

FFTs in Graphics and Vision

Spherical Convolution and Axial Symmetry Detection

Outline



- Math Review
 - Symmetry
 - General Convolution
- Spherical Convolution
- Axial Symmetry Detection



Symmetry:

Given a unitary representation of a group G on a vector space V, we say that a vector $v \in V$ is invariant under the action of G if for all $g \in G$: $\rho_a(v) = v$

The set of G-invariant vector V_G is a vector space.



Symmetry:

The linear map π_G is a <u>projection</u> onto V_G , if:

- $\circ \quad \pi_G(v) \in V_G \text{ for all } v \in V$
- \circ $\pi_G(v) = v$ for all $v \in V_G$
- $\langle v, w \pi_G(w) \rangle = 0$ for all $v \in V_G, w \in V$.

The map π_G is the map sending a vector v to the closest G-invariant vector.



Symmetry:

The measure of symmetry of a vector v with respect to the group G is the size of its projection onto the space of G-invariant vectors:

$$Sym^2(v, G) = ||\pi_G(v)||^2$$



Convolution:

Given two functions f(p) and g(p), we define the convolution of the two functions to be:

$$(f * g)(q) = \int f(p) \cdot g(q - p) dp$$



Convolution:

If we hold the function g fixed we can define a map from the space of functions back into itself:

$$C_g(f) = f * g$$

Claim:

The map C_g is a linear operator.



Convolution:

If we hold the function g fixed we can define a map from the space of functions back into itself:

$$C_g(f) = f * g$$

Claim:

Given functions f and h and scalars α and β :

$$C_g(\alpha f + \beta h)(q) = \int (\alpha \cdot f(p) + \beta \cdot h(p)) \cdot g(q - p) dp$$

$$= \alpha \int f(p) \cdot g(q - p) dp + \beta \int h(p) \cdot g(q - p) dp$$

$$= \alpha \cdot C_g(f) + \beta \cdot C_g(h)$$



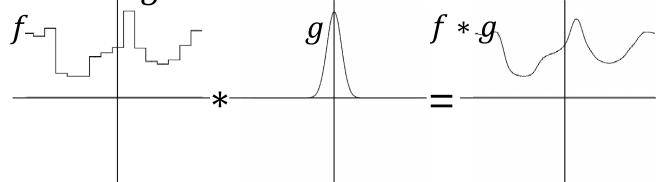
Convolution:

Assume that the function g is real-valued and radial, i.e. the value of g at a point p is completely determined by the distance of p from the origin: $a(p) = \tilde{a}(|p|)$

$$g(p) = \tilde{g}(|p|)$$

Example:

The function g is a Gaussian





Convolution:

Assume that the function *g* is real-valued and radial, i.e. the value of g at a point p is completely determined by the distance of p from the origin:

$$g(p) = \tilde{g}(|p|)$$

Claim:

In this case, C_a is self-adjoint (i.e. symmetric).



Convolution:

Proof:

We need to show that for any functions f and h:

$$\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$$

Expanding the left side, we get:

$$\langle C_g(f), h \rangle = \int (C_g(f))(p) \cdot \overline{h(p)} dp$$



Convolution:

Proof:

We need to show that for any functions f and h:

$$\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$$

Writing out the operator C_g , we get:

$$\langle C_g(f), h \rangle = \int \left(C_g(f) \right) (p) \cdot \overline{h(p)} \, dp$$

$$\downarrow$$

$$\langle C_g(f), h \rangle = \int (f * g)(p) \cdot \overline{h(p)} \, dp$$



Convolution:

Proof:

We need to show that for any functions f and h: $\langle C_a(f), h \rangle = \langle f, C_a(h) \rangle$

Expressing the convolution as an integral gives:

$$\langle C_g(f), h \rangle = \int (f * g)(p) \cdot \overline{h(p)} \, dp$$

$$\downarrow \downarrow$$

$$\langle C_g(f), h \rangle = \int \left(\int f(q) \cdot g(p-q) dq \right) \cdot \overline{h(p)} dp$$



Convolution:

Proof:

We need to show that for any functions f and h: $\langle C_a(f), h \rangle = \langle f, C_a(h) \rangle$

Changing the order of integration, we get:

$$\langle C_g(f), h \rangle = \int \left(\int f(q) \cdot g(p-q) dq \right) \cdot \overline{h(p)} dp$$
 \Downarrow

$$\langle C_g(f), h \rangle = \int \int f(q) \cdot g(p-q) \cdot \overline{h(p)} \, dp \, dq$$



Convolution:

Proof:

We need to show that for any functions f and h: $\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$

Using the fact that g is real-valued and radial:

$$\langle C_g(f), h \rangle = \int \int f(q) \cdot g(p-q) \cdot \overline{h(p)} \, dp \, dq$$

$$\downarrow \downarrow$$

$$\langle C_g(f), h \rangle = \int \int f(q) \cdot \overline{h(p)} \cdot \overline{g(q-p)} \, dp \, dq$$



Convolution:

Proof:

We need to show that for any functions *f* and *h*:

$$\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$$

Using the equation for convolution, we get:

$$\langle C_g(f), h \rangle = \int \int f(q) \cdot \overline{h(p)} \cdot \overline{g(q-p)} \, dp \, dq$$

$$\langle C_g(f), h \rangle = \int f(q) \cdot \overline{(h * g)(q)} \, dq$$



Convolution:

Proof:

We need to show that for any functions f and h:

$$\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$$

Using the equation for C_g , we get:

$$\langle C_g(f), h \rangle = \int f(q) \cdot \overline{(h * g)(q)} \, dq$$

$$\downarrow$$

$$\langle C_g(f), h \rangle = \int f(q) \cdot (C_g(h))(q) dq$$



Convolution:

Proof:

We need to show that for any functions f and h:

$$\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$$

And finally, using the equation for the dot-product:

$$\langle C_g(f), h \rangle = \int f(q) \cdot (C_g(h))(q) dq$$

$$\downarrow \\
\langle C_g(f), h \rangle = \langle f, C_g(h) \rangle$$

Outline



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- Spherical Convolution
- Axial Symmetry Detection

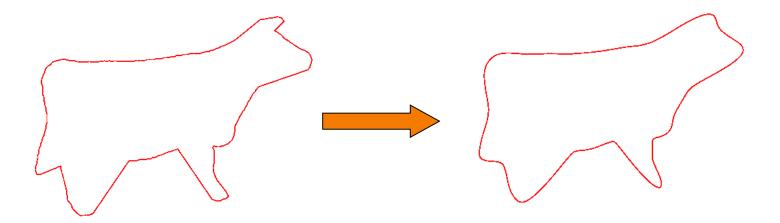


In the case of the circle we used convolution / correlation for two different tasks:



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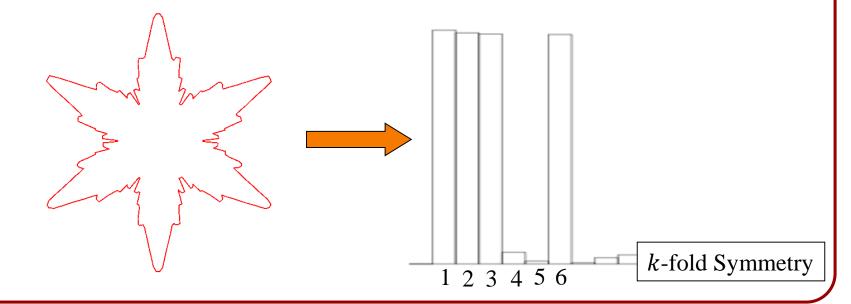
1. We used convolution for operations like smoothing circular functions





In the case of the circle we used convolution / correlation for two different tasks:

- 1. We used convolution for operations like smoothing circular functions
- 2. We used correlation for operations like alignment and symmetry detection





Up to now, we thought of these two operations as essentially the same.

The situation changes as we move to functions on a sphere.

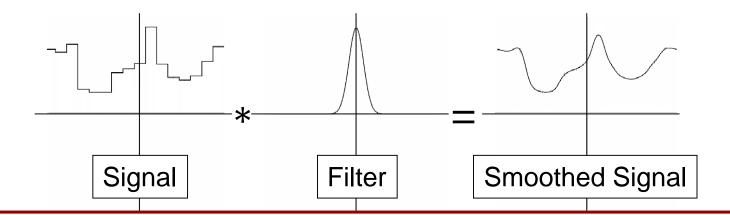


When we perform an operation like smoothing, the input is:

- A function on the circle defining the signal, and
- A function on the circle defining the smoothing filter

The output of the operation is:

A function on the circle



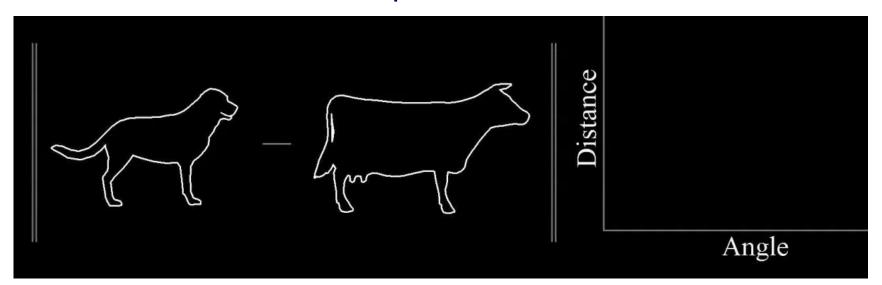


When we perform an operation like alignment, the input is:

Two functions on a circle

The output is:

A function on the space of 2D rotations





In the case of a circle, the situation is simpler because the space of rotations is itself a circle:

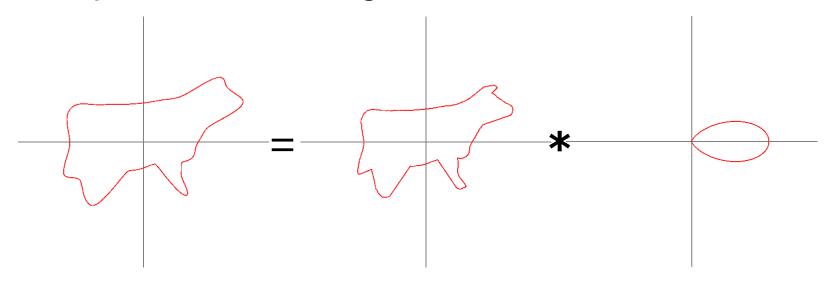
There is a one-to-one mapping from points on a circle to rotations, with a point on a circle with angle θ corresponding to a rotation by an angle of θ .

In the case of the sphere, the situation becomes more complicated:

The sphere is a 2D space while the rotations are a 3D space, so there can't be a one-to-one mapping.

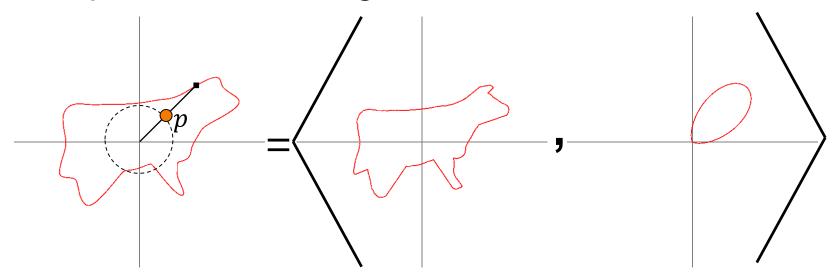


In the case of a circle, we compute the value of the smoothed function at p by rotating the filter so that (1,0) maps to p and then we compute the inner product of the signal with the rotated filter.



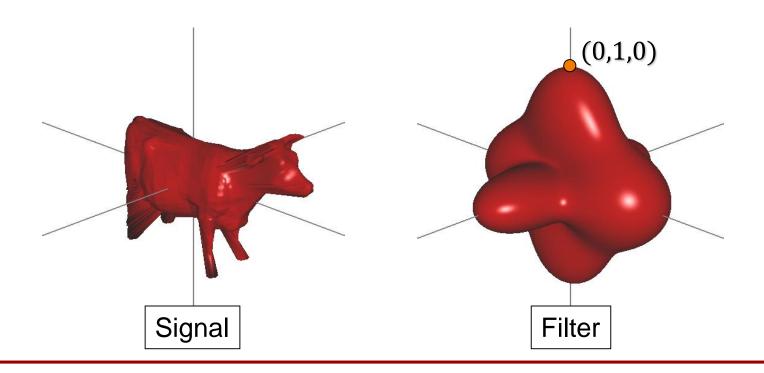


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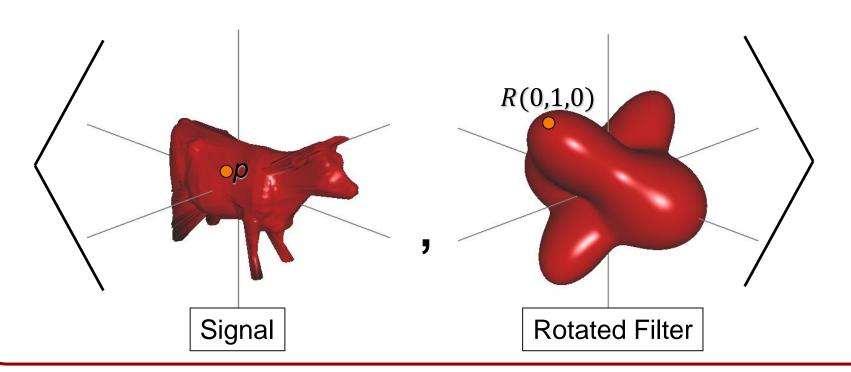
We can try an apply the same type of approach to the case of spherical functions.





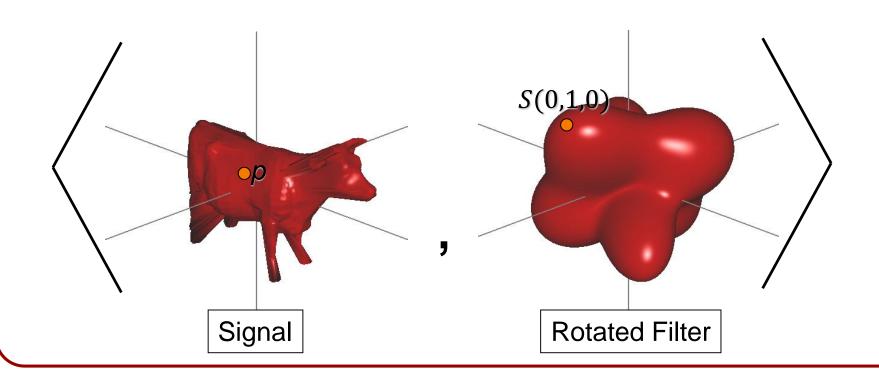
We would like to define a new function on the sphere whose value at the point p is obtained by:

Finding a rotation R that maps the North pole to p and then compute the inner product of the signal with the rotated filter.





The problem is that there are many different rotations that send the North pole to the point p, so this does not lead to a well-defined notion of smoothing.

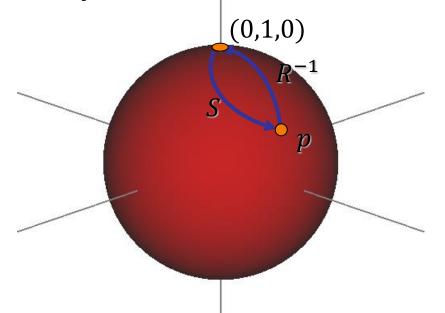




Recall:

If we have two rotations R and S mapping the North pole to the point p, the rotations must differ by an initial rotation about the y-axis:

 $S = R \cdot R_{\mathcal{Y}}(\psi)$





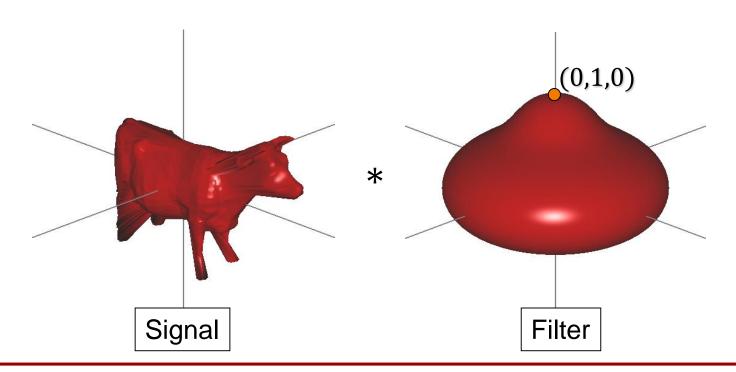
Recall:

Thus, we can make the notion of smoothing well-defined by ensuring that the initial rotation about the y-axis does not change the filter.



Recall:

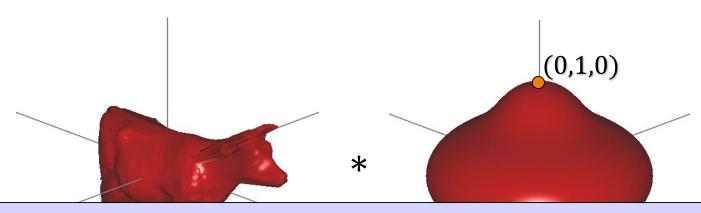
This means that we can extend the circular notion of smoothing to the sphere if we ensure that the filter is symmetric about the *y*-axis:





Recall:

This means that we can extend the circular notion of smoothing to the sphere if we ensure that the filter is symmetric about the *y*-axis:



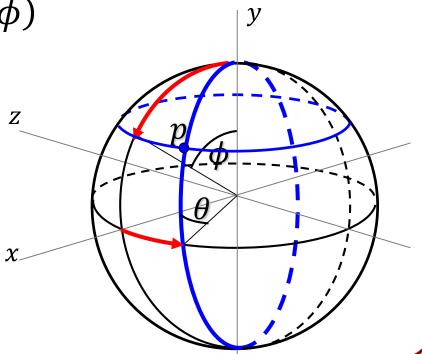
If R and S are rotations mapping the North pole to p, then the rotation of the filter by either R or S will give the same spherical function!



Convolution:

Using the Euler angle representation, we know that the rotation taking the North pole to the point $p = \Phi(\theta, \phi)$ is the rotation:

$$R(\theta, \phi) = R_y(\theta) \cdot R_z(\phi)$$





Convolution:

Thus, given

- A spherical function $f(\theta, \phi)$
- A spherical filter $g(\theta, \phi)$ that is rotationally-symmetric about the y-axis

The convolution of f with g at $p = \Phi(\theta, \phi)$ can be expressed by rotating g so the North pole gets mapped to p and computing the inner product:

$$(f * g)(\theta, \phi) = \langle f, \rho_{R(\theta, \phi)}(g) \rangle$$



Convolution:

Expressing the spherical functions f and g in terms of the spherical harmonic basis, we get:

$$f(\theta,\phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l,m) \cdot Y_{l}^{m}(\theta,\phi)$$
$$g(\theta,\phi) = \sum_{l} \sum_{m=-l}^{l} \hat{g}(l,m) \cdot Y_{l}^{m}(\theta,\phi)$$



Convolution:

Recall that the spherical harmonics can be expressed as a complex exponential in θ times a "polynomial" in $\cos \phi$:

$$Y_l^m(\theta, \phi) = P_l^m(\cos \phi) \cdot e^{im\theta}$$

So a rotation by α degrees about the y-axis acts on the (l, m)-th spherical harmonics by:

$$\rho_{R_{\mathcal{V}}(\alpha)}(Y_l^m) = e^{-im\alpha} \cdot Y_l^m$$



Convolution:

Thus, if the filter g is rotationally symmetric about the y-axis, any rotation about the y-axis must not change g. That is, for all α we must have:

$$\rho_{R_{\mathcal{Y}}(\alpha)}(g) = g$$

Or in terms of the spherical harmonics:

$$\sum_{l} \sum_{m=-l}^{l} \hat{g}(l,m) \cdot Y_{l}^{m}(\theta,\phi) = \sum_{l} \sum_{m=-l}^{l} \hat{g}(l,m) \cdot e^{-im\alpha} \cdot Y_{l}^{m}(\theta,\phi)$$

$$\downarrow \downarrow$$

$$\hat{g}(l,m) = \hat{g}(l,m) \cdot e^{-im\alpha}$$



Convolution:

$$\hat{g}(l,m) = \hat{g}(l,m) \cdot e^{-im\alpha}$$

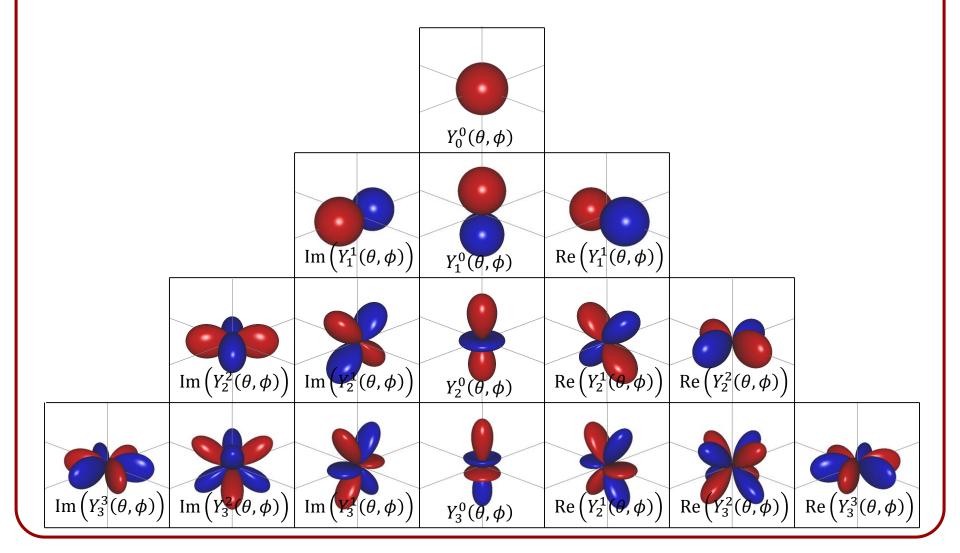
Thus, either:

- $e^{-im\alpha} = 1$ for all $\alpha \Rightarrow m = 0$, or
- \circ $\hat{g}(l,m)=0$

Thus, in terms of the spherical harmonics, we get:

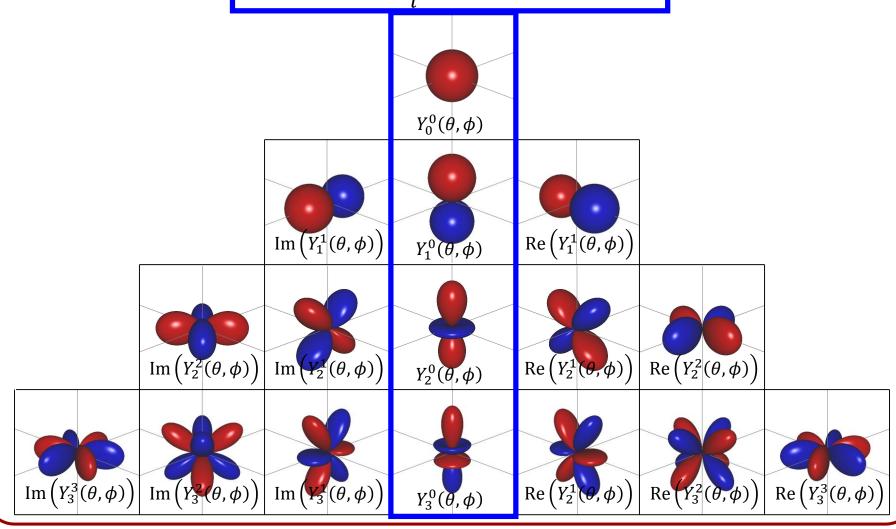
$$g(\theta,\phi) = \sum_{l} \hat{g}(l,0) \cdot Y_{l}^{0}(\theta,\phi)$$







$$g(\theta,\phi) = \sum_{l} \hat{g}(l,0) \cdot Y_{l}^{0}(\theta,\phi)$$





Convolution:

Thus, the expression for the functions in terms of their spherical harmonic decomposition becomes:

$$f(\theta,\phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l,m) \cdot Y_{l}^{m}(\theta,\phi)$$
$$g(\theta,\phi) = \sum_{l} \hat{g}(l,0) \cdot Y_{l}^{0}(\theta,\phi)$$

and we get an expression for the convolution:

$$(f * g)(\theta, \phi) = \left\langle \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot Y_{l}^{m}, \rho_{R(\theta, \phi)} \left(\sum_{l} \hat{g}(l, 0) \cdot Y_{l}^{0} \right) \right\rangle$$



Convolution:

By leveraging the conjugate-linearity of the inner product and using the fact that the transformation ρ_R is linear, we get:

$$(f * g)(\theta, \phi) = \left\langle \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot Y_{l}^{m}, \rho_{R(\theta, \phi)} \left(\sum_{l} \hat{g}(l, 0) \cdot Y_{l}^{0} \right) \right\rangle$$

$$\Downarrow$$

$$(f * g)(\theta, \phi) = \sum_{l,l'} \sum_{m=-l}^{l} \hat{f}(l,m) \cdot \hat{g}(l',0) \langle Y_l^m, \rho_{R(\theta,\phi)}(Y_{l'}^0) \rangle$$



Convolution:

Additionally, we know that:

- A rotation of an *l*-th frequency function will still be an *l*-th frequency function
- The space of l-th frequency functions is orthogonal to the space of l'-th frequency functions (if $l \neq l'$)

Thus, for all $l \neq l'$, we have:

$$\left\langle Y_l^m, \rho_R \left(Y_{l'}^{m'} \right) \right\rangle = 0$$



Convolution:

This lets us simplify the expression for the convolution:

$$(f * g)(\theta, \phi) = \sum_{l,l'} \sum_{m=-l}^{l} \hat{f}(l,m) \cdot \hat{g}(l',0) \langle Y_{l}^{m}, \rho_{R(\theta,\phi)}(Y_{l'}^{0}) \rangle$$

$$\downarrow \downarrow$$

$$(f * g)(\theta, \phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot \hat{g}(l, 0) \langle Y_l^m, \rho_{R(\theta, \phi)}(Y_l^0) \rangle$$



Convolution:

$$(f * g)(\theta, \phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot \hat{g}(l, 0) \langle Y_{l}^{m}, \rho_{R(\theta, \phi)}(Y_{l}^{0}) \rangle$$

To compute the convolution, we need to be able to evaluate the inner product:

$$\langle Y_l^m, \rho_{R(\theta,\phi)}(Y_l^0) \rangle$$



Convolution:

What is the meaning of the function:

$$\langle Y_l^m, \rho_{R(\theta,\phi)}(Y_l^0) \rangle$$

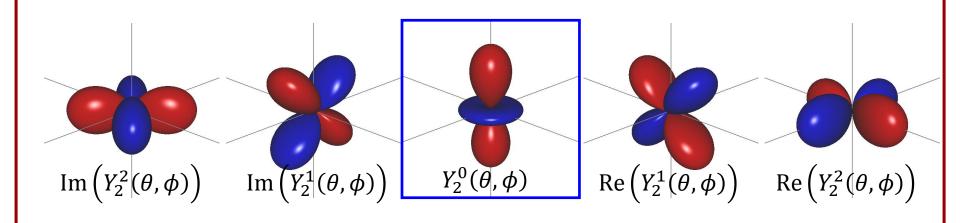
This is a function on the sphere whose value at the point $p = \Phi(\theta, \phi)$ is the inner product of Y_l^m with the rotation of Y_l^0 , where the rotation takes the North pole to p.



Convolution:

We would like to show that this function acts very simply on the spherical harmonics:

$$\langle Y_l^m, \rho_{R(\theta,\phi)}(Y_l^0) \rangle = \lambda_l \cdot Y_l^m(\theta,\phi)$$





Convolution:

Let's consider the operator C_l that maps spherical functions to spherical functions, defined by:

$$(C_l(f))(\theta,\phi) = \langle f, \rho_{R(\theta,\phi)}(Y_l^0) \rangle$$

As before, it turns out this map is a symmetric linear operator on the space of functions.

Thus, there exists an orthonormal basis with respect to which C_l is diagonal.



Convolution:

Let's consider the operator C_l that maps spherical functions to spherical functions, defined by:

$$(C_l(f))(\theta,\phi) = \langle f, \rho_{R(\theta,\phi)}(Y_l^0) \rangle$$

This operator also has the property that it commutes with rotations:

• Rotating a spherical function and then convolving with Y_l^0 is the same as first convolving with Y_l^0 and then rotating.



Convolution:

So, as with the Laplacian, we have a case in which we are given a symmetric operator which commutes with rotations.



L: a symmetric operator

 $R \in SO(3)$: a rotation

 V_{λ} : the space of e. functions of L with e.value λ $R\big(L(f)\big) = L\big(R(f)\big)$



Convolution:

Thus, the subspace of l'-th frequency functions is a space of functions that are eigenvectors of C_l , all with the same eigenvalue:

$$C_l(Y_{l'}^m) = \lambda_{l,l'} \cdot Y_{l'}^m$$

Thus we have:

$$C_l(Y_{l'}^m) = Y_{l'}^m \cdot \begin{cases} 0 & \text{if } l \neq l' \\ \lambda_l & \text{otherwise} \end{cases}$$



Convolution:

Putting this all together, we get:

$$\langle Y_l^m, \rho_{R(\theta, \phi)} Y_l^0 \rangle = \lambda_l \cdot Y_l^m(\theta, \phi)$$

Thus, the equation for the convolution becomes:

$$(f * g)(\theta, \phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot \overline{\hat{g}(l, 0)} \cdot \langle Y_{l}^{m}, \rho_{R(\theta, \phi)}(Y_{l}^{0}) \rangle$$

$$\downarrow \downarrow$$

$$(f * g)(\theta, \phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot \overline{\hat{g}(l, 0)} \cdot \lambda_{l} \cdot Y_{l}^{m}(\theta, \phi)$$



Convolution:

$$(f * g)(\theta, \phi) = \sum_{l} \sum_{m=-l}^{l} \hat{f}(l, m) \cdot \overline{\hat{g}(l, 0)} \cdot \lambda_{l} \cdot Y_{l}^{m}(\theta, \phi)$$

Thus, the convolution of f with g can be obtained by multiplying the (l,m)-th spherical harmonic coefficients of f by $\lambda_l \cdot \overline{\widehat{g}(l,0)}$.

As in the case of functions on a circle, this means that convolution in the spatial domain amounts to multiplication in the frequency domain.



Convolution:

In order to be able to use the convolution theorem for spherical functions, we need to know what the eigenvalues λ_l are.



Convolution:

In order to be able to use the convolution theorem for spherical functions, we need to know what the eigenvalues λ_l are.

It turns out that these are:

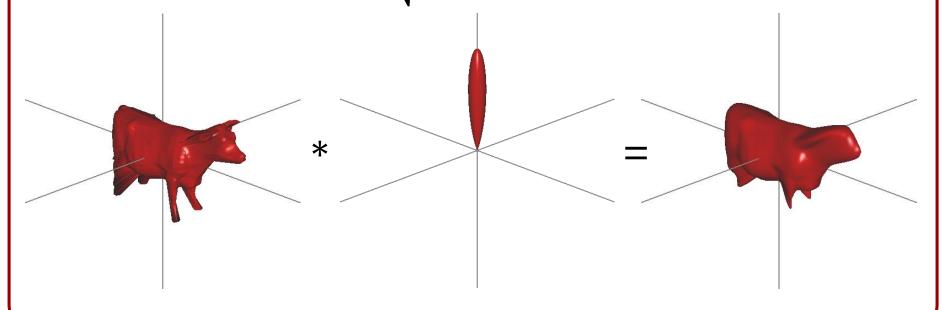
$$\lambda_l = \sqrt{\frac{4\pi}{2l+1}}$$



Convolution:

Which gives us the equation:

$$(\widehat{f * g})(l, m) = \sqrt{\frac{4\pi}{2l+1} \cdot \widehat{f}(l, m) \cdot \overline{\widehat{g}(l, 0)}}$$



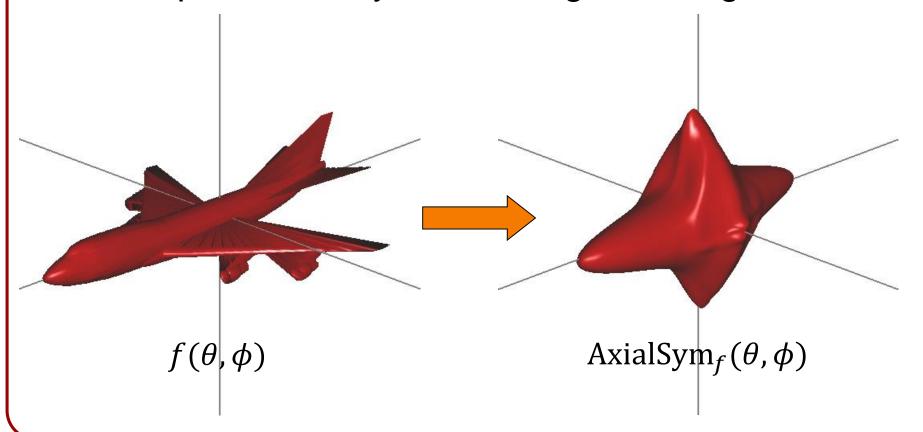
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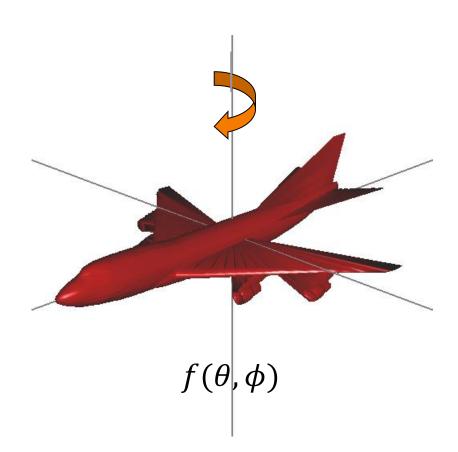


Given a spherical function f, we would like to compute the measure of the axial symmetry of f with respect to every axis through the origin.





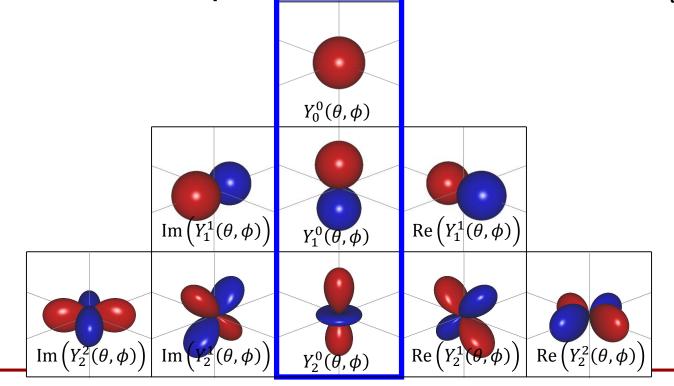
What is the measure of the axial symmetry of f about the y-axis?





What is the measure of the axial symmetry of *f* about the *y*-axis?

We know that f is axially symmetric about the y-axis if it can be expressed as the sum of the Y_i^0 :





What is the measure of the axial symmetry of *f* about the *y*-axis?

We know that f is axially symmetric about the y-axis if it can be expressed as the sum of the Y_l^0 .

We also know that for $m \neq 0$:

$$\langle Y_l^0, Y_l^m \rangle = 0$$

So the projection onto the space of functions that are axially symmetric about the y-axis is:

$$\pi_{\mathcal{Y}}\left(\sum_{l}\sum_{m=-l}^{k}\hat{f}(l,m)\cdot Y_{l}^{m}\right) = \sum_{l}\hat{f}(l,0)\cdot Y_{l}^{0}$$



What is the measure of the axial symmetry of f about the y-axis?

Thus, the measure of the axial symmetry of f about the y-axis is defined as:

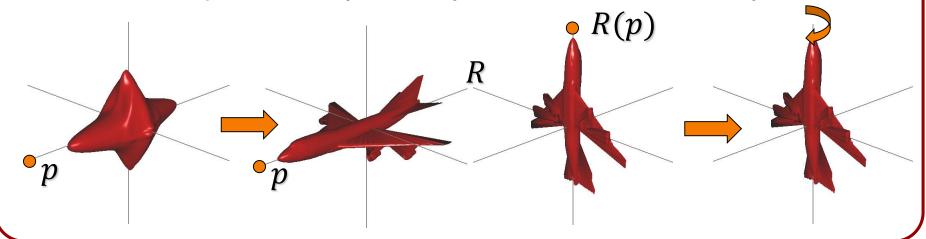
$$YAxialSym2(f) = \left\| \sum_{l} \hat{f}(l,0) \cdot Y_{l}^{0} \right\|^{2}$$
$$= \sum_{l} \left\| \hat{f}(l,0) \right\|^{2}$$



More generally, we would like to be able to compute the measure of the axial symmetry of f with respect to any axis.

To compute the symmetry measure about the line through $p = \Phi(\theta, \phi)$ we:

- Rotate so that p goes to the North pole, and
- \circ Compute the symmetry measure about the y-axis.





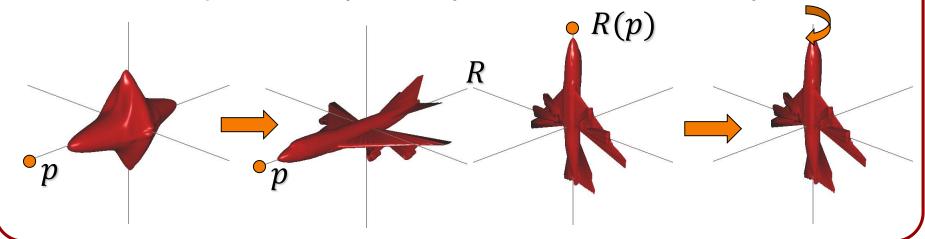
More generally, we would like to be able to compute the measure of the axial symmetry of f

with

Since the rotation $R(\theta, \phi)$ maps the North pole to p, the rotation we are interested in is the inverse, $R^{-1}(\theta, \phi)$.

line

- Rotate so that p goes to the North pole, and
- Compute the symmetry measure about the *y*-axis.





Using the fact that the spherical harmonics form an orthonormal basis, we know that the (l, m)-th harmonic coefficient of f is defined by:

$$\hat{f}(l,m) = \langle f, Y_l^m \rangle$$

Thus, to compute the measure of axial symmetry about the axis through p we need to compute:

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \| \langle \rho_{R^{-1}(\theta, \phi)}(f), Y_l^0 \rangle \|^2$$



Using the fact that ρ is a unitary representation we can re-write this equation as:

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \| \langle \rho_{R^{-1}(\theta, \phi)}(f), Y_{l}^{0} \rangle \|^{2}$$

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \|\langle f, \rho_{R(\theta, \phi)}(Y_l^0) \rangle\|^2$$



Expressing *f* in terms of its spherical harmonic decomposition, we get:

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \|\langle f, \rho_{R(\theta, \phi)}(Y_l^0) \rangle\|^2$$

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \left\| \sum_{m=-l}^{l} \hat{f}(l, m) \langle Y_l^m, \rho_{R(\theta, \phi)}(Y_l^0) \rangle \right\|^2$$



Applying the identity:

$$\langle Y_l^m, \rho_{R(\theta,\phi)} Y_l^0 \rangle = \sqrt{\frac{4\pi}{2l+1}} \cdot Y_l^m(\theta,\phi)$$

we get an expression for the symmetry measure:

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \left\| \sum_{m=-l}^{l} \hat{f}(l, m) \langle Y_l^m, \rho_{R(\theta, \phi)}(Y_l^0) \rangle \right\|^2$$

AxialSym_f²
$$(\theta, \phi) = \sum_{l} \frac{4\pi}{2l+1} \left(\left\| \sum_{m=-l}^{l} \hat{f}(l, m) \cdot Y_{l}^{m}(\theta, \phi) \right\|^{2} \right)$$



AxialSym_f²
$$(\theta, \phi) = \sum_{l} \frac{4\pi}{2l+1} \left(\left\| \sum_{m=-l}^{l} \hat{f}(l, m) \cdot Y_l^m(\theta, \phi) \right\|^2 \right)$$

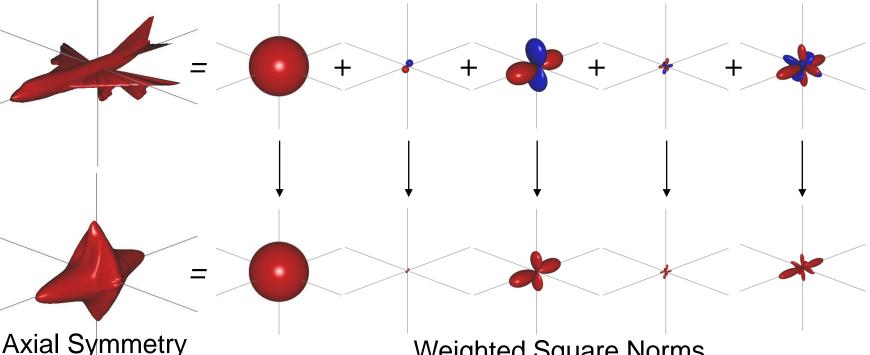
Thus, the measure of axial symmetry can be computed by taking the weighted sum of the squares of the frequency components of f.



AxialSym_f²
$$(\theta, \phi) = \sum_{l} \frac{4\pi}{2l+1} \left(\left\| \sum_{m=-l}^{l} \hat{f}(l, m) \cdot Y_l^m(\theta, \phi) \right\|^2 \right)$$

Initial Function

Frequency Decomposition



Axial Symmetry Descriptor

Weighted Square Norms