

# A nodal-based high-order nonlinear viscosity finite element method for MHD

Tuan Anh Dao

<tuuanh.dao@it.uu.se>

Murtazo Nazarov

<murtazo.nazarov@it.uu.se>

Division of Scientific Computing  
Department of Information Technology  
Uppsala University, Sweden



Vetenskapsrådet



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# Outline

- ① Parameter-free first-order method for scalar conservation laws
- ② Extension to systems and higher-order methods
- ③ Application to the system of MHD equations
- ④ Summary and outlook

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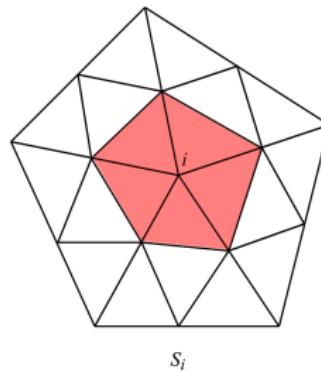
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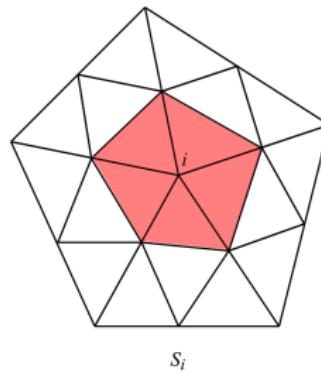
# First-order PP method for $\mathbb{P}_1$



At node  $i$ , define

- $S_i$  – support of the Lagrange basis function  $\varphi_i$

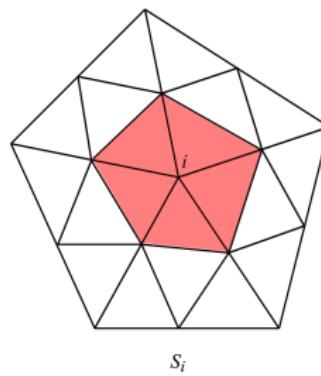
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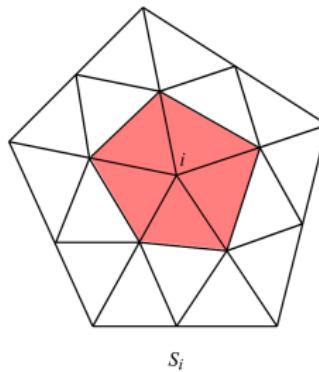
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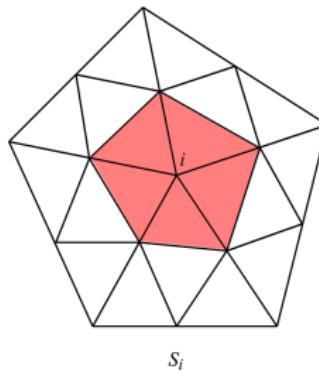
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- $m_i$  – weight of the lumped  $\mathbb{P}_1$  mass matrix at index  $i$

# First-order PP method for $\mathbb{P}_1$

- $C_i := \frac{1}{N_{\text{el}}(S_i)} \max_{K \in S_i, j \neq i, j \in \mathcal{I}(S_i)} \left| \int_K (\mathbb{J}_K^\top \nabla \varphi_j) \cdot (\mathbb{J}_K^\top \nabla \varphi_i) \, d\mathbf{x} \right|^{-1}$

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## Theorem

Define the nodal viscosity as

$$\varepsilon_i := \alpha_i C_i \|f'\|_{L^\infty(S_i)} \max_{i \neq j \in \mathcal{I}(S_i)} |\nabla \varphi_j| m_i.$$

Under a usual CFL condition, the approximation

$$m_i \frac{u_i^{n+1} - u_i^n}{\tau} + \int_{S_i} \nabla \cdot f(u_h^n) \varphi_i \, d\mathbf{x} + \sum_{K \in S_i} \sum_j \varepsilon_j u_j^n b_K(\varphi_j, \varphi_i) = 0$$

is positivity preserving. The form  $b_K(\cdot, \cdot)$  is defined next.

## Alternative $\mathbb{P}_1$ viscous bilinear form in 2D

For each triangle element  $K$ , define  $\Phi_K := \hat{K} \rightarrow K$ , where  $\hat{K}$  is the equilateral triangle such that  $|\hat{K}| = |K|$ .  $\mathbb{J}_K$  is the Jacobian matrix of  $\Phi_K$ .

**Tensor-valued viscosity** ([Guermond & Nazarov, 2014], edge-wise)

$$b(u, v) := \int_{\Omega} (\mathbb{J}^{\top} \nabla u) \cdot (\mathbb{J}^{\top} \nabla v) \, dx, \quad \forall u, v \in V_h.$$

(2D) Local bilinear form, for  $i, j \in \mathcal{I}(K)$ ,

$$b_K(\varphi_j, \varphi_i) := \begin{cases} -\frac{1}{n_K-1} & \text{if } j \neq i, \\ 1 & \text{if } j = i, \end{cases}$$

where  $n_K$  is the number of nodes in  $K$ .

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# Extensions

- High-order viscosity methods: Residual-based, Entropy-based

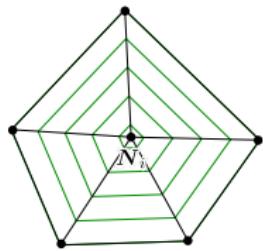
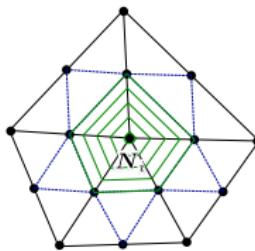
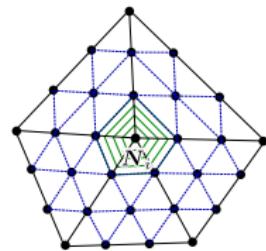
$$\varepsilon_{h,i}^H := C_i \alpha_i \min \left( \|f'\|_{L^\infty(S_i)}, \max_{i \neq j \in \mathcal{I}(S_i)} |\nabla \varphi_j|, |R_i(q_h)| \right) m_i^{(\mathbb{P}_1)}$$

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- Other polynomial spaces: Calculate  $\epsilon_i$  on  $\mathbb{P}_1$  sub-mesh

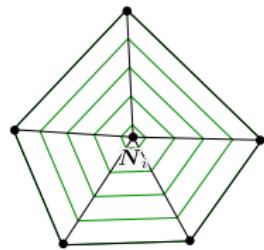
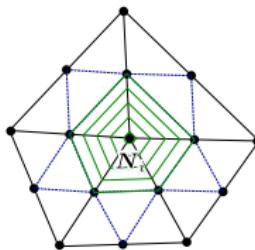
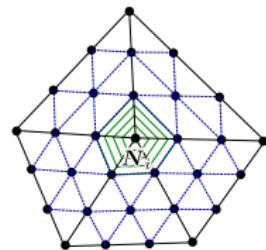
(a)  $\mathbb{P}_1$ (b)  $\mathbb{P}_2$ (c)  $\mathbb{P}_3$

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- Systems: take maximum of the components

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- **The viscosity term does not affect time step limit**

# Convergence test on Burger's equation with shocks

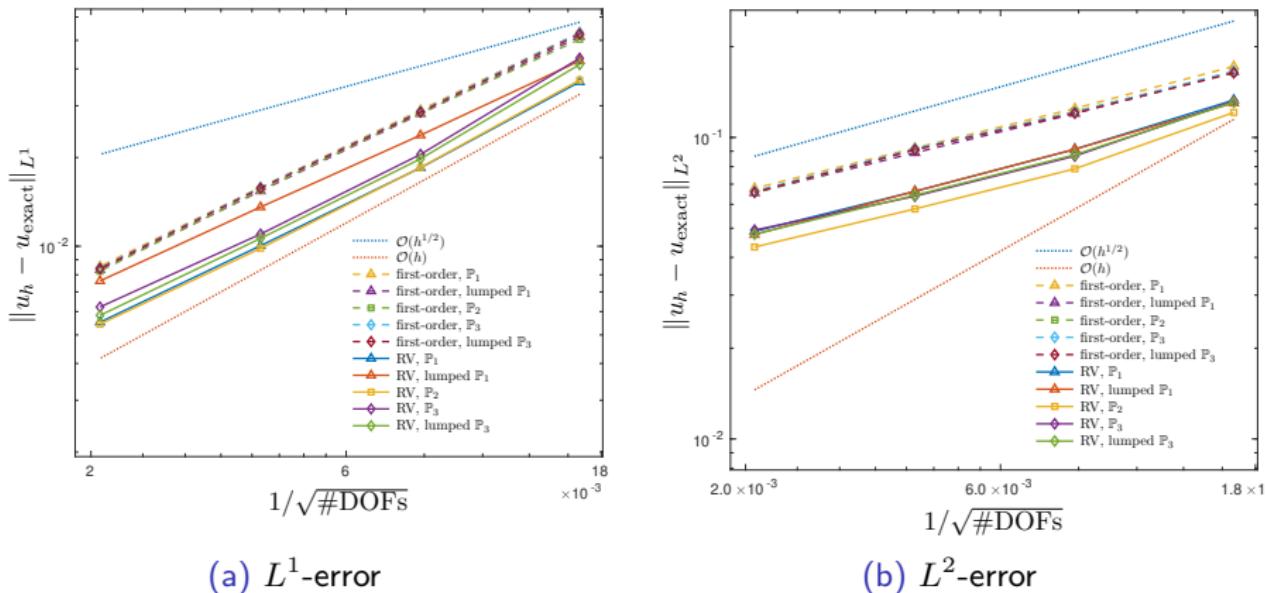
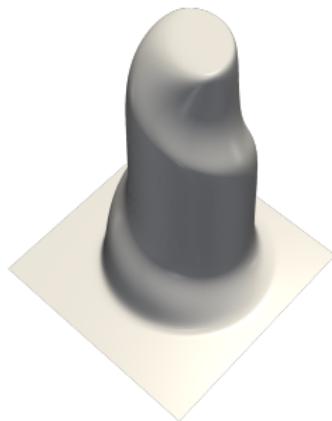
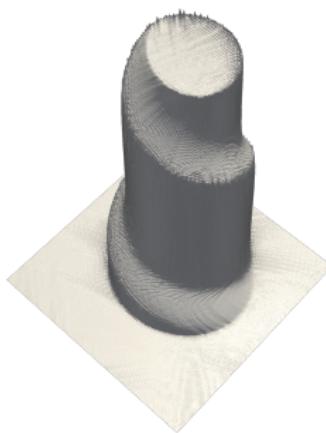


Figure: First-order  $\mathbb{P}_1$ ,  $\mathbb{P}_2$ ,  $\mathbb{P}_3$  shock solutions to the Burger's equation

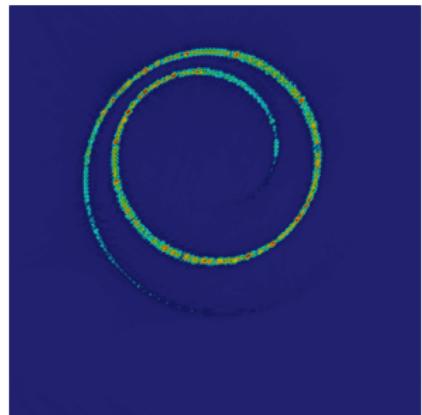
# The KPP problem



(a) First order solution



(b) RV solution

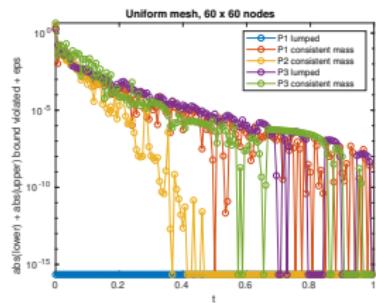
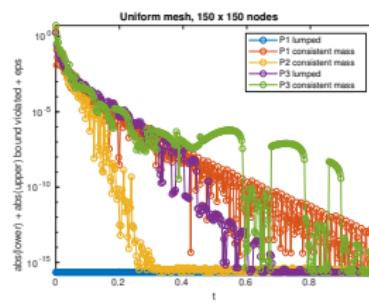
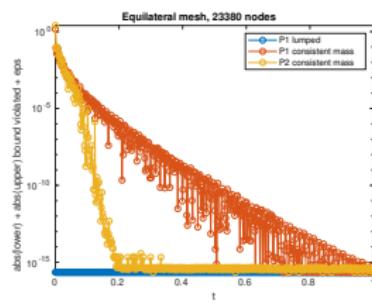


(c) RV viscosity

Figure: KPP solution with  $300 \times 300 \mathbb{P}_3$  nodes

# The KPP problem

## First order solutions

(a)  $60 \times 60$  mesh(b)  $150 \times 150$  mesh

(c) Equilateral mesh

Figure: Discrete maximum principle violation in higher order polynomials.

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# The system of MHD equations

$\mathbf{U} := (\rho(\mathbf{x}, t), \mathbf{m}(\mathbf{x}, t), E(\mathbf{x}, t), \mathbf{B}(\mathbf{x}, t))^{\top}$ ,  $(\mathbf{x}, t) \in \mathbb{R}^d \times \mathbb{R}^+$   
 Density, momentum, total energy, magnetic field

## The ideal MHD equations

$$\partial_t \mathbf{U} + \nabla \cdot (\mathbf{F}_{\mathcal{E}}(\mathbf{U}) + \mathbf{F}_{\mathcal{B}}(\mathbf{U})) = 0, \quad (1)$$

$$\mathbf{F}_{\mathcal{E}} := \begin{pmatrix} \mathbf{m} \\ \mathbf{m} \otimes \mathbf{u} + p\mathbb{I} \\ \mathbf{u}(E + p) \\ 0 \end{pmatrix}, \quad \mathbf{F}_{\mathcal{B}} := \begin{pmatrix} 0 \\ -\boldsymbol{\beta} \\ -\boldsymbol{\beta}\mathbf{u} \\ \mathbf{u} \otimes \mathbf{B} - \mathbf{B} \otimes \mathbf{u} \end{pmatrix},$$

$\boldsymbol{\beta}$  is the Maxwell stress tensor:

$$\boldsymbol{\beta} = \left( -\frac{1}{2}(\mathbf{B} \cdot \mathbf{B})\mathbb{I} + \mathbf{B} \otimes \mathbf{B} \right).$$

# Monolithic parabolic regularization to the MHD equations

Regularize (1) by adding  $\nabla \cdot (\epsilon \nabla \mathbf{U})$ ,

$$\partial_t \mathbf{U} + \nabla \cdot \mathbf{F}_{\mathcal{E}}(\mathbf{U}) + \nabla \cdot \mathbf{F}_{\mathcal{B}}(\mathbf{U}) = \nabla \cdot (\epsilon \nabla \mathbf{U}), \quad (2)$$

where  $\epsilon$  is a small artificial viscosity parameter.

The regularized equation (2) is

- a continuous analogue of the well-known upwind, Lax-Friedrichs schemes;
- suitable to apply the new viscosity method.

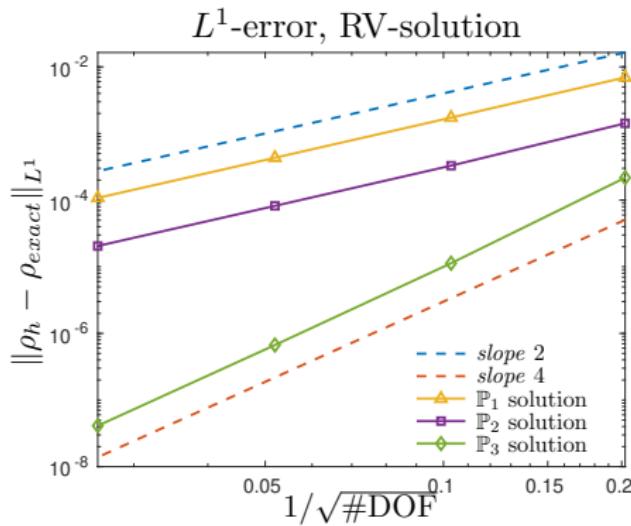
At the PDE level, we have shown that (2)

- is compatible with all the generalized entropies [Harten, 1998];
- fulfills the minimum entropy principles;
- preserves  $\rho > 0, p > 0$ .

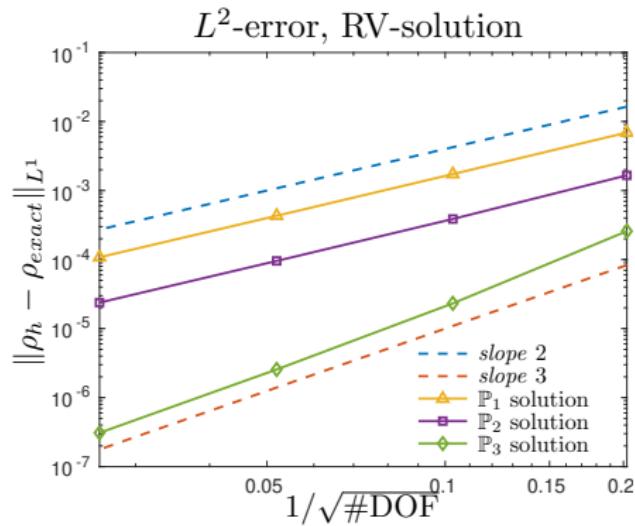
(See [D. & Nazarov, CMAME, 2022])

# Convergence test on a smooth problem

Smooth wave propagation, [Wu and Shu, 2018].



(a)  $L^1$ -error



(b)  $L^2$ -error

Figure:  $\mathbb{P}_1$ ,  $\mathbb{P}_2$ ,  $\mathbb{P}_3$  convergence for smooth solutions.

# The Orszag-Tang problem [Orszag & Tang, 1998]

Domain  $\Omega = [0, 1] \times [0, 1]$ . Gas constant  $\gamma = \frac{5}{3}$ .

Initial solution:

$$(\rho_0, \mathbf{u}_0, p_0, \mathbf{B}_0) = \\ \left( \frac{25}{36\pi}, (-\sin(2\pi y), \sin(2\pi x)), \frac{5}{12\pi}, \left( -\frac{\sin(2\pi y)}{\sqrt{4\pi}}, \frac{\sin(4\pi x)}{\sqrt{4\pi}} \right) \right)$$

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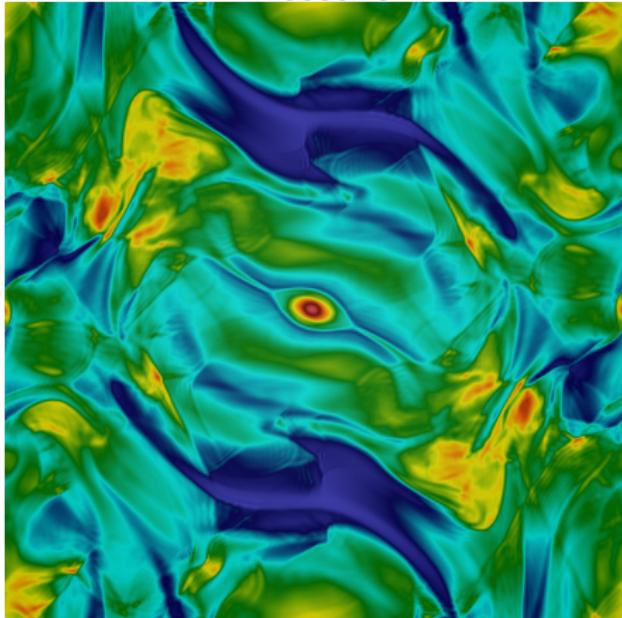
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- at  $t = 1.0$  considered transition to turbulence

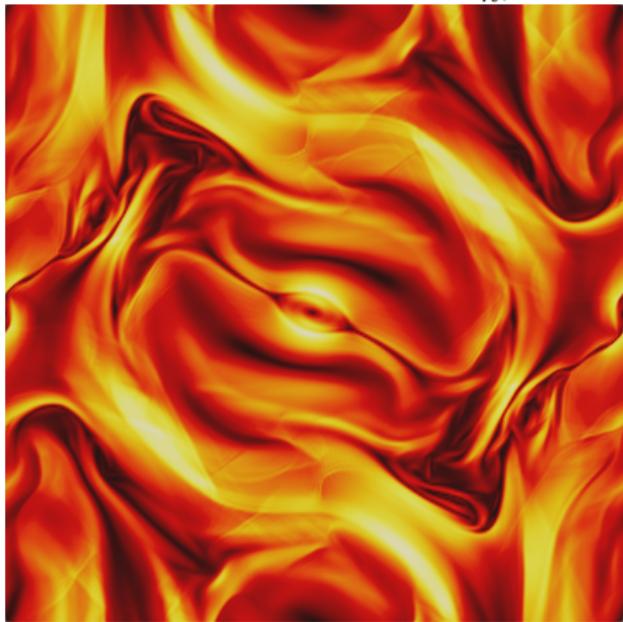
# Orszag-Tang problem [Orszag & Tang, 1998]

$\mathbb{P}_3$  RV solution,  $100 \times 100$  quasi-uniform mesh

Pressure



Magnetic pressure  $B_h^2/2$



# The MHD Rotor problem [Balsara & Spicer, 1998]

Domain  $\Omega = [0, 1] \times [0, 1]$ . Gas constant  $\gamma = 1.4$ .

Initial solution:

$$p_0 = 1, \mathbf{B}_0 = \left( \frac{5}{\sqrt{4\pi}}, 0 \right)$$

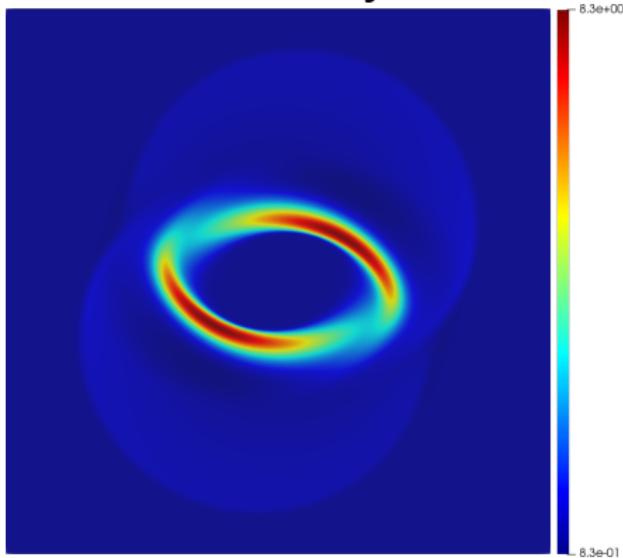
$$(\rho_0, \mathbf{u}_0) = \begin{cases} \left( 10, \left( \frac{2}{r_0} \left( \frac{1}{2} - y \right), \frac{2}{r_0} \left( x - \frac{1}{2} \right) \right)^\top \right) & \text{if } r < r_0, \\ \left( 1 + 9f, \left( f \frac{2}{r} \left( \frac{1}{2} - y \right), f \frac{2}{r} \left( x - \frac{1}{2} \right) \right)^\top \right) & \text{if } r_0 \leq r < r_1, \\ (1, (0, 0)^\top) & \text{otherwise,} \end{cases}$$

where  $r := \sqrt{(x - 0.5)^2 + (y - 0.5)^2}$ .

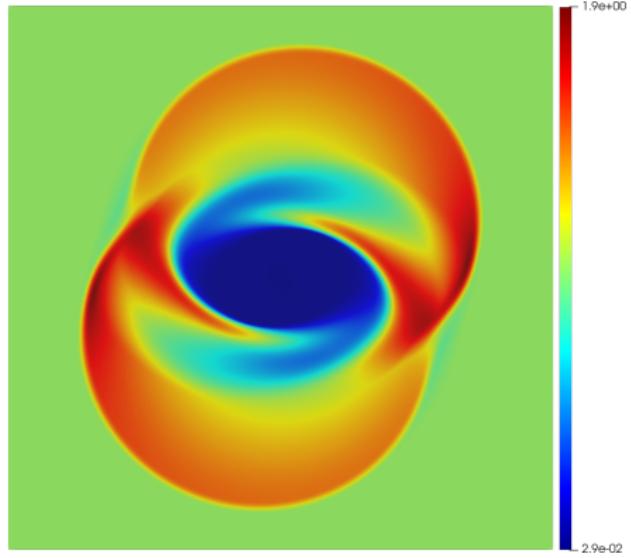
# The MHD Rotor problem [Balsara & Spicer, 1998]

$\mathbb{P}_3$  RV solution,  $100 \times 100$  quasi-uniform mesh

**Density**



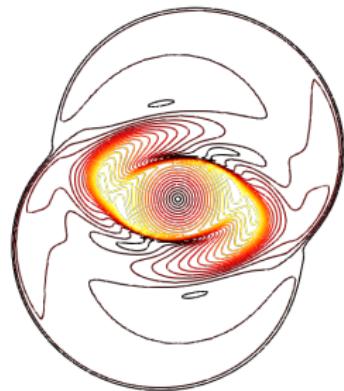
**Pressure**



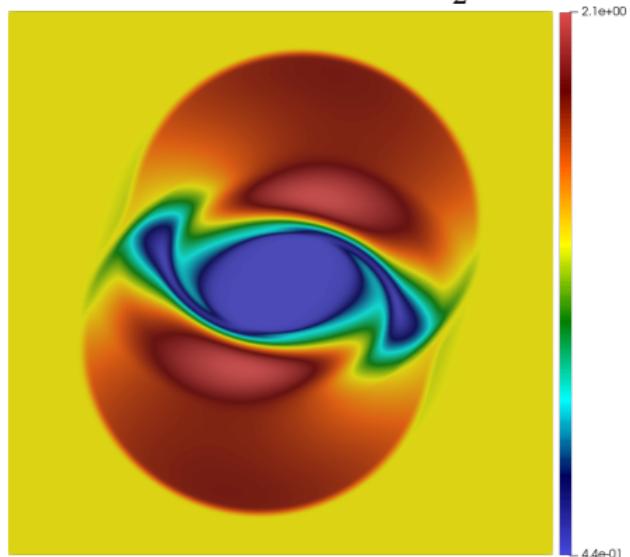
# The MHD Rotor problem [Balsara & Spicer, 1998]

$\mathbb{P}_3$  RV solution,  $100 \times 100$  quasi-uniform mesh

**Mach number**



**Magnetic pressure  $\frac{1}{2} B^2$**



# Astrophysical jets [Wu & Shu, 2018]

Domain  $\Omega = [-0.75, 0.75] \times [0, 1.5]$ . Gas constant  $\gamma = 1.4$ .

Initial solution: ambient

$$p_0 = 1, \mathbf{B}_0 = \left(0, \sqrt{200}\right)^\top, \rho_0 = 0.1\gamma, \mathbf{u}_0 = (0, 0)^\top$$

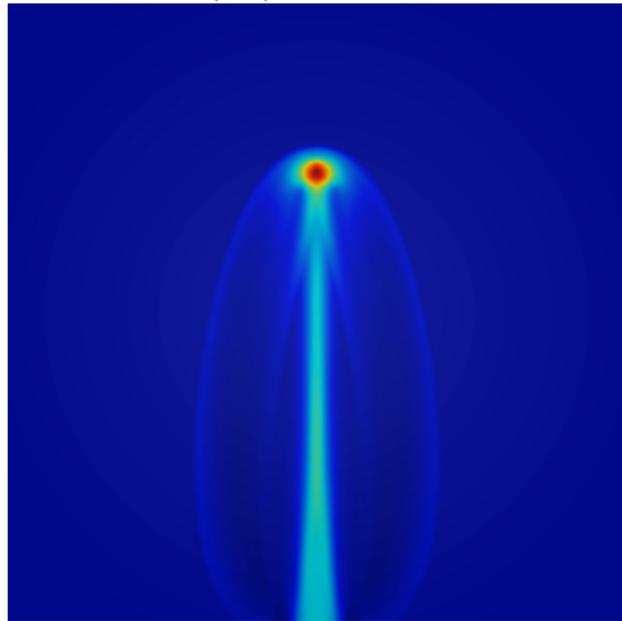
On the inlet  $(x, y) \in [-0.05, 0.05] \times 0$ : Mach 800

$$\mathbf{u}_0 = (0, 800)^\top, \rho = \gamma.$$

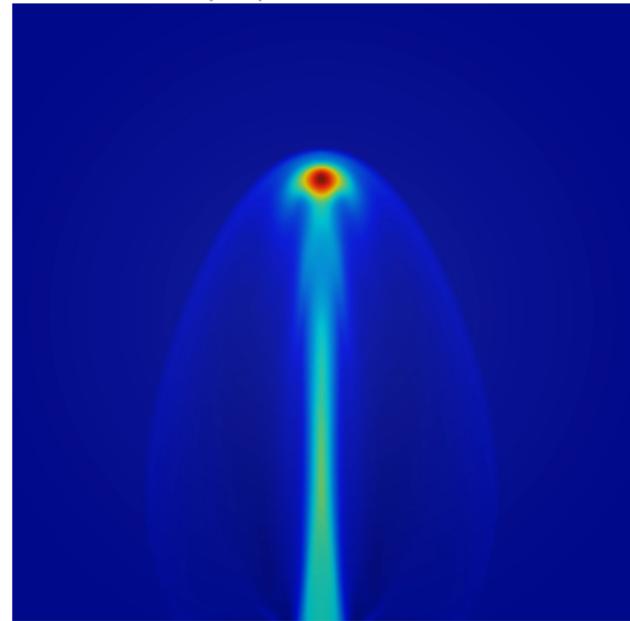
# Astrophysical jets [Wu & Shu, 2018]

RV solution with 371447  $\mathbb{P}_1$  nodes, lumped mass

$$|B| = \sqrt{200}$$



$$|B| = \sqrt{2000}$$



# Outline

1 Parameter-free first-order method for scalar conservation laws

2 Extension to systems and higher-order methods

3 Application to the system of MHD equations

4 Summary and outlook

# Summary

- PP parameter-free first-order FE methods for scalar equations.
- Extensions for systems, higher order methods.
- Tested against several MHD benchmarks.

# Outlook

Future works:

- Invariant-domain properties at the discrete level

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- Invariant-domain properties at the discrete level
- Maximum wave speed for MHD, exact or reasonable upper bound
- Conflicting objectives: conservativeness, divergence, positivity, entropy stability
- High-order invariant-domain preserving FE methods for MHD

Thank you for your listening!