

# Coupled SHM

Prof Dave Spence

- We've looked at
  - undamped motion
  - (decaying) damped motion
  - forced motion
- Now we look at coupled motion
  - Multiple oscillators, 'coupled' in some way so they may exchange energy
- Identify *normal modes*
  - Simple patterns of vibration
  - General solution: superposition of normal modes

# Coupled Oscillators



## Coupled Oscillators

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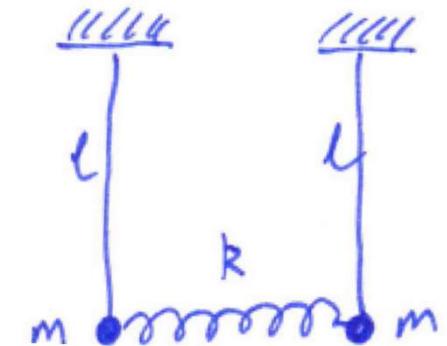
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# Prototype system

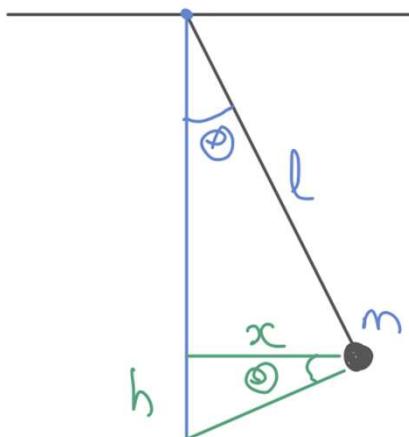
- Consider two pendula coupled by a spring
  - define displacement for each pendulum
  - calculate forces given displacement
  - find equation of motion for each pendulum
- New element: spring force depends on *both* displacements
  - motion of each pendulum depends on the other
  - equation of motion of each depends on position of both
- For small angles, we can neglect the vertical component of the motion.....



# Prototype system: two pendula coupled by a spring

- Use horizontal displacement rather than angle

For small angles, pendulum motion is almost horizontal...



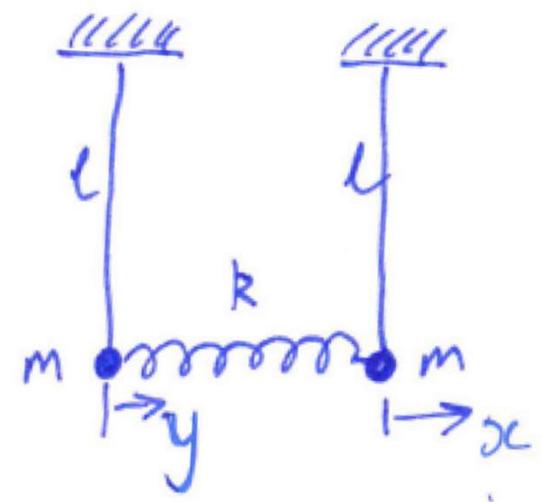
$$x = l \sin \theta \doteq l \theta$$

$$h = x \tan \theta \doteq l \theta^2$$

$$\therefore h \ll x$$

... with restoring force  $-(mg/l)x$

$$\Rightarrow \ddot{x} + (g/l)x = 0$$



Define horizontal displacements  $y$  and  $x$

Extension of spring is  $x - y$

$$F_S = k(x - y)$$

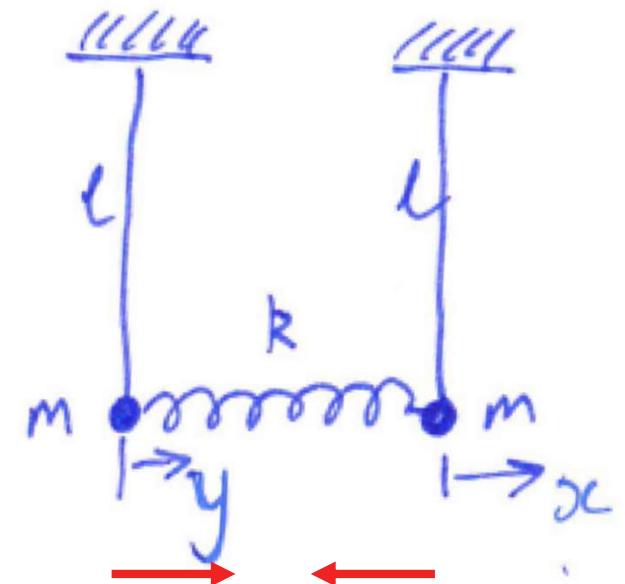
# Prototype system

- calculate forces given displacement
- find equation of motion for each pendulum

$$m\ddot{x} = -(mg/l)x - k(x - y)$$

$$m\ddot{y} = -(mg/l)y + k(x - y)$$

Coupled equations



Spring extension is  $(x - y)$   
Gives +ve force on y mass  
Gives negative force on x mass

## Decoupling the equations

$$m\ddot{x} = -(mg/l)x - k(x - y) \quad (\text{A})$$

$$m\ddot{y} = -(mg/l)y + k(x - y) \quad (\text{B})$$

$$(\text{A}) + (\text{B}) : m(\ddot{x} + \ddot{y}) = -(mg/l)(x + y)$$

$$\therefore \ddot{X} = -(g/l)X \quad \text{where } X = x + y$$

$$(\text{A}) - (\text{B}) : m(\ddot{x} - \ddot{y}) = -(mg/l)(x - y) - 2k(x - y)$$

$$\therefore \ddot{Y} = -(g/l + 2k/m)Y \quad \text{where } Y = x - y$$

- New variables  $X$  and  $Y$  undergo *uncoupled SHM!*

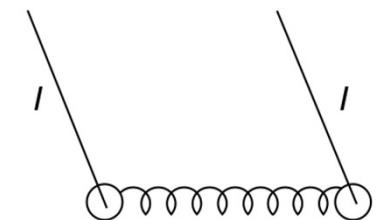
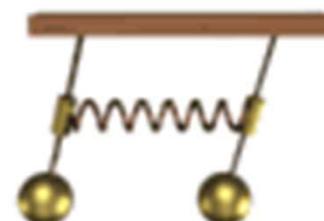
# Normal variables and normal coordinates

Writing  $X = x + y$ :  $\ddot{X} = -(g/l) X$

Writing  $Y = x - y$ :  $\ddot{Y} = -(g/l + 2k/m) Y$

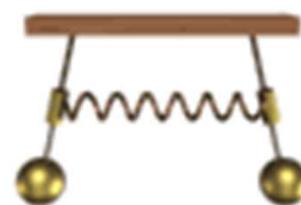
These are the  
*normal modes*.

- If  $Y, \dot{Y} = 0$ , then  $x = y$ , we have a pure oscillation of the  $X = x + y$  normal mode



In-phase  $X$  mode

- If  $X, \dot{X} = 0$ , then  $x = -y$ , and we have a pure oscillation of the  $Y = x - y$  normal mode



Out-of-phase  $Y$  mode

# Normal mode frequencies

$$\ddot{X} = -(g/l) X$$

$$\omega_0^X = (g/l)^{1/2}$$

$$X(t) = A^X \cos(\omega_0^X t + \phi^X)$$

$$\dot{X}(t) = -A^X \omega_0^X \sin(\omega_0^X t + \phi^X)$$

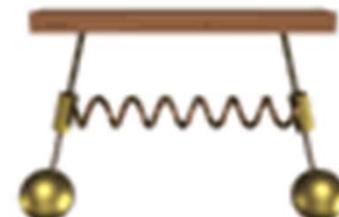
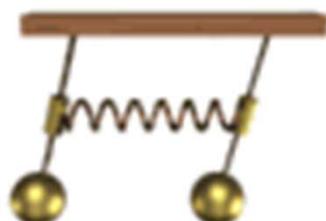
$$\ddot{Y} = -(g/l + 2k/m) Y$$

$$\omega_0^Y = (g/l + 2k/m)^{1/2}$$

$$Y(t) = A^Y \cos(\omega_0^Y t + \phi^Y)$$

$$\dot{Y}(t) = -A^Y \omega_0^Y \sin(\omega_0^Y t + \phi^Y)$$

- Different modes have different frequency
  - The lowest frequency mode has oscillators in phase
  - The highest frequency mode has oscillators out of phase



# Independence of normal modes

$$\ddot{X} = -(g/l) X$$

$$\ddot{Y} = -(g/l + 2k/m) Y$$

- Note that if  $Y, \dot{Y} = 0$  at any instant then this will remain the case forever
- Equally, if  $X, \dot{X} = 0$ , then this will remain the case forever
- Further, the equation for the  $X$  oscillation do not depend on  $Y$  or  $\dot{Y}$ , and vice-versa

**The normal modes are entirely independent  
They do not exchange energy**

- Total energy is  $(aX^2 + b\dot{X}^2) + (cY^2 + d\dot{Y}^2)$

- A pure  $X$  or  $Y$  oscillation continues forever (if no damping!)
- General solution: superposition of  $X$  and  $Y$  normal modes
  - magnitude of the  $X$  and  $Y$  contributions doesn't change with time
  - beating occurs as the normal modes have different frequencies!
  - initial conditions (four!) determine the amplitudes and phases of the  $X$  and  $Y$  contributions
    - express the initial conditions in the normal coordinates ( $X$  and  $Y$  here)
    - find the correct superposition of the normal modes
    - retrieve the corresponding solution for the physical variables

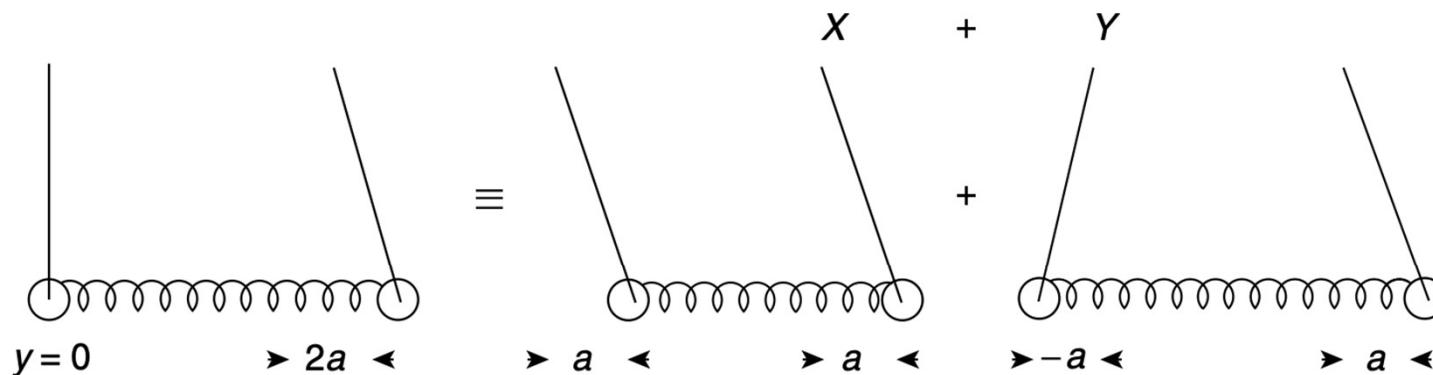
## Example: energy transfer between coupled pendula

Find the behaviour of two coupled pendula if  $x(0) = 2a \quad y(0) = 0$   
 $\dot{x}(0) = 0 \quad \dot{y}(0) = 0$

Find the initial conditions for the normal modes  $X = x + y$  and  $Y = x - y$   
 $X(0) = 2a \quad Y(0) = 2a$   
 $\dot{X}(0) = 0 \quad \dot{Y}(0) = 0$

Find the particular solution for the normal modes

$$X(t) = A^X \cos(\omega_0^X t + \phi^X) \Rightarrow X(t) = 2a \cos(\omega_0^X t)$$
$$Y(t) = A^Y \cos(\omega_0^Y t + \phi^Y) \Rightarrow Y(t) = 2a \cos(\omega_0^Y t)$$



**Figure 4.3** The displacement of one pendulum by an amount  $2a$  is shown as the combination of the two normal coordinates  $X + Y$

## Example (cont)

Translate back to the physical displacements  $x$  and  $y$

$$x = (X+Y)/2$$

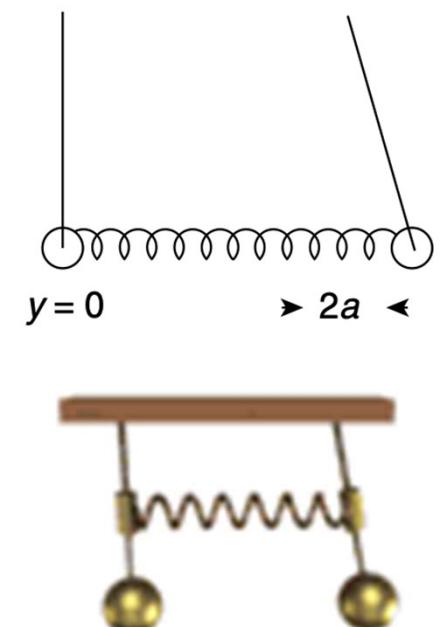
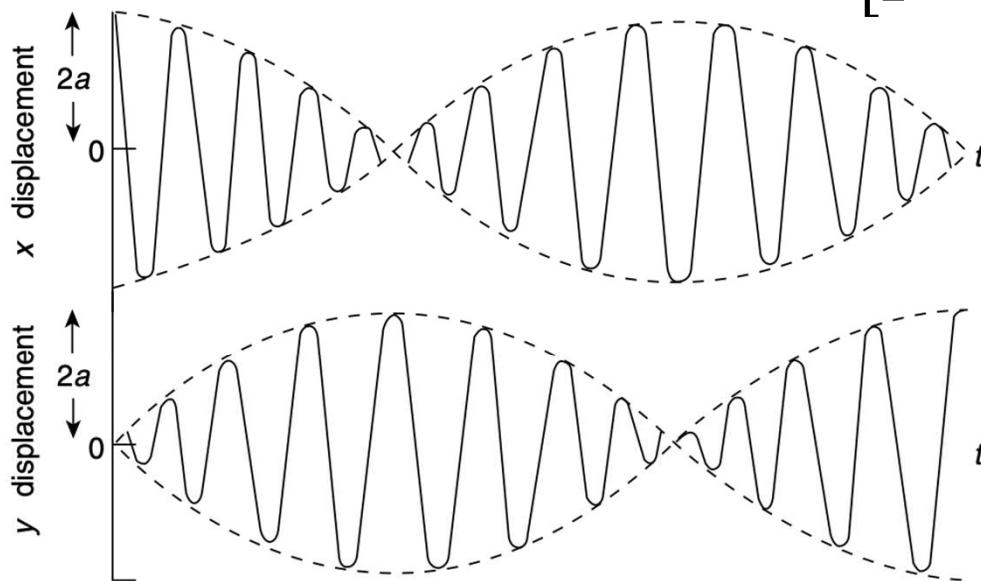
$$X(t) = 2a \cos(\omega_0^X t)$$

$$y = (X-Y)/2$$

$$Y(t) = 2a \cos(\omega_0^Y t)$$

$$x(t) = a\cos(\omega_0^X t) + a\cos(\omega_0^Y t) = 2a\cos\left[\frac{1}{2}(\omega_0^X + \omega_0^Y)t\right]\cos\left[\frac{1}{2}(\omega_0^X - \omega_0^Y)t\right]$$

$$y(t) = a\cos(\omega_0^X t) - a\cos(\omega_0^Y t) = -2a\sin\left[\frac{1}{2}(\omega_0^X + \omega_0^Y)t\right]\sin\left[\frac{1}{2}(\omega_0^X - \omega_0^Y)t\right]$$



**Figure 4.4** Behaviour with time of individual pendulums, showing complete energy exchange between the pendulums as  $x$  decreases from  $2a$  to zero whilst  $y$  grows from zero to  $2a$

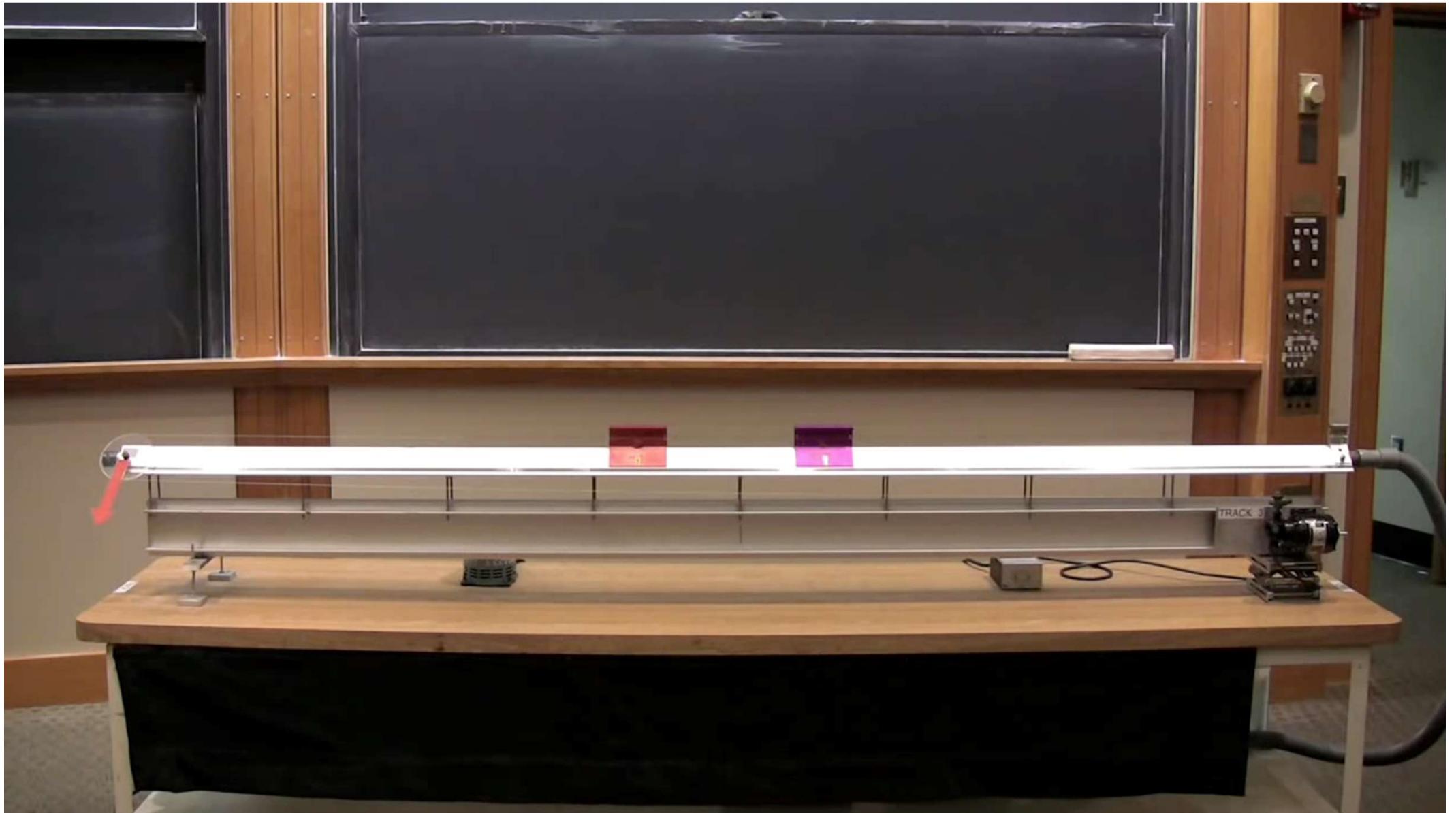
# Properties of normal modes

- For a normal mode
  - Oscillation is independent of other normal modes coordinates
  - Energy is not transferred to other normal modes
  - All oscillators have a fixed phase relation
  - All oscillators have fixed amplitude ratio
  - All oscillators share the same frequency
- Specifying a normal mode
  - Normal mode *dependent variables*  $X, Y$
  - Normal *coordinates*  $X, \dot{X}, Y, \dot{Y}$
  - Normal (mode) frequency  $\omega_0^X, \omega_0^Y$

# How many normal modes?

- The number of normal modes reflects the number of oscillators,  $n$ 
  - $n$  normal modes, each with
    - an associated normal variable  $X, Y, Z\dots$  that is a linear combination of displacements  $x,y,z\dots$
    - a characteristic frequency
    - 2 initial conditions

# Coupled air carts - normal modes



# General method to find normal modes: the coupled pendulum

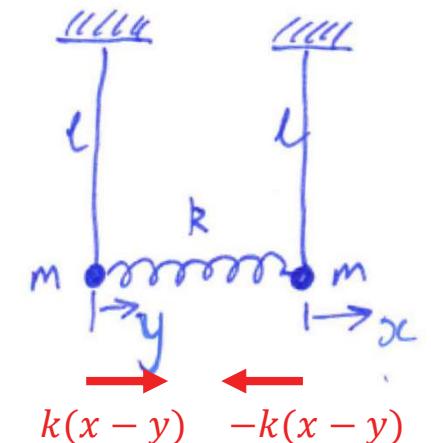
- Difficult/impossible to decouple equations “by eye”
- Instead, seek solutions that oscillate with frequency  $\omega$
- For the coupled pendulum,  $x = A e^{i\omega t}$   $y = B e^{i\omega t}$

$$\ddot{x} = -(g/l)x - (k/m)(x - y)$$

$$\Rightarrow A\omega^2 = (g/l + k/m)A - (k/m)B$$

$$\ddot{y} = -(g/l)y + (k/m)(x - y)$$

$$\Rightarrow B\omega^2 = (g/l + k/m)B - (k/m)A$$



Solve for  $\omega$  to find normal frequencies

# Finding $\omega$

$$A\omega^2 = (g/l + k/m)A - (k/m)B$$

$$B\omega^2 = (g/l + k/m)B - (k/m)A$$

- Two equations, linear in A and B
  - system determines the coefficients on the right hand side
    - $g, l, m, k$  — gravity, pendulum length and mass, spring
  - but there are three unknowns -  $A, B, \omega$
  - random choice of  $\omega \Rightarrow$  only solution is  $A = B = 0 !!$
  - certain choices of  $\omega$ : equations are *not* independent
    - i.e. 2nd equation = multiple of the first
    - either equation  $\Rightarrow$  proportionality between A and B
    - $g, l, m, k$  determine form and frequency of oscillations,  
but *not* the overall amplitude

## Finding $\omega$

$$A\omega^2 = (g/l + k/m)A - (k/m)B \quad B\omega^2 = (g/l + k/m)B - (k/m)A$$

$$(\omega^2 - g/l - k/m)A/B = -k/m \quad (\omega^2 - g/l - k/m) = -(k/m)A/B$$

$$\frac{A}{B} = \frac{-k/m}{\omega^2 - g/l - k/m}$$

$$\frac{\omega^2 - g/l - k/m}{-k/m} = \frac{A}{B}$$

$$\therefore (\omega^2 - g/l - k/m)^2 = (k/m)^2$$

$$\omega^2 - g/l - k/m = \pm k/m$$

$$\omega^2 = g/l, g/l + 2k/m$$

$$\omega = \sqrt{g/l}, \sqrt{g/l + 2k/m} \quad (\text{require } \omega \geq 0)$$

## Solving with matrices

$$A\omega^2 = (g/l + k/m)A - (k/m)B$$

$$B\omega^2 = (g/l + k/m)B - (k/m)A$$

Rewrite in matrix form as

$$\begin{bmatrix} g/l + k/m & -k/m \\ -k/m & g/l + k/m \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \omega^2 \begin{bmatrix} A \\ B \end{bmatrix}$$

(The general form  $\mathbf{M}\mathbf{v} = \lambda\mathbf{v}$  is called an *eigenvalue equation*.)

$$\begin{bmatrix} g/l + k/m - \omega^2 & -k/m \\ -k/m & g/l + k/m - \omega^2 \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

# "Reminder": Matrices

$$\underbrace{\begin{pmatrix} \cdot & \cdot \end{pmatrix}}_M \begin{pmatrix} A \\ B \end{pmatrix} = 0$$

$$M = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

$$\Delta = ad - bc$$

$$M^{-1} = \frac{1}{\Delta} \begin{pmatrix} d & -b \\ -c & a \end{pmatrix}$$

$$MM^{-1} = M^{-1}M = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$Mv = 0$$

$$\underbrace{M^{-1}M}_ {\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}} v = M^{-1}0 = 0$$

$$\therefore v = \underline{0}$$

## General solution with matrices...

$$\begin{bmatrix} g/l + k/m - \omega^2 & -k/m \\ -k/m & g/l + k/m - \omega^2 \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

If determinant is nonzero, matrix has an inverse and  $A=B=0$ , so instead we want

$$\begin{vmatrix} g/l + k/m - \omega^2 & -k/m \\ -k/m & g/l + k/m - \omega^2 \end{vmatrix} = 0$$

$$(g/l + k/m - \omega^2)^2 - (k/m)^2 = 0$$

$$(g/l + k/m - \omega^2) = \pm(k/m)$$

$$\Rightarrow \omega^2 = g/l, \text{ or } (g/l + 2k/m)$$

## General solution with matrices...

$$\omega^2 = g/l, \text{ or } (g/l + 2k/m)$$

...substituting  $\omega^2$  back

$$A\omega^2 = (g/l + k/m)A - (k/m)B$$

$$B\omega^2 = (g/l + k/m)B - (k/m)A$$

$$A\omega^2 = (g/l + k/m)A - k/m \cdot B$$

$$A = B \text{ for } \omega^2 = g/l, \text{ and } A = -B \text{ for } \omega^2 = (g/l + 2k/m)$$

Eigenvectors describe relationship between normal and physical variables

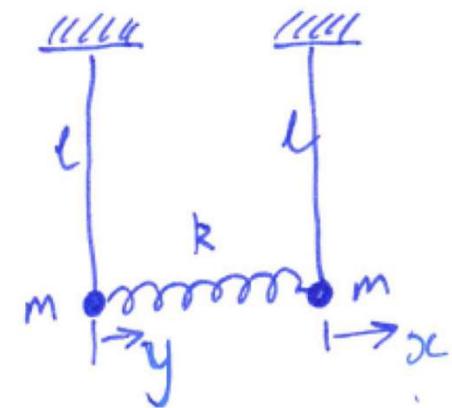
# Summary: Finding normal modes with matrices

$$m\ddot{x} = -(mg/l)x - k(x - y)$$

$$m\ddot{y} = -(mg/l)y + k(x - y)$$

- Seek solutions that oscillate with frequency  $\omega$

$$x = A e^{i\omega t} \quad y = B e^{i\omega t}$$



$$\begin{bmatrix} g/l + k/m & -k/m \\ -k/m & g/l + k/m \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \omega^2 \begin{bmatrix} A \\ B \end{bmatrix}$$

Eigenvalue equation

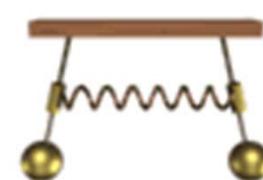
eigenvector

eigenvalue

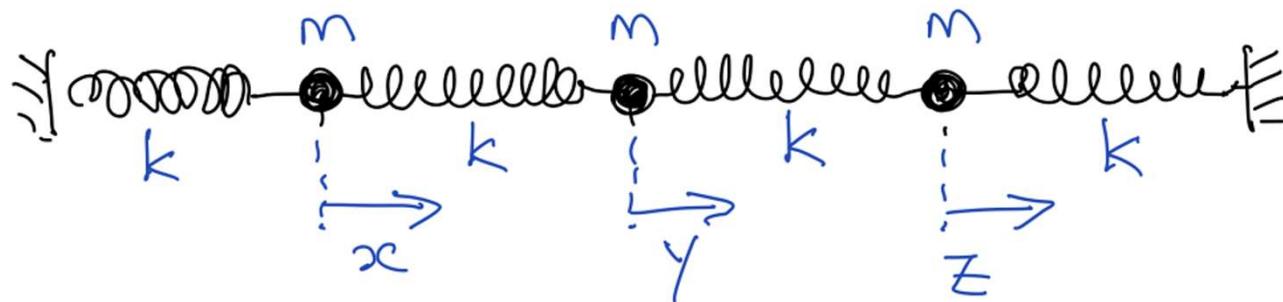


$$\omega^2 = g/l$$
  
$$A = B$$

$$g/l + 2k/m$$
  
$$A = -B$$



# Example — two walls, three masses and four springs



$$m\ddot{x} = -kx + k(y-x) = -2kx + ky$$

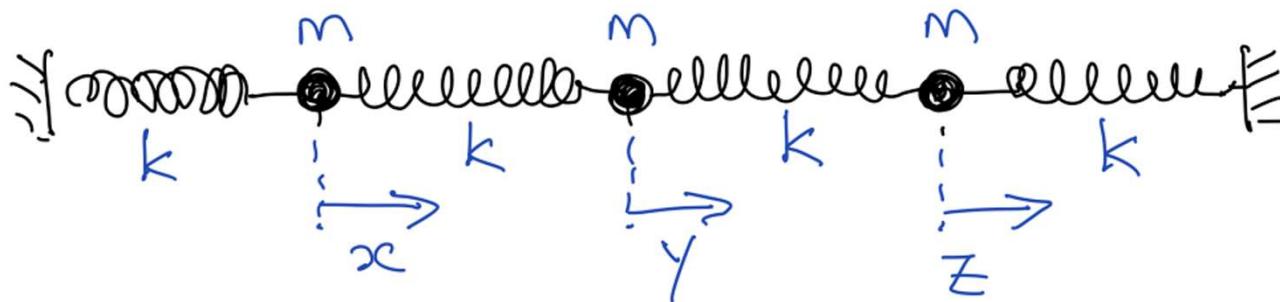
$$m\ddot{y} = -k(y-x) + k(z-y) = kx - 2ky + kz$$

$$m\ddot{z} = -k(z-y) - kz = ky - 2kz$$

Now assume  $x = A e^{i\omega t}$   
 $y = B e^{i\omega t}$   
 $z = C e^{i\omega t}$

Then  $\ddot{x} = -\omega^2 x$   
 $\ddot{y} = -\omega^2 y$   
 $\ddot{z} = -\omega^2 z$

# Example — two walls, three masses and four springs



$$m\ddot{x} = -2kx + ky$$

$$m\ddot{y} = kx - 2ky + kz$$

$$m\ddot{z} = ky - 2kz$$

$$\ddot{x} = -\omega^2 x$$

$$\ddot{y} = -\omega^2 y$$

$$\ddot{z} = -\omega^2 z$$

$$x = A e^{i\omega t}$$

$$y = B e^{i\omega t}$$

$$z = C e^{i\omega t}$$

$$-\omega^2 \begin{pmatrix} A \\ B \\ C \end{pmatrix} = \frac{k}{m} \begin{pmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ 0 & 1 & -2 \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix}$$

Put  $\omega^2 = \lambda(k/m)$   $\Rightarrow$   $\begin{pmatrix} \lambda-2 & 1 & 0 \\ 1 & \lambda-2 & 1 \\ 0 & 1 & \lambda-2 \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$

# Example — two walls, three masses and four springs

$$m\ddot{x} = -2kx + ky$$

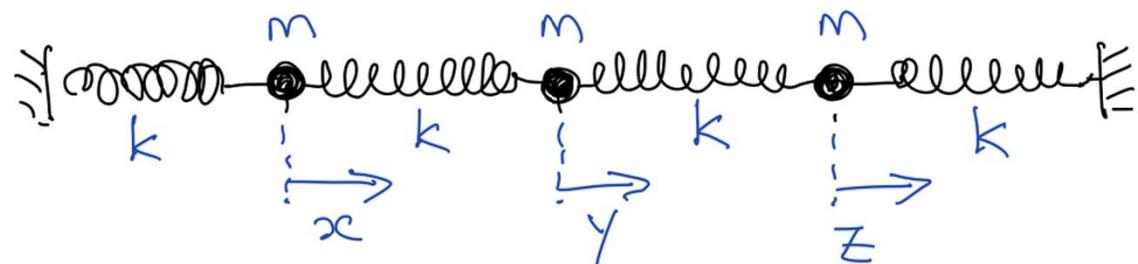
$$m\ddot{y} = kx - 2ky + kz$$

$$m\ddot{z} = ky - 2kz$$

$$x = A e^{i\omega t}$$

$$y = B e^{i\omega t} \quad \lambda = \omega^2/(k/m)$$

$$z = C e^{i\omega t}$$



$$\begin{pmatrix} \lambda-2 & 1 & 0 \\ 1 & \lambda-2 & 1 \\ 0 & 1 & \lambda-2 \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$$

$$0 = \begin{vmatrix} \lambda-2 & 1 & 0 \\ 1 & \lambda-2 & 1 \\ 0 & 1 & \lambda-2 \end{vmatrix} = (\lambda-2)[(\lambda-2)^2 - 1] - (\lambda-2) = (\lambda-2)[(\lambda-2)^2 - 2]$$

$$\therefore \lambda = 2, 2 \pm \sqrt{2} \quad \text{i.e. } \omega^2 = 2k/m, (2 \pm \sqrt{2})k/m$$

# Example — two walls, three masses and four springs

$$m\ddot{x} = -2kx + ky$$

$$m\ddot{y} = kx - 2ky + kz$$

$$m\ddot{z} = ky - 2kz$$

$$x = A e^{i\omega t}$$

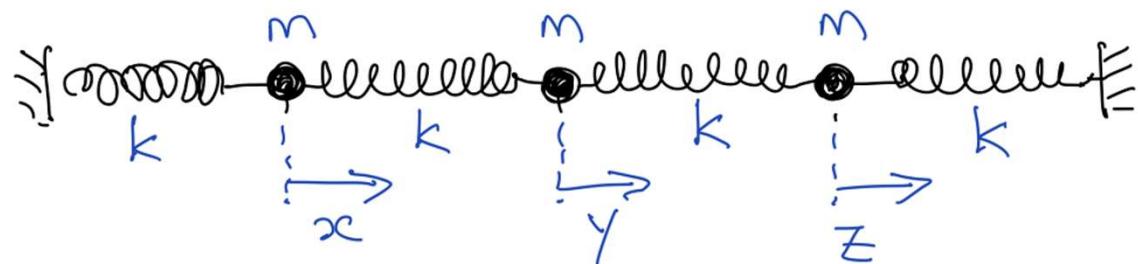
$$y = B e^{i\omega t} \quad \lambda = \omega^2/(k/m)$$

$$z = C e^{i\omega t}$$

$$\begin{pmatrix} \lambda-2 & 1 & 0 \\ 1 & \lambda-2 & 1 \\ 0 & 1 & \lambda-2 \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$$

$$\therefore \lambda = 2, 2 \pm \sqrt{2}$$

$$\text{i.e. } \omega^2 = 2k/m, (2 \pm \sqrt{2})k/m$$

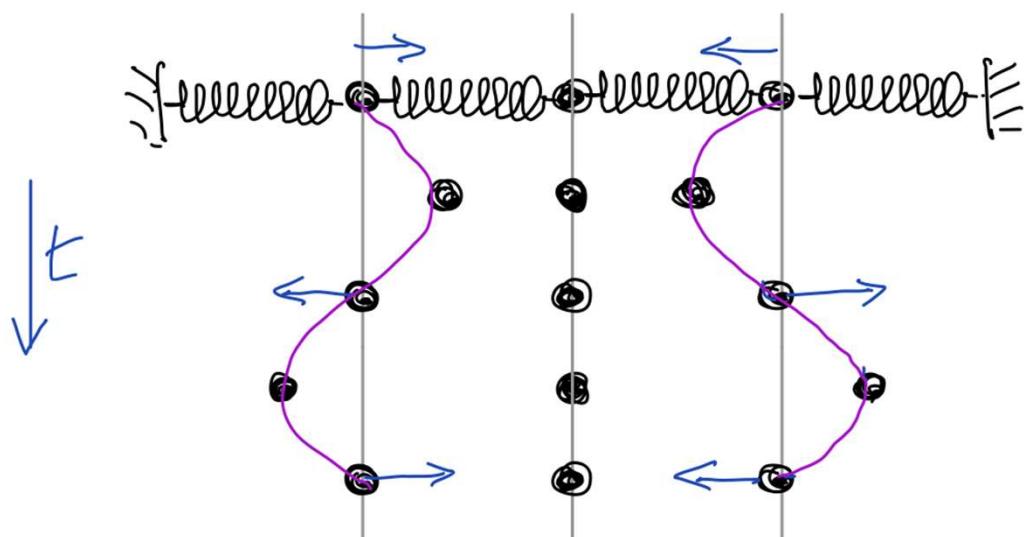


$$\underline{\lambda = 2}$$

$$\begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$$

$$\omega^2 = 2k/m$$

$$B=0 \quad C=-A$$



# Example — two walls, three masses and four springs

$$m\ddot{x} = -2kx + ky$$

$$m\ddot{y} = kx - 2ky + kz$$

$$m\ddot{z} = ky - 2kz$$

$$x = A e^{i\omega t}$$

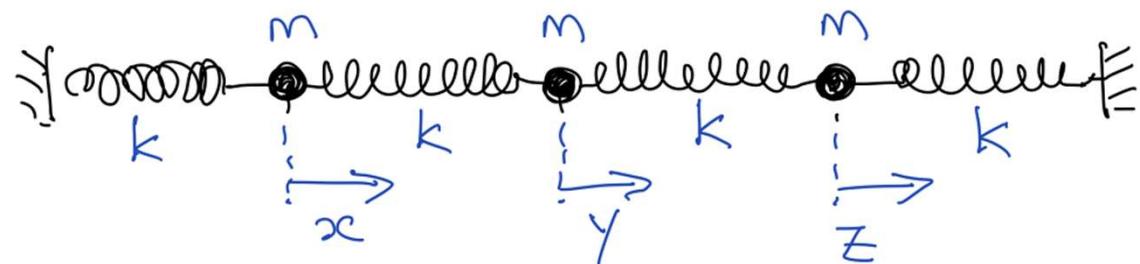
$$y = B e^{i\omega t} \quad \lambda = \omega^2/(k/m)$$

$$z = C e^{i\omega t}$$

$$\begin{pmatrix} x-2 & 1 & 0 \\ 1 & x-2 & 1 \\ 0 & 1 & \lambda-2 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} = 0$$

$$\lambda = 2, 2 \pm \sqrt{2}$$

$$\omega^2 = 2\omega_0^2, (2 \pm \sqrt{2})\omega_0^2$$



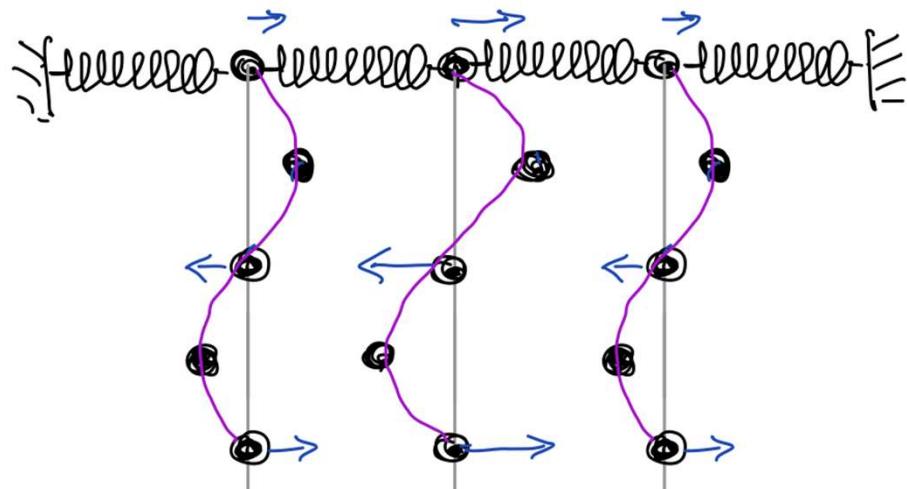
$$\lambda = 2 - \sqrt{2}$$

$$\begin{pmatrix} -\sqrt{2} & 1 & 0 \\ 1 & -\sqrt{2} & 1 \\ 0 & 1 & -\sqrt{2} \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$$

$$\omega^2 = (2 - \sqrt{2}) k/m$$

$$B = \sqrt{2} A$$

$$C = A$$



# Example — two walls, three masses and four springs

$$m\ddot{x} = -2kx + ky$$

$$m\ddot{y} = kx - 2ky + kz$$

$$m\ddot{z} = ky - 2kz$$

$$x = A e^{i\omega t}$$

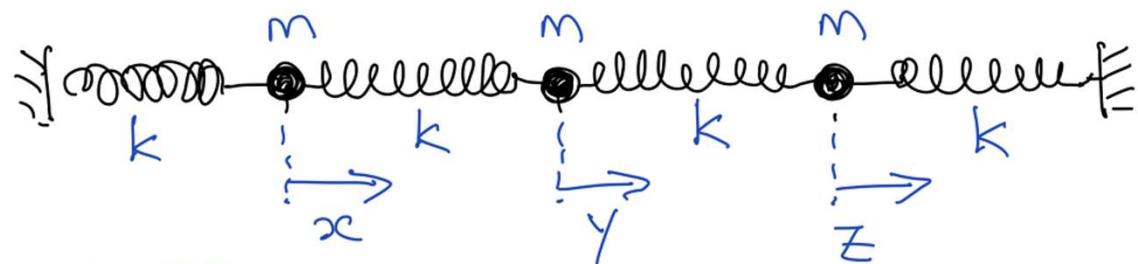
$$y = B e^{i\omega t} \quad \lambda = \omega^2/(k/m)$$

$$z = C e^{i\omega t}$$

$$\begin{pmatrix} x-2 & 1 & 0 \\ 1 & x-2 & 1 \\ 0 & 1 & \lambda-2 \end{pmatrix} \begin{pmatrix} x \\ y \\ z \end{pmatrix} = 0$$

$$\lambda = 2, 2 \pm \sqrt{2}$$

$$\omega^2 = 2\omega_0^2, (2 \pm \sqrt{2})\omega_0^2$$



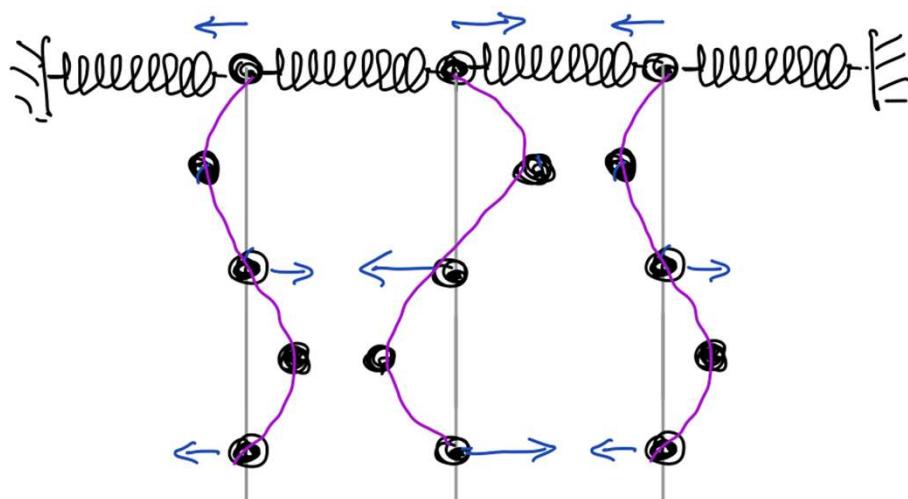
$$\underline{\lambda = 2 + \sqrt{2}}$$

$$\begin{pmatrix} \sqrt{2} & 1 & 0 \\ 1 & \sqrt{2} & 1 \\ 0 & 1 & \sqrt{2} \end{pmatrix} \begin{pmatrix} A \\ B \\ C \end{pmatrix} = 0$$

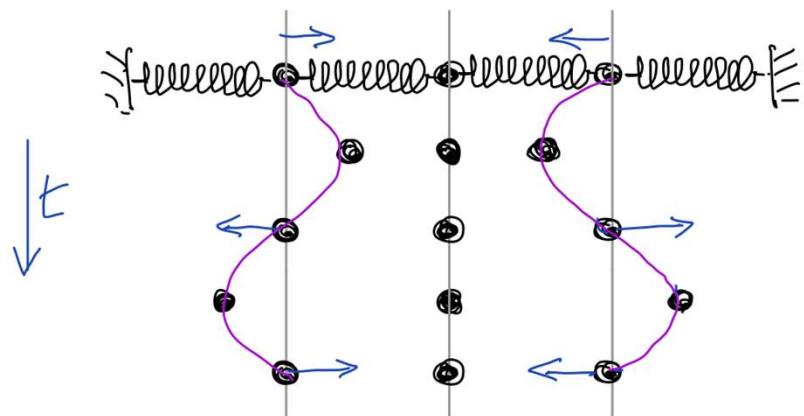
$$\omega^2 = (2 + \sqrt{2}) k/m$$

$$B = -\sqrt{2}A$$

$$C = A$$



# Example — two walls, three masses and four springs



$$\omega^2 = 2 \text{ k/m}$$

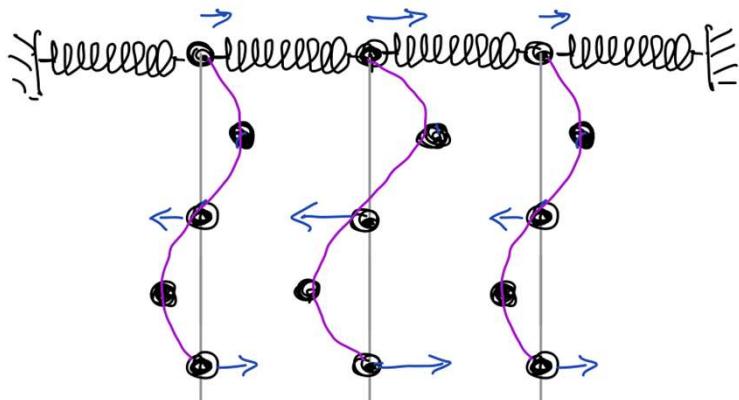
$$B = 0$$

$$C = -A$$

$$x = A e^{i\omega t}$$

$$y = B e^{i\omega t}$$

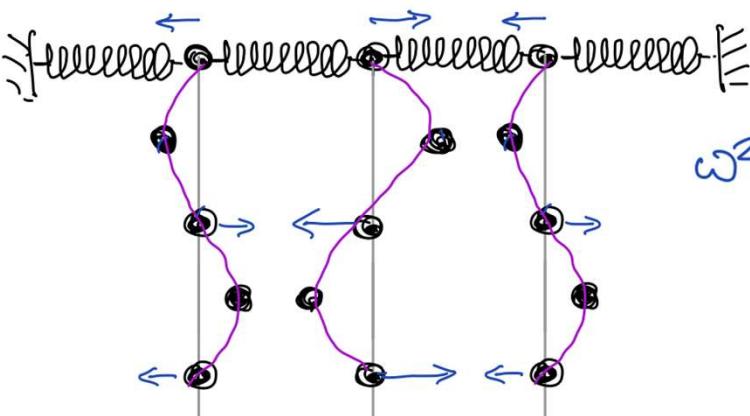
$$z = C e^{i\omega t}$$



$$\omega^2 = (2 - \sqrt{2}) \text{ k/m}$$

$$B = \sqrt{2} A$$

$$C = A$$



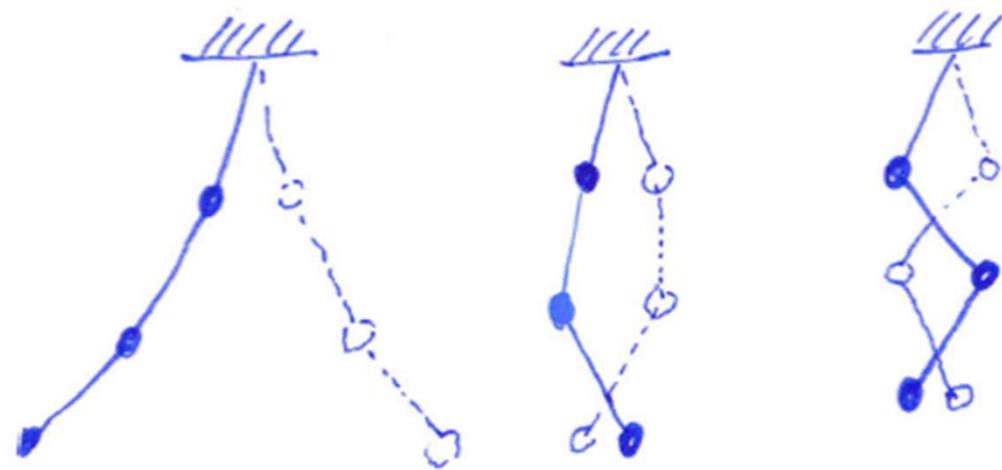
$$\omega^2 = (2 + \sqrt{2}) \text{ k/m}$$

$$B = -\sqrt{2} A$$

$$C = A$$

- Number of modes = number of oscillators
- Lowest frequency mode: all masses are in-phase
- Highest frequency mode: adjacent masses are out-of-phase

## Examples – more!



- Number of modes = number of oscillators
- Lowest frequency mode: all masses are in-phase
- Highest frequency mode: adjacent masses are out-of-phase

