Electricity and Magnetism



Overview Magnetism

- 28-5: Introduction, magnetism: field and force
- 1-6: Magnetism: Biot-Savart, Ampere
- 4-6: Electromagnetic induction
- 8-6: Electromagnetic induction
- 11-6: Maxwell's equations and electromagnetic waves
- 15-6: Local laws for the magnetostatic field, local laws for the electromagnetic field, magnetic field intensity H
- 18-6: available for answering questions, exercises



Equations of electromagnetism

- The four complete laws of electromagnetism are collectively called **Maxwell's equations**. They describe all electromagnetic fields in the universe, outside the realm of quantum physics.
- Gauss for the electric field E
 - Charge produces electric field
 - Field lines begin and end on charges or close

$$\oint \vec{E} \cdot d\vec{A} = \frac{q}{\varepsilon_0}$$

 $\oint \vec{B} \cdot d\vec{A} = 0$

- Gauss for the magnetic field B
 - No magnetic charges
 - Magnetic field lines close
- Faraday
 - Changing magnetic flux produces electric field

$$\oint \vec{E} \cdot d\vec{r} = -\frac{d\Phi_B}{dt}$$

- Ampere
 - Electric currents and changing electric flux produce magnetic field $\oint \vec{B} \cdot d\vec{r} = \mu_0 I + \mu_0 \varepsilon_0 \frac{d\Phi_E}{dt}$



Magnetostatic field: magnetic materials

- Learning objectives
- Introducing the magnetic field strength, the magnetic susceptibility, and the magnetization
- Magnetic materials
 - Diamagnetic/paramagnetic materials
 - Soft magnetic materials
 - Hard magnetic materials (permanent magnets)
- Application to toroid and solenoid
- Rewriting Maxwell's equations including dielectric and magnetic material properties
- Local form of Maxwell's equations
- Boundary conditions



Learning objectives

- Know
 - The magnetic field strength H,
 - The magnetic susceptibility χ_m ,
 - The relative magnetic permeability μ_{rm}
 - The magnetization M
- Know and separate between different magnetic materials
- Know the local form of Maxwell's equations
- Know the boundary conditions



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Dielectric material properties

$$\oint \vec{E} \cdot d\vec{A} = \frac{q}{\varepsilon_0} = \frac{q_f + q_b}{\varepsilon_0}$$

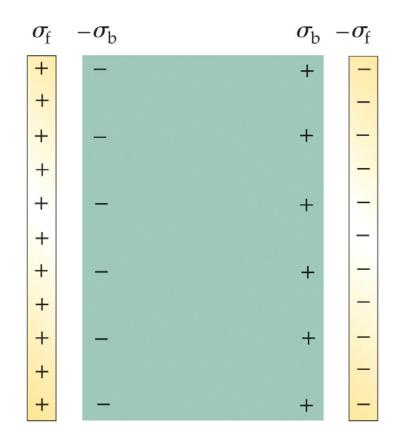
$$\oint \vec{D} \cdot d\vec{A} = q_f$$

$$\oint \vec{D} \cdot d\vec{A} = q_f$$

$$\vec{D} = \varepsilon_0 \vec{E} + \vec{P} = \varepsilon_0 \varepsilon_r \vec{E} = \varepsilon_0 \kappa \vec{E}$$

$$\vec{P} = \varepsilon_0(\varepsilon_r - 1)\vec{E} = \varepsilon_0(\kappa - 1)\vec{E} = \varepsilon_0 \chi_e \vec{E}$$

- D is electric flux density
- P is the polarisation
- χ_e is the electric susceptibility
- ε_0 is the permittivity in vacuum
- $\varepsilon_r = \kappa$ is the relative permittivity
- Electric flux density increases with polarisation
- Electric field strength E reduces due to polarisation
- Assuming linearity!!!



Magnetic material properties (static)

$$\oint \vec{B} \cdot d\vec{r} = \mu_0 I = \mu_0 (I_f + I_a)$$

$$\oint \vec{H} \cdot d\vec{r} = I_f$$

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$$

$$\vec{B} = \mu_0(\vec{H} + \vec{M}) = \mu_0(1 + \chi_m)\vec{H} = \mu_0^I \mu_r \vec{H}$$



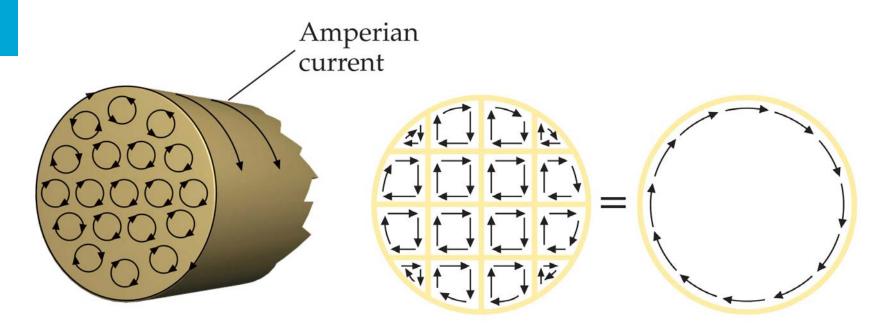
- M is magnetisation
- χ_m is the magnetic susceptibility
- μ_0 is the magnetic permeability in vacuum
- μ_r is the relative permeability
- Magnetic flux density B increases due to magnetisation
- Magnetic field strength H decreases due to magnetisation
- Assuming linearity!!!



- What is the unit of the magnetization M?

 - B. A/m
 - C. T/m²
 - D. Am²

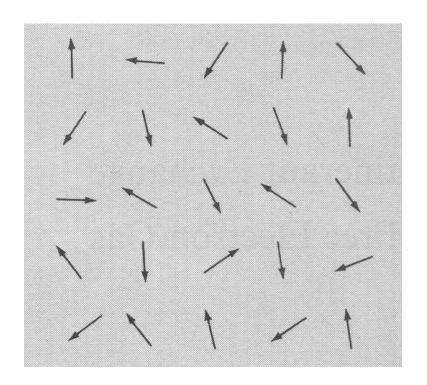
Amperian current

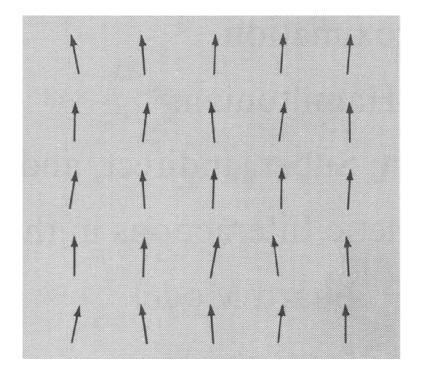


- Electrons circling around atoms
- Electron's intrinsic magnetic dipole moment



Magnetism in materials







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Magnetic susceptibility diamagnetic and paramagnetic materials

 Magnetisation often negligible

$$\vec{B} = \mu_0 (\vec{H} + \vec{M})$$

$$\vec{B} = \mu_0 (1 + \chi_m) \vec{H} = \mu_0 \mu_r \vec{H}$$

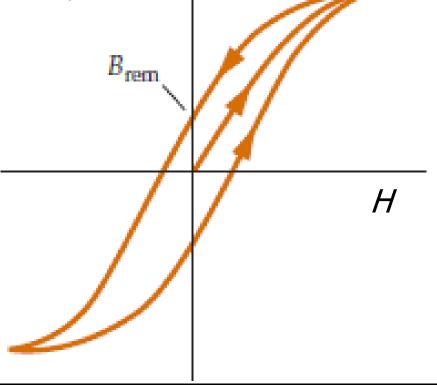
Material	Xm
Aluminium	2.3 · 10 ⁻⁵
Copper	-0.98 · 10-5
Gold	-3.6 · 10 ⁻⁵
Silver	-2.6 · 10 ⁻⁵
Titanium	7.06 · 10 ⁻⁵
Hydrogen (1atm)	-9.9 · 10 ⁻⁹
Carbon dioxide (1 atm)	-2.3 · 10 ⁻⁹
Nitrogen (1 atm)	-5 · 10 ⁻⁹
Oxygen (1 atm)	2090 · 10 ⁻⁹



Magnetic susceptibility soft ferromagnetic materials

$$\vec{B} = \mu_0(\vec{H} + \vec{M}) = \mu_0(1 + \chi_m)\vec{H} = \mu_0\mu_r\vec{H}$$

- Hysteresis characteristic:
 - Saturation
 - Losses



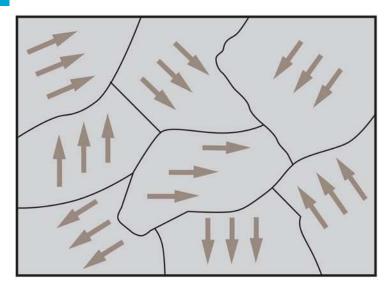
Magnetic susceptibility soft ferromagnetic materials

$$\vec{B} = \mu_0(\vec{H} + \vec{M}) = \mu_0(1 + \chi_m)\vec{H} = \mu_0\mu_r\vec{H}$$

Material	$\mu_0 M_s (T)$	$\chi_{\rm m} (\approx \mu_{\rm r})$
Iron (annealed)	2.16	5500
Iron-silicon (96% - 4%)	1.95	7000
Permalloy (FeNi)	1.6	25000
Mu-metal (NiFeCuCr)	0.65	100000

- Susceptibility can be very large
- Applied in transformers, motors, generators, coils

Magnetic domains in soft ferromagnetic materials



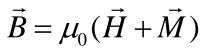
schematic example for domains netto resulting in M=0

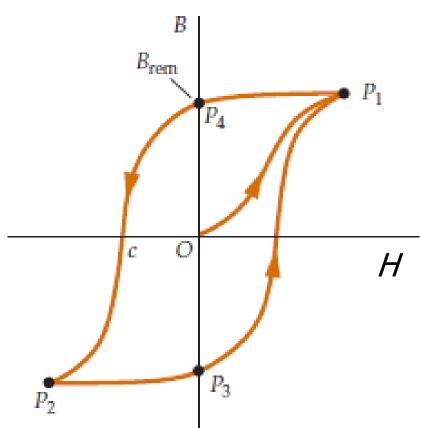


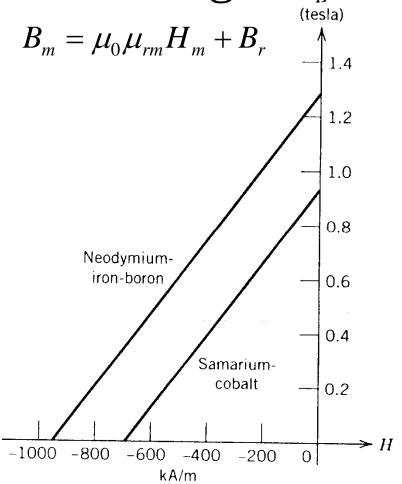
SEM picture of 4-colors of domain orientations in Fe_{0,97}Si_{0,03}



Magnetisation hard ferromagnetic materials: permanent magnet _B









Magnetisation hard ferromagnetic materials: permanent magnet

$$\vec{B} = \mu_0(\vec{H} + \vec{M})$$

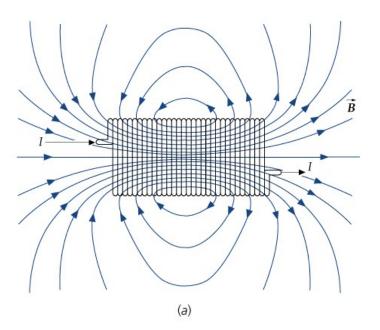
	$B_r(T)$	$H_c(A/m)$
Superpermalloy	0,8	Ca. 0,3
Chrome steel	1,0	4000
Ferrite	0,4	250000
Alnico 5	1,2	44000
Platinum cobalt	0,6	290000
SmCo	1,0	750000
NdFeB	1,4	1000000

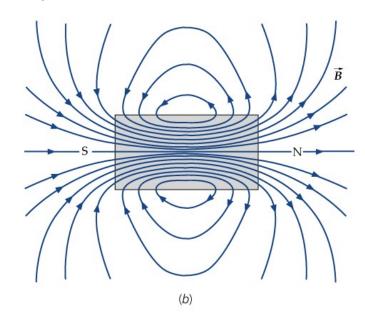


Magnetic field of a solenoid and a magnet

 If a magnet has a constant magnetisation, the external field is the same as for a solenoid.

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$$





- A strong permanent magnet has
 - A. A high saturation flux density and a narrow BH curve
 - A high remanent flux density and a wide BH curve
 - C. A low remanent flux density and a high coercive force



- The magnetic circuit of a transformer has
 - A. A high saturation flux density and a narrow BH curve
 - A high remanent flux density and a wide BH curve
 - C. A low remanent flux density and a high coercive force



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Toroid

$$\oint \vec{B} \cdot d\vec{r} = \mu_0 (I_f + I_a)$$

$$\oint \vec{H} \cdot d\vec{r} = I_{\it encircled}$$

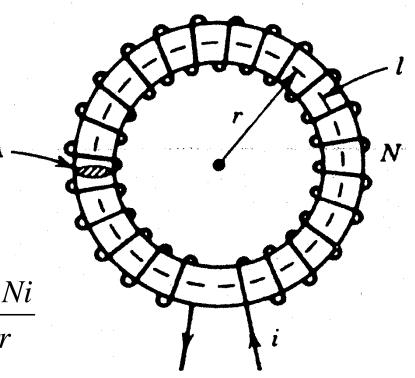
$$2\pi rH = Ni$$

$$B = \mu_0(H + M) = \mu_0 \mu_r H = \frac{\mu_0 \mu_r Ni}{2\pi r}$$

$$\Phi = \int \vec{B} \cdot d\vec{A} \approx NAB = \frac{\mu_0 \mu_r A N^2 I}{2\pi r}$$

$$L = \frac{\Phi}{I} = \frac{\mu_0 \mu_r A N^2}{2\pi r}$$

$$L = \frac{\Phi}{I} = \frac{\mu_0 \mu_r A N^2}{2\pi r}$$



Assumptions:

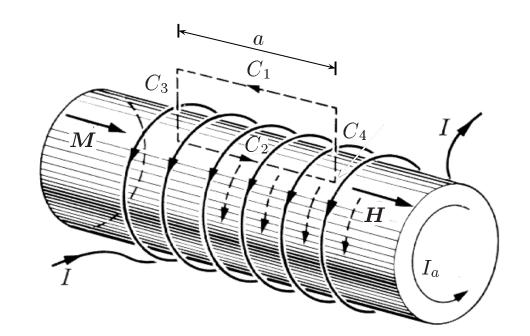
- **Symmetry**
- B constant over A

Solenoid

$$\oint \vec{B} \cdot d\vec{r} = \mu_0 (I_f + I_a)$$

$$\oint \vec{H} \cdot d\vec{r} = I_{\it encircled}$$

$$lH = Ni$$



$$B = \mu_0(H + M) = \mu_0 \mu_r H = \frac{\mu_0 \mu_r Ni}{l}$$

$$\Phi = \int \vec{B} \cdot d\vec{A} \approx NAB = \frac{\mu_0 \mu_r A N^2 I}{l}$$

$$L = \frac{\Phi}{I} = \frac{\mu_0 \mu_r A N^2}{l}$$

$$L = \frac{\Phi}{I} = \frac{\mu_0 \mu_r A N^2}{l}$$

Assumption:

- H outside negligible
- Problematic. Why?

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Equations of electromagnetism

- The four complete laws of electromagnetism are collectively called Maxwell's equations.
- Gauss for the electric field E
 - Charge produces electric field
 - Field lines begin and end on charges
- Gauss for the magnetic field B
 - No magnetic charges
 - Magnetic field lines close
- Faraday
 - Changing magnetic flux produces electric field

• Electric currents and changing electric flux produce magnetic field $\oint \vec{B} \cdot d\vec{r} = \mu_0 I + \mu_0 \varepsilon_0 \frac{d\Phi_E}{dt}$

$$\oint \vec{E} \cdot d\vec{A} = \frac{q}{\varepsilon_0}$$

$$\oint \vec{B} \cdot d\vec{A} = 0$$

$$\oint \vec{E} \cdot d\vec{r} = -\frac{d\Phi_B}{dt}$$



Equations of electromagnetism rewritten

- Gauss for the electric field E
 - Charge produces electric field
 - Field lines begin and end on charges
- Gauss for the magnetic field B
 - No magnetic charges
 - Magnetic field lines close
- Faraday
 - Changing magnetic flux produces electric field
- Ampere
 - Electric currents and changing electric flux produce magnetic field $\oint \vec{H} \cdot \vec{e}$
- Constitutive relations

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$$

$$\oint \vec{D} \cdot d\vec{A} = \int \rho_f dV$$

$$\oint \vec{B} \cdot d\vec{A} = 0$$

$$\oint \vec{E} \cdot d\vec{r} = -\frac{d}{dt} \int \vec{B} \cdot d\vec{A}$$

$$\oint \vec{H} \cdot d\vec{r} = \int \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int \vec{D} \cdot d\vec{A}$$

$$\vec{D} = \varepsilon_0 \vec{E} + \vec{P}$$

Maxwell's equations in local form

From integral form to local form using Gauss and Stokes

$$\oint \vec{D} \cdot d\vec{A} = \int \vec{\nabla} \cdot \vec{D} dV = \int \rho_f dV$$

$$\oint \vec{B} \cdot d\vec{A} = \int \vec{\nabla} \cdot \vec{B} dV = 0$$

$$\oint \vec{E} \cdot d\vec{r} = \int (\vec{\nabla} \times \vec{E}) \cdot d\vec{A} = -\frac{d}{dt} \int \vec{B} \cdot d\vec{A}$$

$$\oint \vec{H} \cdot d\vec{r} = \int (\vec{\nabla} \times \vec{H}) \cdot d\vec{A} = \int \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int \vec{D} \cdot d\vec{A} \qquad \vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$$

$$\vec{\nabla} \cdot \vec{D} = \rho_{\scriptscriptstyle f}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$ec{
abla} imes ec{H} = ec{J} + rac{\partial D}{\partial t}$$

$$\vec{D} = \varepsilon_0 \vec{E} + \vec{P}$$

Maxwell's equations in local form

From integral form to local form using Gauss and Stokes

$$\oint \vec{E} \cdot \hat{\tau} \, \mathrm{d} \, s = \iint (\vec{\nabla} \times \vec{E}) \cdot \hat{n} \, \mathrm{d} \, A = -\frac{\mathrm{d}}{\mathrm{d} \, t} \iint \vec{B} \cdot \hat{n} \, \mathrm{d} \, A \qquad \vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\oint \vec{H} \cdot \hat{\tau} \, \mathrm{d}s = \iint (\vec{\nabla} \times \vec{H}) \cdot \hat{n} \, \mathrm{d}A = \iint \vec{J} \cdot \hat{n} \, \mathrm{d}A + \frac{d}{dt} \iint \vec{D} \cdot \hat{n} \, \mathrm{d}A$$

$$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$
• The constitutive relations

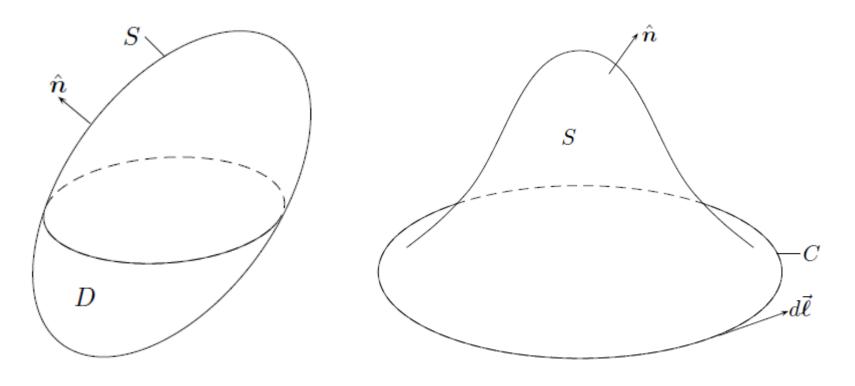
The constitutive relations

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$$

$$\vec{D} = \varepsilon_0 \vec{E} + \vec{P}$$

 $\vec{\nabla} \cdot \vec{B} = 0$

Closed and open surfaces



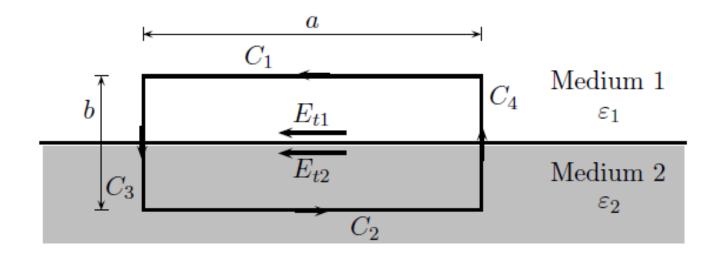


Magnetostatic field: magnetic materials

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Boundary conditions for the electric field: tangential

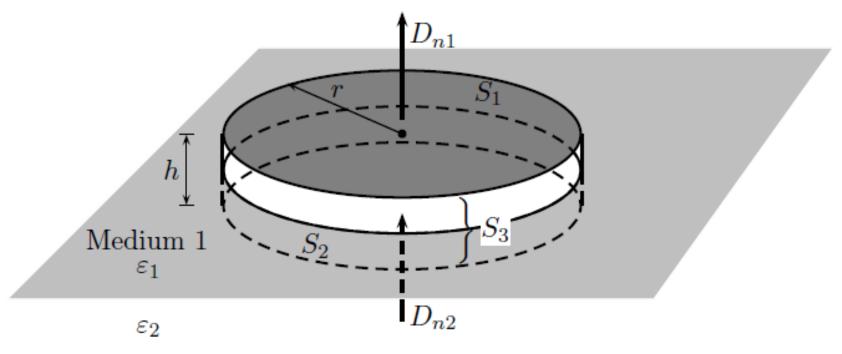


$$\oint \vec{E} \cdot d\vec{r} = -\frac{d}{dt} \int \vec{B} \cdot d\vec{A}$$

• If *b* goes to zero, Faraday says

$$\oint \vec{E} \cdot d\vec{r} = 0 \implies E_{t1} = E_{t2}$$

Boundary conditions for the electric field: normal

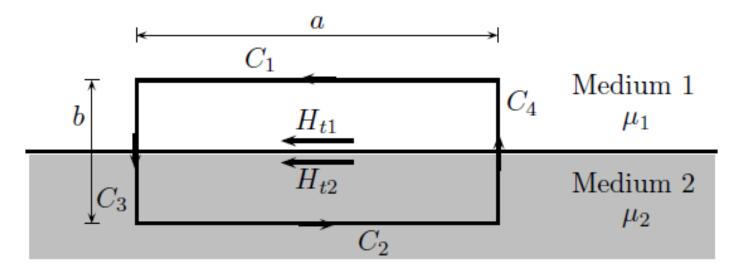


Medium 2

• If *h* goes to zero, Gauss for the electric displacement says

$$\oint \vec{D} \cdot d\vec{A} = Q_f \quad \Rightarrow \quad D_{n1} - D_{n2} = \sigma_f$$

Boundary conditions for the magnetic field: tangential



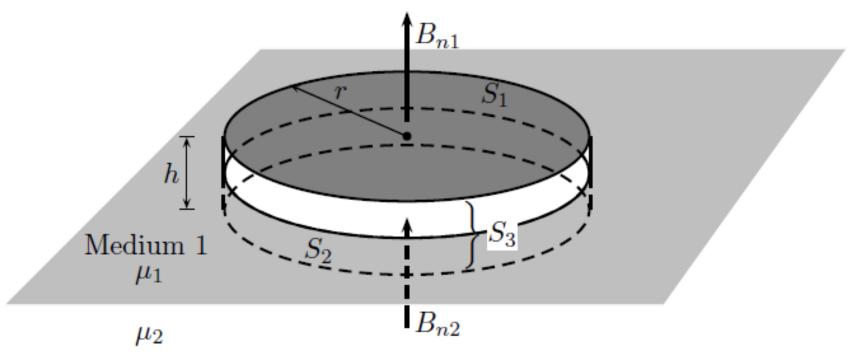
$$\oint \vec{H} \cdot d\vec{r} = \int \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int \vec{D} \cdot d\vec{A}$$

• If *b* goes to zero, Ampere says

$$\oint \vec{H} \cdot d\vec{r} = 0 \implies H_{t1} = H_{t2}$$



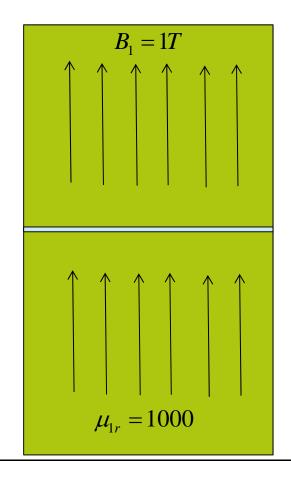
Boundary conditions for the electric field: normal



Medium 2

• If h goes to zero, Gauss for the magnetic flux density says

$$\oint \vec{B} \cdot d\vec{A} = 0 \implies B_{n1} = B_{n2}$$

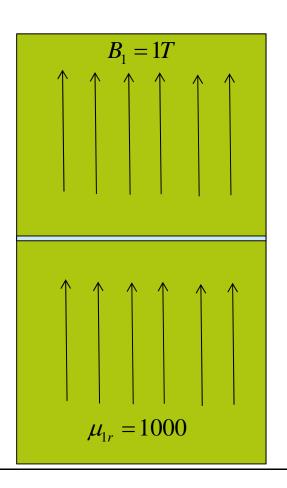


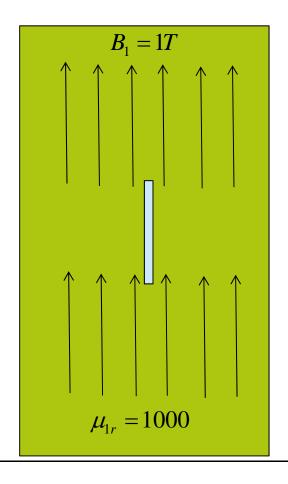
$$\oint \vec{B} \cdot d\vec{A} = 0 \implies B_{n1} = B_{n2} = 1T$$

$$H_n = \frac{B_n}{\mu_0 \mu_r}$$

$$H_{ng} = \frac{B_{ng}}{\mu_0} = 796 \text{ kA/m}$$

$$H_{nFe} = \frac{B_{nFe}}{\mu_0 \mu_r} = 796 \text{ A/m}$$

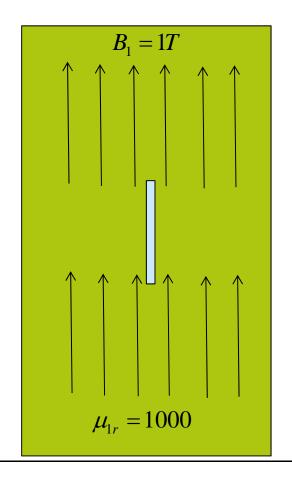




$$\oint \vec{H} \cdot d\vec{r} = 0 \implies H_{t1} = H_{t2}$$

$$H_{tFe} = \frac{B_{tFe}}{\mu_0 \mu_r} = 796 \,\text{A/m}$$

$$B_{tg} = \mu_0 H_{tg} = \mu_0 H_{tFe} = 1 \text{mT}$$



Maxwell's equations and materials

From integral form to local form using Gauss and Stokes

$$\begin{split} &\oint \vec{D} \cdot d\vec{A} = \int \vec{\nabla} \cdot \vec{D} dV = \int \rho_f dV \\ &\oint \vec{B} \cdot d\vec{A} = \int \vec{\nabla} \cdot \vec{B} dV = 0 \end{split}$$

$$\oint \vec{E} \cdot d\vec{r} = \int (\vec{\nabla} \times \vec{E}) \cdot d\vec{A} = -\frac{d}{dt} \int \vec{B} \cdot d\vec{A}$$

$$\oint \vec{H} \cdot d\vec{r} = \int (\vec{\nabla} \times \vec{H}) \cdot d\vec{A} = \int \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int \vec{D} \cdot d\vec{A} \qquad \vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

$$\oint \vec{H} \cdot d\vec{r} = 0 \implies H_{t1} = H_{t2}$$

$$\oint \vec{B} \cdot d\vec{A} = 0 \implies B_{n1} = B_{n2}$$

$$\vec{\nabla} \cdot \vec{D} = \rho_f$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial D}{\partial t}$$

$$\vec{B} = \mu_0 \vec{H} + \mu_0 \vec{M}$$
 $\vec{D} = \varepsilon_0 \vec{E} + \vec{P}$

$$\oint \vec{E} \cdot d\vec{r} = 0 \implies E_{t1} = E_{t2}$$

$$\oint \vec{D} \cdot d\vec{A} = Q_f \quad \Rightarrow \quad D_{n1} - D_{n2} = \sigma_f$$

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- 1-6: Magnetism: Biot-Savart, Ampere
- 4-6: Electromagnetic induction
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- 11-6: Maxwell's equations and electromagnetic waves
- 15-6: Local laws for the magnetostatic field, local laws for the electromagnetic field, magnetic field intensity H
- 18-6: available for answering questions, exercises



Summary magnetism force and field 26

- Magnetism involves moving electric charge.
- Magnetic fields exert forces on moving electric charges:

• For a moving charge:
$$\vec{F} = q\vec{v} \times \vec{B}$$

• For a current:
$$\vec{F} = I\vec{L} \times \vec{B}$$

Magnetic fields arise from moving electric charge, as described by

• Biot-Savart law:
$$\vec{R} - \int d\vec{R} = \int \mu_0$$

Biot-Savart law:
$$\vec{B} = \int d\vec{B} = \int \frac{\mu_0}{4\pi} \frac{I \, dL \times \hat{r}}{r^2}$$

• Ampère's law:
$$\oint \vec{B} \cdot d\vec{r} = \mu_0 I_{encircled}$$

- Magnetic fields encircle the currents and moving charges that are their sources.
 - Magnetic field lines don't begin or end.
 - This is expressed in Gauss's law for magnetism:

$$\oint \vec{B} \cdot d\vec{A} = 0$$



Summary electromagnetic induction 27

 Faraday's law describes electromagnetic induction, most fundamentally the phenomenon whereby a changing magnetic field produces an electric field:

$$\oint \vec{E} \cdot d\vec{r} = -\frac{d\Phi_B}{dt}$$

- This induced electric field is nonconservative and its field lines have no beginnings or endings.
- In the presence of a circuit, the induced electric field gives rise to an induced emf and an induced current.
 - Lenz's law states that the direction of the induced current is such that the magnetic field it produces acts to *oppose* the change that gives rise to it.
 - Self-inductance is a circuit property whereby changing current in a circuit results in an induced emf that opposes the change.
- Consideration of current buildup in an inductor shows that all magnetic fields store energy, with energy density $B^2/2\mu_0$.



Equations of electromagnetism 29

- The four complete laws of electromagnetism are collectively called **Maxwell's equations**. They describe all electromagnetic fields in the universe, outside the realm of quantum physics.
- Gauss for the electric field E
 - Charge produces electric field
 - Field lines begin and end on charges or close

$$\oint \vec{E} \cdot d\vec{A} = \frac{q}{\varepsilon_0}$$

 $\oint \vec{B} \cdot d\vec{A} = 0$

- Gauss for the magnetic field B
 - No magnetic charges
 - Magnetic field lines close
- Faraday
 - Changing magnetic flux produces electric field

$$\oint \vec{E} \cdot d\vec{r} = -\frac{d\Phi_B}{dt}$$

- Ampere
 - Electric currents and changing electric flux $\oint \vec{B} \cdot d\vec{r} = \mu_0 I + \mu_0 \varepsilon_0 \frac{d\Phi_E}{dt}$ produce magnetic field



Maxwell's equations and materials

From integral form to local form using Gauss and Stokes

$$\begin{split} \oint \vec{D} \cdot d\vec{A} &= \int \vec{\nabla} \cdot \vec{D} dV = \int \rho_f dV \\ \oint \vec{B} \cdot d\vec{A} &= \int \vec{\nabla} \cdot \vec{B} dV = 0 \\ \end{aligned} \qquad \qquad \vec{\nabla} \cdot \vec{D} = \rho_f$$

$$\oint \vec{E} \cdot d\vec{r} = \int (\vec{\nabla} \times \vec{E}) \cdot d\vec{A} = -\frac{d}{dt} \int \vec{B} \cdot d\vec{A}$$

$$\oint \vec{H} \cdot d\vec{r} = \int (\vec{\nabla} \times \vec{H}) \cdot d\vec{A} = \int \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int \vec{D} \cdot d\vec{A} \qquad \vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

$$\frac{1}{dt} \int D \cdot dA \qquad \forall \times H = J + \frac{1}{\partial t}$$

$$\oint \vec{H} \cdot d\vec{r} = 0 \implies H_{t1} = H_{t2}$$

$$\oint \vec{B} \cdot d\vec{A} = 0 \implies B_{n1} = B_{n2}$$

$$ec{B} = \mu_0 ec{H} + \mu_0 ec{M}$$
 $ec{D} = arepsilon_0 ec{E} + ec{P}$

$$\oint \vec{E} \cdot d\vec{r} = 0 \implies E_{t1} = E_{t2}$$

 $\vec{\nabla} \times \vec{E} = -\frac{\partial B}{\partial x}$

$$\oint \vec{D} \cdot d\vec{A} = Q_f \quad \Rightarrow \quad D_{n1} - D_{n2} = \sigma_f$$