

AMATH 503: Homework 3
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(1)

(a) Using the formula for a complex Fourier series in the notes we have the following in a domain of $-2 < x < 2$ with $n = 0, \pm 1, \pm 2, \pm 3 \dots$

$$|x| = \sum_{-\infty}^{\infty} c_n e^{\frac{-in\pi x}{2}} \quad (1)$$

Similarly, we obtain the coefficient c_n with the following:

$$c_n = \frac{1}{4} \int_{-2}^2 |x| e^{\frac{in\pi x}{2}} dx \quad (2)$$

To solve this, we split the integral into two domains and integration by parts:

$$\begin{aligned}
 c_n &= \frac{1}{4} \left[\int_0^2 x e^{\frac{in\pi x}{2}} dx - \int_{-2}^0 x e^{\frac{in\pi x}{2}} dx \right] \\
 u &= x, \quad du = 1 \\
 dv &= e^{\frac{in\pi x}{2}} \\
 v &= \frac{2}{in\pi} e^{\frac{in\pi x}{2}} \\
 uv &= \int v du = \frac{2x}{in\pi} e^{\frac{in\pi x}{2}} - \frac{4}{n^2\pi^2} e^{\frac{in\pi x}{2}} \quad (3) \\
 c_n &= \frac{1}{4} \left[\left[\frac{2xe^{\frac{in\pi x}{2}}}{in\pi} + \frac{4e^{\frac{in\pi x}{2}}}{n^2\pi^2} \right] \Big|_0^2 - \left[\frac{2xe^{\frac{in\pi x}{2}}}{in\pi} + \frac{4e^{\frac{in\pi x}{2}}}{n^2\pi^2} \right] \Big|_{-2}^0 \right] \\
 c_n &= \frac{1}{4} \left[\left[\frac{4e^{in\pi}}{in\pi} + \frac{4e^{in\pi}}{n^2\pi^2} - \frac{4}{n^2\pi^2} \right] - \left[\frac{4}{n^2\pi^2} + \frac{4e^{-in\pi}}{in\pi} + \frac{4e^{-in\pi}}{n^2\pi^2} \right] \right] \\
 c_n &= \frac{-in\pi e^{in\pi} + in\pi e^{-in\pi} + e^{in\pi} + e^{-in\pi} - 2}{n^2\pi^2}
 \end{aligned}$$

Now we observe that:

$$\begin{aligned}
 e^{in\pi} &= \cos(n\pi) + i\sin(n\pi) = \cos(n\pi) = (-1)^n \\
 e^{-in\pi} &= \cos(n\pi) - i\sin(n\pi) = \cos(n\pi) = (-1)^n \\
 e^{in\pi} &= e^{-in\pi}
 \end{aligned} \quad (4)$$

This allows for the following simplification:

$$c_n = \frac{2(-1)^n - 2}{n^2\pi^2} \quad (5)$$

This is obviously not defined for $n = 0$ however, and so we calculate this independently:

$$\begin{aligned}
 c_0 &= \frac{1}{4} \int_{-2}^2 |x| dx \\
 c_0 &= \frac{1}{4} \left[\left[\frac{x^2}{2} \right] \Big|_0^2 - \left[\frac{x^2}{2} \right] \Big|_{-2}^0 \right] = 1
 \end{aligned} \quad (6)$$

Putting all this together, we arrive at the following Fourier series, noting that the even modes are zero (except c_0), and that $n = \pm 1, \pm 2, \pm 3 \dots$

$$|x| = 1 + \frac{2}{\pi^2} \sum_{n=-\infty}^{\infty} \frac{(-1)^n - 1}{n^2} e^{-\frac{in\pi x}{2}} \quad (7)$$

(b) Using Euler's Formula we can rewrite the exponential above in the following way:

$$e^{-\frac{in\pi x}{2}} = \cos\left(\frac{n\pi x}{2}\right) - i\sin\left(\frac{n\pi x}{2}\right) \quad (8)$$

If we observe that \sin is an odd function, it follows that $i\sin\left(\frac{-n\pi x}{2}\right) = -i\sin\left(\frac{n\pi x}{2}\right)$.

Additionally, the n term in the divisor is squared, so it is no different for $\pm n$.

This means that when $n < 0$, all the \sin values in the sum will cancel with their corresponding values from $n > 0$, for example when $n = 1$ we have $-i\sin\left(\frac{\pi x}{2}\right)$ and when $n = -1$ we have $i\sin\left(\frac{n\pi x}{2}\right)$.

Similarly, we can observe that cosine is an even function and $\cos(nx) = \cos(-nx)$.

Thus we can change our sum to be from $0 < n < \infty$ and multiply the sum by 2.

Lastly, we can also observe that the even modes are zero, substituting $n = 2k - 1$ and define our sum with $k = 1, 2, 3 \dots$. This means $c_n = \frac{-4}{n^2\pi^2}$ as well. The resulting cosine series after discarding the sines, changing the range of n , eliminating even modes, and multiplying by 2, is:

$$|x| = 1 - \frac{8}{\pi^2} \sum_{k=1}^{\infty} \frac{\cos\left(\frac{(2k-1)\pi x}{2}\right)}{(2k-1)^2} \quad (9)$$

(c)

$$\frac{d}{dx} \left[1 - \frac{8}{\pi^2} \sum_{k=1}^{\infty} \frac{\cos\left(\frac{(2k-1)\pi x}{2}\right)}{(2k-1)^2} \right] = \frac{4}{\pi} \sum_{k=1}^{\infty} \frac{\sin\left(\frac{n\pi x}{2}\right)}{2k-1} \quad (10)$$

This is the sine series for $f(x) = 1$ in a domain of $0 < x < 2$. To show this, we simply observe that if $f(x) = 1$, our sine series coefficient is:

$$\begin{aligned} a_n &= \int_0^2 \sin\left(\frac{n\pi x}{2}\right) \\ a_n &= \left[-\frac{2}{n\pi} \cos\left(\frac{n\pi x}{2}\right) \right]_0^2 \\ a_n &= \frac{2 - 2(-1)^n}{n\pi} \end{aligned} \tag{11}$$

Even modes are zero, thus:

$$\begin{aligned} f(x) &= \sum_{n=1}^{\infty} \frac{4}{n\pi} \sin\left(\frac{n\pi x}{2}\right) \\ f(x) &= \frac{4}{\pi} \sum_{k=1}^{\infty} \frac{\sin\left(\frac{(2k-1)\pi x}{2}\right)}{2k-1} \end{aligned}$$

This is the same as the result of taking our derivative, which makes sense since we know the derivative of $|x|$ is -1 from $-\infty < x < 0$ and 1 from $0 < x < \infty$. Furthermore, we know that a fourier sine series for a non-periodic function will only be accurate in the domain $0 < x < L$, and because it's an odd function, will be reflected accordingly outside that domain. Thus we would expect the Fourier series for the derivative of $|x|$ to be accurate from $-L < x < L$ and be $2L$ periodic outside this. The resulting function would be the sine series of $f(x) = 1$ and that is precisely what we got.

(2) We first define the fourier transforms in the different dimensions:

$$\begin{aligned} \mathcal{F}_1[u(x, y, z, t)] &= \int_{-\infty}^{\infty} u(x, y, z, t) e^{i\omega_1 x} dx \\ \mathcal{F}_2[u(x, y, z, t)] &= \int_{-\infty}^{\infty} u(x, y, z, t) e^{i\omega_2 y} dy \\ \mathcal{F}_3[u(x, y, z, t)] &= \int_{-\infty}^{\infty} u(x, y, z, t) e^{i\omega_3 z} dz \end{aligned} \tag{12}$$

We can then fully define u in fourier space:

$$\mathcal{F}_1[\mathcal{F}_2[\mathcal{F}_3[u(x, y, z, t)]]] = U(\omega_1, \omega_2, \omega_3, t) \quad (13)$$

Now defining the terms in the PDE, in particular using the definition of a derivative in Fourier space from the notes:

$$\begin{aligned} \mathcal{F}_1\mathcal{F}_2\mathcal{F}_3[u_t] &= \frac{d}{dt}U(\omega_1, \omega_2, \omega_3, t) \\ \mathcal{F}_1\mathcal{F}_2\mathcal{F}_3[u_x x] &= (i\omega_1)^2 U(\omega_1, \omega_2, \omega_3, t) \\ \mathcal{F}_1\mathcal{F}_2\mathcal{F}_3[u_y y] &= (i\omega_2)^2 U(\omega_1, \omega_2, \omega_3, t) \\ \mathcal{F}_1\mathcal{F}_2\mathcal{F}_3[u_z z] &= (i\omega_3)^2 U(\omega_1, \omega_2, \omega_3, t) \end{aligned} \quad (14)$$

We can now write out the full PDE in Fourier space:

$$\frac{d}{dt}U = -D(\omega_1^2 + \omega_2^2 + \omega_3^2)U \quad (15)$$

We then separate variables and integrate with respect to t , letting $D = \alpha^2$:

$$\begin{aligned} \frac{dU}{U} &= -\alpha^2(\omega_1^2 + \omega_2^2 + \omega_3^2)dt \\ \ln U &= -\alpha^2(\omega_1^2 + \omega_2^2 + \omega_3^2)t + C \\ U &= Ae^{-\alpha^2(\omega_1^2 + \omega_2^2 + \omega_3^2)t} \end{aligned} \quad (16)$$

Applying the initial condition, we have:

$$\begin{aligned} A &= U(\omega_1, \omega_2, \omega_3, 0) \\ U(\omega_1, \omega_2, \omega_3, 0) &= \mathcal{F}_1\mathcal{F}_2\mathcal{F}_3\delta(x)\delta(y)\delta(z) \end{aligned} \quad (17)$$

As we learned in class, the integral across a delta function is equal to 1 because it is summing the area under an infinitely small Riemann rectangle. Thus we

have:

$$U = e^{-\alpha^2(\omega_1^2 + \omega_2^2 + \omega_3^2)t} \quad (18)$$

To apply the reverse transform, we can split U up into three exponentials and apply the reverse transform to each one:

$$\begin{aligned} U &= e^{-\alpha^2\omega_1^2t} e^{-\alpha^2\omega_2^2t} e^{-\alpha^2\omega_3^2t} \\ u(x, y, z, t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\alpha^2\omega_1^2t} e^{i\omega_1x} d\omega_1 \\ &\quad * \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\alpha^2\omega_2^2t} e^{i\omega_2y} d\omega_2 \\ &\quad * \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\alpha^2\omega_3^2t} e^{i\omega_3z} d\omega_3 \end{aligned} \quad (19)$$

These integrals can be re-written in the form:

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\alpha^2\omega_1^2t + i\omega_1x} d\omega_1 \quad (20)$$

Using the general formula for Euler's Integral, taken from the class notes section (8.2), and noting that t is being treated as a constant here since we are integrating in respect to the omegas. The formula is as follows:

$$\int_{-\infty}^{\infty} e^{p^2\omega^2 \pm q\omega} d\omega = \frac{\sqrt{\pi}}{p} e^{\frac{q^2}{4p^2}} \quad (21)$$

Applying this to each integral, we get the answer:

$$\begin{aligned} u(x, y, z, t) &= \left(\frac{1}{2\pi}\right)^3 \left(\frac{\sqrt{\pi}}{\alpha\sqrt{t}}\right)^3 e^{\frac{-x^2}{4\alpha^2t}} e^{\frac{-y^2}{4\alpha^2t}} e^{\frac{-z^2}{4\alpha^2t}} \\ u(x, y, z, t) &= \frac{1}{8\pi^3} \left(\frac{\pi^{\frac{3}{2}}}{\alpha^3 t^{\frac{3}{2}}}\right) e^{-\frac{(x^2+y^2+z^2)}{4\alpha^2t}} \\ u(x, y, z, t) &= \frac{1}{8\alpha^3 t^{\frac{3}{2}} \pi^{\frac{3}{2}}} e^{-\frac{(x^2+y^2+z^2)}{4\alpha^2t}} \end{aligned} \quad (22)$$