

Optimizing acceptance rates and drawing initial states using the density of states

This document gives mathematical details about the density of states-based methods used in RESAAS to optimize Replica Exchange schedules and to draw good initial states for consecutive simulation runs. All credit for these ideas goes to Prof. Michael Habeck at Jena University / Germany.

General theory

Two replicas are sampling distributions $p_1(x)$ and $p_2(x)$ and are currently at states x_1 and x_2 . The probability for accepting an exchange is then given by

$$p_{\text{acc}}(x_1, x_2 | p_1, p_2) = \min \left\{ 1, \frac{p_1(x_2) p_2(x_1)}{p_1(x_1) p_2(x_2)} \right\}. \quad (1)$$

Assume p_1 and p_2 depend on x only through a common function $E(x)$, the “energy”, that is, we have $p_i(x) = p_i(E(x))$.

To calculate the average acceptance rate \bar{p}_{acc} , we calculate the expectation value of p_{acc} with respect to the joint distribution $p(x_1, x_2) = p_1(x_1)p_2(x_2)$ and we have

$$\bar{p}_{\text{acc}} = \iint dx_1 dx_2 p_{\text{acc}}(x_1, x_2 | p_1, p_2) p_1(x_1) p_2(x_2). \quad (2)$$

This is a double integral over two potentially high-dimensional variables. To make progress, we introduce the density of states (DOS) $g(E)$, which counts the multiplicity of energies in an interval $[E, E + dE]$:

$$g(E) = \int dx \delta(E - E(x))$$

Using the DOS, we can write the expectation value of any function $f(x)$ which depends on x only via the energy E (meaning, again, $f(x) = f(E(x))$), as an integral over E instead of x :

$$\int dx f(E(x)) = \int dE g(E) f(E)$$

We further note that often, we only know distributions p up to a normalization constant Z , and we have

$$p(x) = \frac{1}{Z} q(x)$$

Using the DOS, above equation and the assumption that $q(x) = q(E(x))$, we can now rewrite the expected acceptance rate as an integral over energies:

$$\begin{aligned}\bar{p}_{\text{acc}} &= \frac{1}{Z_1 Z_2} \iint dx_1 dx_2 \min \left\{ 1, \frac{q_1(E(x_2))}{q_1(E(x_1))} \frac{q_2(E(x_1))}{q_2(E(x_2))} \right\} q_1(E(x_1)) q_2(E(x_2)) \\ &= \frac{1}{Z_1 Z_2} \iint dx_1 dx_2 \min \{ q_1(E(x_1)) q_2(E(x_2)), q_1(E(x_2)) q_2(E(x_1)) \} \\ &= \frac{1}{Z_1 Z_2} \iint dE_1 dE_2 g(E_1) g(E_2) \min \{ q_1(E_1) q_2(E_2), q_1(E_2) q_2(E_1) \}\end{aligned}\tag{3}$$

Note that the normalization constants Z_1, Z_2 can be written in terms of the energies E , too:

$$Z_i = \int dx q_i(E(x)) = \int dE g(E) q_i(E)\tag{4}$$

Long story short: if we have access to $g(E)$, we can estimate the expected acceptance rate by means of simple numeric approximation of one-dimensional and two-dimensional integrals over energies.

We can get a good estimate of $g(E)$ via multiple histogram reweighting[1]; which results in an estimate of $g(E)$ at every sampled value and thus the integrals collapse into sums.

Example: Boltzmann ensemble

The only family of tempered distributions currently supported in RESAAS is the Boltzmann ensemble. It is given by

$$p_\beta(E) = \frac{\exp(-\beta E)}{Z(\beta)}$$

and we thus define

$$\begin{aligned}p_1(x) &= \frac{1}{Z(\beta_1)} \exp(-\beta_1 E(x)), \\ p_2(x) &= \frac{1}{Z(\beta_2)} \exp(-\beta_2 E(x)).\end{aligned}$$

The expected acceptance rate is then

$$\begin{aligned}\bar{p}_{\text{acc}} &= \frac{1}{Z(\beta_1) Z(\beta_2)} \times \\ &\quad \iint dE_1 dE_2 g(E_1) g(E_2) \min \{ \exp(-\beta_1 E_1 - \beta_2 E_2), \exp(-\beta_1 E_2 - \beta_2 E_1) \}\end{aligned}\tag{5}$$

and the normalization constants are given by

$$Z(\beta_i) = \int dE g(E) \exp(-\beta_i E).$$

Assuming we have an estimate of $g(E)$ from multiple histogram reweighting, this means we can calculate the expected acceptance rate for any two pairs of inverse temperatures β_1, β_2 . Note that in the source code, the acceptance rate and the normalization constants are calculated in log-space for numerical stability.

Determining an inverse temperature schedule with constant acceptance rates

Say now that $\beta_0 = 1$ gives us the distribution we actually are interested in and $\beta_N = \epsilon \approx 0$ is an almost uniform distribution. The above now suggests an iterative algorithm to get a sequence $\beta_0 > \beta_1 > \dots > \beta_N$ for which the acceptance rates between simulations at consecutive temperatures β_i, β_{i+1} attain a constant value $\bar{p}_{\text{acc}}^{\text{target}}$ [2]. we start with $\beta = 1$ and lower β in increments of $\Delta\beta$,

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params ← [1.0]
Δ ← δ
while params[-1] - Δ > βmin do
  if  $\bar{p}_{\text{acc}}(\text{params}[-1] - \Delta) \leq p_{\text{acc}}^{\text{target}}$  then
    append params[-1] - Δ to params
    Δ ← δ
  else
    Δ ← Δ + δ
  end if
end while
append params[-1] - Δ to params

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until the expected acceptance rate drops below $p_{\text{acc}}^{\text{target}}$. We then save the β value before the last decrement as β_1 . We then calculate acceptance rates between β_1 and steadily decreasing β values, until the acceptance reate drops below $p_{\text{acc}}^{\text{target}}$ again and we save the previous β value as β_2 and so on and so forth. At one point, we will hit a predefined β_{min} . We then terminate and have obtained the desired sequence of β values as our optimized schedule.

Drawing initial states using the density of states

To shorten the burnin period of simulations other than the very first one, RESAAS uses the density of states to reweight samples from a previous simulation and then to draw “good” initial states in high-probability regions for a Replica Exchange simulation.

Given an unnormalized probability density $q(x)$ which, as above, depends on x

only via an “energy” E , meaning $q(x) = q(E_x)$ (sorry for the abuse of notation), we observe that

$$q(x) = \frac{g(E_x)q(E_x)}{Z_i}$$

Remember that the relation between states x and energies E_x is not one-to-one. The x -ses are samples obtained from a previous simulation and for each x we have an estimate $g(E_x)$. Using above relation we can thus calculate, for each previous sample x , its probability weight under the new distribution q , which can be any member of the set of tempered distributions used in the next Replica Exchange run. Approximately drawing from q then amounts to draw samples from the categorical distribution of x -ses with probabilities given by the sample weights. This idea is also due to Prof. Michael Habeck[2].

References

- [1] Habeck, Michael. “Evaluation of marginal likelihoods via the density of states.” *AISTATS* (2012).
- [2] Habeck, Michael. “Ensemble annealing of complex physical systems.” *arXiv: Computational Physics* (2015).