

# Advanced Quantum Theory

## Homework 5

Tim Lawson

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### 4.2 Wavefunction

Consider a bosonic, spinless system with three single-particle states. Write down the normalised wavefunction of the multiple-particle state  $|3, 0, 1\rangle$  in occupation-number representation.

For bosons, we have that:

$$\psi_{i_1, \dots, i_N}(\vec{r}_1, \dots, \vec{r}_N) = C \sum_{\pi} \psi_{i_{\pi(1)}}(\vec{r}_1) \dots \psi_{i_{\pi(N)}}(\vec{r}_N), \quad C = \frac{1}{\sqrt{N! \prod_i n_i!}} \quad (4.2.1)$$

For  $N = 4$ :

$$C = \frac{1}{\sqrt{4! 3! 0! 1!}} = \frac{1}{12}$$

There are  $4! = 24$  permutations of the indices  $i_1, \dots, i_4$ . Hence:

$$\begin{aligned} \psi_{3,0,1}(\vec{r}_1, \vec{r}_2, \vec{r}_3, \vec{r}_4) &= \frac{1}{12} \sum_{\pi} \psi_{i_{\pi(1)}}(\vec{r}_1) \psi_{i_{\pi(2)}}(\vec{r}_2) \psi_{i_{\pi(3)}}(\vec{r}_3) \psi_{i_{\pi(4)}}(\vec{r}_4) \\ &= \frac{1}{12} \left( \psi_1(\vec{r}_1) \psi_1(\vec{r}_2) \psi_1(\vec{r}_3) \psi_3(\vec{r}_4) + \psi_1(\vec{r}_1) \psi_1(\vec{r}_2) \psi_1(\vec{r}_4) \psi_3(\vec{r}_3) + \right. \\ &\quad \psi_1(\vec{r}_1) \psi_1(\vec{r}_3) \psi_1(\vec{r}_2) \psi_3(\vec{r}_4) + \psi_1(\vec{r}_1) \psi_1(\vec{r}_3) \psi_1(\vec{r}_4) \psi_3(\vec{r}_2) + \\ &\quad \psi_1(\vec{r}_1) \psi_1(\vec{r}_4) \psi_1(\vec{r}_2) \psi_3(\vec{r}_3) + \psi_1(\vec{r}_1) \psi_1(\vec{r}_4) \psi_1(\vec{r}_3) \psi_3(\vec{r}_2) + \\ &\quad \psi_1(\vec{r}_2) \psi_1(\vec{r}_1) \psi_1(\vec{r}_3) \psi_3(\vec{r}_4) + \psi_1(\vec{r}_2) \psi_1(\vec{r}_1) \psi_1(\vec{r}_4) \psi_3(\vec{r}_3) + \\ &\quad \psi_1(\vec{r}_2) \psi_1(\vec{r}_3) \psi_1(\vec{r}_1) \psi_3(\vec{r}_4) + \psi_1(\vec{r}_2) \psi_1(\vec{r}_3) \psi_1(\vec{r}_4) \psi_3(\vec{r}_1) + \\ &\quad \psi_1(\vec{r}_2) \psi_1(\vec{r}_4) \psi_1(\vec{r}_1) \psi_3(\vec{r}_3) + \psi_1(\vec{r}_2) \psi_1(\vec{r}_4) \psi_1(\vec{r}_3) \psi_3(\vec{r}_1) + \\ &\quad \psi_1(\vec{r}_3) \psi_1(\vec{r}_1) \psi_1(\vec{r}_2) \psi_3(\vec{r}_4) + \psi_1(\vec{r}_3) \psi_1(\vec{r}_1) \psi_1(\vec{r}_4) \psi_3(\vec{r}_2) + \\ &\quad \psi_1(\vec{r}_3) \psi_1(\vec{r}_2) \psi_1(\vec{r}_1) \psi_3(\vec{r}_4) + \psi_1(\vec{r}_3) \psi_1(\vec{r}_2) \psi_1(\vec{r}_4) \psi_3(\vec{r}_1) + \\ &\quad \psi_1(\vec{r}_3) \psi_1(\vec{r}_4) \psi_1(\vec{r}_1) \psi_3(\vec{r}_2) + \psi_1(\vec{r}_3) \psi_1(\vec{r}_4) \psi_1(\vec{r}_2) \psi_3(\vec{r}_1) + \\ &\quad \psi_1(\vec{r}_4) \psi_1(\vec{r}_1) \psi_1(\vec{r}_2) \psi_3(\vec{r}_3) + \psi_1(\vec{r}_4) \psi_1(\vec{r}_1) \psi_1(\vec{r}_3) \psi_3(\vec{r}_2) + \\ &\quad \psi_1(\vec{r}_4) \psi_1(\vec{r}_2) \psi_1(\vec{r}_1) \psi_3(\vec{r}_3) + \psi_1(\vec{r}_4) \psi_1(\vec{r}_2) \psi_1(\vec{r}_3) \psi_3(\vec{r}_1) + \\ &\quad \left. \psi_1(\vec{r}_4) \psi_1(\vec{r}_3) \psi_1(\vec{r}_1) \psi_3(\vec{r}_2) + \psi_1(\vec{r}_4) \psi_1(\vec{r}_3) \psi_1(\vec{r}_2) \psi_3(\vec{r}_1) \right) \end{aligned}$$

## 4.5 Bose-Hubbard model

Consider a system with an arbitrary number of indistinguishable bosonic particles. It has two sites and the Hamiltonian:

$$\hat{H} = -\hat{a}_1^\dagger \hat{a}_2 - \hat{a}_2^\dagger \hat{a}_1 + \frac{1}{2} \sum_i \hat{n}_i (\hat{n}_i - 1), \quad \hat{n}_i = \hat{a}_i^\dagger \hat{a}_i \quad (4.5.1)$$

(a) Show that applying  $\hat{H}$  to a state does not change its particle number.

Heuristically, each term in the Hamiltonian contains an equal number of creation and annihilation operators. Apply  $\hat{H}$  to a state  $|n_1, n_2\rangle$  in occupation-number representation:

$$\begin{aligned} \hat{H} &= \left( -\hat{a}_1^\dagger \hat{a}_2 - \hat{a}_2^\dagger \hat{a}_1 + \frac{1}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) \right) |n_1, n_2\rangle \\ &= -\hat{a}_1^\dagger \hat{a}_2 |n_1, n_2\rangle - \hat{a}_2^\dagger \hat{a}_1 |n_1, n_2\rangle + \left( \frac{1}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) \right) |n_1, n_2\rangle \end{aligned}$$

The three terms are:

$$\begin{aligned} -\hat{a}_1^\dagger \hat{a}_2 |n_1, n_2\rangle &= -\sqrt{n_2} \hat{a}_1^\dagger |n_1, n_2 - 1\rangle \\ &= -\sqrt{n_1 + 1} \sqrt{n_2} |n_1 + 1, n_2 - 1\rangle \end{aligned}$$

$$\begin{aligned} -\hat{a}_2^\dagger \hat{a}_1 |n_1, n_2\rangle &= -\sqrt{n_1} \hat{a}_2^\dagger |n_1 - 1, n_2\rangle \\ &= -\sqrt{n_1} \sqrt{n_2 + 1} |n_1 - 1, n_2 + 1\rangle \end{aligned}$$

$$\left( \frac{1}{2} \sum_i \hat{n}_i (\hat{n}_i - 1) \right) |n_1, n_2\rangle = \frac{1}{2} (n_1(n_1 - 1) + n_2(n_2 - 1)) |n_1, n_2\rangle$$

In each case, the particle number is unchanged.

(b) Write down the basis states of the system with two particles.

In occupation-number representation, the basis states are  $|2, 0\rangle, |1, 1\rangle, |0, 2\rangle$ .

(c) Represent the Hamiltonian in matrix form and determine its eigenvalues.

Apply  $\hat{H}$  to the basis states:

$$\hat{H}|2, 0\rangle = -\sqrt{2} \hat{a}_2^\dagger |1, 0\rangle = -\sqrt{2} |1, 1\rangle$$

$$\hat{H}|1, 1\rangle = -\hat{a}_1^\dagger |1, 0\rangle - \hat{a}_2^\dagger |0, 1\rangle = -\sqrt{2} |2, 0\rangle - \sqrt{2} |0, 2\rangle$$

$$\hat{H}|0, 2\rangle = -\sqrt{2} \hat{a}_1^\dagger |0, 1\rangle = -\sqrt{2} |1, 1\rangle$$

In matrix form:

$$\begin{aligned} H(i, j) &= \begin{bmatrix} \langle 2, 0 | \hat{H} | 2, 0 \rangle & \langle 2, 0 | \hat{H} | 1, 1 \rangle & \langle 2, 0 | \hat{H} | 0, 2 \rangle \\ \langle 1, 1 | \hat{H} | 2, 0 \rangle & \langle 1, 1 | \hat{H} | 1, 1 \rangle & \langle 1, 1 | \hat{H} | 0, 2 \rangle \\ \langle 0, 2 | \hat{H} | 2, 0 \rangle & \langle 0, 2 | \hat{H} | 1, 1 \rangle & \langle 0, 2 | \hat{H} | 0, 2 \rangle \end{bmatrix} \\ &= \begin{bmatrix} 0 & -\sqrt{2} & 0 \\ -\sqrt{2} & 0 & -\sqrt{2} \\ 0 & -\sqrt{2} & 0 \end{bmatrix} \end{aligned}$$

The eigenvalues of  $A$  are the roots of its characteristic polynomial,  $\det(A - \lambda I) = 0$ .

$$\begin{aligned}\det(A - \lambda I) &= \begin{vmatrix} -\lambda & -\sqrt{2} & 0 \\ -\sqrt{2} & -\lambda & -\sqrt{2} \\ 0 & -\sqrt{2} & -\lambda \end{vmatrix} = -\lambda \begin{vmatrix} -\lambda & -\sqrt{2} \\ -\sqrt{2} & -\lambda \end{vmatrix} + \sqrt{2} \begin{vmatrix} -\sqrt{2} & -\sqrt{2} \\ 0 & -\lambda \end{vmatrix} \\ &= -\lambda (\lambda^2 - 2) + 2\lambda = -\lambda (4 - \lambda^2) = 0\end{aligned}$$

Hence, the eigenvalues are  $\lambda = 0, \pm 2$ .