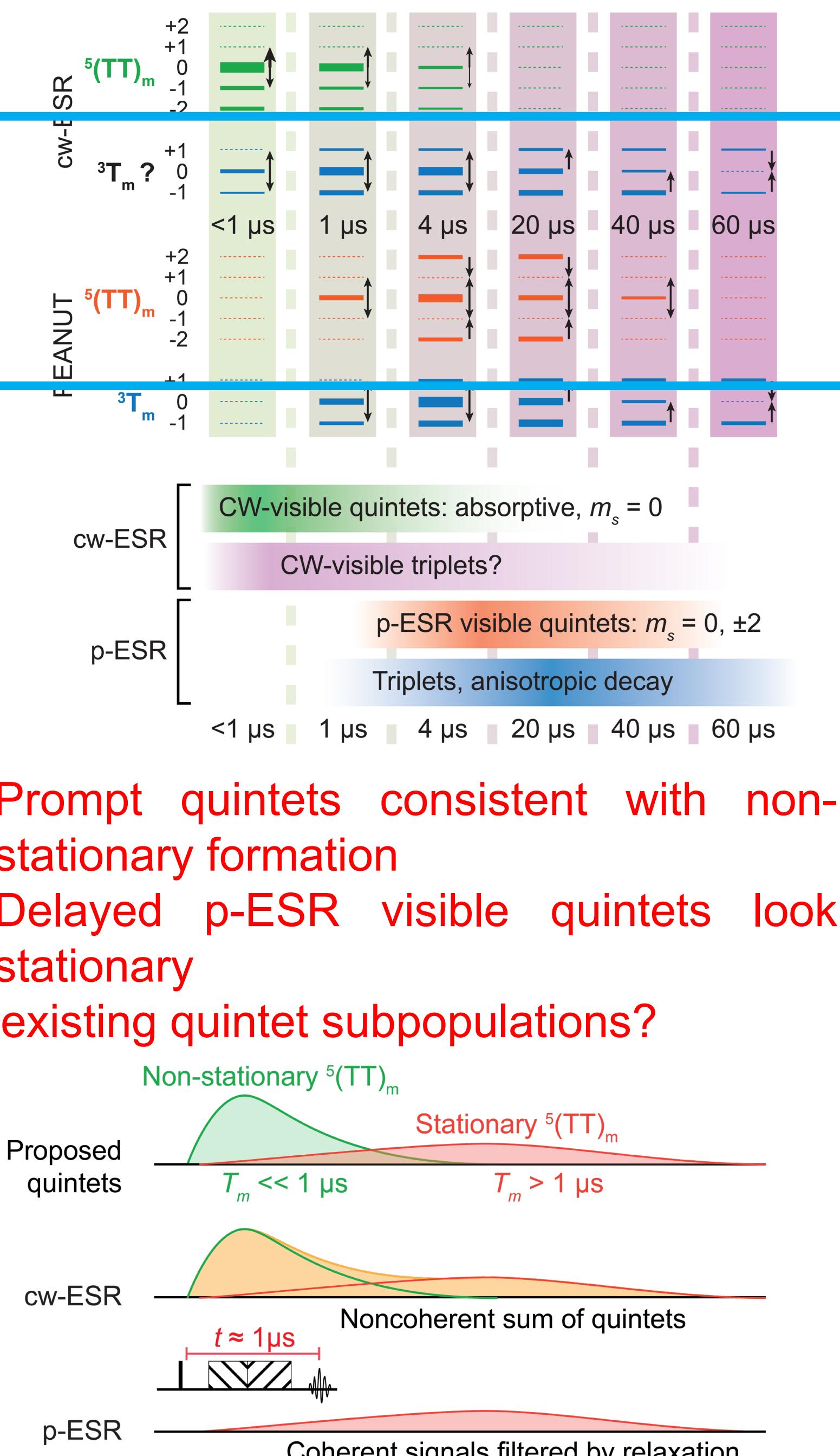
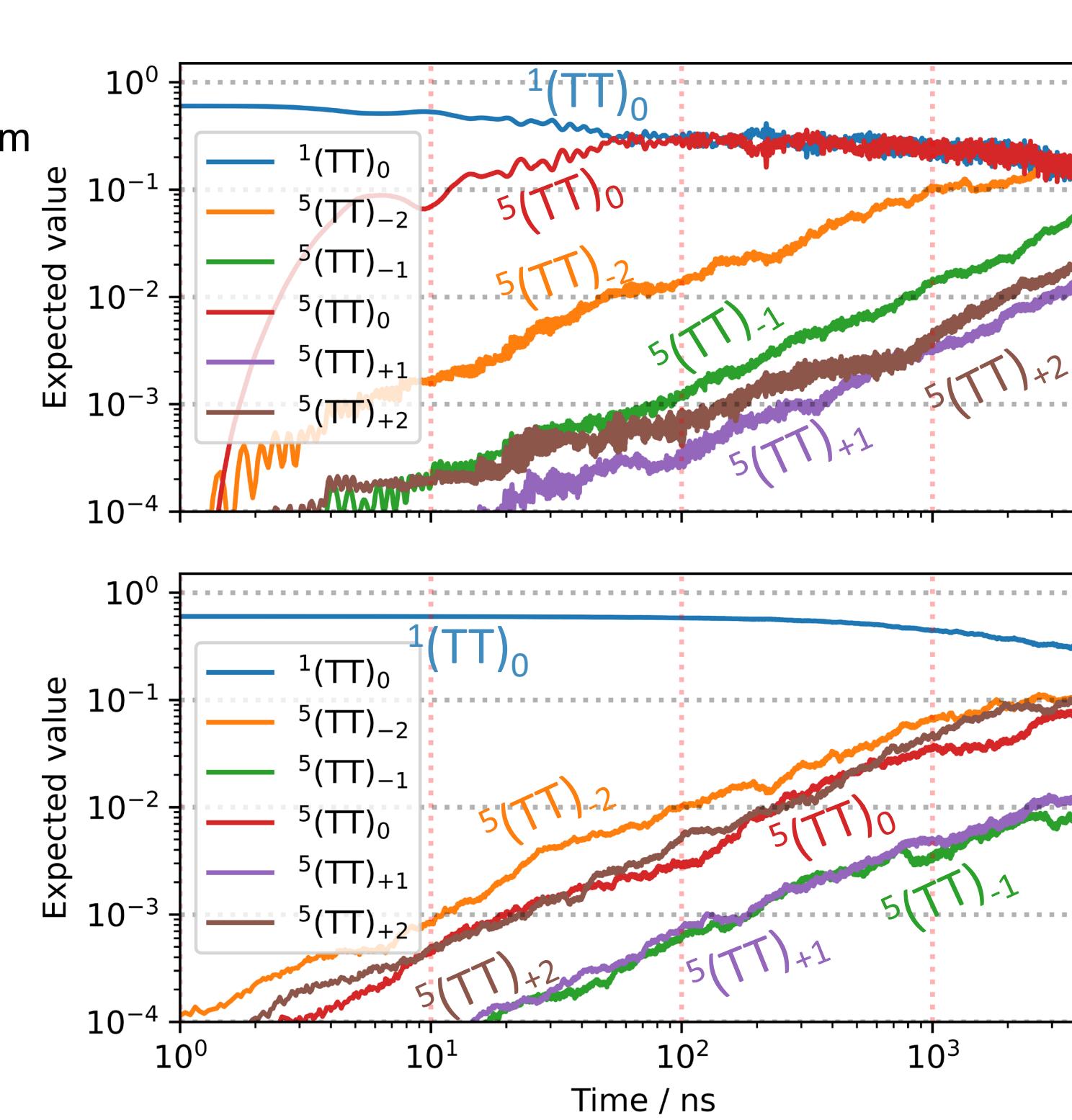


- ‘Non-stationary’  $1(TT)_0 \leftrightarrow 5(TT)_m$
- Fast mixing at low  $J$
  - Limited coherence times?
  - Adiabatic transit  $\rightarrow 5(TT)_{m<0}$
  - Slow, large-scale dynamics?

- ‘Stationary’  $1(TT)_0 \leftrightarrow 5(TT)_m$
- Slower mixing while  $J(t) \gg 0$
  - Longer coherence times?
  - Diabatic transfer  $\rightarrow 5(TT)_{\pm 2}$
  - Fast, small-scale dynamics?

## Reconciling cw-ESR and p-ESR results

We see different populations and dynamics by cw-ESR and p-ESR.  $J(t)$  reaches zero. Fast coherent mixing, possibly limited spin coherence times, adiabatic transit back to  $J_{high}$  biases toward low-lying quintets



## For Further Information

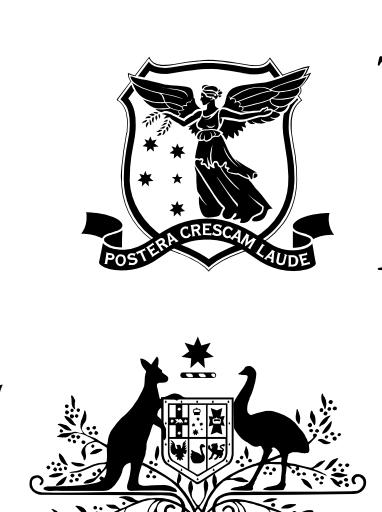
a  
b

## References

- Tayebjee, M. J. Y. et al. *Nature Physics* **13**, 182–188 (2017).
- Collins, M. I., et al. *J. Chem. Phys.* **151**, 164104 (2019).

## Acknowledgements

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