



# Anyons

## Quantum particles with fractional statistics

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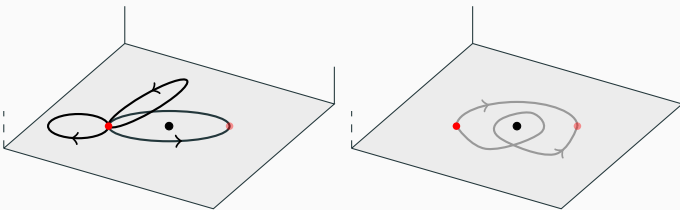
# Bosons and Fermions

$M_N^{(d)}$  : configuration space of  $N$  particles in  $d$  spatial dimensions.

$$M_2^{(d)} = (\mathbb{R}^d \setminus \{0\}) / S_2$$

In  $d = 3$  its **fundamental group** is  $S_2$ . Be  $\theta$  the **statistical angle**:

$$e^{i\theta} = \pm 1 \implies \theta = 0 \text{ (boson)}, \pi \text{ (fermion)}$$

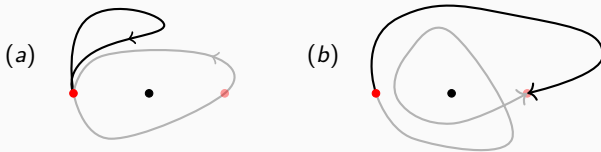


A particle remains fixed at  $\bullet$ , whereas the other moves. Antipodal points with respect to  $\bullet$ , like  $\bullet$  and  $\circ$ , are identified. There are two homotopy classes.

# Anyons and fractional statistics

In  $d = 2$  the fundamental group is the **braid group**  $B_2$ . **Anyons** appear, with **fractional statistical angles**:

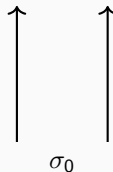
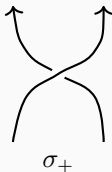
$$\theta \in [0, 2\pi)$$



For  $d = 2$ , **winding** is well defined, thus homotopy classes are infinite in number.

(a)  $\sigma_0, \sigma_+^2$

(b)  $\sigma_-, \sigma_+^3$



# Topological phases and flux tube model

If the configurational space is not simply connected,  $\psi$  is **polydromous**.

$$\psi \xrightarrow{\gamma(t) \in g} a_g(t) \psi \quad \text{with} \quad a_g(t) = \exp \left( \frac{i}{\pi} \sum_{i < j} \theta_{ij} \int_{\gamma} \int_0^t dt' \dot{\varphi}_{ij}(t') \right)$$

$\theta_{ij}$ ,  $\varphi_{ij}$  being the  $(i, j)$ -couple's and **statistical** and **azimuthal angles**.

$$K(\mathbf{x}, t; \mathbf{x}_0, 0) = \sum_{g \in \pi_1} a_g(t) \overbrace{K_g(\mathbf{x}, t; \mathbf{x}_0, 0)}^{\text{partial propagators}} = \int_{(\mathbf{x}_0, 0)}^{(\mathbf{x}, t)} \mathcal{D}[\gamma(t)] e^{\frac{i}{\hbar} S}$$

$$L(\mathbf{x}, \dot{\mathbf{x}}, t) = \underbrace{L_0(\mathbf{x}, \dot{\mathbf{x}}, t)}_{\text{proper term}} + \underbrace{\frac{\hbar}{\pi} \sum_{i < j} \theta_{ij} \dot{\varphi}_{ij}(t)}_{L_{int} \text{ (topological term)}} \quad \text{effective lagrangian}$$

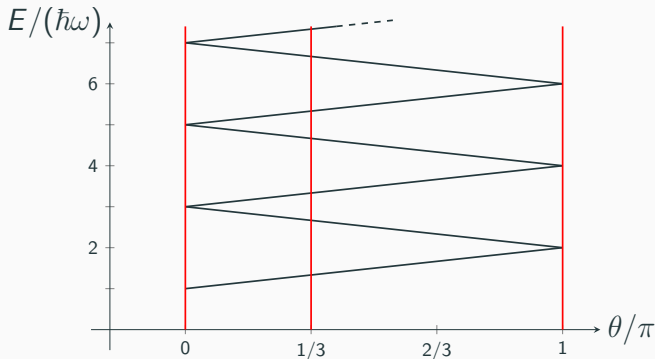
**Flux tube model:** free anyons are treated like bosons with lagrangian  $L$ , carrying **fictitious charge** and **magnetic flux**.  $L_{int}$  emerges as flux-charge interaction.

# Two harmonically interacting anyons

- **Ansatz**  $\psi(r, \phi) = R(r)e^{il\phi}$
- **Anyon statistics** with angle  $\theta$ :  $\psi(r, \phi + \pi) = e^{i\theta}\psi(r, \phi)$

$$\left[ \frac{d}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{l^2}{r^2} - \frac{2\mu}{\hbar^2} (U(r) - E) \right] R(r) = 0, \quad \begin{cases} l = \theta/\pi + 2k \\ k \in \mathbb{Z} \end{cases}$$

Hence the spectrum is  $E_{nk} = \hbar\omega (2n + 1 + |\theta/\pi + 2k|)$



# There and back again

*Practical applications of these phenomena seem remote . . .*

*Wilczek, 1982*

- **Anyons exist as quasi-particles** in systems that show **fractional quantum Hall effect** (Laughlin states).
- Significant research in **Topological Quantum Computing**. Anyonic statistics could indeed make available a **fault-tolerant** quantum computation:  
information is stored globally  $\longrightarrow$  no local perturbation can alter it

**Thanks for your attention**

# References

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