Let's check:

$$S^{2}|\pm\pm\rangle = 2|\pm\pm\rangle$$

$$S_{2}|\pm\pm\rangle = \pm|\pm\pm\rangle$$

$$S^{2}|\pm(|+-\rangle+|-+\rangle) = \frac{1}{\sqrt{2}}[(|+-\rangle+|-+\rangle)+(|-+\rangle+|+-\rangle)]$$

$$= 2 \times \frac{1}{\sqrt{2}}(|+-\rangle+|-+\rangle)$$

$$S_{2}|\pm(|+-\rangle+|-+\rangle) = \frac{1}{\sqrt{2}}(0+0) = 0$$

We can label these three states by  $S^2$ :  $2 = s(s+1) \Rightarrow s=1$ 

5z: m=1,0,-1

15,m> -> 11,1>, 11,0>, 11,-1>

We're missing one state, which must be orthogonal to the other 3:

$$5^{2} = (|+-\rangle - |-+\rangle) = \frac{1}{12} \left[ (|+-\rangle + |-+\rangle) - (|-+\rangle + |+-\rangle) \right]$$

$$= 0$$

15,m>= 10,0>

So adding Z spin 1/2 particles yields a spin one eigenket W sz projections ± 1,0 and a spin zero eigenket. tlow can we form eigenkets w/ simultaneous eigenvalues? recall ladder operators for 11/2:

$$5 \pm = 5 \times \pm i \leq y$$
  
 $5 - 1 + 0 = 1 - 0$   
 $5 + 1 - 0 = 1 + 0$   
 $5 - 1 - 0 = 0 + 1 + 0 = 0$ 

So for the addition of two spin 1/2 particles we have:

Since  $1\pm\pm$  has simultaneous eigenvalues of  $5^2$   $\pm$  5z we can start "at the top (bottom) of the ladder and lower (raise) to the bottom"

$$S_{-} | + + \rangle = S_{-} \otimes 1 | + + \rangle + 1 \otimes S_{-} | + + \rangle$$

$$= | - + \rangle + | + - \rangle$$

$$S_{-} (| - + \rangle + | + - \rangle) = S_{-} \otimes 1 (| - + \rangle + | + - \rangle) + 1 \otimes S_{-} (| - + \rangle + | + - \rangle)$$

$$= 0 + | - - \rangle + | - - \rangle + 0$$

So we have 3 potential states, but we want them nicely normalized:

$$|++\rangle \rightarrow |++\rangle$$

$$|-+\rangle + |+-\rangle \rightarrow \frac{1}{6}(|-+\rangle + |+-\rangle)$$

$$2|--\rangle \rightarrow |--\rangle$$

## Gaussian Integrals

$$I = \int_{-\infty}^{\infty} d\rho e^{-\frac{i}{2}a\rho^2 + J\rho}$$

complete the square: 
$$-\frac{1}{2}ap^2+Jp=-\frac{1}{2}a(p-\frac{1}{2})^2+X$$

$$=-\frac{1}{2}ap^2+Jp-\frac{1}{2}\frac{J^2}{a}+X \Rightarrow X=\frac{1}{2}\frac{J^2}{a}$$

$$I = \int_{-\infty}^{\infty} d\rho e^{-\frac{1}{2}\alpha(\rho - \frac{\omega}{\alpha})^2 + \frac{\omega^2}{2\alpha}}$$

let 
$$p \rightarrow p + \frac{1}{a} \Rightarrow dp \rightarrow dp$$

$$I = \int_{-\infty}^{\infty} e^{-\frac{1}{2}ap^2 + \frac{1}{2a}} = e^{\frac{\sqrt{2}}{2a}} \int_{-\infty}^{\infty} dp e^{-\frac{1}{2}ap^2}$$

let 
$$p \Rightarrow ka \Rightarrow dp \Rightarrow \frac{d}{ds}$$

$$I = \frac{1}{4}e^{3ka}\int dp e^{-ka}p^{2}$$

## Consider:

Wer:  

$$I^{2} \propto \left[ \int d\rho e^{-\frac{1}{2}\rho^{2}} \right]^{2} = \int dx \int dy e^{-\frac{1}{2}(x^{2}+y^{2})}$$

$$= \int r dr d\theta e^{-\frac{1}{2}r^{2}}$$

$$= 2\pi \int \left( \frac{1}{2} dr^{2} \right) e^{-\frac{1}{2}r^{2}}$$

$$= \pi \left( -2e^{-\frac{1}{2}r^{2}} \right) r^{2} = 0$$

$$= 2\pi$$

$$I = \int_{a}^{2\pi} e^{\int_{2a}^{2}}$$

For a multidimensional integral:
$$ap^2 \rightarrow \vec{p}^{\dagger} A \vec{p} = p_i^* A_{ij} p_i$$

$$I = \int_{-\infty}^{\infty} d\hat{p} e^{-\frac{1}{2}p_i^* A_{ij} p_j + J_{ij}^* p_i}$$

$$= \int_{-\infty}^{\infty} d\vec{p}' e^{-\frac{1}{p}} p_{i}^{i*} D_{ii} p_{i}' + J_{i}^{i*} M_{ij} p_{i}'$$

$$= \int_{-\infty}^{\infty} d\vec{p}' e^{-\frac{1}{p}} p_{i}^{i*} D_{ii} p_{i}' + J_{i}^{i*} M_{ij} p_{i}'$$

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$$= \int_{-\infty}^{\infty} d\vec{p}' e^{-\frac{1}{p}} p_{i}' + J_{i}^{i*} D_{ii} p_{i}' + J_{i}^{i*} D_{ii}' + J_{i}^{i*} D_$$

$$=\prod_{i}\frac{2\pi}{p_{ii}}e^{\frac{1}{2}J_{i}^{*}}D_{ij}J_{i}^{*}$$

since Dij is the diagonalized Aij det Ai; = [] Zi = [] Dii
eigenvalues

$$\frac{1}{D_{ij}} = D_{ij}^{-1}, \quad A_{i\ell}^{-1} = (M_{ij} D_{jk} M_{k\ell})^{-1}$$

$$= M_{ij} D_{jk} M_{k\ell}$$

$$D_{ij}^{*} D_{ij}^{-1} J_{i}^{*} = J_{i}^{*} M_{ij} D_{jk} M_{k\ell} J_{\ell} = J_{i}^{*} A_{ij} J_{j}$$

PK = Mkepe

$$I = \int \frac{(Z\pi)^{n}}{\det A} e^{\frac{1}{2}J^{\dagger}A^{\dagger}J} \qquad \text{for } A \text{ nxn matrix}$$