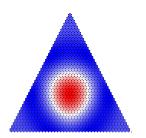
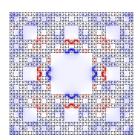
Plasmonic Properties of QM 2D Fractals.

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TCM











Plasmons

Definition (Plasmon)

Quantum of plasma oscillation (think quasiparticle).

Definition (Plasma Oscillation)

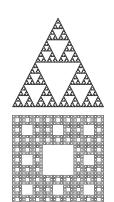
Longitudinal waves of electron charge-density.



- Features: nowhere differentiable, have fractal dimensions, self-similar, etc.
- Ramification: minimum number of links that one needs to remove to separate a macroscopic part of infinite fractal.

Finite ramification: Solved analytically by L.P. Kadanoff (1982).

Infinite ramification: Unsolvable analytically (?). Numerical calculations exist (done at Radboud), but not of plasmonic properties.





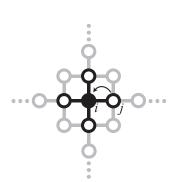
- atomic basis {|a⟩}
- Hamiltonian of the system is

$$\hat{\mathcal{H}} = -t \, \sum_{\langle a,b
angle} (\hat{c}_a^\dagger \hat{c}_b + ext{h.c.}) \; ,$$

 \hat{c}_a^{\dagger} and \hat{c}_a are creation and annihilation operators, t is the hopping value.

• Hopping:

$$\hat{c}_a^{\dagger} \hat{c}_b | \psi \rangle$$
"take" from b
"move" to a



System

- Electron system is described (in one-particle approximation) by
 - eigenenergies $\{E_i\}$ and
 - eigenstates $\{|i\rangle\}$.
- Grand-canonical ensemble, i.e. temperature T and chemical potential μ are fixed.
- Occupational numbers are given by the Fermi-Dirac distribution.

Problem

- All plasmonic properties are obtained from the dielectric function $\hat{\varepsilon}$.
 - \implies Goal: calculate $\hat{\varepsilon}$.
- There exist good methods for calculation of the dielectric function in momentum space, i.e. they rely on translational symmetry of the system.
- But in fractals, there is no translational symmetry!
- Calculations have to be done in real space.
 - ⇒ New method.

Dielectric Function in Real Space

Within the Random Phase Approximation (RPA) we obtain:

$$\langle \mathbf{r}|\hat{V}_{\mathrm{ext}}(\omega)|\mathbf{r}
angle = \int\!\!\mathrm{d}^3r'\;\langle \mathbf{r}|\hat{arepsilon}(\omega)|\mathbf{r}'
angle\langle \mathbf{r}'|\hat{V}|\mathbf{r}'
angle\;,\;\;\mathrm{where}$$

$$\langle \mathbf{r} | \hat{\varepsilon}(\omega) | \mathbf{r}' \rangle = \langle \mathbf{r} | \mathbf{r}' \rangle - \lim_{\eta \to 0+} \sum_{i,j} \frac{n_i - n_j}{E_i - E_j - \hbar(\omega + i\eta)} \int d^3 r'' \frac{e^2}{\|\mathbf{r} - \mathbf{r}''\|} \times \langle j | \mathbf{r}'' \rangle \langle \mathbf{r}'' | i \rangle \langle i | \mathbf{r}' \rangle \langle \mathbf{r}' | j \rangle .$$

Algorithm

Calculating $\hat{\varepsilon}$ for a single ω is an $\mathcal{O}(N^4)$ problem, where N is the number of particles.

- Hack: rewrite everything in terms of matrix operations and use BLAS (Basic Linear Algebra Subroutines).
- Use a supercomputer for calculations.
- Hack: use symmetries of the system.

Code

- Written C++ package to calculate $\hat{\varepsilon}$ for a general system described by TB.
- \approx 8.5 thousands lines of code (excluding data analysis scripts).
- \approx 400000 computational hours spent, that is \approx 45 years of calculations.

Electron Energy Loss Spectroscopy (EELS)

- Shoot electrons on the sample.
 - Electron beam with a well-defined wavevector k.
 - Inelastic scattering results in energy loss $E=\hbar\omega$ and momentum transfer $\hbar {\bf q}$.
- Look at (usually) transmitted electrons.

$$\frac{\mathrm{d}^2 \sigma}{\mathrm{d}\Omega \mathrm{d}E} \propto -\mathrm{Im} \left[\frac{1}{\langle \mathbf{q} | \hat{\varepsilon}(\omega) | \mathbf{q} \rangle} \right] \ .$$
 loss function

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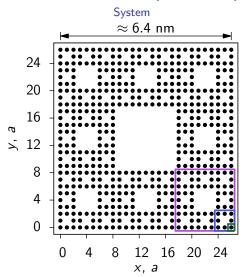
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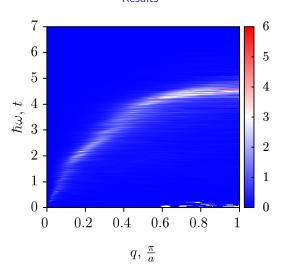
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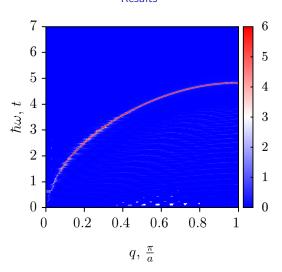
Third Iteration Sierpinski Carpet



Third Iteration Sierpinski Carpet



Square Lattice Results



Summary

- It is possible to do calculations of $\hat{\varepsilon}$ in real space.
- There are well defined plasmons in Sierpinski carpets, which is unexpected for a system without translational invariance.
- The created package can be reused to study plasmonic properties of other systems.

Outlook

- Publish the results, i.e. write a paper.
- Look at other systems with no translational symmetry.

Any questions?

Simple example of plasma oscillation

- Free electron gas with positively charged atom cores in the background.
- Oscillating electric field $\mathbf{E}(t) = \mathbf{E_0} \exp(-i\omega t)$.
- Equation of motion:

$$\ddot{\mathbf{x}} + \gamma \dot{\mathbf{x}} = -\frac{e}{m} \mathbf{E}(t) .$$

- Displaced electrons generate a polarization $\mathbf{P} = -Ne\mathbf{x}$ (N is the electron density).
- From $\epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon_0 \varepsilon(\omega) \mathbf{E}$ we get $\varepsilon(\omega)$.
- In high frequency limit $\varepsilon(\omega) \approx 1 \frac{\omega_p^2}{\omega^2}$, ω_p plasma frequency.
- If $\omega=\omega_p$, we get longitudinal oscillation of electron-charge density.



Random Phase Approximation

Perturbation of the form

$$\hat{V}e^{-i(\omega+i\eta)t}$$
, with $\eta \to +0$.

$$(\hat{V} \text{ is diagonal in r-basis, i.e. } \langle \mathbf{r}|\hat{V}|\mathbf{r}'
angle = V(\mathbf{r})\delta^3(\mathbf{r}-\mathbf{r}'))$$

- Results in an induced charge density variation $\delta \hat{N}(t)$
- Self-consistency equation

$$\langle \mathbf{r}|\hat{V}_{\text{tot}}(t)|\mathbf{r}\rangle = \langle \mathbf{r}|\hat{V}_{\text{ext}}(t)|\mathbf{r}\rangle + \int\!\!\mathrm{d}^3r'\;\langle \mathbf{r}|\hat{V}_{\text{Coulomb}}|\mathbf{r}'\rangle\langle \mathbf{r}'|\delta\hat{N}(t)|\mathbf{r}'\rangle\;.$$