

For QED we had $G = \partial_\mu A_\mu$

Now we choose $G = \partial_\mu A_\mu - \frac{1}{\sqrt{2}} e v \pi$

$$\delta A_\mu = -\partial_\mu \alpha$$

$$\left. \begin{aligned} \delta \varphi &= +i e \alpha \varphi \\ \delta \varphi^* &= -i e \alpha \varphi^* \end{aligned} \right\} \pi = -i \sqrt{2} \ln(\varphi) \rightarrow \delta \pi = -i \sqrt{2} \frac{\delta \varphi - \delta \varphi^*}{2}$$

$$= \frac{-i \sqrt{2}}{2} (+i e \alpha \varphi + i e \alpha \varphi^*)$$

$$= + \frac{e \sqrt{2}}{2} \alpha \frac{2(v+\sigma)}{\sqrt{2}} = +e(v+\sigma)\alpha$$

$$\delta G = -\partial_\mu \partial^\mu \alpha + \frac{1}{\sqrt{2}} e v (-e(v+\sigma)) \alpha$$

$$= [-\square - \frac{1}{2} e^2 v (v+\sigma)] \alpha$$

$$\frac{\delta G}{\delta \alpha} = -\square - \frac{1}{2} e^2 v (v+\sigma)$$

$$\text{So } \mathcal{L}_{gh} = +\bar{c}^a (-\partial_\mu \partial^\mu - \frac{1}{2} e^2 v (v+\sigma)) c^a$$

$$= (\partial_\mu \bar{c}^a)(\partial^\mu c^a) - \underbrace{\frac{1}{2} e^2 v^2 \bar{c}^a c^a}_{\text{ghost mass } m_{gh}^2 = \frac{1}{2} e^2 v^2} - \underbrace{\frac{1}{2} e^2 v \sigma \bar{c}^a c^a}_{\text{ghost-}\sigma \text{ interaction}}$$

$$\text{And } \mathcal{L}_{GF} = \frac{-1}{2\xi} G^2 = \frac{-1}{2\xi} (\partial_\mu A_\mu - \frac{1}{\sqrt{2}} e v \pi)^2$$

$$= \frac{-1}{2\xi} [(\partial_\mu A_\mu)^2 - 2 \frac{1}{\sqrt{2}} e v \pi \partial_\mu A_\mu + \frac{1}{2} e^2 v^2 \pi^2]$$

$$= \underbrace{\frac{-1}{2\xi} (\partial_\mu A_\mu)^2}_{\text{normal GF term}} + \underbrace{e v \pi \partial_\mu A_\mu}_{\text{IBP} \rightarrow -e v A_\mu \partial_\mu \pi \text{ exactly cancel } A-\pi \text{ mixing}} - \frac{1}{2} e^2 v^2 \pi^2$$

$m_\pi^2 = \frac{1}{2} e^2 v^2$ (π is IR so $\frac{1}{2}$ is not part of m_π^2)

We have:

$$\delta A_\mu = -\partial_\mu \alpha$$

$$\delta \phi = ie\alpha \phi$$

$$\delta \phi^* = -ie\alpha \phi^*$$

So taking

$$\begin{aligned} \delta \Gamma = 0 &= \frac{\delta \Gamma}{\delta A_\mu} (-\partial_\mu \alpha) + \frac{\delta \Gamma}{\delta \phi} ie\alpha \phi + \frac{\delta \Gamma}{\delta \phi^*} (-ie\alpha \phi^*) \\ &= \underbrace{\left[\partial_\mu \frac{\delta \Gamma}{\delta A_\mu} + ie \left(\frac{\delta \Gamma}{\delta \phi} \phi - \frac{\delta \Gamma}{\delta \phi^*} \phi^* \right) \right]}_{=0 \text{ bc } \alpha \text{ arbitrary}} \alpha \end{aligned}$$

Taking a variation w.r to A_μ , then setting all fields to their vevs:

$$\partial_\mu \frac{\delta^2 \Gamma}{\delta A_\mu \delta A_\mu} + ie \left(\frac{\delta^2 \Gamma}{\delta A_\mu \delta \phi} \langle \phi \rangle - \frac{\delta^2 \Gamma}{\delta A_\mu \delta \phi^*} \langle \phi^* \rangle \right) = 0$$

$\uparrow \quad \quad \quad \uparrow$
 both $\frac{v}{\sqrt{2}}$

$$\phi = \frac{v + \sigma + i\pi}{\sqrt{2}} \rightarrow \frac{\delta}{\delta \phi} = \underbrace{\frac{\delta \phi}{\delta \sigma}}_{1/\sqrt{2}} \frac{\delta}{\delta \sigma} + \underbrace{\frac{\delta \phi}{\delta \pi}}_{i/\sqrt{2}} \frac{\delta}{\delta \pi}, \quad \frac{\delta}{\delta \phi^*} = \frac{1}{\sqrt{2}} \frac{\delta}{\delta \sigma} - \frac{i}{\sqrt{2}} \frac{\delta}{\delta \pi}$$

$$\partial_\mu \frac{\delta^2 \Gamma}{\delta A_\mu \delta A_\mu} + \frac{iev}{2} \left(\frac{\delta^2 \Gamma}{\delta A_\mu \delta \sigma} + i \frac{\delta^2 \Gamma}{\delta A_\mu \delta \pi} - \frac{\delta^2 \Gamma}{\delta A_\mu \delta \sigma} + i \frac{\delta^2 \Gamma}{\delta A_\mu \delta \pi} \right) = 0$$

$$\rightarrow \partial_\mu \frac{\delta^2 \Gamma}{\delta A_\mu \delta A_\mu} = ev \frac{\delta^2 \Gamma}{\delta A_\mu \delta \pi}$$

\uparrow
 Longitudinal $2p+1$ is prop to $A-\pi$ mixing! $\partial_\mu (m \otimes m) = ev (m \otimes -)$

Returning to the master Ward ID \dot{z} varying w/r to π :

$$\partial_\mu \frac{\delta^2 \Gamma}{\delta A_\mu \delta \pi} + ie \left(\frac{\delta^2 \Gamma}{\delta \phi \delta \pi} \langle \phi \rangle + \frac{\delta \Gamma}{\delta \phi} \frac{\delta \phi}{\delta \pi} - \frac{\delta^2 \Gamma}{\delta \phi^* \delta \pi} \langle \phi^* \rangle - \frac{\delta \Gamma}{\delta \phi^*} \frac{\delta \phi^*}{\delta \pi} \right) = 0$$

$$= \partial_\mu \frac{\delta^2 \Gamma}{\delta A_\mu \delta \pi} + ie \left(\frac{1}{\sqrt{2}} \frac{\delta^2 \Gamma}{\delta \phi \delta \pi} \frac{v}{\sqrt{2}} + \frac{i}{\sqrt{2}} \frac{\delta^2 \Gamma}{\delta \pi \delta \pi} \frac{v}{\sqrt{2}} + \frac{1}{\sqrt{2}} \frac{\delta \Gamma}{\delta \phi} \frac{i}{\sqrt{2}} + \frac{i}{\sqrt{2}} \frac{\delta \Gamma}{\delta \pi} \frac{i}{\sqrt{2}} \right. \\ \left. - \frac{1}{\sqrt{2}} \frac{\delta^2 \Gamma}{\delta \phi \delta \pi} \frac{v}{\sqrt{2}} + \frac{i}{\sqrt{2}} \frac{\delta^2 \Gamma}{\delta \pi \delta \pi} \frac{v}{\sqrt{2}} - \frac{1}{\sqrt{2}} \frac{\delta \Gamma}{\delta \phi} \left(-\frac{i}{\sqrt{2}}\right) + \frac{i}{\sqrt{2}} \frac{\delta \Gamma}{\delta \pi} \left(-\frac{i}{\sqrt{2}}\right) \right) = 0$$

$$= \partial_\mu \frac{\delta^2 \Gamma}{\delta A_\mu \delta \pi} - e \left(\frac{\delta^2 \Gamma}{\delta \pi \delta \pi} v + \frac{\delta \Gamma}{\delta \phi} \right) = 0$$

$$\partial_\mu \left(\text{tadpole diagram} \right) - e v \left(\text{tadpole diagram} \right) - e \left(\text{tadpole diagram} \right) = 0$$

Setting all fields in the master Ward ID to zero we see:

$$\underbrace{\partial_\mu \frac{\delta \Gamma}{\delta A_\mu}}_{=0, \text{ no } A_\mu \text{ tadpole}} + ie \frac{v}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} \frac{\delta \Gamma}{\delta \phi} + \frac{i}{\sqrt{2}} \frac{\delta \Gamma}{\delta \pi} - \frac{1}{\sqrt{2}} \frac{\delta \Gamma}{\delta \phi} + \frac{i}{\sqrt{2}} \frac{\delta \Gamma}{\delta \pi} \right) = 0$$

$$\Rightarrow \frac{\delta \Gamma}{\delta \pi} v = 0, \text{ since } v \neq 0 \text{ we conclude } \frac{\delta \Gamma}{\delta \pi} = 0$$

$\bar{Q} \tilde{H} U_R$ invar

$$\tilde{H} \equiv i\sigma_2 H^*$$

$$H^* \rightarrow e^{-ig_1/2 \theta_1 + \frac{-ig_2}{2} \theta_2 \sigma^*} \quad (\text{recall } H \rightarrow e^{ig_1/2 \theta_1 + \frac{ig_2}{2} \theta_2 \sigma})$$

$$\bar{Q} \rightarrow e^{-ig_1/6 \theta_1 - \frac{i}{2} g_2 \theta_2 \sigma - \frac{i}{2} g_3 \theta_3 \lambda}$$

$$U_R \rightarrow e^{ig_1 \theta_1 \frac{2}{3} + \frac{i}{2} g_3 \theta_3 \lambda}$$

So $U(1)_Y$ is trivial: $e^{-ig_1/2 \theta_1} \cdot e^{-ig_1/6 \theta_1} e^{ig_1 \theta_1 \frac{2}{3}} = e^{ig_1 (\frac{1}{2} - \frac{1}{6} + \frac{2}{3}) \theta_1} = 1$

$U(1)_C$ is too: $e^{-\frac{i}{2} g_3 \theta_3 \lambda} e^{\frac{i}{2} g_3 \theta_3 \lambda} = 1$

For $SU(2)_L$ recall: $(\sigma^1)^* = \sigma^1$, $(\sigma^2)^* = -\sigma^2$, $(\sigma^3)^* = \sigma^3$

$$\sigma^i \sigma^2 + \sigma^2 \sigma^{i*} = \begin{cases} i=1: \sigma^1 \sigma^2 + \sigma^2 \sigma^{1*} = \{\sigma^1, \sigma^2\} = 0 \\ 2: \sigma^2 \sigma^2 + \sigma^2 \sigma^{2*} = 0 \\ 3: \sigma^3 \sigma^2 + \sigma^2 \sigma^{3*} = \{\sigma^3, \sigma^2\} = 0 \end{cases}$$

$$\begin{aligned} \text{So } e^{-\frac{i}{2} g_2 \theta_2 \sigma} i\sigma_2 e^{\frac{i}{2} g_2 \theta_2 \sigma^*} &\rightarrow (1 - \frac{ig_2}{2} \theta_2 \sigma) i\sigma_2 (1 - \frac{i}{2} g_2 \theta_2 \sigma^*) \\ &= i\sigma_2 + \frac{g_2^2}{2} \theta_2^2 \sigma^i \sigma^i \sigma_2 + \frac{g_2^2}{2} \sigma_2 \theta_2^2 (\sigma^i)^* + O(\theta_2^4) \\ &= i\sigma_2 + \frac{g_2^2}{2} \theta_2^2 \underbrace{(\sigma^i \sigma_2 + \sigma_2 \sigma^{i*})}_{=0} \\ &= i\sigma_2 \end{aligned}$$

So $\bar{Q} \tilde{H} U_R$ is $U(1)_Y \otimes SU(2)_L \otimes SU(3)_C$ invar.