# VE320 Introduction to Semiconductor Physics and Devices

#### VE320 Teaching Group SU2022

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#### **Notations**

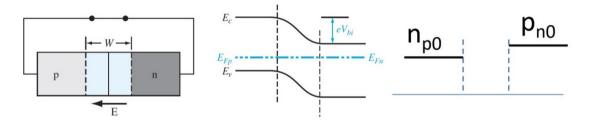
N <sub>a</sub>
$N_d$
$n_{n0} = N_d$
$p_{p0} = N_a$
$n_{p0} = n_i^2/N_a$
$p_{n0} = n_i^2/N_d$
$n_p$
$p_n$
$n_p\left(-x_p\right)$
$p_n(x_n)$
$\delta n_p = n_p - n_{p0}$
$\delta p_n = p_n - p_{n0}$

Acceptor concentration in the p region of the pn junction Donor concentration in the n region of the pn junction Thermal-equilibrium majority carrier electron concentration in the n region Thermal-equilibrium majority carrier hole concentration in the p region Thermal-equilibrium minority carrier electron concentration in the p region Thermal-equilibrium minority carrier hole concentration in the n region Total minority carrier electron concentration in the p region Total minority carrier hole concentration in the n region Minority carrier electron concentration in the p region at space charge edge Minority carrier hole concentration in the n region at space charge edge Excess minority carrier electron concentration in the p region Excess minority carrier hole concentration in the n region

#### Ideal Assumptions

- The abrupt depletion layer approximation applies. The space charge regions have abrupt boundaries, and the semiconductor is neutral outside of the depletion region.
- The Maxwell-Boltzmann approximation applies to carrier statistics.
- The concepts of low injection and complete ionization apply.
- The total current is a constant throughout the entire pn structure.
- The individual electron and hole currents are continuous functions through the pn structure.
- The individual electron and hole currents are constant throughout the depletion region.

#### Concentration Relation on Two Sides



#### Concentration Relation on Two Sides

$$V_{bi} = \frac{kT}{e} \ln \left( \frac{N_a N_d}{n_i^2} \right)$$

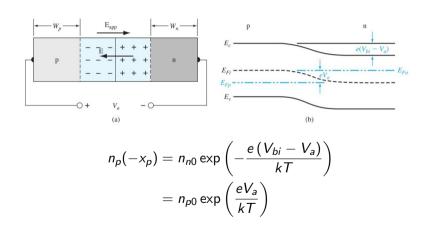
$$n_{n0} = N_d$$

$$n_{p0} = \frac{n_i^2}{N_a}$$

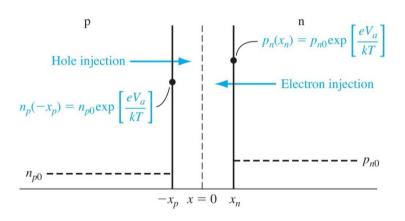
$$\Rightarrow n_p(-x_p) = n_{p0} = n_{n0} \exp \left( -\frac{eV_{bi}}{kT} \right)$$

Relates the minority carrier electron concentration on the p side of the junction to the majority carrier electron concentration on the n side of the junction in thermal equilibrium.

#### Forward Biased



#### Forward Biased



### Minority Carrier Distribution

$$D_{p} \frac{\mathrm{d}^{2} (\delta p_{n})}{\mathrm{d}x^{2}} - \mu_{p} \left( E \frac{\mathrm{d} (\delta p_{n})}{\mathrm{d}x} + p \frac{\mathrm{d}E}{\mathrm{d}x} \right) + g' - \frac{\delta p_{n}}{\tau_{pt}} = 0$$

Assumption: the electric field is zero in both the neutral p and n regions. In n region for  $x > x_n$ , we have g' = 0. The equation becomes

$$\frac{\mathrm{d}^{2}(\delta p_{n})}{\mathrm{d}x^{2}} - \frac{\delta p_{n}}{L_{p}^{2}} = 0, \quad (x > x_{n}), \quad L_{n}^{2} = D_{n}\tau_{n0}$$

Solve it with boundary conditions

$$p_n(x_n) = p_{n0} \exp\left(\frac{eV_a}{kT}\right)$$
  
 $p_n(x \to +\infty) = p_{n0}$ 

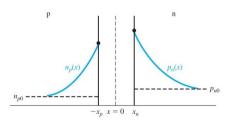
### Minority Carrier Distribution

The solution is

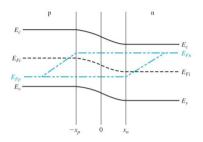
$$\delta p_n(x) = p_n(x) - p_{n0} = p_{n0} \left[ \exp\left(\frac{eV_a}{kT}\right) - 1 \right] \exp\left(\frac{x_n - x}{L_p}\right), \quad x \ge x_n$$

Similarly,

$$\delta n_p(x) = n_p(x) - n_{p0} = n_{p0} \left[ \exp\left(\frac{eV_a}{kT}\right) - 1 \right] \exp\left(\frac{x_p + x}{L_n}\right), \quad x \le -x_p$$



#### Quasi-Fermi Level



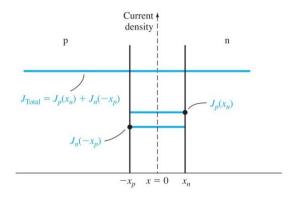
$$p = p_0 + \delta p = n_i \exp\left(\frac{E_{Fi} - E_{Fp}}{kT}\right)$$

Therefore the quasi-Fermi levels are linear functions of distance in the neutral p and n regions. Also,

$$np = n_i^2 \exp\left(\frac{E_{Fn} - E_{Fp}}{kT}\right)$$

### Ideal pn Junction Current

Assumption 4(a): The total current is a constant throughout the entire pn structure.



$$J_{p}(x_{n}) = -eD_{p}\frac{\mathrm{d}\left(\delta p_{n}(x)\right)}{\mathrm{d}x}\bigg|_{x=x}$$

### Ideal pn Junction Current

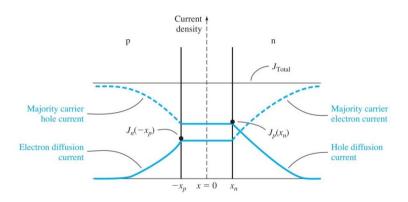
$$J_{p}(x_{n}) = \frac{eD_{p}p_{n0}}{L_{p}} \left[ \exp\left(\frac{eV_{a}}{kT}\right) - 1 \right]$$

$$J_{n}(-x_{p}) = \frac{eD_{n}n_{p0}}{L_{n}} \left[ \exp\left(\frac{eV_{a}}{kT}\right) - 1 \right]$$

Then the total current density

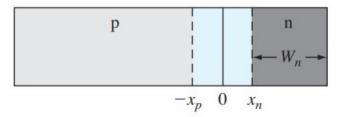
$$J=J_{s}\left[\exp\left(rac{eV_{a}}{kT}
ight)-1
ight]$$
 where  $J_{s}=\left[rac{eD_{p}p_{n0}}{L_{p}}+rac{eD_{n}n_{p0}}{L_{n}}
ight]$ 

#### Ideal pn Junction Current



The above figure shows idea electron and hole current components through a pn junction under forward bias.

We assumed in the previous analysis that both p and n regions were long compared with the minority carrier diffusion lengths. In many pn junction structures, one region may, in fact, be short compared with the minority carrier diffusion length. Assume the length  $W_n$  is much smaller than the minority carrier hole diffusion length,  $L_p$ . Assume an infinite surface recombination velocity and therefore an excess minority carrier concentration of zero. Calculate the minority carrier hole diffusion current density  $J_p$ .



The steady-state excess minority carrier hole concentration in the n region:

$$\frac{d^2(\delta p_n)}{dx^2} - \frac{\delta p_n}{L_p^2} = 0$$

Boundary conditions:

$$p_n(x_n) = p_{n0} \exp\left(\frac{eV_a}{kT}\right)$$
$$p_n(x = x_n + W_n) = p_{n0}$$

The general solution is then

$$\delta p_n(x) = p_n(x) - p_{n0} = Ae^{x/L_p} + Be^{-x/L_p} \quad (x \ge x_n)$$

We get the solution

$$\delta p_n(x) = p_{n0} \left[ \exp\left(\frac{eV_a}{kT}\right) - 1 \right] \frac{\sinh\left[\left(x_n + W_n - x\right)/L_p\right]}{\sinh\left[W_n/L_p\right]}$$

If  $W_n \ll L_p$ , we can approximate the hyperbolic sine terms by

$$\sinh\left(\frac{x_n+W_n-x}{L_p}\right)\approx\left(\frac{x_n+W_n-x}{L_p}\right)$$

and

$$\sinh\left(\frac{W_n}{L_p}\right) \approx \left(\frac{W_n}{L_p}\right)$$

Then

$$\delta p_n(x) = p_{n0} \left[ \exp\left(\frac{eV_a}{kT}\right) - 1 \right] \left(\frac{x_n + W_n - x}{W_n}\right)$$

The minority carrier hole diffusion current density is given by

$$J_{p} = -eD_{p} \frac{d \left[\delta p_{n}(x)\right]}{dx}$$

so that in the short n region, we have

$$J_p(x) = \frac{eD_p p_{n0}}{W_n} \left[ \exp\left(\frac{eV_a}{kT}\right) - 1 \right]$$

## Questions?