VE230 Electromagnetics

Chapter 3

September 11, 2022



Exercise 3.1

Refer to Fig. 3-4.

- a) Find the relation between the angle of arrival, α , of the electron beam at the screen and the deflecting electric field intensity $E_{\rm d}$.
 - b) Find the relation between w and L such that $d_1 = d_0/20$.

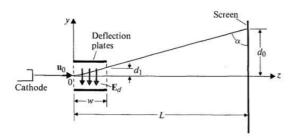


FIGURE 3-4
Electrostatic deflection system of a cathode-ray oscilloscope (Example 3-3)

Answer:

(a)

$$\alpha = \tan^{-1}\left(\frac{L-w}{d_0 - d_1}\right) = \tan^{-1}\left(\frac{L-w}{d_2}\right) = \tan^{-1}\left(\frac{mu_0^2}{ewE_d}\right)$$

(b)

$$d_1 = \frac{d_0}{20} \longrightarrow \frac{eE_d}{2m} \left(\frac{w}{u_0}\right)^2 = \frac{1}{20} \frac{eE_d}{mu_0^2} w \left(L - \frac{w}{2}\right)$$
$$L/w = 10.5$$

Exercise 3.2

The cathode-ray oscilloscope (CRO) shown in Fig. 3-4 is used to measure the voltage applied to the parallel deflection plates.

- a) Assuming no breakdown in insulation, what is the maximum voltage that can be measured if the distance of separation between the plates is h?
 - b) What is the restriction on L if the diameter of the screen is D?
- c) What can be done with a fixed geometry to double the CRO's maximum measurable voltage?

a) Max. voltage V max. will make
$$d_1 = h/2$$
 at $z = w$. $\frac{h}{2} = \frac{e}{2m} \left(\frac{V_{max}}{h} \right) \left(\frac{w}{u_0} \right)^2 \longrightarrow V_{max} = \frac{m}{e} \left(\frac{u_0 h}{v_v} \right)^2$.

b) At the screen, $(d_0)_{\text{max}} = D/2$. L mast be $\leq L_{\text{max}}$, Where

$$L_{\text{max}} = \frac{1}{2} \left(w + \frac{m u_0^2 Dh}{e w V_{\text{max}}} \right).$$

c) Double V_{max} by doubling u_0^2 , or doubling the anode accelerating voltage.

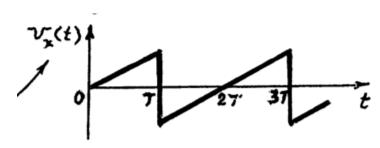
Exercise 3.3

The deflection system of a cathode-ray oscilloscope usually consists of two pairs of parallel plates producing orthogonal electric fields. Assume the presence of another set of plates in Fig. 3-4 that establishes a uniform electric field $\mathbf{E}_x = \mathbf{a}_x E_x$ in the deflection region. Deflection voltages $v_x(t)$ and $v_y(t)$ are applied to produce \mathbf{E}_x and \mathbf{E}_y , respectively. Determine the types of waveforms that $v_x(t)$ and $v_y(t)$ should have if the electrons are to trace the following graphs on the fluorescent screen:

- a) a horizontal line,
- b) a straight line having a negative unity slope,
- c) a circle,
- d) two cycles of a sine wave.

Answer:

- a) $v_x(t) = V_0 \cos \omega t$ (or $V_0 \sin \omega t$), $v_y = 0$.
- b) $v_x(t) = V_0 \cos \omega t$, $v_y(t) = -V_0 \cos \omega t$.
- c) $v_x(t) = V_0 \cos \omega t$, $v_y(t) = V_0 \sin \omega t$.
- d) $v_y(t) = V_0 \sin \frac{2\pi}{T} t$,
- $v_x(t)$: Periodic saw-tooth wave with period 2T.



Exercise 3.4

Write a short article explaining the principle of operation of xerography. (Use library resources if needed.)

Answer:

略

Exercise 3.5

Two point charges, Q_1 and Q_2 , are located at (1,2,0) and (2,0,0), respectively. Find the relation between Q_1 and Q_2 such that the total force on a test charge at the point P(-1,1,0) will have

- a) no x-component,
- b) no y-component.

Answer:

$$\overrightarrow{Q_1P} = -\overline{a}_x 2 - \overline{a}_y; \quad \overrightarrow{Q_2P} = -\overline{a}_x 3 + \overline{a}_y$$

$$\overline{E}_{p_1} = \frac{Q_1}{4\pi\epsilon_0(\sqrt{5})^3} \left(-\overline{a}_x 2 - \overline{a}_y \right); \overline{E}_{p_2} = \frac{Q_2}{4\pi\epsilon_0(\sqrt{10})^3} \left(-\overline{a}_x 3 + \overline{a}_y \right)$$

- a) No x-component of $\bar{E}_p: -\frac{2Q_1}{(\sqrt{5})^3} \frac{3Q_2}{(\sqrt{10})^3} = 0$, or $\frac{Q_1}{Q_2} = -\frac{3}{4\sqrt{2}}$. b) No y-component of $\bar{E}_p: -\frac{Q_1}{(\sqrt{5})^3} + \frac{Q_2}{(\sqrt{10})^3} = 0$, or $\frac{Q_1}{Q_2} = \frac{1}{2\sqrt{2}}$.

Exercise 3.6

Two very small conducting spheres, each of a mass 1.0×10^{-4} (kg), are suspended at a common point by very thin nonconducting threads of a length 0.2(m). A charge Q is placed on each sphere. The electric force of repulsion separates the spheres, and an equilibrium is reached when the suspending threads make an angle of 10°. Assuming a gravitational force of 9.80 (N/kg) and a negligible mass for the threads, find Q

Answer:

At equilibrium, electric force F_e and gravitational force F_9 must add to give a resultant along the thread.

$$\frac{F_e}{F_g} = \tan 5^\circ = 0.0875$$

$$F_g = mg = 9.80 \times 10^{-4} (\text{N})$$

$$F_e = \frac{Q^2}{4\pi\varepsilon_0 (2 \times 0.2 \sin 5^\circ)^2} = 7.41 \times 10^{-12} Q^2 (N)$$

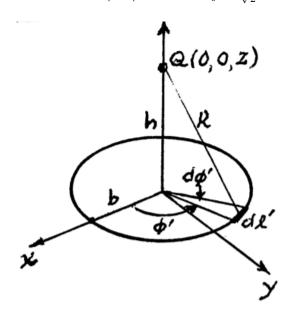
$$\longrightarrow Q = 3.40 (\text{nC})$$

Exercise 3.7

Find the force between a charged circular loop of radius b and uniform charge density ρ_{ℓ} and a point charge Q located on the loop axis at a distance h from the plane of the loop. What is the force when $h \gg b$, and when h = 0? Plot the force as a function of h.

$$\begin{split} dV_Q &= \frac{\rho_l b d\phi'}{4\pi E_0 \left(z^2 + b^2\right)^{1/2}}.\\ V_Q &= \frac{\rho_l b}{4\pi \varepsilon_0 \left(z^2 + b^2\right)^{1/2}} \int_0^{2\pi} d\phi' = \frac{\rho_l b}{2E_0 \left(z^2 + b^2\right)^{1/2}}\\ E_Q &= -\bar{\nabla} V = -\bar{a}_z \frac{dV_a}{dz} = a_z \frac{\rho_l bz}{2\varepsilon_0 \left(z^2 + b^2\right)^{3/2}}\\ \bar{F}_Q &= Q\bar{E}_Q = \bar{a}_z \frac{Q\rho_l bz}{2E_0 \left(z^2 + b^2\right)^{3/2}}. \end{split}$$

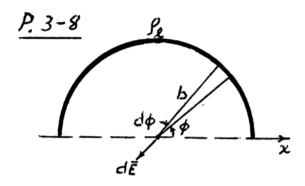
(Direction depending on whether Q is above or below the loop.) When $h \gg b$, $|\bar{F}_Q| \cong \frac{QP_lb}{2\epsilon_0h^2}$ (inverse-square law). When h = 0, $\bar{F}_Q = 0$, Max. $|\bar{F}_Q|$ occurs at $\frac{z}{b} = \frac{1}{\sqrt{2}}$.



Exercise 3.8

A line charge of uniform density ρ_{ℓ} in free space forms a semicircle of radius b. Determine the magnitude and direction of the electric field intensity at the center of the semicircle.

$$\begin{split} dE_y &= -\frac{\rho_l(bd\phi)}{4 \cdot \pi \epsilon_0 b^2} \sin \phi_1 \\ E &= \bar{a}_y E_y = -\bar{a}_y \frac{\rho_l}{4\pi \epsilon_0 b} \int_0^{\pi} \sin \phi d\phi \\ &= -\bar{a}_y \frac{\rho_l}{2\pi \epsilon_0 b} \end{split}$$

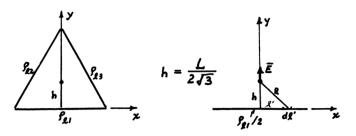


Three uniform line charges $-\rho_{\ell 1}$, $\rho_{\ell 2}$, and $\rho_{\ell 3}$, each of length L- form an equilateral triangle. Assuming that $\rho_{\ell 1}=2\rho_{\ell 2}=2\rho_{\ell 3}$, determine the electric field intensity at the center of the triangle.

Answer:

E at the center of the triangle would be zero it all three line charges were of the same charge density. The problem is equivalent to that of a single line charge of density P_l , 1/2. By symmetry, there will only be a y-component

$$\bar{E} = \bar{a}_y E_y = \bar{a}_y \int_{-L/2}^{L/2} \frac{(P_{l_1}/2) \, dl'}{4\pi\epsilon_0 R^3} \left(\frac{h}{R}\right) = \bar{a}_y \int_{-L/2}^{L/2} \frac{\rho_{l_1} h dl}{8\pi\epsilon_0 \left(h^2 + l'^2\right)^{3/2}}
= \bar{a}_y \frac{3\rho_{\ell_1}}{4\pi\epsilon_0 L} = \bar{a}_y \frac{3\rho_{\ell_2}}{2\pi\epsilon_0 L}.$$



Exercise 3.10

Assuming that the electric field intensity is $\mathbf{E} = \mathbf{a}_x 100x(\text{ V/m})$, find the total electric charge contained inside

- a) a cubical volume 100(mm) on a side centered symmetrically at the origin,
- b) a cylindrical volume around the z-axis having a radius 50 (mm) and a height 100 (mm) centered at the origin.

Answer:

Use Gauss's law: $\oint_s \bar{E} \cdot d\bar{s} = Q/\epsilon_0$.

a) E is normal to the two faces at $x = \pm 0.05$ (m), where $\bar{E} = \pm \bar{a}_x 5$ and $\bar{a}_n = \pm \bar{a}_x$ respectively.

$$Q = 2\epsilon_0 (5 \times 0.1^2) = 0.1\epsilon_0 = 8.84 \times 10^{-13} (C)$$

b) $E = \bar{a}_r(100x)\cos\phi - \bar{a}_\phi(100x)\sin\phi = \bar{a}_r(100r\cos^2\phi) - \bar{a}_\phi(100r\sin^2\phi)$.

$$\oint_{s} E \cdot \bar{a}_{n} ds = \int_{0}^{2\pi} \left(1 \cos \times 0.05 \cos^{2} \phi \right) (0.1 \times 0.05) d\phi = 0.025\pi.$$

$$Q = 0.025\pi \epsilon_{0} = 6.44 \times 10^{-13} (C).$$

Exercise 3.11

A spherical distribution of charge $\rho = \rho_0 \left[1 - (R^2/b^2)\right]$ exists in the region $0 \le R \le b$. This charge distribution is concentrically surrounded by a conducting shell with inner radius $R_i(>b)$ and outer radius R_0 . Determine **E** everywhere.

Answer:

Spherical symmetry:
$$\bar{E} = \bar{a}_R E_R$$
. Apply Gauss's law.
1) $0 \leqslant R \leqslant b$. $4\pi R^2 E_{R1} = \frac{\rho_0}{\epsilon_0} \int_0^R \left(1 - \frac{R^2}{b^2}\right) + \pi R^2 dR = \frac{4\pi\rho_0}{\epsilon_0} \left(\frac{R^3}{3} - \frac{R^5}{5b^2}\right)$.

$$E_{R1} = \frac{\rho_0}{\epsilon_0} R \left(\frac{1}{3} - \frac{R^2}{5b^2} \right)$$

2)
$$b \geqslant R < R_i$$
. $4\pi R^2 = \frac{\rho_0}{\epsilon_0} \int_0^b \left(1 - \frac{R^2}{b^2}\right) 4\pi R^2 dR = \frac{5\pi\rho_0}{15\epsilon_0} b^3$,

$$E_{R_2} = \frac{2\rho_0 b^3}{15 \cdot \epsilon_0 R^2}$$

$$R_i < R < R_0$$
. $E_{R3} = 0$

4)
$$R > R_0$$
.

$$E_{R4} = \frac{2\rho_0 b^3}{15\epsilon_0 R^2}$$

Exercise 3.12

Two infinitely long coaxial cylindrical surfaces, r = a and r = b(b > a), carry surface charge densities ρ_{sa} and ρ_{sb} , respectively.

- a) Determine **E** everywhere.
- b) What must be the relation between a and b in order that **E** vanishes for r > b?

Answer:

Cylindrical symmetry: $\bar{E} = \bar{a}_r E_r$. Apply Gauss's law.

a)

$$r < a, E_r = 0$$

$$a < r < b, E_r = a\rho_{sa}/\epsilon_0 r$$

$$r > b, E_r = \frac{a\rho_{sa} + b\rho_{sb}}{\epsilon_0 r}$$

b)

$$\frac{b}{a} = -\frac{\rho_{sa}}{\rho_{sb}}$$

Exercise 3.13

Determine the work done in carrying a $-2(\mu C)$ charge from $P_1(2,1,-1)$ to $P_2(8,2,-1)$ in the field $\mathbf{E} = \mathbf{a}_x y + \mathbf{a}_y x$

- a) along the parabola $x = 2y^2$
- b) along the straight line joining P_1 and P_2 .

 $W = -q \int \bar{E} \cdot d\bar{l} = -q \int (ydx + xdy).$

a) Along the parabola $x = 2y^2$, dx = 4ydy.

$$W = -q \int_{1}^{2} 6y^{2} dy = -14q = 28(\mu J).$$

b) Along the straight line x = 6y - 4, dx = 6dy.

$$W = -q \int_{1}^{2} (12y - 4) dy = -14q = 28(\mu J).$$

Exercise 3.14

At what values of θ does the electric field intensity of a z-directed dipole have no z-component?

Answer:

Refer to

$$\mathbf{E} = -\nabla V = -\mathbf{a}_R \frac{\partial V}{\partial R} - \mathbf{a}_\theta \frac{\partial V}{R \partial \theta}$$
$$= \frac{p}{4\pi\epsilon_0 R^3} \left(\mathbf{a}_R 2 \cos \theta + \mathbf{a}_\theta \sin \theta \right)$$

E will have no z - component if $E_R \cos \theta = E_\theta \sin \theta$, or $2 \cos^2 \theta = \sin^2 \theta$. $\longrightarrow \theta = 54.7^\circ$ and 125.3° .

Exercise 3.15

Three charges (+q, -2q, and +q) are arranged along the z-axis at z = d/2, z = 0, and z = -d/2, respectively.

- a) Determine V and E at a distant point $P(R, \theta, \phi)$.
- b) Find the equations for equipotential surfaces and streamlines.
- c) Sketch a family of equipotential lines and streamlines.

(Such an arrangement of three charges is called a linear electrostatic quadrupole.)

Answer:

$$V = \frac{q}{4\pi\epsilon_0 R} \left(\frac{R}{R_1} + \frac{R}{R_2} - 2\right) \cdot$$

$$R_1^2 = R^2 + \left(\frac{d}{2}\right)^2 - Rd\cos\theta$$

$$\frac{R}{R_1} \cong \left[1 + \left(\frac{d}{2R}\right)^2 - \frac{d}{R}\cos\theta\right]^{-1/2}$$

$$\cong 1 + \frac{d}{2R}\cos\theta + \frac{d^2}{4R^2} \frac{3\cos^2\theta - 1}{2},$$

$$\frac{R}{R_2} \cong 1 - \frac{d}{2R}\cos\theta + \frac{d^2}{4R^2} \frac{3\cos^2\theta - 1}{2}$$

$$V = \frac{gd^2}{16\pi\epsilon_0 R^3} \left(3\cos^2\theta - 1\right), \quad R^3 \gg d^3.$$
a)
$$\bar{E} = -\bar{\nabla}V = -\bar{a}_R \frac{\partial V}{\partial R} - \bar{a}_\theta \cdot \frac{\partial V}{R\theta}$$

$$= \frac{3qd^2}{16\pi\epsilon_0 R^4} \left[\bar{a}_R \left(3\cos^2\theta - 1\right) + \bar{a}_\theta \sin 2\theta\right].$$

b) Equation for equipotential surfaces: $R^3 = c_V (3\cos^2\theta - 1)$ Equation for streamlines:

$$\frac{dR}{E_R} = \frac{Rd\theta}{E_{\theta}}$$

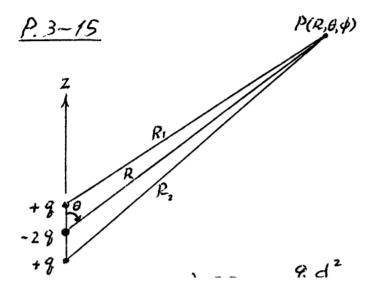
$$\frac{dR}{3\cos^2\theta - 1} = \frac{Rd\theta}{\sin 2\theta}$$

$$\frac{dR}{R} = \frac{3}{2}\frac{d(\sin\theta)}{\sin\theta} - \frac{d\theta}{\sin 2\theta}$$

$$\longrightarrow R = c_E' \frac{\sin^{3/2}\theta}{\sqrt{|\tan\theta|}}$$

$$R^2 = c_E \sin^2\theta \cos\theta$$

or



Exercise 3.16

A finite line charge of length L carrying uniform line charge density ρ_{ℓ} is coincident with the x-axis.

- a) Determine V in the plane bisecting the line charge.
- b) Determine **E** from ρ_{ℓ} directly by applying Coulomb's law.
- c) Check the answer in part (b) with $-\nabla V$.

Answer:

a)

$$\begin{split} V = & 2 \int_0^{L/2} \frac{\rho_l dx}{4\pi\epsilon_0 R} \\ = & \frac{\rho_l}{2\pi\epsilon_0} \int_0^{L/2} \frac{dx}{\sqrt{x^2 + y^2}} \\ = & \frac{\rho_l}{2\pi\epsilon_0} \{ \ln[\sqrt{(\frac{L}{2})^2 + y^2} + \frac{L}{2}] - \ln y \}. \end{split}$$

b) From Coulomb's law:

$$\bar{E} = \bar{a}_y E_y = 2 \int_0^{L/2} \bar{\alpha}_y \frac{\rho_l y dx}{4\pi \epsilon_0 R^3} := \bar{a}_y \frac{\rho_l}{2\pi \epsilon_0 y} \left[\frac{L/2}{\sqrt{(L/2)^2 + y^2}} \right]$$

c) $\bar{E} = -\bar{\nabla}V$ gives the same answer as in b).

Exercise 3.17

In Example 3-5 we obtained the electric field intensity around an infinitely long line charge of a uniform charge density in a very simple manner by applying Gauss's law. Since $|\mathbf{E}|$ is a function of r only, any coaxial cylinder around the infinite line charge is an equipotential surface. In practice, all conductors are of finite length. A finite line charge carrying a constant charge density ρ_{ℓ} along the axis, however, does not produce a constant potential on a concentric cylindrical surface. Given the finite line charge ρ_{ℓ} of length L in Fig. 3-40, find the potential on the cylindrical surface of radius b as a function of x and plot it.

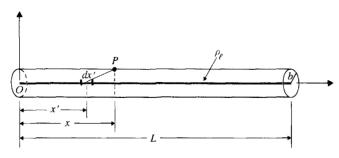
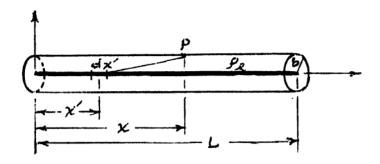


FIGURE 3-40 A finite line charge (Problem P.3-17).

$$dV_P(x) = \frac{\rho_l dx'}{4\pi\epsilon_0 \sqrt{(x-x')^2 + b^2}}$$

$$V_P(x) = \frac{\rho_l}{4\pi\epsilon_0} \int_0^L \frac{dx}{\sqrt{(x-x')^2 + b^2}}$$

$$= \frac{\rho_l}{4\pi\epsilon_0} \left[\sinh^{-1} \left(\frac{L-x}{b} \right) + \sinh^{-1} \left(\frac{x}{b} \right) \right]$$

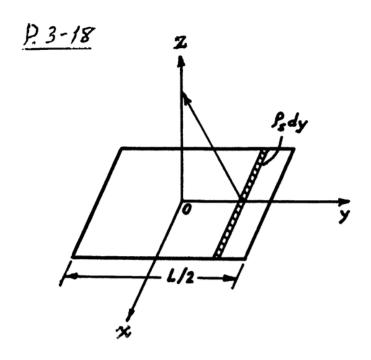


A charge Q is distributed uniformly over an $L \times L$ square plate. Determine V and \mathbf{E} at a point on the axis perpendicular to the plate and through its center.

$$V = 2 \cdot \frac{\rho_s}{2\pi\epsilon_0} \int_0^{L/2} \left\{ \ln \left[\sqrt{\left(\frac{L}{2}\right)^2 + y^2 + z^2} + \left(\frac{L}{2}\right) \right] - \ln \sqrt{y^2 + z^2} \right\} dy$$

$$= \frac{Q}{\pi\epsilon_0 L^2} \left\{ \frac{L}{2} \ln \left[\frac{\sqrt{2\left(\frac{L}{2}\right)^2 + z^2} + \frac{L}{2}}{\sqrt{2\left(\frac{L}{2}\right)^2 + z^2} - \frac{L}{2}} \right] - z \tan^{-1} \left[\frac{\left(\frac{L}{2}\right)^2}{z\sqrt{2\left(\frac{L}{2}\right)^2 + z^2}} \right] \right\}$$

$$\bar{E} = -\bar{\nabla}V = \bar{a}_z \frac{Q}{\pi\epsilon_0 L^2} \tan^{-1} \left[\frac{\left(\frac{L}{2}\right)^2}{2\sqrt{2\left(\frac{L}{2}\right)^2 + z^2}} \right]$$



A charge Q is distributed uniformly over the wall of a circular tube of radius b and height h. Determine V and \mathbf{E} on its axis

- a) at a point outside the tube, then
- b) at a point inside the tube.

Answer:

Assume the circular tube sits on the xy-plane with its axis coinciding with the z-axis. The surface charge on the tube wall is $\rho_s = Q/2\pi bh$. First find the potential along the axis at z due to a circular line, charge of density situated at z'.

$$V = \oint \frac{\rho_l dl}{4\pi\epsilon_0 R} = \int_0^{2\pi} \frac{\rho_l b d\phi'}{4\pi\epsilon_0 \sqrt{b^2 + (z - z')^2}} = \frac{\rho_l b}{2\epsilon_0 \sqrt{b^2 + (z - z')^2}}.$$

a) The expression above is the contribution dV due to a circular line charge of density $\rho_l = \rho_s dz^{\text{in}}$ at z

$$dV = \frac{\rho_s b dz'}{2\epsilon_0 \sqrt{b^2 + (z - z')^2}}.$$

At a point outside the tube, z > h:

$$V_0 = \int_{z'=0}^{z'=h} dV = \frac{b\rho_s}{2\epsilon_0} \ln \frac{z + \sqrt{b^2 + z^2}}{(z - h) + \sqrt{b^2 + (z - h)^2}}.$$
$$\bar{E}_0 = -\bar{a}_z \frac{dV}{dz} = \bar{a}_z \frac{b\rho_s}{2\epsilon_0} \left[\frac{1}{\sqrt{b^2 + (z - h)^2}} - \frac{1}{\sqrt{b^2 + z^2}} \right].$$

b) At a point inside the tube, z < h:

$$V_{i} = \frac{\rho_{s}b}{2\epsilon_{0}} \left[\int_{0}^{z} \frac{dz'}{\sqrt{b^{2} + (z - z')^{2}}} + \int_{z}^{h} \frac{dz'}{\sqrt{b^{2} + (z' - z)^{2}}} \right]$$
$$= \frac{\rho_{s}b}{2\epsilon_{0}} \ln \frac{1}{b^{2}} \left(z + \sqrt{b^{2} + z^{2}} \right) \left[(h - z) + \sqrt{b^{2} + (h - z)^{2}} \right].$$

 \bar{E}_i has the-same expression as \bar{E}_0 .

Exercise 3.20

An early model of the atomic structure of a chemical element was that the atom was a spherical cloud of uniformly distributed positive charge Ne, where N is the atomic number and e is the magnitude of electronic charge. Electrons, each carrying a negative charge -e, were considered to be imbedded in the cloud. Assuming the spherical charge cloud to have a radius R_0 and neglecting collision effects,

- a) find the force experienced by an imbedded electron at a distance r from the center;
- b) describe the motion of the electron;
- c) explain why this atomic model is unsatisfactory.

Positive charge Ne uniformly distributed over a sphere of radius $R_0: \rho = \frac{Ne}{\frac{4}{3}\pi R_0^3} = \frac{3Ne}{4\pi R_0^3}$. Inside the sphere (applying Gauss's law): $\bar{E} = \bar{a}_R E_R = \bar{a}_R \frac{Ner}{4\pi\epsilon_0 R_0^3}$.

- a) Force experienced by an electron $-e: \bar{F} = -\overline{a_R} \frac{Ne^2 r}{4\pi\epsilon_0 R_0^3}$.
- b) Equation of motion for an electron with mass m:

$$m\frac{d^2r}{dt^2}=-\frac{Ne^2r}{4\pi\epsilon_0R_0^3},$$
 or
$$\frac{d^2r}{dt^2}+\left(\frac{Ne^2}{4\pi\epsilon_0mR_0^3}\right)r=0$$
 or
$$\frac{d^2r}{dt^2}+\omega_e^2r=0$$
 where
$$\omega_e=\sqrt{\frac{Ne^2}{4\pi\epsilon_0mR_0^3}}$$

Hence the electrons would undergo a simple harmonic motion with an angular frequency ω_e .

c) The oscillating electrons would lase power through radiation and lead to an unstable atomic model.

Exercise 3.21

A simple classical model of an atom consists of a nucleus of a positive charge Ne surrounded by a spherical electron cloud of the same total negative charge. (N is the atomic number and e is the magnitude of electronic charge.) An external electric field \mathbf{E}_o will cause the nucleus to be displaced a distance r_o from the center of the electron cloud, thus polarizing the atom. Assuming a uniform charge distribution within the electron cloud of radius b, find r_o

Answer:

Applied E_0 causes a displacement r_0 .

Force of separation: qE_0

Restoring force(attraction): qE_x

 E_x at q due to spherical volume of electrons- of radius r_0 is (by Gauss's law)

$$E_x = \frac{\rho r_0}{3\epsilon_0} = -\frac{r_0}{3\epsilon_0} |\rho|$$

where

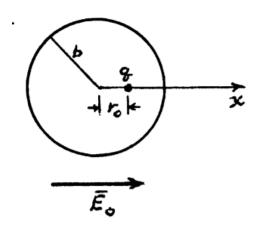
$$|\rho| = \frac{q}{\frac{4}{3}\pi b^3} = \frac{3\text{Ne}}{4\pi b^3}$$

At equilibrium:

$$E_0 = |E_x| = \frac{r_0 N e}{4\pi\epsilon_0 b^3}$$

or

$$r_0 = \frac{4\pi\epsilon_0 b^3}{Ne} E_0$$



The polarization in a dielectric cube of side L centered at the origin is given by P = $P_o\left(\mathbf{a}_x x + \mathbf{a}_y y + \mathbf{a}_z z\right)$

- a) Determine the surface and volume bound-charge densities.
- b) Show that the total bound charge is zero.

a) $\rho_{ps} = \bar{p} \cdot \bar{a}_n = P_0 \frac{L}{2}$ on all six faces of the-cube. $\rho_p = -\bar{\nabla} \cdot \bar{p} = -3P_0$.

$$\rho_p = -\bar{\nabla} \cdot \bar{p} = -3P_0.$$

b) $Q_{s'}=(6L^2) \rho_{ps}=3P_0L^3, \quad Q_v=(L^3) \rho_{p,}=-3P_0L^3.$ Total bound charge $=Q_s+Q_v=0.$

Exercise 3.23

Determine the electric field intensity at the center of a small spherical cavity cut out of a large block of dielectric in which a polarization **P** exists.

Answer:

Assume $\bar{\rho} = \bar{a}_z P$. Surface charge density

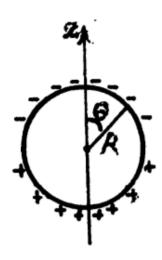
$$\rho_{ps} = \bar{p} \cdot \bar{a}_n = (\bar{a}_z p) \cdot (-\bar{a}_R)$$
$$= -P \cos \theta.$$

The z-component of the electric field intensity due to a ring of ρ_{ps} contained in width $Rd\theta$ at θ is

$$dE_z = \frac{P\cos\theta}{4\pi\varepsilon_0 R^2} (2\pi R\sin\theta) (Rd\theta) \cos\theta$$
$$= \frac{P}{2\varepsilon_0} \cos^2\theta \sin\theta d\theta.$$

At the center of the cavity: $\bar{E} = \bar{a}_z E_z = \bar{a}_z \frac{\rho}{2\epsilon_0} \int_0^{\pi} \cos^2 \theta \sin \theta d\theta$

$$=\bar{P}/3\epsilon_0.$$



Solve the following problems:

- a) Find the breakdown voltage of a parallel-plate capacitor, assuming that conducting plates are 50(mm) apart and the medium between them is air.
- b) Find the breakdown voltage if the entire space between the conducting plates is filled with plexiglass, which has a dielectric constant 3 and a dielectric strength 20(kV/mm)
- c) If a 10 (mm) thick plexiglass is inserted between the plates, what is the maximum voltage that can be applied to the plates without a breakdown?

Answer:

a: air

p: plexigiass

b: breakdown

a)
$$V_b = E_{ba} d_a = 3 \times 50 = 150 (kV)$$

b)
$$V_b = E_{bp}d_p = 20 \times 50 = 1,000(kV)$$

c)
$$V_b = E_a d_a + E_n d_n = E_a (50 - d_n) + E_n d_n$$

c)
$$V_b = E_{bp}d_p = 20 \times 30 = 1,000 (kV)$$

c) $V_b = E_a d_a + E_p d_p = E_a (50 - d_p) + E_p d_p$
Now, $D_a = D_p \to \epsilon_0 E_a = \epsilon_0 \epsilon_{rp} E_p \longrightarrow E_p = \frac{E_a}{\epsilon_{rp}} = \frac{E_u}{3} < E_a$

 $E_{ba} < E_{bp} \longrightarrow \text{Breakdown occurs in air-region first}$

$$\therefore V_b = E_{ba}(50 - 10) + \frac{E_{ba}}{3} \times 10 = 3\left(40 + \frac{1}{3} \times 10\right) = 130(\text{kV}).$$

Exercise 3.25

Assume that the z=0 plane separates two lossless dielectric regions with $\epsilon_{r1}=2$ and $\epsilon_{r2}=3$. If we know that \mathbf{E}_1 in region 1 is $\mathbf{a}_x 2y - \mathbf{a}_y 3x + \mathbf{a}_z (5+z)$, what do we also know about \mathbf{E}_2 and \mathbf{D}_2 in region 2? Can we determine \mathbf{E}_2 and \mathbf{D}_2 at any point in region 2? Explain.

At the
$$z=0$$
 plane: $\bar{E}_1 = \bar{a}_{x^2}{}^2y - \bar{a}_y \cdot 3x + \bar{a}_y 5$.

$$\bar{E}_{1t}(z=0) = \bar{E}_{2t}(z=0) = \bar{a}_x 2y - \bar{a}_y 3x.$$

$$\bar{D}_{1n}(z=0) = \bar{D}_{2n}(z=0) \longrightarrow 2\bar{E}_{1n}(z=0) = 3\bar{E}_{2n}(z=0),$$

$$\therefore E_{2n}(z=0) = \frac{2}{3} (\bar{a}_z 5) = \bar{a}_z \frac{10}{3}.$$

$$\bar{E}_2(z=0) = \bar{a}_x 2y - \bar{a}_y 3x + \bar{a}_z \frac{10}{3}.$$

$$\bar{D}_2(z=0) = \left(\bar{a}_x 2y - \bar{a}_y 3x + \bar{a}_z \frac{10}{3}\right) 3\epsilon_0.$$

Determine the boundary conditions for the tangential and the normal components of **P** at an interface between two perfect dielectric media with dielectric constants ϵ_{r1} and ϵ_{r2} .

Answer:

$$\bar{P}_t = \epsilon_0 \left(\epsilon_r - 1 \right) \bar{E}_t. \quad \bar{E}_{t1} = \bar{E}_{t2} \quad \longrightarrow \frac{1}{\epsilon_{r1} - 1} p_{t1} = \frac{1}{\epsilon_{r2} - 1} p_{t2}.$$

$$\epsilon_{r1} \bar{E}_{n1} = \epsilon_{r2} \bar{E}_{n_2} \longrightarrow \frac{\epsilon_{r1}}{\epsilon_{r1} - 1} p_{n1} = \frac{\epsilon_{r2}}{\epsilon_{r_2} - 1} p_{n_2}.$$

Exercise 3.27

What are the boundary conditions that must be satisfied by the electric potential at an interface between two perfect dielectrics with dielectric constants ϵ_{r1} and ϵ_{r2} ?

Answer:

$$\epsilon_1 \frac{\partial V_1}{\partial n} = \epsilon_2 \frac{\partial V_2}{\partial n}$$
 and $V_1 = V_2$

Exercise 3.28

Dielectric lenses can be used to collimate electromagnetic fields. In Fig. 3-41 the left surface of the lens is that of a circular cylinder, and the right surface is a plane. If \mathbf{E}_1 at point $P(r_o, 45^\circ, z)$ in region 1 is $\mathbf{a}_r 5 - \mathbf{a}_\phi 3$, what must be the dielectric constant of the lens in order that \mathbf{E}_3 in region 3 is parallel to the x-axis?

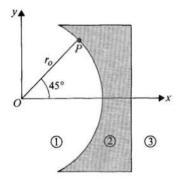


FIGURE 3-41 A dielectric lens (Problem P.3-28).

Answer:

Assume $E_2 = \bar{a}_r E_{2r} + \bar{a}_\phi E_{2\phi}$.

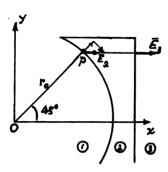
Boundary condition: $\bar{a}_n \times \bar{E}_1 = \bar{a}_n \times \bar{E}_2 \longrightarrow E_{2\phi} = -3$.

For E_3 , and hence E_2 , to be parallel to the x-axis, $E_{2\phi} = -E_{2r}$

$$\rightarrow E_{2r} = 3.$$

Boundary condition: $\bar{a}_n \cdot \bar{D}_1 = \bar{a}_n \cdot \bar{D}_2$

$$\longrightarrow \epsilon_1 E_{r1} = \epsilon_2 E_2 \longrightarrow \epsilon_0 5 = \epsilon_0 \epsilon_{r2} 3$$
$$\epsilon_{r2} = \frac{5}{3} = 1.667$$



Exercise 3.29

Refer to Example 3-16. Assuming the same r_i and r_o and requiring the maximum electric field intensities in the insulating materials not to exceed 25% of their dielectric strengths, determine the voltage rating of the coaxial cable

- a) if $r_p = 1.75r_i$
- b) if $r_p = 1.35r_i$
- c) Plot the variations of E_r and V versus r for both part (a) and part (b).

Answer:

Given: $r_i = 0.4$ (cm) = 0.004 (m), $r_0 = 0.832$ (cm) = 2.08 r_i .

25% of the dielectric strength of rubber- ($\epsilon_{rr}=3.2$) is $0.25\times25\times10^6=6.25\times10^6({\rm v/m})$. 25% of the dielectric strength of polystyrene ($\epsilon_{r-p}=2.6$) is $0.25\times20\times10^6=5\times10^6({\rm v/m})$. a) $r_p=1.75r_i, r_0=1.189r_p$.

$$\begin{aligned} &\text{Max. } E_r = \frac{P_l}{2\pi\epsilon_0} \left(\frac{1}{3.2r_i}\right) \longrightarrow \frac{\rho_l}{2\pi\epsilon_b} = 6.25 \times 10^6 \times 3.2r_i = 20 \times 10^6 r_i, \\ &\text{Max. } E_p = \frac{P_l}{2\pi\epsilon_0} \left(\frac{1}{2.6r_p}\right) = \frac{\rho_l}{2\pi\epsilon_0} \left(\frac{1}{4.55r_i}\right) \longrightarrow \frac{P_l}{2\pi t_0} = 5 \times 10^6 \times 4.55r_i = 22.75 \times 10^6 r_i. \\ &\frac{\text{Max. } E_r}{\text{Max. } E_p} = \frac{4.55}{3.2} > \frac{6.25}{5} \longrightarrow \text{Max. } E_r \text{ determines the allowable } P_l/2\pi\epsilon_0. \end{aligned}$$

$$\therefore \frac{\rho_l}{2\pi\epsilon_0} = 20 \times 10^6 \times 0.004 = 8 \times 10^4.$$

Voltage rating $V_{\text{max.}} = 8 \times 10^4 \left(\frac{1}{2.6} \ln \frac{r_0}{r_p} + \frac{1}{3.2} \ln \frac{r_p}{r_i} \right)$

$$= 8 \times 10^4 \left(\frac{1}{2.6} \ln 1.189 + \frac{1}{3.2} \ln 1.75 \right) = 19.3(kV).$$

b) $r_p = 1.35 r_i, r_0 = 1.54 r_p$ Max, $r_p = \frac{\rho_l}{2\pi\epsilon_0} \left(\frac{1}{2.6r_p}\right) = \frac{\rho}{2\pi\epsilon_0} \left(\frac{1}{3.51r_i}\right)$, which determines the allowable $\frac{\rho_l}{2\pi\epsilon_0} = 5 \times 10^6 \times 3.51 r_i = 1.76 \times 10^6 r_i = 7 \times 10^3$.

$$V_{\text{max.}} = 7 \times 10^3 \left(\frac{1}{2.6} \ln 1.54 + \frac{1}{3.2} \ln 1.35 \right) = 1.82 (\text{ kV}).$$

Exercise 3.30

The space between a parallel-plate capacitor of area S is filled with a dielectric whose permittivity varies linearly from ϵ_1 at one plate (y = 0) to ϵ_2 at the other plate (y = d). Neglecting fringing effect, find the capacitance.

Answer:

$$\epsilon = \frac{\epsilon_2 - \epsilon_1}{d}y + \epsilon_1$$
Assume Q on plate at $y = d$. $\bar{E} = -\bar{a}_y \frac{\rho_s}{\epsilon} = \frac{\rho_s}{s\left(\frac{\epsilon_2 - \epsilon_1}{d}y + t_1\right)}$

$$V = -\int_{y=0}^{y=d} \bar{E} \cdot d\bar{l} = \frac{Qd \ln\left(\epsilon_2/\epsilon_1\right)}{S\left(\epsilon_2 - \epsilon_1\right)}$$

$$C = \frac{Q}{V} = \frac{S\left(\epsilon_2 - \epsilon_1\right)}{d \ln\left(\epsilon_2/\epsilon_1\right)}$$

Exercise 3.31

Assume that the outer conductor of the cylindrical capacitor in Example 3 – 18 is grounded and that the inner conductor is maintained at a potential V_0 .

- a) Find the electric field intensity, $\mathbf{E}(a)$, at the surface of the inner conductor.
- b) With the inner radius, b, of the outer conductor fixed, find a so that E(a) is minimized.
 - c) Find this minimum E(a).
 - d) Determine the capacitance under the conditions of part (b).

EXAMPLE 3-18 A cylindrical capacitor consists of an inner conductor of radius a and an outer conductor whose inner radius is b. The space between the conductors is filled with a dielectric of permittivity ϵ , and the length of the capacitor is L. Determine the capacitance of this capacitor.

Solution We use cylindrical coordinates for this problem. First we assume charges +Q and -Q on the surface of the inner conductor and the inner surface of the outer conductor, respectively. The **E** field in the dielectric can be obtained by applying Gauss's law to a cylindrical Gaussian surface within the dielectric a < r < b. (Note that Eq. (3-122) gives only the normal component of the **E** field at a conductor surface. Since the conductor surfaces are not planes here, the **E** field is not constant in the dielectric and Eq. (3-122)

cannot be used to find **E** in the a < r < b region.) Referring to Fig. 3-29 and applying Gauss's law, we have

$$\mathbf{E} = \mathbf{a}_r E_r = \mathbf{a}_r \frac{Q}{2\pi\epsilon Lr}$$

Again we neglect the fringing effect of the field near the edges of the conductors. The potential difference between the inner and outer conductors is

$$V_{ab} = -\int_{r=b}^{r=a} \mathbf{E} \cdot d\ell = -\int_{b}^{a} \left(\mathbf{a}_{r} \frac{Q}{2\pi \epsilon L r} \right) \cdot (\mathbf{a}_{r} dr)$$
$$= \frac{Q}{2\pi \epsilon L} \ln \left(\frac{b}{a} \right)$$

Therefore, for a cylindrical capacitor,

$$C = \frac{Q}{V_{ab}} = \frac{2\pi\epsilon L}{\ln\left(\frac{b}{a}\right)}$$

We could not solve this problem from an assumed V_{ab} because the electric field is not uniform between the inner and outer conductors. Thus we would not know how to express **E** and Q in terms of V_{ab} until we learned how to solve such a boundary value problem.

Answer:

Let ρ_l be the-lineal charge density on the inner conductor. $E = \bar{a}_r \frac{\rho_R}{2\pi\epsilon_r}$

$$V_0 = -\int_b^a E \cdot d\bar{r} = \frac{\rho_l}{2\pi\epsilon} \ln\left(\frac{b}{a}\right) \longrightarrow \rho_l = \frac{2\pi\epsilon V_0}{\ln(b/a)}.$$

- a) $E(a) = \bar{a}_r \frac{V_0}{a \ln(b/a)}$.
- b) For a fixed b, the function to be minimized is. $(x = b/a) f(x) = \frac{V_0 x}{b \cdot \ln x}$. Setting $\begin{array}{l} \frac{df(x)}{dx} = 0 \text{ yields } \ln x = 1, \text{ or } x = \frac{b}{a} = e = 2.718. \\ \text{c) } \min E(a) = eV_0/b. \\ \text{d) } (C' = \frac{\rho_l}{V_0} = \frac{2\pi\epsilon}{\ln(b/a)} = 2\pi\epsilon(F/m) \end{array}$

Exercise 3.32

The radius of the core and the inner radius of the outer conductor of a very long coaxial transmission line are r_i and r_o , respectively. The space between the conductors is filled with two coaxial layers of dielectrics. The dielectric constants of the dielectrics are ϵ_{r1} for $r_i < r < b$ and ϵ_{r2} for $b < r < r_o$. Determine its capacitance per unit length.

$$\bar{D} = \bar{a}_r \frac{\rho_l}{2\pi r} \cdot \bar{E}_i = \bar{a}_r \frac{\rho_l}{2\pi \epsilon_0 \epsilon_{r1} r}, \quad r_i < r < b$$

$$\bar{E}_2 = \bar{a}_r \frac{\rho_l}{2\pi \epsilon_0 \epsilon_{r2} r}, \quad b < r < r_c$$

$$V = -\int_{r_0}^{r_i} \bar{E} \cdot d\bar{r} = \frac{\rho_l}{2 \cdot \pi \epsilon_0} \left[\frac{1}{\epsilon_{r1}} \ln \left(\frac{b}{r_i} \right) + \frac{1}{\epsilon_{r2}} \ln \left(\frac{r_0}{b} \right) \right]$$

$$C' = \frac{\rho_l}{V} = \frac{2\pi \epsilon_0}{\frac{1}{\epsilon_{r1}} \ln \left(\frac{b}{r_i} \right) + \frac{1}{\epsilon_{r2}} \ln \left(\frac{r_c}{b} \right)} \quad (F/m).$$

A cylindrical capacitor of length L consists of coaxial conducting surfaces of radii r_i and r_o . Two dielectric media of different dielectric constants ϵ_{r1} and ϵ_{r2} fill the space between the conducting surfaces as shown in Fig. 3-42. Determine its capacitance.

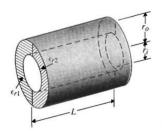


FIGURE 3-42 A cylindrical capacitor with two dielectric media (Problem P.3-33).

Answer:

Gauss's law:
$$\oint_{s} \bar{D} \cdot d\bar{s} = \rho_{l}L$$

$$\bar{E}_{1} = \bar{E}_{2} = \bar{a}_{r}E_{r}, \quad \pi r L \left(\epsilon_{0}\epsilon_{r1} + \epsilon_{0}\epsilon_{r2}\right) E_{r} = \rho_{l}L$$

$$\longrightarrow E_{r} = \frac{\rho_{l}}{\pi r \epsilon_{0} \left(\epsilon_{r1} + \epsilon_{r2}\right)}; V = -\int_{r_{o}}^{r_{i}} E_{r} dr = \frac{\rho_{l}}{\epsilon_{0} \left(\epsilon_{r1} + \epsilon_{r2}\right)} \ln \left(\frac{r_{o}}{r_{i}}\right)$$

$$\therefore C = \frac{\rho_{l}L}{V} = \frac{\pi \epsilon_{0} \left(\epsilon_{r1} + \epsilon_{r2}\right) L}{\ln \left(r_{o}/r_{i}\right)}.$$

Exercise 3.34

A capacitor consists of two coaxial metallic cylindrical surfaces of a length 30(mm) and radii 5(mm) and 7(mm). The dielectric material between the surfaces has a relative permittivity $\epsilon_r = 2 + (4/r)$, where r is measured in mm. Determine the capacitance of the capacitor.

Answer:

$$\bar{E} = \bar{a}_r \frac{\rho_l}{2\pi\epsilon r} = \bar{a}_r \frac{\rho_l}{2\pi\epsilon_0 \left(2 + \frac{4}{r}\right) r} = \bar{a}_r \frac{\rho_l}{4\pi\epsilon_0 (r+2)}$$

$$V = -\int_{r_0}^{r_i} \bar{E} \cdot d\bar{r} = \frac{\rho_l}{4\pi\epsilon_0} \ln(r+2) \Big|_{5}^{7} = \frac{\rho_l}{4\pi\epsilon_0} \ln\left(\frac{9}{7}\right)$$

$$C = \frac{\rho_l L}{V} = \frac{4\pi\epsilon_0 L}{\ln(9/7)} = 1500\epsilon_0 = 13.26(\mu F).$$

Exercise 3.35

Assuming the earth to be a large conducting sphere (radius = 6.37×10^3 km) surrounded by air, find

- a) the capacitance of the earth;
- b) the maximum charge that can exist on the earth before the air breaks down.

a)
$$C = 4\pi\epsilon_0 R = \frac{1}{9} \times 10^{-9} \times (6.37 \times 10^6) = 7.08 \times 10^{-4} (\text{ F})$$

b)
$$E_b = 3 \times 10^6 = \frac{Q_{\text{max}}}{4\pi\epsilon_0 R^2} \longrightarrow Q_{\text{max}} = 1.35 \times 10^{10} (C)$$

Determine the capacitance of an isolated conducting sphere of radius b that is coated with a dielectric layer of uniform thickness d. The dielectric has an electric susceptibility χ_e .

Answer:

Assume charge Q on conducting sphere,

$$b < R < b + d : \quad \bar{E}_1 = \bar{a}_R \frac{Q}{4\pi\epsilon_0 (1 + \chi_e) R^2}$$

$$R > b + d \quad \bar{E}_2 = \bar{a}_R \frac{Q}{4\pi\epsilon_0 R^2}.$$

$$V = -\int_{\infty}^b \bar{E} \cdot d\bar{l} = -\int_{\infty}^{b+d} E_2 dR - \int_{b+d}^b E_1 dR = \frac{Q}{4\pi\epsilon_0 (1 + \chi_e)} \left(\frac{\chi_e}{b+d} + \frac{1}{b}\right)$$

$$C = \frac{Q}{V} = \frac{4\pi\epsilon_0 (1 + \chi_e)}{\frac{\chi_e}{b+d} + \frac{1}{b}}.$$

Exercise 3.37

A capacitor consists of two concentric spherical shells of radii R_i and R_o . The space between them is filled with a dielectric of relative permittivity ϵ_r from R_i to b ($R_i < b < R_o$) and another dielectric of relative permittivity $2\epsilon_r$ from b to R_o .

- a) Determine **E** and **D** everywhere in terms of an applied voltage V.
- b) Determine the capacitance.

Answer:

Assume charge Q in inner shall and -Q on inner-shell. $R_i < R < R_0 : \bar{D} = \bar{a}_R \frac{Q}{4\pi R^2}$.

$$R_{i} < R < b : E_{1} = \frac{\bar{D}}{\epsilon_{0} \cdot \epsilon_{r}}; \quad b < R < R_{0} : \bar{E}_{2} = \frac{\bar{D}}{2\epsilon_{0}\epsilon_{r}}$$

$$V = -\int_{R_{0}}^{R_{i}} \bar{E} \cdot d\bar{R} = -\int_{b}^{R_{i}} E_{1} dR - \int_{R_{0}}^{b} E_{2} dR = \frac{Q}{4\pi\epsilon_{0}\epsilon_{r}} \left(\frac{1}{R_{i}} - \frac{1}{2b} - \frac{1}{2R_{0}}\right)$$
a) $\bar{D} = \bar{a}_{R} \frac{\epsilon_{0}\epsilon_{r}V}{R^{2}\left(\frac{1}{R_{i}} - \frac{1}{2b} - \frac{1}{2R_{0}}\right)}, R_{i} < R_{0}. \quad \bar{D} = 0, \bar{E} = 0 \text{ for. } R < R_{i} \text{ and } R > R_{0}$

$$\bar{E}_{1} = \bar{a}_{R} \frac{V}{R^{2}\left(\frac{1}{R_{i}} - \frac{1}{2b} - \frac{1}{2R_{0}}\right)}; \quad \bar{E}_{2} = \bar{a}_{R} \frac{V}{R^{2}\left(\frac{2}{R_{i}} - \frac{1}{b} - \frac{1}{R_{0}}\right)}.$$
b) $C = \frac{4\pi\epsilon_{0}\epsilon_{r}}{\frac{1}{R_{i}} - \frac{1}{2b} - \frac{1}{2R_{0}}}$

Exercise 3.38

The two parallel conducting wires of a power transmission line have a radius a and are spaced at a distance d apart. The wires are at a height h above the ground. Assuming the ground to be perfectly conducting and both d and h to be much larger than a, find the expressions for the mutual and self-partial capacitances per unit length.

Image wires 1' and 2'.

The potential differences referring to the ground are:

$$V_{10} = \frac{1}{2} V_{11'} = \frac{\rho_{l1}}{2\pi\epsilon_0} \ln \frac{2h}{a} + \frac{\rho_{l2}}{2\pi\epsilon_0} \ln \frac{D}{d}$$
$$V_{20} = \frac{1}{2} V_{22'} = \frac{\rho_{l1}}{2\pi\epsilon_0} \ln \frac{D}{d} + \frac{\rho_{l2}}{2\pi\epsilon_0} \ln \frac{2h}{d}$$

where $D = (4h^2 + d^2)^{1/2}$

Eqs. (1)and (2) can be solved for ρ_{l1} and ρ_{l2} in terms of V_{10} and V_{20}

$$\rho_{l1} = \Delta_0 \left(V_{10} \ln \frac{2h}{a} - V_{20} \ln \frac{D}{a} \right) = c_{11} V_{10} + c_{12} V_{20}$$

$$\rho_{l2} = \Delta_0 \left(-V_{10} \ln \frac{D}{a} + V_{10} \ln \frac{2h}{a} \right) = c_{21} V_{10} + c_{22} V_{20}$$

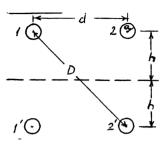
where

$$\Delta_0 = \frac{2\pi\epsilon_0}{\left(\ln\frac{2h}{a}\right)^2 - \left(\ln\frac{D}{a}\right)^2}$$

From Eqs. (3-157a, b, c):

$$C_{12} = -c_{12} = \Delta_0 \ln \frac{D}{a} = \frac{2\pi\epsilon_0 \ln(D/a)}{\left(\ln \frac{2h}{a}\right)^2 - \left(\ln \frac{D}{a}\right)^2}.$$

$$C_{10} = C_{20} = c_{11} + c_{12} = \frac{2\pi\epsilon_0}{\ln(2h/a) + \ln(D/a)}.$$



Exercise 3.39

An isolated system consists of three very long parallel conducting wires. The axes of all three wires lie in a plane. The two outside wires are of a radius b and both are at a distance d = 500b from a center wire of a radius 2b. Determine the partial capacitances per unit length.

Assuming line charge densities we have (see Example 3-21)

$$2\pi\epsilon_0 V_{10} = \rho_{l0} \ln \frac{b}{d} + \rho_{l1} \ln \frac{d}{2b} + \rho_{l2} \ln \frac{2d}{d}$$
$$= \rho_{l0} \ln \frac{1}{500} + \rho_{l1} \ln 250 + \rho_{l2} \ln 2$$
$$2\pi\epsilon_0 V_{20} = \rho_{l0} \ln \frac{b}{2d} + \rho_{l1} \ln \frac{d}{d} + \rho_{l2} \ln \frac{2d}{b}$$
$$= \rho_{l0} \ln \frac{1}{1000} + 0 + \rho_{l2} \ln 1000$$

For an isolated system:

$$\rho_{l0} = -(\rho_{l1} + \rho_{l2})$$

Eqs. (1) and (2)become
$$2\pi\epsilon_0 V_{10} = \rho_{l1} \left(\ln 250 + \ln 500\right) + \rho_{l2} (\ln 2 + \ln 500)$$

$$= \rho_{l1} \ln 125 \times 10^3 + \rho_{l2} \ln 10^3$$

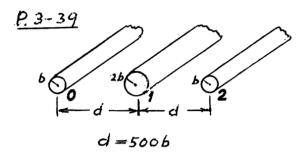
$$2\pi\epsilon_0 V_{20} = \rho_{l_1} \ln 1000 + \rho_{\ell_2} (\ln 1000 + \ln 1000) = \rho_{\ell_1} \ln 10^3 + \rho_{\ell_2} \ln 10^6$$
Solving (3) and (4):
$$\rho_{l1} = \Delta_0 \left(V_{10} \ln 10^6 - V_{20} \ln 10^3 \right),$$

$$\rho_{l2} = \Delta_0 \left(-V_{10} \ln 10^3 + V_{20} \ln 125 \times 10^3 \right),$$
where
$$\Delta_0 = \frac{2\pi\epsilon_0}{\left(\ln 10^6\right) \left(\ln 125 \times 10^3\right) - \left(\ln 10^3\right)^2}.$$

$$C_{12} = -c_{12} = \Delta_0 \ln 10^3 = 3.36 (\text{pF/m}),$$

$$C_{10} = c_{11} + c_{12} = \Delta_0 \left(\ln 10^6 - \ln 10^3\right) = 3.36 (\text{pF/m}),$$

$$C_{20} = c_{22} + c_{12} = \Delta_0 \left(\ln 125 \times 10^3 - \ln 10^3\right) = 2.35 (\text{pF/m}).$$



Exercise 3.40

Calculate the amount of electrostatic energy of a uniform sphere of charge with radius b and volume charge density ρ stored in the following regions:

- a) inside the sphere,
- b) outside the sphere.

Check your results with those in Example 3-22.

$$\bar{D} = \bar{a}_R \frac{R}{3} \rho, R < b; \quad \bar{D} = \overline{a}_R \frac{b^3 \rho}{3R^2}, R > b; \quad \bar{E} = \frac{\bar{D}}{\epsilon_0}$$
 a)
$$W_i = \frac{1}{2} \int_V \bar{D} \cdot \bar{E} dv = \frac{1}{2} \int_0^b \frac{1}{\epsilon_0} \left(\frac{R}{3} \rho \right)^2 4\pi R^2 dR = \frac{2\pi b^5 \rho^2}{45\epsilon_0}$$
 b)
$$W_o = \frac{1}{2} \int_b^\infty \frac{1}{\epsilon_0} \left(\frac{b^3 \rho}{3R^2} \right)^2 4\pi R^2 dR = \frac{2\pi b^5 p^2}{9\epsilon_0}$$
 Total
$$W = W_i + W_o = \frac{4\pi b^5 p^2}{15\epsilon_0}$$

Einstein's theory of relativity stipulates that the work required to assemble a charge is stored as energy in the mass and is equal to mc^2 , where m is the mass and $c \cong 3 \times 10^8 (\text{ m/s})$ is the velocity of light. Assuming the electron to be a perfect sphere, find its radius from its charge and mass $(9.1 \times 10^{-31} \text{ kg})$.

Answer:

From Eq. (3 – 169)
$$W = \frac{3Q^2}{20\pi\epsilon_0 b}$$

$$W = \frac{3e^2}{20\pi\epsilon_0 b} = mc^2$$

$$b = \frac{3e^2}{20\pi\epsilon_0 mc^2} = 1.69 \times 10^{-15} (\text{ m}).$$

Exercise 3.42

Find the electrostatic energy stored in the region of space R>b around an electric dipole of moment p.

Answer:

$$\bar{E} = \frac{p}{4\pi\epsilon_0 R^3} \left(\bar{a}_R 2 \cos \theta + \bar{a}_\theta \sin \theta \right)$$

$$W = \frac{1}{2}\epsilon_0 \int_V E^2 dV = \frac{\epsilon_0}{2} \left(\frac{p}{4\pi\epsilon_0} \right)^2 \int_0^{2\pi} d\phi \int_0^{\pi} d\theta \int_b^{\infty} \frac{1}{R^6} \left(4 \cos^2 \theta + \sin^2 \theta \right) R^2 \sin \theta dR$$

$$= \frac{p^2}{12\pi\epsilon_0 b^3}$$

Exercise 3.43

Prove that Eqs. (3 - 180)

$$W_e = \frac{1}{2}CV^2$$

for stored electrostatic energy hold true for any two-conductor capacitor.

Answer:

Two conductors at potentials V_1 and V_2 carrying charges Q and -Q:

$$W_e = \frac{1}{2}V_1 \int_{s_1} \rho_{s1} ds + \frac{1}{2}V_2 \int_{s_2} \rho_{s2} ds = \frac{1}{2}Q (V_1 - V_2)$$
$$= \frac{1}{2}CV^2, \quad V = V_1 - V_2.$$

Exercise 3.44

A parallel-plate capacitor of width w, length L, and separation d is partially filled with a dielectric medium of dielectric constant ϵ_r , as shown in Fig. 3-43. A battery of V_0 volts is connected between the plates.

- a) Find \mathbf{D}, \mathbf{E} , and ρ_s in each region.
- b) Find distance x such that the electrostatic energy stored in each region is the same.

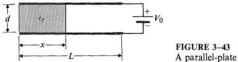


FIGURE 3-43
A parallel-plate capacitor (Problem P.3-44).

Answer:

Region 1-dielectric; region 2 -air.

a)
$$\bar{E}_1 = -\bar{a}_y \frac{V_0}{d}, \bar{D}_1 = -\bar{a}_y \epsilon_0 \epsilon_r \frac{V_0}{d}, \rho_{s1} = \epsilon_0 \epsilon_r \frac{V_0}{d} \text{(top plate)}$$

$$\bar{E}_2 = -\bar{a}_y \frac{V_0}{d}, \bar{D}_2 = -\bar{a}_y \epsilon_0 \frac{V_0}{d}, \rho_{s2} = \epsilon_0 \frac{V_0}{d} \text{(top plate)}$$
(b)
$$\frac{W_{e1}}{W_{e2}} = \frac{\epsilon_r x}{L - x} = 1 \longrightarrow x = \frac{L}{\epsilon_r + 1}$$

Exercise 3.45

Using the principle of virtual displacement, derive an expression for the force between two point charges +Q and -Q separated by a distance x in free space.

From Eqs.
$$(3 - 165)$$

$$W_e = \frac{1}{2} \sum_{k=1}^{N} Q_k V_k$$

and
$$(3 - 166)$$

$$V_k = \frac{1}{4\pi\epsilon_0} \sum_{\substack{j=1\\(j\neq k)}}^N \frac{Q_j}{R_{jk}}.$$

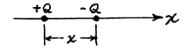
$$W_e = -\frac{Q^2}{4\pi\epsilon_0 x}, \quad x = x_- - x_+.$$

Force on

$$+Q: \bar{F}_{Q+} = -\bar{\nabla}_{+}W_{e} = -\bar{a}_{x}\frac{\partial W_{e}}{\partial x_{+}} = \bar{a}_{x}\frac{Q^{2}}{4\pi\epsilon_{0}x^{2}}$$

Force on

$$-Q: \bar{F}_{Q-} = -\bar{\nabla}_{-}W_{e} = -\bar{a}_{x}\frac{\partial W_{e}}{\partial x_{-}} = -\bar{a}_{x}\frac{Q^{2}}{4\pi\epsilon_{0}x^{2}} = -\bar{F}_{a+}$$



Exercise 3.46

A constant voltage V_0 is applied to a partially filled parallel-plate capacitor shown in Fig. 3-44. The permittivity of the dielectric is ϵ , and the area of the plates is S. Find the force on the upper plate.

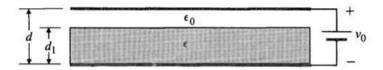


FIGURE 3-44 A parallel-plate capacitor (Problem P.3-46).

Answer:

Denoting the plate separation by y and using Eq (3-192)

$$dW + dW_e = dW_s$$

$$\bar{F}_v = \bar{a}_y \frac{\partial W_e}{\partial y} = \bar{a}_y \frac{\partial}{\partial y} \left(\frac{1}{2} C V_0^2 \right) = \bar{a}_y \frac{V_0^2}{2} \frac{dC}{dy},$$

where C is the capacitance of the series connection of C_d and $C_a: C_d = \frac{\epsilon S}{d_1}; \quad C_a = \frac{\epsilon_0 S'}{y-d_1}$.

$$C = \frac{c_d c_a}{c_d + c_a} = \frac{\epsilon \epsilon_0 S}{\epsilon (y - d_1) + \epsilon_0 d_1}$$

$$\therefore \bar{F}_V = -\bar{a}_y \frac{\epsilon^2 \epsilon_0 S V_0}{2 \left[\epsilon (d - d_1) + \epsilon_0 d_1\right]^2} \text{ (attractive force)}.$$

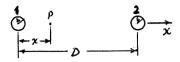
The conductors of an isolated two-wire transmission line, each of radius b, are spaced at a distance D apart. Assuming $D \gg b$ and a voltage V_0 between the lines, find the force per unit length on the lines.

Answer:

Assume line charge densities ρ_l and $-\rho_l$ on conductors 1 and 2 respectively. At point P,

$$\begin{split} \bar{E}_p &= \bar{E}_1 + \bar{E}_2 = \bar{a}_x \left[\frac{\rho_l}{2\pi\epsilon_0 x} + \frac{\rho_l}{2\pi\epsilon_0 (D-x)} \right] \\ V_0 &= V_1 - V_2 = \int_b^{D-b} \bar{E}_p \cdot d\bar{x} = \frac{\rho_l}{2\pi\epsilon_0} \int_b^{D-b} \left(\frac{1}{x} + \frac{1}{D-x} \right) dx \\ &= \frac{\rho_l}{2\pi\epsilon_0} \left(\ln \frac{D-b}{b} - \ln \frac{b}{D-b} \right) = \frac{\rho_l}{\pi\epsilon_0} \ln \frac{D-b}{b} \cong \frac{\rho_l}{\pi\epsilon_0} \ln \frac{D}{b} \\ C' &= \frac{\rho_l}{V_0} = \frac{\pi\epsilon_0}{\ln(D/b)} \quad (F/m). \\ \bar{F}' &= \bar{\nabla} W_e = \bar{a}_x \frac{V_0}{2} \frac{\partial C'}{\partial D} = -\bar{a}_x \frac{\pi\epsilon_0 V_0^2}{2D[\ln(D/b)]^2} \end{split}$$

in the direction of decreasing D.



Exercise 3.48

A parallel-plate capacitor of width w, length L, and separation d has a solid dielectric slab of permittivity ϵ in the space between the plates. The capacitor is charged to a voltage V_0 by a battery, as indicated in Fig. 3-45. Assuming that the dielectric slab is withdrawn to the position shown, determine the force acting on the slab

- a) with the switch closed,
- b) after the switch is first opened.

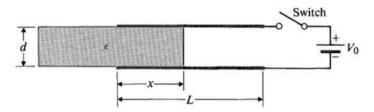


FIGURE 3-45 A partially filled parallel-plate capacitor (Problem P.3-48).

a) Switch closed $\longrightarrow V = V_0 = \text{constant.}$

$$W_e = \frac{1}{2}CV_0^2, \quad C = \frac{w}{d} \left[\epsilon x + \epsilon_0 (L - x) \right].$$
$$\bar{F}_v = \bar{\nabla} W_e = \bar{a}_x \frac{V_0^2}{2} \frac{\partial C}{\partial x} = \bar{a}_x \frac{V_0 w}{2d} \left(\epsilon - \epsilon_0 \right).$$

b) Switch open $\longrightarrow Q = \text{constant} = CV_0$.

$$\begin{split} W_e &= \frac{Q^2}{2C}, \\ \bar{F}_Q &= -\bar{\nabla} W_e = -\bar{a}_x \frac{Q^2}{2} \frac{\partial}{\partial x} \left(\frac{1}{C} \right) \\ &= \bar{a}_x \cdot \frac{Q^2 d}{2w} \cdot \frac{\epsilon - \epsilon_0}{\left[\epsilon x + \epsilon_0 (L - x) \right]^2} = \bar{a}_x \frac{V_0 w}{2d} \left(\epsilon - \epsilon_0 \right). \end{split}$$

Reference

1. Cheng, David Keun. Field and wave electromagnetics. Pearson Education India, 1989.