

## Grover's algorithm

①

Having not sorted array of elements, the task is to find certain element (which we will call the "winner").

Illustrative example: having phone number find the name of the person in phone book.

$$\{x_i\} \quad i=0, 1, 2, \dots, N-1$$

N elements

~~w~~ - winner, which we look for

We have also function (so called oracle) which for given element answers if it is a winner or not:

$$f(x) = \begin{cases} 0 & \text{for } x \neq w \\ 1 & \text{for } x = w \end{cases} \quad (1)$$

- Number of calls to function  $f$  is the complexity of the algorithm.
- Classically, since the array is unsorted, statistically we need  $N/2$  calls, which means complexity is  $O(N)$

- such oracle functions are very popular in quantum algorithms
- having function  $f(x)$  we can use the following form of oracle as an operator:

$$\hat{O}|x\rangle \equiv (-1)^{f(x)}|x\rangle \quad (2)$$

Let's consider the situation where states  $|x_0\rangle, |x_1\rangle, \dots, |x_{N-1}\rangle$  are basis states in some Hilbert space of dimension  $N$ .

We can construct such Hilbert space using  $n$  qubits:  ~~$N$~~   $N = \underline{2^n}$

Let's consider following state:

$$|\alpha\rangle = \frac{1}{\sqrt{N-1}} \sum_{x \neq w} |x\rangle \Rightarrow |x\rangle \perp |w\rangle \quad (3)$$

↑  
as  $|x\rangle$  does not contain  $|w\rangle$

Let's consider also state  $|y\rangle$  being a linear combination of all basis states:

$$|y\rangle \equiv \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle = H^{\otimes n} |0^{\otimes n}\rangle \quad (4)$$

e have

(2)

$$|\psi\rangle = \frac{1}{\sqrt{N}} \left( \sum_{x \neq w} |x\rangle + |w\rangle \right) = \frac{1}{\sqrt{N}} \sum_{x \neq w} |x\rangle + \frac{1}{\sqrt{N}} |w\rangle =$$

} from (3):

$$= \left\{ \sum_{x \neq w} |x\rangle = \sqrt{N-1} |\alpha\rangle \right\} = \frac{\sqrt{N-1}}{\sqrt{N}} |\alpha\rangle + \frac{1}{\sqrt{N}} |w\rangle =$$

$$= \sqrt{1 - \frac{1}{N}} |\alpha\rangle + \sqrt{\frac{1}{N}} |w\rangle \quad (5)$$

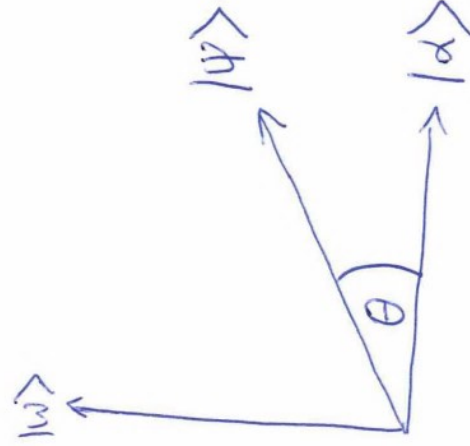
Now by analogy to  $\sin\theta$  and  $\cos\theta$ :

$$\sin^2\theta + \cos^2\theta = 1 \Rightarrow \sqrt{1 - \sin^2\theta} = \sqrt{\cos^2\theta} \quad (6)$$

We see, that there is always such  $\theta$ , that

$$|\psi\rangle = \cos\theta |\alpha\rangle + \sin\theta |w\rangle$$

Graphically:





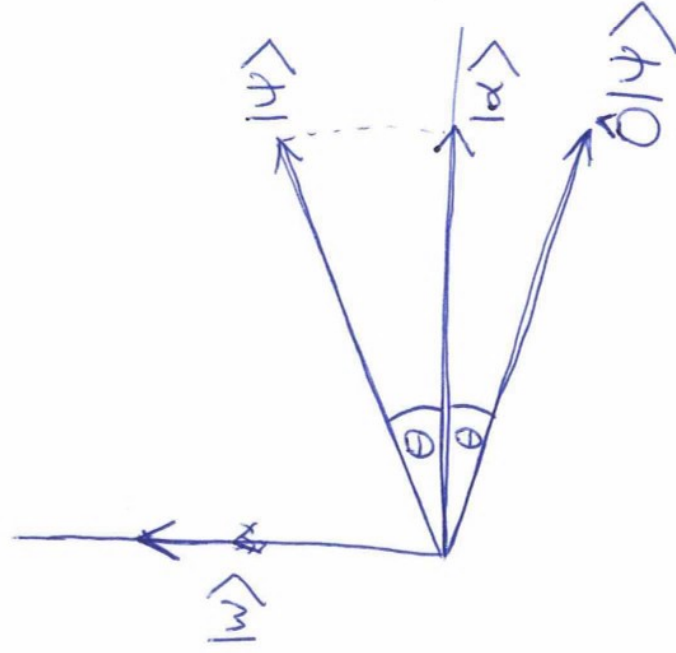
$$\boxed{\sin \theta = \left[ \frac{1}{N} \right]} \Rightarrow \text{for large } N, \theta \text{ is small}$$

If we have state  $|\psi\rangle$  let's see what effect has acting on it with oracle operator  $\hat{O}$ :

$$\begin{aligned} \hat{O}|\psi\rangle &= \hat{O} \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle = \frac{1}{\sqrt{N}} \hat{O}(|x_0\rangle + |x_1\rangle + \dots + |w\rangle + \dots + |x_{N-1}\rangle) \\ &= \frac{1}{\sqrt{N}} (\hat{O}|x_0\rangle + \hat{O}|x_1\rangle + \dots + \hat{O}|x_{N-1}\rangle) = \\ &= \left\{ \begin{array}{l} \text{based on (2)} \\ \text{and (1)} \end{array} \right\} = \frac{1}{\sqrt{N}} (|x_0\rangle + |x_1\rangle + \dots + \underbrace{-|w\rangle}_{\text{flipped}} + \dots + |x_{N-1}\rangle) \end{aligned}$$

$\Downarrow$

$\hat{O}|\psi\rangle$  does not change any of the amplitudes except the one of  $|w\rangle$  which is flipped

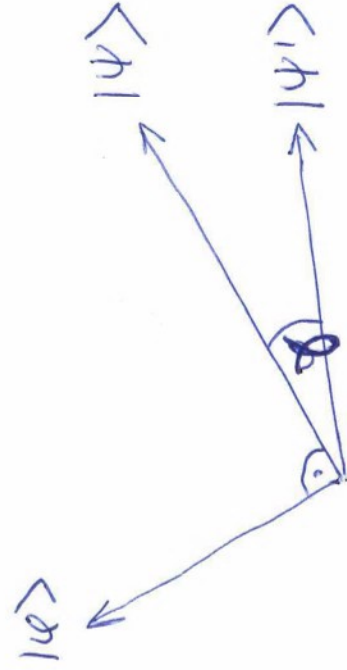


③ Now, let's consider the following operator:

$$\hat{D} \equiv 2|\psi\rangle\langle\psi| - \mathbb{1}$$

Reflection operator (or diffuser)

and consider following states in our Hilbert space:



where  $|\psi\rangle$  is state orthogonal to  $|\psi\rangle$  and so  $|\psi\rangle$  and  $|\psi\rangle$  project a hyperplane which means they are basis states for any other state on this plane.

For example state  $|\psi'\rangle$  can be expressed as:

$$|\psi'\rangle = \eta_1 |\psi\rangle + \eta_2 |\psi\rangle$$



Now, let's act with operator  $\hat{D}$  on state

$|\psi'\rangle$ :

$$\hat{D}|\psi'\rangle = \hat{D}(\eta_1|\psi\rangle + \eta_2|\varphi\rangle) =$$

$$= (2|\psi\rangle\langle\psi| - \mathbb{I})(\eta_1|\psi\rangle + \eta_2|\varphi\rangle) =$$

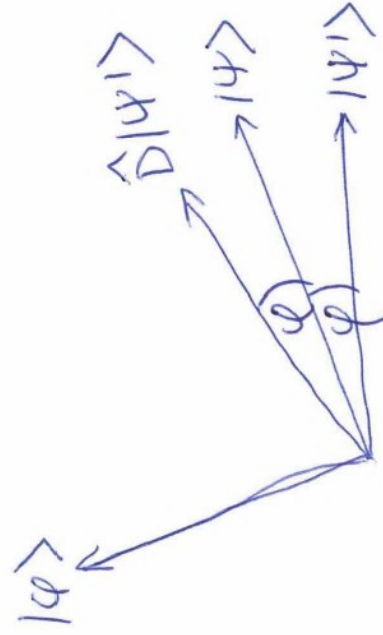
$$= 2\eta_1|\psi\rangle\langle\psi|\psi\rangle + \underbrace{2\eta_2|\psi\rangle\langle\psi|\varphi\rangle}_{=0 \text{ (based on definition of } |\varphi\rangle)} - \eta_1|\psi\rangle - \eta_2|\varphi\rangle$$

$$= 2\eta_1|\psi\rangle - \eta_1|\psi\rangle - \eta_2|\varphi\rangle =$$

$$= \eta_1|\psi\rangle - \eta_2|\varphi\rangle$$



$\hat{D}$  flips the " $|\varphi\rangle$  coordinate" of any state to opposite one:



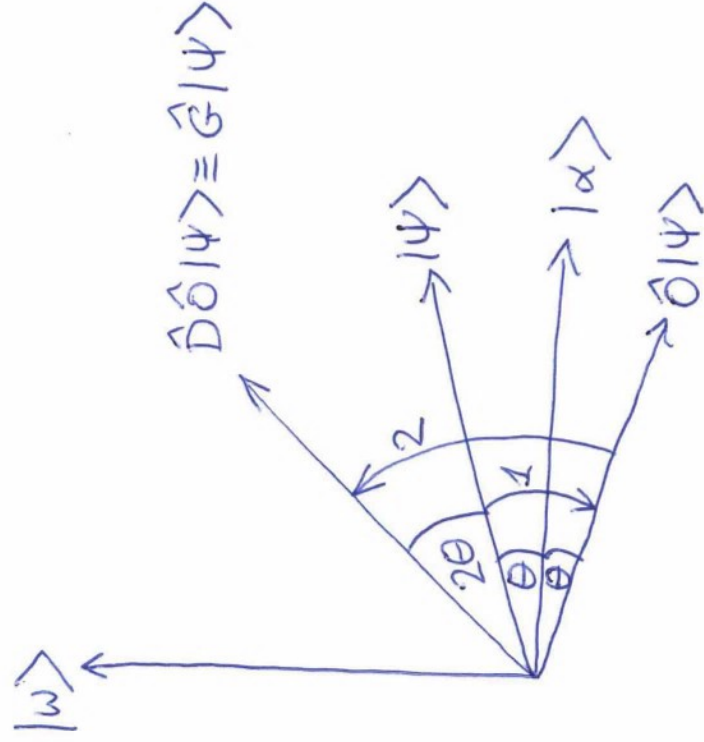
low, our final Grover operator is defined as:

$$\hat{G} \equiv \underbrace{(2|\psi\rangle\langle\psi| - \mathbb{I})}_{\hat{D}} \hat{O}$$

or

$$\hat{G} \equiv \hat{D} \hat{O}$$

and its effect can be graphically represented as:



- we only know  $|\psi\rangle$  at the beginning (superposition of all states - H gates on all qubits), and have operator  $\hat{O}$
- $\hat{D}$  is a constant operator, not depending on the problem



- having just this, we can perform operations ( $\hat{O}$  and  $\hat{D}$ , so effectively  $\hat{G}$ ) moving initial state  $| \psi \rangle$  to state  $\hat{G}| \psi \rangle$  which we know is closer to our searched state  $| w \rangle$
- being "closer" to  $| w \rangle$  means, that when we measure the state we will have bigger probability of getting  $| w \rangle$  as the result.

- now, we can perform Grover operation  $\hat{G}$  multiple times getting us closer and closer to  $| w \rangle$

How many times should we apply

Grover operator?

- one application of  $\hat{G}$  rotates our state by  $2\theta$  in direction of  $| w \rangle$
- from (5) and (6) we see, that

$$\sin \theta = \frac{1}{\sqrt{N}}$$

- for large  $N$ ,  $\frac{1}{\sqrt{N}}$  is small, which means, that  $\downarrow$



$$\sin \theta \approx \theta$$



$$\theta \approx \frac{1}{\sqrt{N}} / \cdot 2\sqrt{N}$$

$$2\theta \cdot \sqrt{N} \approx 2$$

let's assume, that we need to

rotate by  $\frac{\pi}{2}$   $\leftarrow$  as  $|4\rangle = \cos\theta|2\rangle + \sin\theta|w\rangle$   
 we want  $|4\rangle = |w\rangle$   
 $\cos\theta = 0 \Rightarrow \theta = \frac{\pi}{2}$

$$2\theta \cdot \sqrt{N} \approx 2 / \cdot \frac{\pi}{4}$$

$$\frac{\pi\sqrt{N}}{4} \cdot 2\theta \approx \frac{\pi}{2}$$

$\uparrow$  rotation in one iteration  
 $\uparrow$  total rotation angle  
 $\uparrow$  number of iterations



We need approximately

$\frac{\pi\sqrt{N}}{4}$  iterations

to align our initial state  $|4\rangle$  with  $|w\rangle$  which we look for.





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# Grover algorithm example on 3 qubits

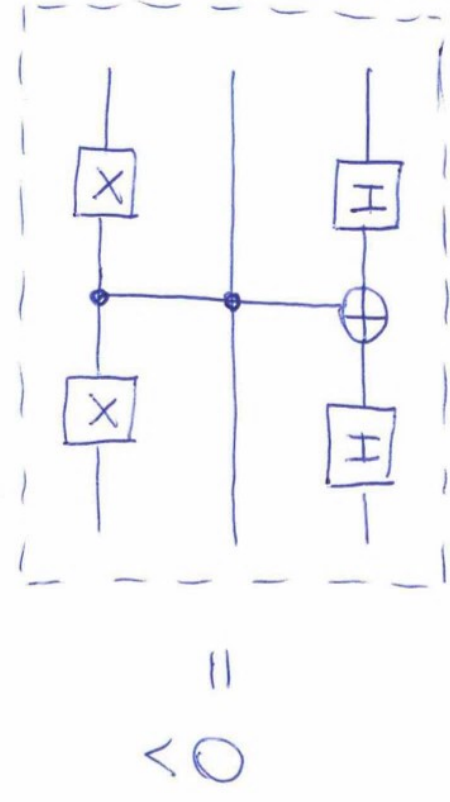
3 qubits can represent 8 states:

$|000\rangle, |001\rangle, |010\rangle, \dots, |111\rangle$

Let's choose for our winner  $|w\rangle$  the following state:

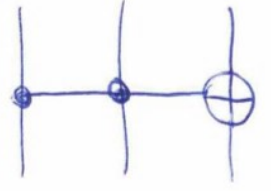
$$|w\rangle \equiv |110\rangle = |6\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 1 \\ 0 \\ 0 \end{bmatrix}$$

How to construct Oracle?



→ it can be proven experimentally:)

where



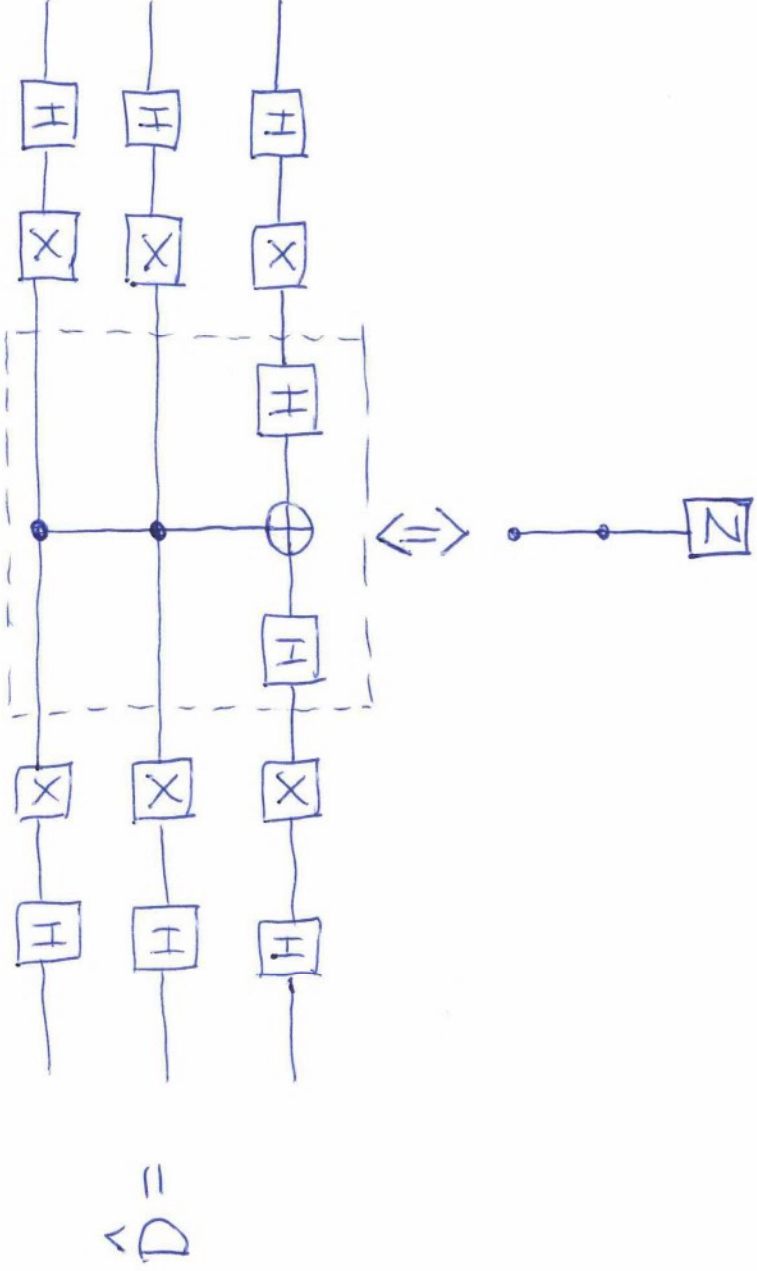
is so called Toffoli gate, which is a CNOT but with one additional control qubit





Toffoli gate flips last qubit only if first two are  $|1\rangle$

Now, how to construct diffuser?



• the diffuser is fixed, not dependent on the problem