Advanced Quantum Theory Homework 1

Tim Lawson

November 17, 2023

2.1 Path integral in phase space

Show that the propagator of a quantum system can be written as:

$$K(\vec{r}_f, \vec{r}_0, t) = \int D[\vec{r}] D[\vec{p}] \exp\left(\frac{i}{\hbar} \int_0^t \left(\vec{p}(t') \cdot \dot{\vec{r}}(t') - H(\vec{r}(t'), \vec{p}(t'))\right) dt'\right)$$
(2.1.1)

We have that:

$$K(\vec{r}_f, \vec{r}_0, t) = \int d^n r_1 \dots \int d^n r_{N-1} \prod_{i=0}^{N-1} \langle \vec{r}_{j+1} | e^{-\frac{i}{\hbar} \hat{H} \tau} | \vec{r}_j \rangle$$
 (2.1.2)

Assume that the Hamiltonian has the form $\hat{H} = \hat{T}(\vec{p}) + \hat{U}(\vec{r})$. Replace the exponential by the first-order Taylor polynomial and vice versa:

$$\begin{split} \langle \, \vec{r}_{j+1} \, | e^{-\frac{i}{\hbar} \hat{H} \tau} | \, \vec{r}_{j} \, \rangle &\approx \langle \, \vec{r}_{j+1} \, | \, \left(1 - \frac{i}{\hbar} \hat{H} \tau \right) | \, \vec{r}_{j} \, \rangle \\ &\approx \langle \, \vec{r}_{j+1} \, | \, \left(1 - \frac{i}{\hbar} \hat{T} \tau \right) \left(1 - \frac{i}{\hbar} \hat{U} \tau \right) | \, \vec{r}_{j} \, \rangle \\ &\approx \langle \, \vec{r}_{j+1} \, | \, e^{-\frac{i}{\hbar} \hat{T} \tau} e^{-\frac{i}{\hbar} \hat{U} \tau} | \, \vec{r}_{j} \, \rangle \end{split}$$

Apply \hat{U} to the position eigenstate $|\vec{r_i}\rangle$:

$$\langle \vec{r}_{i+1} | e^{-\frac{i}{\hbar}\hat{H}\tau} | \vec{r}_i \rangle \approx \langle \vec{r}_{i+1} | e^{-\frac{i}{\hbar}\hat{T}\tau} | \vec{r}_i \rangle e^{-\frac{i}{\hbar}U(\vec{r}_j)\tau}$$

Insert a resolution of the identity in terms of momentum eigenfunctions:

$$\langle \vec{r}_{j+1} | e^{-\frac{i}{\hbar}\hat{H}\tau} | \vec{r}_{j} \rangle = \langle \vec{r}_{j+1} | e^{-\frac{i}{\hbar}\hat{T}\tau} \int d^{n}p_{j} | \vec{p}_{j} \rangle \langle \vec{p}_{j} | \vec{r}_{j} \rangle e^{-\frac{i}{\hbar}U(\vec{r}_{j})\tau}$$

$$= \int d^{n}p_{j}e^{-\frac{i}{\hbar}U(\vec{r}_{j})\tau} \langle \vec{r}_{j+1} | e^{-\frac{i}{\hbar}\hat{T}\tau} | \vec{p}_{j} \rangle \langle \vec{p}_{j} | \vec{r}_{j} \rangle$$

Apply \hat{T} to the momentum eigenstate $|\vec{p}_i\rangle$ and substitute:

$$\begin{split} \langle\,\vec{r}_{j+1}\,|e^{-\frac{i}{\hbar}\hat{H}\tau}|\,\vec{r}_{j}\,\rangle &= \int\mathrm{d}^{n}p_{j}\exp\left(-\frac{i}{\hbar}\left(T(\vec{p}_{j})+U(\vec{r}_{j})\right)\tau\right)\langle\,\vec{r}_{j+1}\,|\,\vec{p}_{j}\,\rangle\langle\,\vec{p}_{j}\,|\,\vec{r}_{j}\,\rangle\\ &= \int\mathrm{d}^{n}p_{j}\exp\left(-\frac{i}{\hbar}H(\vec{r}_{j},\vec{p}_{j})\tau\right)\langle\,\vec{r}_{j+1}\,|\,\vec{p}_{j}\,\rangle\langle\,\vec{p}_{j}\,|\,\vec{r}_{j}\,\rangle\\ &= \int\mathrm{d}^{n}p_{j}\exp\left(-\frac{i}{\hbar}H(\vec{r}_{j},\vec{p}_{j})\tau\right)\frac{1}{(2\pi\hbar)^{n/2}}e^{\frac{i}{\hbar}\vec{p}_{j}\cdot\vec{r}_{j+1}}\frac{1}{(2\pi\hbar)^{n/2}}e^{-\frac{i}{\hbar}\vec{p}_{j}\cdot\vec{r}_{j}}\\ &= \frac{1}{(2\pi\hbar)^{n}}\int\mathrm{d}^{n}p_{j}\exp\left(\frac{i}{\hbar}\left(\vec{p}_{j}\cdot\frac{\vec{r}_{j+1}-\vec{r}_{j}}{\tau}-H(\vec{r}_{j},\vec{p}_{j})\right)\tau\right) \end{split}$$

The product in equation 2.1.2 is then:

$$\begin{split} & \prod_{j=0}^{N-1} \frac{1}{(2\pi\hbar)^n} \int \mathrm{d}^n p_j \exp\left(\frac{i}{\hbar} \left(\vec{p}_j \cdot \frac{\vec{r}_{j+1} - \vec{r}_j}{\tau} - H(\vec{r}_j, \vec{p}_j)\right) \tau\right) \\ &= \frac{1}{(2\pi\hbar)^{Nn}} \int \mathrm{d}^n p_0 \dots \int \mathrm{d}^n p_{N-1} \exp\left(\frac{i}{\hbar} \sum_{j=0}^{N-1} \left(\vec{p}_j \cdot \frac{\vec{r}_{j+1} - \vec{r}_j}{\tau} - H(\vec{r}_j, \vec{p}_j)\right) \tau\right) \end{split}$$

Equation 2.1.2 becomes:

$$K(\vec{r}_f, \vec{r}_0, t) = \frac{1}{(2\pi\hbar)^{Nn}} \int d^n r_1 \dots \int d^n r_{N-1} \int d^n p_0 \dots \int d^n p_{N-1}$$
$$\exp\left(\frac{i}{\hbar} \sum_{j=0}^{N-1} \left(\vec{p}_j \cdot \frac{\vec{r}_{j+1} - \vec{r}_j}{\tau} - H(\vec{r}_j, \vec{p}_j)\right) \tau\right)$$

Replace the sum by an integral in the limit $N \to \infty, \, \tau \to 0$ to obtain:

$$K(\vec{r}_f, \vec{r}_0, t) = \int D\left[\vec{r}\right] D\left[\vec{p}\right] \exp\left(\frac{i}{\hbar} \int_0^t \left(\vec{p}(t') \cdot \dot{\vec{r}}(t') - H(\vec{r}(t'), \vec{p}(t'))\right) dt'\right)$$

where

$$\int D\left[\vec{r}\right]D\left[\vec{p}\right]\ldots = \lim_{N\to\infty}\frac{1}{(2\pi\hbar)^{Nn}}\int d^nr_1\ldots\int d^nr_{N-1}\int d^np_0\ldots\int d^np_{N-1}\ldots$$