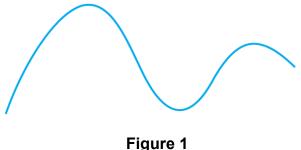
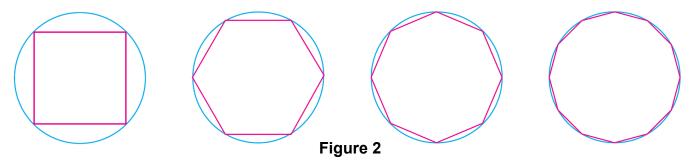
What do we mean by the length of a curve? We might think of fitting a piece of string to the curve in Figure 1 and then measuring the string against a ruler. But that might be difficult to do with much accuracy if we have a complicated curve.

We need a precise definition for the length of an arc of a curve, in the same spirit as the definitions we developed for the concepts of area and volume.



If the curve is a polygon, we can easily find its length; we just add the lengths of the line segments that form the polygon. (We can use the distance formula to find the distance between the endpoints of each segment.) We are going to define the length of a general curve by first approximating it by a polygon and then taking a limit as the number of segments of the polygon is increased. This process is familiar for the case of a circle, where the circumference is the limit of lengths of inscribed polygons (see Figure 2).



Suppose that a curve C is defined by the equation y = f(x) where f is continuous and  $a \le x \le b$ .

We obtain a polygonal approximation to C by dividing the interval [a, b] into n subintervals with endpoints  $x_0, x_1, \ldots, x_n$  and equal width  $\Delta x$ . If  $y_i = f(x_i)$ , then the point  $P_i(x_i, y_i)$  lies on C and the polygon with vertices  $P_0, P_1, \ldots, P_n$ , illustrated in Figure 3, is an approximation to C.

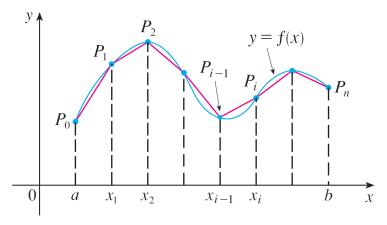


Figure 3

The length L of C is approximately the length of this polygon and the approximation gets better as we let n increase. (See Figure 4, where the arc of the curve between  $P_{i-1}$  and  $P_i$  has been magnified and approximations with successively smaller values of  $\Delta x$  are shown.)

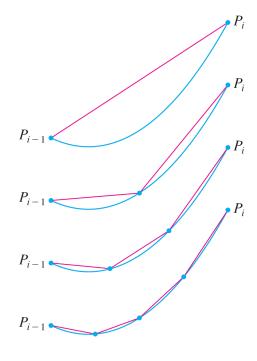


Figure 4

Therefore we define the **length** L of the curve C with equation, y = f(x),  $a \le x \le b$  as the limit of the lengths of these inscribed polygons (if the limit exists):

1

$$L = \lim_{n \to \infty} \sum_{i=1}^{n} |P_{i-1}P_i|$$

Notice that the procedure for defining arc length is very similar to the procedure we used for defining area and volume: We divided the curve into a large number of small parts. We then found the approximate lengths of the small parts and added them. Finally, we took the limit as  $n \to \infty$ 

The definition of arc length given by Equation 1 is not very convenient for computational purposes, but we can derive an integral formula for L in the case where f has a continuous derivative. [Such a function is called **smooth** because a small change in x produces a small change in f'(x).]

If we let  $\Delta y_i = y_i - y_{i-1}$ , then

$$|P_{i-1}P_i| = \sqrt{(x_i - x_{i-1})^2 + (y_i - y_{i-1})^2} = \sqrt{(\Delta x)^2 + (\Delta y_i)^2}$$

By applying the Mean Value Theorem to f on the interval,  $[x_i - x_{i-1}]$ , we find that there is a number  $x_i^*$  between  $x_{i-1}$  and  $x_i$  such that

$$f(x_i) - f(x_{i-1}) = f'(x_i^*)(x_i - x_{i-1})$$

that is

$$\Delta y_i = f'(x_i^*) \Delta x$$

#### Thus we have

$$\begin{aligned} |P_{i-1}P_i| &= \sqrt{(\Delta x)^2 + (\Delta y_i)^2} = \sqrt{(\Delta x)^2 + [f'(x_i^*) \, \Delta x]^2} \\ &= \sqrt{1 + [f'(x_i^*)]^2} \, \sqrt{(\Delta x)^2} = \sqrt{1 + [f'(x_i^*)]^2} \, \Delta x \quad \text{(since } \Delta x > 0) \end{aligned}$$

Therefore, by Definition 1,

$$L = \lim_{n \to \infty} \sum_{i=1}^{n} |P_{i-1}P_i| = \lim_{n \to \infty} \sum_{i=1}^{n} \sqrt{1 + [f'(x_i^*)]^2} \Delta x$$

We recognize this expression as being equal to

$$\int_{a}^{b} \sqrt{1 + [f'(x)]^2} dx$$

by the definition of a definite integral. This integral exists because the function  $g(x) = \sqrt{1 + [f'(x)]^2}$  is continuous. Thus we have proved the following theorem:

**The Arc Length Formula** If f' is continuous on [a, b], then the length of the curve y = f(x),  $a \le x \le b$ , is

$$L = \int_{a}^{b} \sqrt{1 + [f'(x)]^{2}} \, dx$$

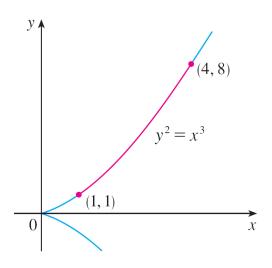
If we use Leibniz notation for derivatives, we can write the arc length formula as follows:

3

$$L = \int_{a}^{b} \sqrt{1 + \left(\frac{dy}{dx}\right)^2} \, dx$$

### Example 1

Find the length of the arc of the semicubical parabola  $y^2 = x^3$  between the points (1, 1) and (4, 8). (See Figure 5.)



### Example 1 – Solution

For the top half of the curve we have

$$y = x^{3/2}$$
  $\frac{dy}{dx} = \frac{3}{2}x^{1/2}$ 

So the arc length formula gives

$$L = \int_{1}^{4} \sqrt{1 + \left(\frac{dy}{dx}\right)^{2}} \, dx = \int_{1}^{4} \sqrt{1 + \frac{9}{4}x} \, dx$$

If we substitute  $u = 1 + \frac{9}{4}x$ , then  $du = \frac{9}{4} dx$ . When x = 1,  $u = \frac{13}{4}$ ; when x = 4, u = 10.

## Example 1 – Solution

#### **Therefore**

$$L = \frac{4}{9} \int_{13/4}^{10} \sqrt{u} \ du = \frac{4}{9} \cdot \frac{2}{3} u^{3/2} \Big]_{13/4}^{10}$$
$$= \frac{8}{27} \Big[ 10^{3/2} - \Big( \frac{13}{4} \Big)^{3/2} \Big]$$
$$= \frac{1}{27} \Big( 80 \sqrt{10} - 13 \sqrt{13} \Big)$$

If a curve has the equation x = g(y),  $c \le y \le d$ , and g'(y) is continuous, then by interchanging the roles of x and y in Formula 2 or Equation 3, we obtain the following formula for its length:

$$L = \int_{c}^{d} \sqrt{1 + [g'(y)]^{2}} \, dy = \int_{c}^{d} \sqrt{1 + \left(\frac{dx}{dy}\right)^{2}} \, dy$$

We will find it useful to have a function that measures the arc length of a curve from a particular starting point to any other point on the curve. Thus if a smooth curve C has the equation y = f(x),  $a \le x \le b$  let s(x) be the distance along C from the initial point  $P_0(a, f(a))$  to the point Q(x, f(x)). Then s is a function, called the **arc length function**, and, by Formula 2,

$$s(x) = \int_{a}^{x} \sqrt{1 + [f'(t)]^2} dt$$

(We have replaced the variable of integration by so that does not have two meanings.) We can use Part 1 of the Fundamental Theorem of Calculus to differentiate Equation 5 (since the integrand is continuous):

$$\frac{ds}{dx} = \sqrt{1 + [f'(x)]^2} = \sqrt{1 + \left(\frac{dy}{dx}\right)^2}$$

Equation 6 shows that the rate of change of s with respect to x is always at least 1 and is equal to 1 when f'(x), the slope of the curve, is 0.

The differential of arc length is

7

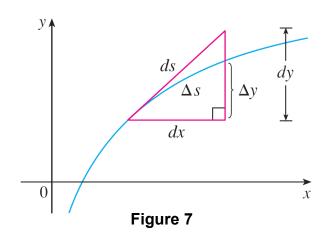
$$ds = \sqrt{1 + \left(\frac{dy}{dx}\right)^2} \, dx$$

and this equation is sometimes written in the symmetric form

8

$$(ds)^2 = (dx)^2 + (dy)^2$$

The geometric interpretation of Equation 8 is shown in Figure 7. It can be used as a mnemonic device for remembering both of the Formulas 3 and 4.



If we write  $L = \int ds$ , then from Equation 8 either we can solve to get (7), which gives (3), or we can solve to get

$$ds = \sqrt{1 + \left(\frac{dx}{dy}\right)^2} \, dy$$

### Example 4

Find the arc length function for the curve  $y = x^2 - \frac{1}{8} \ln x$  taking  $P_0(1, 1)$  as the starting point.

#### Solution:

If 
$$f(x) = x^2 - \frac{1}{8} \ln x$$
, then
$$f'(x) = 2x - \frac{1}{8x}$$

$$1 + [f'(x)]^2 = 1 + \left(2x - \frac{1}{8x}\right)^2 = 1 + 4x^2 - \frac{1}{2} + \frac{1}{64x^2}$$

$$= 4x^2 + \frac{1}{2} + \frac{1}{64x^2} = \left(2x + \frac{1}{8x}\right)^2$$

$$\sqrt{1 + [f'(x)]^2} = 2x + \frac{1}{8x}$$

### Example 4 – Solution

Thus the arc length function is given by

$$s(x) = \int_{1}^{x} \sqrt{1 + [f'(t)]^2} dt$$

$$= \int_{1}^{x} \left(2t + \frac{1}{8t}\right) dt = t^{2} + \frac{1}{8} \ln t \Big]_{1}^{x}$$

$$= x^2 + \frac{1}{8} \ln x - 1$$

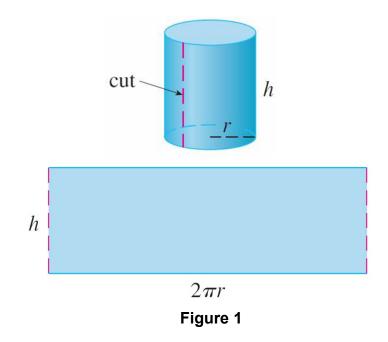
For instance, the arc length along the curve from (1, 1) to (3, f(3)) is

$$s(3) = 3^2 + \frac{1}{8} \ln 3 - 1 = 8 + \frac{\ln 3}{8} \approx 8.1373$$

A surface of revolution is formed when a curve is rotated about a line. Such a surface is the lateral boundary of a solid of revolution.

We want to define the area of a surface of revolution in such a way that it corresponds to our intuition. If the surface area is *A*, we can imagine that painting the surface would require the same amount of paint as does a flat region with area *A*.

Let's start with some simple surfaces. The lateral surface area of a circular cylinder with radius r and height h is taken to be  $A = 2\pi rh$  because we can imagine cutting the cylinder and unrolling it (as in Figure 1) to obtain a rectangle with dimensions  $2\pi r$  and h.



Likewise, we can take a circular cone with base radius r and slant height l, cut it along the dashed line in Figure 2, and flatten it to form a sector of a circle with radius l and central angle  $\theta = 2\pi r/l$ .

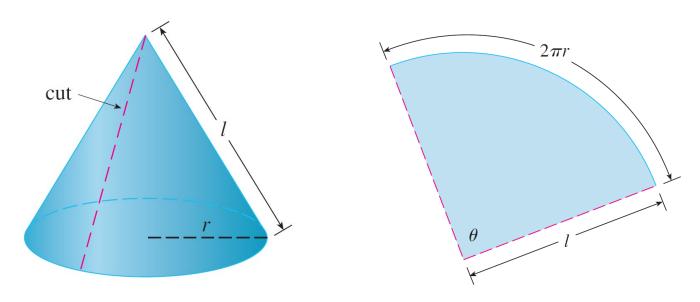


Figure 2

We know that, in general, the area of a sector of a circle with radius I and angle  $\theta$  is  $\frac{1}{2}I^2\theta$  and so in this case the area is

$$A = \frac{1}{2}I^2\theta = \frac{1}{2}l^2\left(\frac{2\pi r}{l}\right) = \pi rI$$

Therefore we define the lateral surface area of a cone to be  $A = \pi rI$ .

What about more complicated surfaces of revolution? If we follow the strategy we used with arc length, we can approximate the original curve by a polygon.

When this polygon is rotated about an axis, it creates a simpler surface whose surface area approximates the actual surface area.

By taking a limit, we can determine the exact surface area.

The approximating surface, then, consists of a number of bands, each formed by rotating a line segment about an axis.

To find the surface area, each of these bands can be considered a portion of a circular cone, as shown in Figure 3.

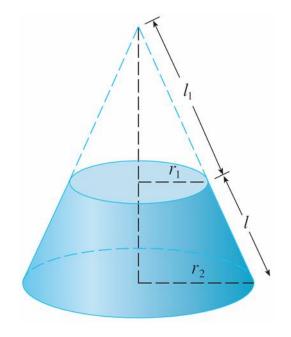


Figure 3

The area of the band (or frustum of a cone) with slant height l and upper and lower radii  $r_1$  and  $r_2$  is found by subtracting the areas of two cones:

$$A = \pi r_2 (l_1 + l) - \pi r_1 l_1 = \pi [(r_2 - r_1) l_1 + r_2 l]$$

From similar triangles we have

$$\frac{l_1}{r_1} = \frac{l_1 + l}{r_2}$$

which gives

$$r_2I_1 = r_1I_1 + r_1I$$
 or  $(r_2 - r_1)I_1 = r_1I$ 

Putting this in Equation 1, we get

$$A = \pi \left( r_1 I + r_2 I \right)$$

or

2

$$A = 2\pi rl$$

where  $r = \frac{1}{2}(r_1 + r_2)$  is the average radius of the band.

Now we apply this formula to our strategy. Consider the surface shown in Figure 4, which is obtained by rotating the curve y = f(x),  $a \le x \le b$ , about the *x*-axis, where *f* is positive and has a continuous derivative.

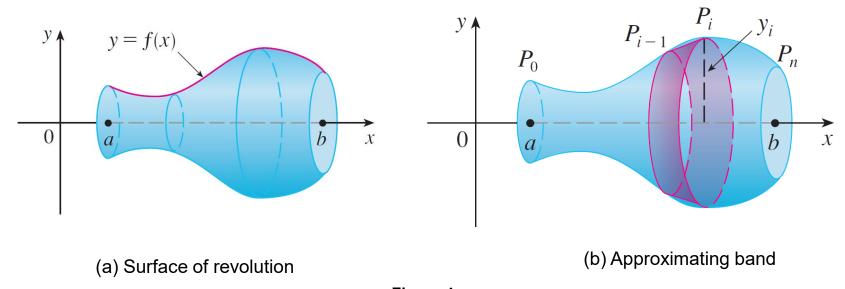


Figure 4

In order to define its surface area, we divide the interval [a, b] into n subintervals with endpoints  $x_0, x_1, \ldots, x_n$  and equal width  $\Delta x$ .

If  $y_i = f(x_i)$ , then the point  $P_i(x_i, y_i)$  lies on the curve.

The part of the surface between  $x_{i-1}$  and  $x_i$  is approximated by taking the line segment  $P_i - {}_1P_i$  and rotating it about the x-axis.

The result is a band with slant height  $I = |P_{i-1}P_i|$  and average radius  $r = \frac{1}{2}(y_{i-1} + y_i)$  so, by Formula 2, its surface area is

$$2\pi \frac{y_{i-1} + y_i}{2} |P_{i-1}P_i|$$

We have

$$|P_{i-1}P_i| = \sqrt{1 + [f'(x_i^*)]^2} \Delta x$$

where  $x_i^*$  is some number in  $[x_{i-1}, x_i]$ .

When  $\Delta x$  is small, we have  $y_i = f(x_i) \approx f(x_i^*)$  and also  $Y_{i-1} = f(x_{i-1}) \approx f(x_i^*)$ , since f is continuous.

#### **Therefore**

$$2\pi \frac{y_{i-1} + y_i}{2} |P_{i-1}P_i| \approx 2\pi f(x_i^*) \sqrt{1 + [f'(x_i^*)]^2} \Delta x$$

and so an approximation to what we think of as the area of the complete surface of revolution is

$$\sum_{i=1}^{n} 2\pi f(x_i^*) \sqrt{1 + [f'(x_i^*)]^2} \Delta x$$

This approximation appears to become better as  $n \to \infty$  and, recognizing as 3 Riemann sum for the function

$$g(x) = 2\pi f(x) \sqrt{1 + [f'(x)]^2}$$
, we have

$$\lim_{n \to \infty} \sum_{i=1}^{n} 2\pi f(x_i^*) \sqrt{1 + [f'(x_i^*)]^2} \Delta x = \int_a^b 2\pi f(x) \sqrt{1 + [f'(x)]^2} dx$$

Therefore, in the case where f is positive and has a continuous derivative, we define the **surface area** of the surface obtained by rotating the curve y = f(x),  $a \le x \le b$ , about the x-axis as



$$S = \int_{a}^{b} 2\pi f(x) \sqrt{1 + [f'(x)]^{2}} dx$$

With the Leibniz notation for derivatives, this formula becomes

$$S = \int_a^b 2\pi y \sqrt{1 + \left(\frac{dy}{dx}\right)^2} \, dx$$

If the curve is described as x = g(y),  $c \le y \le d$ , then the formula for surface area becomes

$$S = \int_{c}^{d} 2\pi y \sqrt{1 + \left(\frac{dx}{dy}\right)^{2}} \, dy$$

#### Area of a Surface of Revolution

Now both Formulas 5 and 6 can be summarized symbolically, using the notation for arc length, as

$$S = \int 2\pi y \, ds$$

For rotation about the *y*-axis, the surface area formula becomes

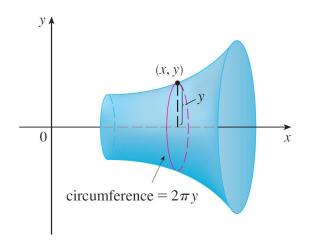
$$S = \int 2\pi x \, ds$$

where, as before, we can use either

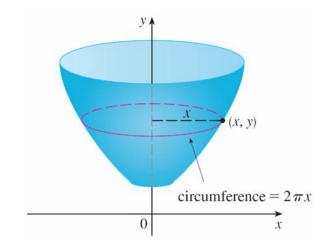
$$ds = \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$
 or  $ds = \sqrt{1 + \left(\frac{dx}{dy}\right)^2} dy$ 

### Area of a Surface of Revolution

These formulas can be remembered by thinking of  $2\pi y$  or  $2\pi x$  as the circumference of a circle traced out by the point (x, y) on the curve as it is rotated about the x-axis or y-axis, respectively (see Figure 5).



(a) Rotation about x-axis:  $S = \int 2\pi y \, ds$ 



(b) Rotation about *y*-axis:  $S = \int 2\pi x \, ds$ 

Figure 5

### Example 1

The curve  $y = \sqrt{4 - x^2}$ ,  $-1 \le x \le 1$ , is an arc of the circle  $x^2 + y^2 = 4$ .

Find the area of the surface obtained by rotating this arc about the *x*-axis. (The surface is a portion of a sphere of radius 2. See Figure 6.)

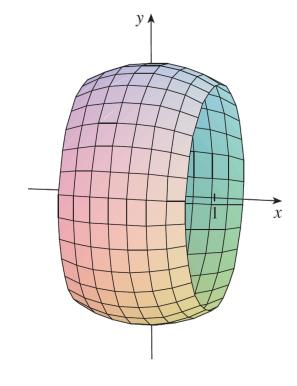


Figure 6

We have

$$\frac{dy}{dx} = \frac{1}{2}(4 - x^2)^{-1/2}(-2x) = \frac{-x}{\sqrt{4 - x^2}}$$

and so, by Formula 5, the surface area is

$$S = \int_{-1}^{1} 2\pi y \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$

$$=2\pi\int_{-1}^{1}\sqrt{4-x^2}\sqrt{1+\frac{x^2}{4-x^2}}\,dx$$

$$= 2\pi \int_{-1}^{1} \sqrt{4 - x^2} \frac{2}{\sqrt{4 - x^2}} dx$$

$$= 4\pi \int_{-1}^{1} 1 dx$$

$$= 4\pi (2)$$

$$= 8\pi$$

## Example 3

Find the area of the surface generated by rotating the curve  $y = e^x$ ,  $0 \le x \le 1$ , about the x-axis.

#### Solution:

Using Formula 5 with

$$y = e^x$$
 and  $\frac{dy}{dx} = e^x$ 

we have

$$S = \int_0^1 2\pi y \sqrt{1 + \left(\frac{dy}{dx}\right)^2} \, dx$$
$$= 2\pi \int_0^1 e^x \sqrt{1 + e^{2x}} \, dx$$

$$=2\pi\int_{1}^{e}\sqrt{1+u^{2}}\,du$$

(where 
$$u = e^x$$
)

$$=2\pi\int_{\pi/4}^{\alpha}\sec^3\theta\,d\theta$$

(where  $u = \tan \theta$  and  $\alpha = \tan^{-1} e$ )

$$= 2\pi \cdot \frac{1}{2} \left[ \sec \theta \tan \theta + \ln |\sec \theta + \tan \theta| \right]_{\pi/4}^{\alpha}$$

= 
$$\pi[\sec \alpha \tan \alpha + \ln(\sec \alpha + \tan \alpha) - \sqrt{2} - \ln(\sqrt{2} + 1)]$$

Since  $\tan \alpha = e$ , we have  $\sec^2 \alpha = 1 + \tan^2 \alpha = 1 + e^2$  and

$$S = \pi \left[ e + \sqrt{1 + e^2} + \ln(e + \sqrt{1 + e^2}) - \sqrt{2} - \ln(\sqrt{2} + 1) \right]$$

8.3

### Applications to Physics and Engineering

### Applications to Physics and Engineering

Among the many applications of integral calculus to physics and engineering, we consider two here: force due to water pressure, and centers of mass.

Deep-sea divers realize that water pressure increases as they dive deeper. This is because the weight of the water above them increases.

In general, suppose that a thin horizontal plate with area A square meters is submerged in a fluid of density  $\rho$  kilograms per cubic meter at a depth d meters below the surface of the fluid as in Figure 1.

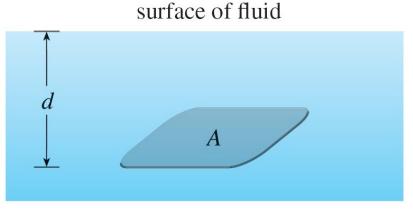


Figure 1

The fluid directly above the plate has volume V = Ad, so its mass is  $m = \rho V = \rho Ad$ . The force exerted by the fluid on the plate is therefore

$$F = mg = \rho gAd$$

where *g* is the acceleration due to gravity. The **pressure** *P* on the plate is defined to be the force per unit area:

$$P = \frac{F}{A} = \rho g d$$

The SI unit for measuring pressure is newtons per square meter, which is called a pascal (abbreviation:  $1 \text{ N/m}^2 = 1 \text{ Pa}$ ). Since this is a small unit, the kilopascal (kPa) is often used.

For instance, because the density of water is  $\rho$  = 1000 kg/m<sup>3</sup>, the pressure at the bottom of a swimming pool 2 m deep is

$$P = \rho gd = 1000 \text{ kg/m}^3 \times 9.8 \text{ m/s}^2 \times 2 \text{ m}$$
  
= 19,600 Pa = 19.6 kPa

An important principle of fluid pressure is the experimentally verified fact that at any point in a liquid the pressure is the same in all directions. (A diver feels the same pressure on nose and both ears.)

Thus the pressure in *any* direction at a depth d in a fluid with mass density  $\rho$  is given by

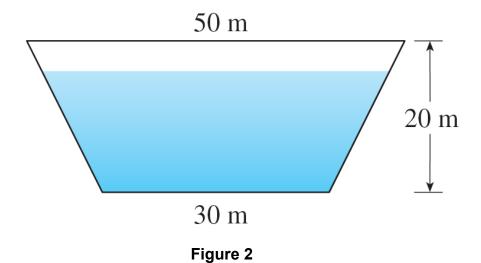
$$P = \rho gd = \delta d$$

This helps us determine the hydrostatic force against a vertical plate or wall or dam in a fluid.

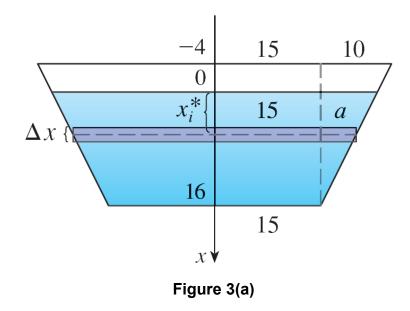
This is not a straightforward problem because the pressure is not constant but increases as the depth increases.

### Example 1

A dam has the shape of the trapezoid shown in Figure 2. The height is 20 m and the width is 50 m at the top and 30 m at the bottom. Find the force on the dam due to hydrostatic pressure if the water level is 4 m from the top of the dam.



We choose a vertical *x*-axis with origin at the surface of the water as in Figure 3(a).



The depth of the water is 16 m, so we divide the interval [0, 16] into subintervals of equal length with endpoints  $x_i$  and we choose  $x_i^* \in [x_{i-1}, x_i]$ .

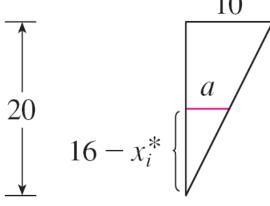
The *i*th horizontal strip of the dam is approximated by a rectangle with height  $\Delta x$  and width  $w_i$ , where, from similar triangles in Figure 3(b),

$$\frac{a}{16 - x_i^*} = \frac{10}{20}$$

or 
$$a = \frac{16 - x_i^*}{2} = 8 - \frac{x_i^*}{2}$$

and so

$$W_i = 2(15 + a)$$



$$= 2(15 + 8 - \frac{1}{2}x_i^*)$$

$$= 46 - x_i^*$$

If  $A_i$  is the area of the *i*th strip, then

$$A_i \approx w_i \, \Delta x$$
$$= (46 - x_i^*) \, \Delta x$$

If  $\Delta x$  is small, then the pressure  $P_i$  on the *i*th strip is almost constant and we can use Equation 1 to write

$$P_i \approx 1000gx_i^*$$

The hydrostatic force  $F_i$  acting on the ith strip is the product of the pressure and the area:

$$F_i = P_i A_i$$

$$\approx 1000 g x_i^* (46 - x_i^*) \Delta x$$

Adding these forces and taking the limit as  $n \to \infty$ , we obtain the total hydrostatic force on the dam:

$$F = \lim_{n \to \infty} \sum_{i=1}^{n} 1000gx_i^* (46 - x_i^*) \Delta x$$

$$= \int_0^{16} 1000gx (46 - x) dx$$

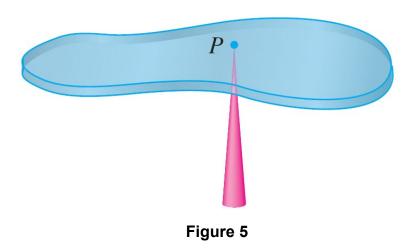
$$= 1000(9.8) \int_0^{16} (46x - x^2) dx$$

$$= 9800 \left[ 23x^2 - \frac{x^3}{3} \right]_0^{16}$$

$$\approx 4.43 \times 10^7 \text{ N}$$

Our main objective here is to find the point *P* on which a thin plate of any given shape balances horizontally as in Figure 8.

This point is called the **center of mass** (or center of gravity) of the plate.



We first consider the simpler situation illustrated in Figure 6, where two masses  $m_1$  and  $m_2$  are attached to a rod of negligible mass on opposite sides of a fulcrum and at distances  $d_1$  and  $d_2$  from the fulcrum.

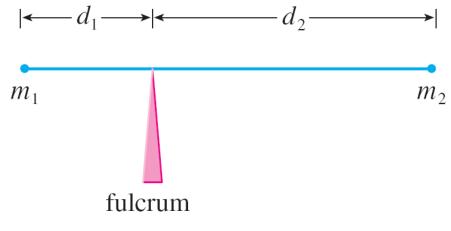


Figure 6

The rod will balance if

2

$$m_1d_1=m_2d_2$$

This is an experimental fact discovered by Archimedes and called the Law of the Lever. (Think of a lighter person balancing a heavier one on a seesaw by sitting farther away from the center.)

Now suppose that the rod lies along the *x*-axis with  $m_1$  at  $x_1$  and  $m_2$  at  $x_2$  and the center of mass at  $\overline{x}$ .

If we compare Figures 6 and 7, we see that  $d_1 = \bar{x} - x_1$  and  $d_2 = x_2 - \bar{x}$  and so Equation 2 gives

$$m_1(\bar{x} - x_1) = m_2(x_2 - \bar{x})$$
 $m_1\bar{x} + m_2\bar{x} = m_1x_1 + m_2x_2$ 
 $\bar{x} = \frac{m_1x_1 + m_2x_2}{m_1 + m_2}$ 

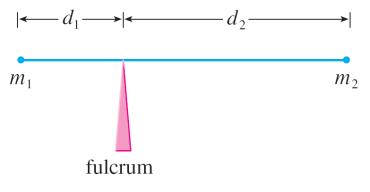


Figure 6

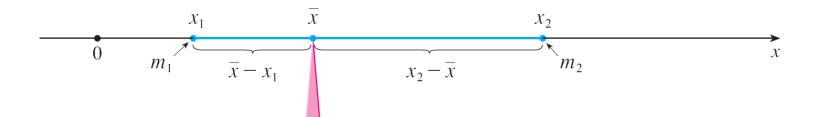


Figure 7

The numbers  $m_1x_1$  and  $m_2x_2$  are called the **moments** of the masses  $m_1$  and  $m_2$  (with respect to the origin), and Equation 3 says that the center of mass  $\overline{x}$  is obtained by adding the moments of the masses and dividing by the total mass  $m = m_1 + m_2$ .

In general, if we have a system of n particles with masses  $m_1, m_2, \ldots, m_n$  located at the points  $x_1, x_2, \ldots, x_n$  on the x-axis, it can be shown similarly that the center of mass of the system is located at

$$\overline{x} = \frac{\sum_{i=1}^{n} m_i x_i}{\sum_{i=1}^{n} m_i} = \frac{\sum_{i=1}^{n} m_i x_i}{m}$$

where  $m = \sum m_i$  is the total mass of the system, and the sum of the individual moments

$$M = \sum_{i=1}^n m_i x_i$$

is called the moment of the system about the origin.

Then Equation 4 could be rewritten as  $m\bar{x} = M$ , which says that if the total mass were considered as being concentrated at the center of mass  $\bar{x}$ , then its moment would be the same as the moment of the system.

Now we consider a system of n particles with masses  $m_1, m_2, \ldots, m_n$  located at the points  $(x_1, y_1), (x_2, y_2), \ldots, (x_n, y_n)$  in the xy-plane as shown in Figure 8.

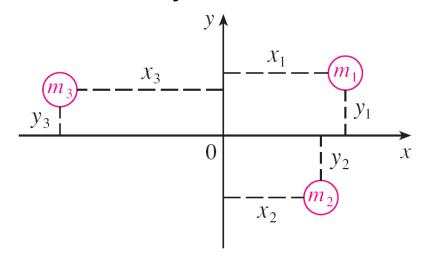


Figure 8

By analogy with the one-dimensional case, we define the **moment of the system about the** *y***-axis** to be

$$M_{y} = \sum_{i=1}^{n} m_{i} x_{i}$$

and the moment of the system about the x-axis as

$$M_x = \sum_{i=1}^n m_i y_i$$

Then  $M_y$  measures the tendency of the system to rotate about the *y*-axis and  $M_x$  measures the tendency to rotate about the *x*-axis.

As in the one-dimensional case, the coordinates  $(\bar{x}, \bar{y})$  of the center of mass are given in terms of the moments by the formulas

$$\overline{x} = \frac{M_y}{m}$$
  $\overline{y} = \frac{M_x}{m}$ 

where  $m = \sum m_i$  is the total mass. Since  $m\overline{x} = M_y$  and  $m\overline{y} = M_x$ , the center of mass  $(\overline{x}, \overline{y})$  is the point where a single particle of mass m would have the same moments as the system.

### Example 3

Find the moments and center of mass of the system of objects that have masses 3, 4, and 8 at the points (–1, 1), (2, –1), and (3, 2), respectively.

#### Solution:

We use Equations 5 and 6 to compute the moments:

$$M_y = 3(-1) + 4(2) + 8(3) = 29$$
  
 $M_x = 3(1) + 4(-1) + 8(2) = 15$ 

Since m = 3 + 4 + 8 = 15, we use Equations 11 to obtain

$$\overline{x} = \frac{M_y}{m}$$

$$= \frac{29}{15}$$

$$\overline{y} = \frac{M_x}{m}$$

$$= \frac{15}{15} = 1$$

Thus the center of mass is  $(1\frac{14}{15}, 1)$ . (See Figure 9.)

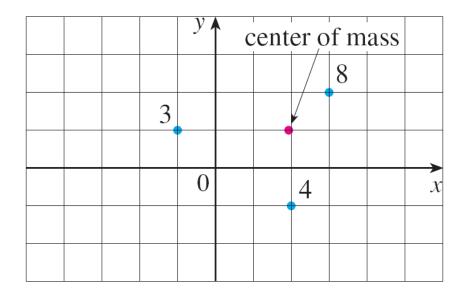


Figure 9

Next we consider a flat plate (called a *lamina*) with uniform density  $\rho$  that occupies a region  $\Re$  of the plane.

We wish to locate the center of mass of the plate, which is called the **centroid** of  $\Re$ .

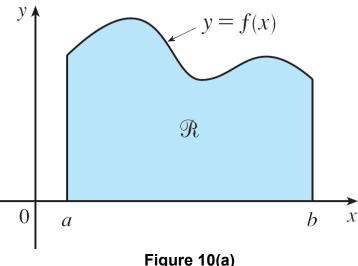
In doing so we use the following physical principles: The **symmetry principle** says that if  $\Re$  is symmetric about a line I, then the centroid of  $\Re$  lies on I. (If  $\Re$  is reflected about I, then  $\Re$  remains the same so its centroid remains fixed. But the only fixed points lie on I.)

Thus the centroid of a rectangle is its center.

Moments should be defined so that if the entire mass of a region is concentrated at the center of mass, then its moments remain unchanged.

Also, the moment of the union of two nonoverlapping regions should be the sum of the moments of the individual regions.

Suppose that the region  $\Re$  is of the type shown in Figure 10(a); that is,  $\Re$  lies between the lines x = a and x = b, above the x-axis, and beneath the graph of f, where f is a continuous function.



We divide the interval [a, b] into n subintervals with endpoints  $x_0, x_1, \ldots, x_n$  and equal width  $\Delta x$ . We choose the sample point  $x_i^*$  to be the midpoint  $\bar{x}_i$  of the ith subinterval, that is,  $\bar{x}_i = (x_{i-1} + x_i)/2$ .

This determines the polygonal approximation to R shown in

Figure 10(b).

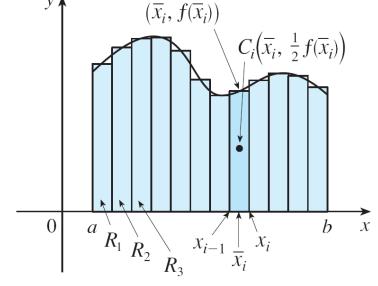


Figure 10(b)

The centroid of the *i*th approximating rectangle  $R_i$  is its center  $C_i(\bar{x}_i, \frac{1}{2}f(\bar{x}_i))$ . Its area is  $f(\bar{x}_i) \Delta x$ , so its mass is

$$\rho f(\overline{x}_i) \Delta x$$

The moment of  $R_i$  about the *y*-axis is the product of its mass and the distance from  $C_i$  to the *y*-axis, which is  $\bar{x}_i$ . Thus

$$M_{y}(R_{i}) = [\rho f(\overline{x}_{i}) \Delta x] \overline{x}_{i} = \rho \overline{x}_{i} f(\overline{x}_{i}) \Delta x$$

Adding these moments, we obtain the moment of the polygonal approximation to  $\Re$ , and then by taking the limit as  $n \to \infty$  we obtain the moment of  $\Re$  itself about the *y*-axis:

$$M_{y} = \lim_{n \to \infty} \sum_{i=1}^{n} \rho \overline{x}_{i} f(\overline{x}_{i}) \Delta x = \rho \int_{a}^{b} x f(x) dx$$

In a similar fashion we compute the moment of  $R_i$  about the x-axis as the product of its mass and the distance from  $C_i$  to the x-axis:

$$M_{x}(R_{i}) = \left[\rho f(\overline{x}_{i}) \Delta x\right] \frac{1}{2} f(\overline{x}_{i}) = \rho \cdot \frac{1}{2} \left[f(\overline{x}_{i})\right]^{2} \Delta x$$

Again we add these moments and take the limit to obtain the moment of  $\Re$  about the *x*-axis:

$$M_{x} = \lim_{n \to \infty} \sum_{i=1}^{n} \rho \cdot \frac{1}{2} [f(\bar{x}_{i})]^{2} \Delta x = \rho \int_{a}^{b} \frac{1}{2} [f(x)]^{2} dx$$

Just as for systems of particles, the center of mass of the plate is defined so that  $m\bar{x} = M_y$  and  $m\bar{y} = M_x$ . But the mass of the plate is the product of its density and its area:

$$m = \rho A = \rho \int_a^b f(x) \, dx$$

and so

$$\overline{x} = \frac{M_y}{m} = \frac{\rho \int_a^b x f(x) dx}{\rho \int_a^b f(x) dx} = \frac{\int_a^b x f(x) dx}{\int_a^b f(x) dx}$$

$$\overline{y} = \frac{M_x}{m} = \frac{\rho \int_a^b \frac{1}{2} [f(x)]^2 dx}{\rho \int_a^b f(x) dx} = \frac{\int_a^b \frac{1}{2} [f(x)]^2 dx}{\int_a^b f(x) dx}$$

Notice the cancellation of the  $\rho$ 's. The location of the center of mass is independent of the density.

In summary, the center of mass of the plate (or the centroid of  $\Re$ ) is located at the point  $(\bar{x}, \bar{y})$ , where

$$\overline{x} = \frac{1}{A} \int_a^b x f(x) \, dx \qquad \overline{y} = \frac{1}{A} \int_a^b \frac{1}{2} [f(x)]^2 \, dx$$

If the region  $\Re$  lies between two curves y = f(x) and y = g(x), where  $f(x) \ge g(x)$ , as illustrated in Figure 13, then the same sort of argument that led to Formulas 8 can be used to show that the centroid of  $\Re$  is  $(\bar{x}, \bar{y})$ , where

9

$$\overline{x} = \frac{1}{A} \int_a^b x [f(x) - g(x)] dx$$

$$\overline{y} = \frac{1}{A} \int_a^b \frac{1}{2} \{ [f(x)]^2 - [g(x)]^2 \} dx$$

We end this section by showing a surprising connection between centroids and volumes of revolution.

**Theorem of Pappus** Let  $\mathcal{R}$  be a plane region that lies entirely on one side of a line l in the plane. If  $\mathcal{R}$  is rotated about l, then the volume of the resulting solid is the product of the area A of  $\mathcal{R}$  and the distance d traveled by the centroid of  $\mathcal{R}$ .

8.4

#### Applications to Economics and Biology

### Applications to Economics and Biology

In this section we consider some applications of integration to economics (consumer surplus) and biology (blood flow, cardiac output).

Recall that the demand function p(x) is the price that a company has to charge in order to sell x units of a commodity.

Usually, selling larger quantities requires lowering prices, so the demand function is a decreasing function. The graph of a typical demand function, called a demand curve, is shown in Figure 1.

If X is the amount of the commodity that is currently available, then P = p(X) is the current selling price.

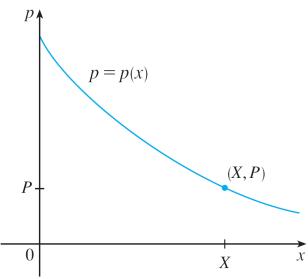
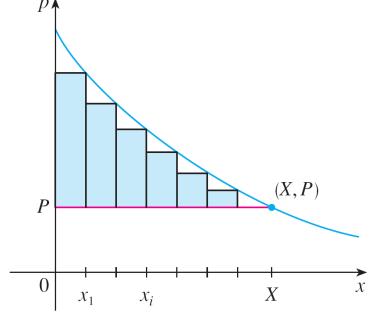


Figure 1
A typical demand curve

We divide the interval [0, X] into n subintervals, each of length  $\Delta x = X/n$ , and let  $x_i^* = x_i$  be the right endpoint of the ith subinterval, as in Figure 2.

If, after the first  $x_{i-1}$  units were sold, a total of only  $x_i$  units had been available and the price per unit had been set at  $p(x_i)$  dollars, then the additional  $\Delta x$  units could have been sold (but no more).



The consumers who would have paid  $p(x_i)$  dollars placed a high value on the product; they would have paid what it was worth to them.

So, in paying only P dollars they have saved an amount of (savings per unit) (number of units) =  $[p(x_i) - P] \Delta x$ 

Considering similar groups of willing consumers for each of the subintervals and adding the savings, we get the total savings:

$$\sum_{i=1}^{n} [p(x_i) - P] \Delta x$$

(This sum corresponds to the area enclosed by the rectangles in Figure 2.)

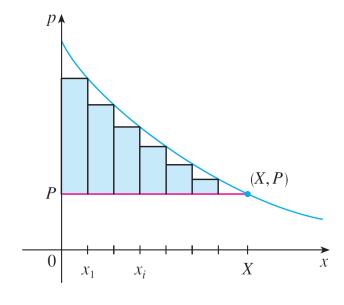


Figure 2

If we let  $n \to \infty$ , this Riemann sum approaches the integral

$$\int_0^X \left[ p(x) - P \right] dx$$

which economists call the **consumer surplus** for the commodity.

The consumer surplus represents the amount of money saved by consumers in purchasing the commodity at price *P*, corresponding to an amount demanded of *X*.

Figure 3 shows the interpretation of the consumer surplus as the area under the demand curve and above the line p = P.

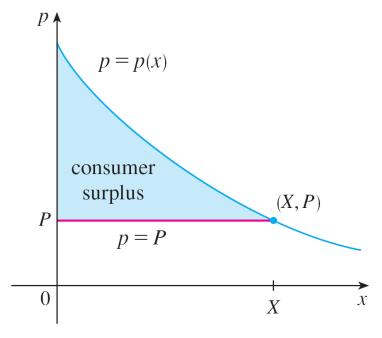


Figure 3

# Example 1

The demand for a product, in dollars, is

$$p = 1200 - 0.2x - 0.0001x^2$$

Find the consumer surplus when the sales level is 500.

#### Solution:

Since the number of products sold is X = 500, the corresponding price is

$$P = 1200 - (0.2)(500) - (0.0001)(500)^2$$
$$= 1075$$

# Example 1 – Solution

Therefore, from Definition 1, the consumer surplus is

$$\int_0^{500} [p(x) - P] dx = \int_0^{500} (1200 - 0.2x - 0.0001x^2 - 1075) dx$$

$$= \int_0^{500} (125 - 0.2x - 0.0001x^2) dx$$

$$= 125x - 0.1x^2 - (0.0001) \left(\frac{x^3}{3}\right) \Big]_0^{500}$$

$$= (125)(500) - (0.1)(500)^2 - \frac{(0.0001)(500)^3}{3}$$

$$= $33,333.33$$

We have discussed the law of laminar flow:

$$v(r) = \frac{P}{4\eta l} \left( R^2 - r^2 \right)$$

which gives the velocity v of blood that flows along a blood vessel with radius R and length I at a distance r from the central axis, where P is the pressure difference between the ends of the vessel and  $\eta$  is the viscosity of the blood.

Now, in order to compute the rate of blood flow, or *flux* (volume per unit time), we consider smaller, equally spaced radii  $r_1, r_2, \ldots$ 

The approximate area of the ring (or washer) with inner radius  $r_{i-1}$  and outer radius  $r_i$  is

 $2\pi r_i \Delta r$ 

where

$$\Delta r = r_i - r_{i-1}$$

(See Figure 4.)

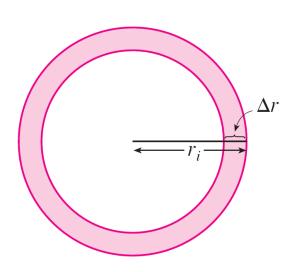


Figure 4

If  $\Delta r$  is small, then the velocity is almost constant throughout this ring and can be approximated by  $v(r_i)$ .

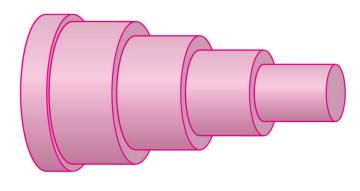
Thus the volume of blood per unit time that flows across the ring is approximately

$$(2\pi r_i \Delta r) \ v(r_i) = 2\pi r_i \ v(r_i) \Delta r$$

and the total volume of blood that flows across a cross-section per unit time is about

$$\sum_{i=1}^{n} 2\pi r_i \ v(r_i) \ \Delta r$$

This approximation is illustrated in Figure 5.



Notice that the velocity (and hence the volume per unit time) increases toward the center of the blood vessel.

The approximation gets better as *n* increases.

When we take the limit we get the exact value of the **flux** (or *discharge*), which is the volume of blood that passes a cross-section per unit time:

$$F = \lim_{n \to \infty} \sum_{i=1}^{n} 2\pi r_i v(r_i) \Delta r$$
$$= \int_{0}^{R} 2\pi r v(r) dr$$

$$= \int_{0}^{R} 2\pi r \frac{P}{4\eta l} (R^{2} - r^{2}) dr$$

$$= \frac{\pi P}{2\eta l} \int_{0}^{R} (R^{2}r - r^{3}) dr$$

$$= \frac{\pi P}{2\eta l} \left[ R^{2} \frac{r^{2}}{2} - \frac{r^{4}}{4} \right]_{r=0}^{r=R}$$

$$= \frac{\pi P}{2\eta l} \left[ \frac{R^{4}}{2} - \frac{R^{4}}{4} \right]$$

$$\frac{\pi P R^{4}}{2}$$

The resulting equation

$$F = \frac{\pi P R^{+}}{8\eta l}$$

is called **Poiseuille's Law**; it shows that the flux is proportional to the fourth power of the radius of the blood vessel.

Figure 6 shows the human cardiovascular system.

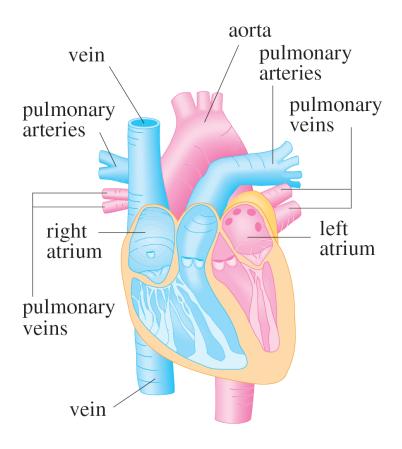


Figure 6

Blood returns from the body through the veins, enters the right atrium of the heart, and is pumped to the lungs through the pulmonary arteries for oxygenation.

It then flows back into the left atrium through the pulmonary veins and then out to the rest of the body through the aorta.

The **cardiac output** of the heart is the volume of blood pumped by the heart per unit time, that is, the rate of flow into the aorta.

The *dye dilution method* is used to measure the cardiac output.

Dye is injected into the right atrium and flows through the heart into the aorta. A probe inserted into the aorta measures the concentration of the dye leaving the heart at equally spaced times over a time interval [0, T] until the dye has cleared.

Let c(t) be the concentration of the dye at time t. If we divide [0, T] into subintervals of equal length  $\Delta t$ , then the amount of dye that flows past the measuring point during the subinterval from  $t = t_{i-1}$  to  $t = t_i$  is approximately

(concentration) (volume) =  $c(t_i)$  ( $F \Delta t$ )

where *F* is the rate of flow that we are trying to determine.

Thus the total amount of dye is approximately

$$\sum_{i=1}^{n} c(t_i) F \Delta t = F \sum_{i=1}^{n} c(t_i) \Delta t$$

and, letting  $n \to \infty$ , we find that the amount of dye is

$$A = F \int_0^T c(t) dt$$

Thus the cardiac output is given by

$$F = \frac{A}{\int_0^T c(t) dt}$$

where the amount of dye A is known and the integral can be approximated from the concentration readings.

### Example 2

A 5-mg bolus of dye is injected into a right atrium. The concentration of the dye (in milligrams per liter) is measured in the aorta at one-second intervals as shown in the chart. Estimate the cardiac output.

| t | c(t) | t  | c(t) |
|---|------|----|------|
| 0 | 0    | 6  | 6.1  |
| 1 | 0.4  | 7  | 4.0  |
| 2 | 2.8  | 8  | 2.3  |
| 3 | 6.5  | 9  | 1.1  |
| 4 | 9.8  | 10 | 0    |
| 5 | 8.9  |    |      |

### Example 2 – Solution

Here A = 5,  $\Delta t = 1$ , and T = 10. We use Simpson's Rule to approximate the integral of the concentration:

$$\int_0^{10} c(t) dt \approx \frac{1}{3} \left[ 0 + 4(0.4) + 2(2.8) + 4(6.5) + 2(9.8) + 4(8.9) + 2(6.1) + 4(4.0) + 2(2.3) + 4(1.1) + 0 \right]$$

 $\approx 41.87$ 

Thus Formula 3 gives the cardiac output to be

$$F = \frac{A}{\int_0^{10} c(t) dt} \approx \frac{5}{41.87}$$
$$\approx 0.12 \text{ L/s}$$
$$= 7.2 \text{ L/min}$$

8.5

# Probability

Calculus plays a role in the analysis of random behavior. Suppose we consider the cholesterol level of a person chosen at random from a certain age group, or the height of an adult female chosen at random, or the lifetime of a randomly chosen battery of a certain type.

Such quantities are called **continuous random variables** because their values actually range over an interval of real numbers, although they might be measured or recorded only to the nearest integer.

We might want to know the probability that a blood cholesterol level is greater than 250, or the probability that the height of an adult female is between 60 and 70 inches, or the probability that the battery we are buying lasts between 100 and 200 hours.

If X represents the lifetime of that type of battery, we denote this last probability as follows:

$$P(100 \le X \le 200)$$

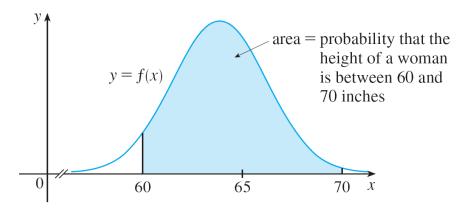
According to the frequency interpretation of probability, this number is the long-run proportion of all batteries of the specified type whose lifetimes are between 100 and 200 hours. Since it represents a proportion, the probability naturally falls between 0 and 1.

Every continuous random variable *X* has a **probability density function** *f*. This means that the probability that *X* lies between *a* and *b* is found by integrating *f* from *a* to *b*:

$$P(a \le X \le b) = \int_a^b f(x) \, dx$$

For example, Figure 1 shows the graph of a model for the probability density function f for a random variable X defined to be the height in inches of an adult female in the United States (according to data from the National Health Survey).

The probability that the height of a woman chosen at random from this population is between 60 and 70 inches is equal to the area under the graph of *f* from 60 to 70.



Probability density function for the height of an adult female

Figure 1

In general, the probability density function f of a random variable X satisfies the condition  $f(x) \ge 0$  for all x.

Because probabilities are measured on a scale from 0 to 1, it follows that

$$\int_{-\infty}^{\infty} f(x) \ dx = 1$$

## Example 1

Let f(x) = 0.006x(10 - x) for  $0 \le x \le 10$  and f(x) = 0 for all other values of x.

- (a) Verify that f is a probability density function.
- **(b)** Find  $P(4 \le X \le 8)$ .

#### Solution:

(a) For  $0 \le x \le 10$  we have  $0.006x(10 - x) \ge 0$ , so  $f(x) \ge 0$  for all x. We also need to check that Equation 2 is satisfied:

$$\int_{-\infty}^{\infty} f(x) dx = \int_{0}^{10} 0.006x(10 - x) dx$$
$$= 0.006 \int_{0}^{10} (10x - x^{2}) dx$$

## Example 1 – Solution

$$= 0.006 \left[ 5x^2 - \frac{1}{3}x^3 \right]_0^{10}$$
$$= 0.006 \left( 500 - \frac{1000}{3} \right)$$
$$= 1$$

Therefore *f* is a probability density function.

## Example 1 – Solution

**(b)** The probability that X lies between 4 and 8 is

$$P(4 \le X \le 8) = \int_{4}^{8} f(x) dx$$

$$= 0.006 \int_{4}^{8} (10x - x^{2}) dx$$

$$= 0.006 \left[ 5x^{2} - \frac{1}{3}x^{3} \right]_{4}^{8}$$

$$= 0.544$$

Suppose you're waiting for a company to answer your phone call and you wonder how long, on average, you can expect to wait.

Let f(t) be the corresponding density function, where t is measured in minutes, and think of a sample of N people who have called this company.

Most likely, none of them had to wait more than an hour, so let's restrict our attention to the interval  $0 \le t \le 60$ .

Let's divide that interval into n intervals of length  $\Delta t$  and endpoints 0,  $t_1$ ,  $t_2$ , . . . ,  $t_{60}$ . (Think of  $\Delta t$  as lasting a minute, or half a minute, or 10 seconds, or even a second.)

The probability that somebody's call gets answered during the time period from  $t_{i-1}$  to  $t_i$  is the area under the curve y = f(t) from  $t_{i-1}$  to  $t_i$ , which is approximately equal to  $f(\bar{t}_i) \Delta t$ . (This is the area of the approximating rectangle in Figure 3, where  $\bar{t}_i$  is the midpoint of the interval.)

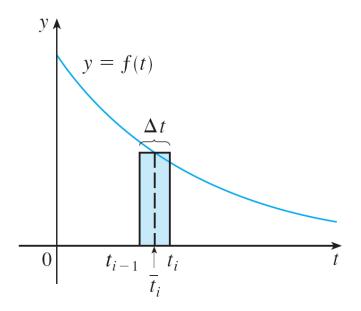


Figure 3

Since the long-run proportion of calls that get answered in the time period from  $t_{i-1}$  to  $t_i$  is  $f(\bar{t}_i) \Delta t$ , we expect that, out of our sample of N callers, the number whose call was answered in that time period is approximately  $N f(\bar{t}_i) \Delta t$  and the time that each waited is about  $\bar{t}_i$ .

Therefore the total time they waited is the product of these numbers: approximately  $\bar{t}_i$  [ $N f(\bar{t}_i) \Delta t$ ].

Adding over all such intervals, we get the approximate total of everybody's waiting times:

$$\sum_{i=1}^{n} N \overline{t}_{i} f(\overline{t}_{i}) \Delta t$$

If we now divide by the number of callers N, we get the approximate average waiting time:

$$\sum_{i=1}^n \overline{t}_i f(\overline{t}_i) \ \Delta t$$

We recognize this as a Riemann sum for the function t f(t). As the time interval shrinks (that is,  $\Delta t \to 0$  and  $n \to \infty$ ), this Riemann sum approaches the integral

$$\int_0^{60} t f(t) dt$$

This integral is called the *mean waiting time*.

In general, the **mean** of any probability density function *f* is defined to be

$$\mu = \int_{-\infty}^{\infty} x f(x) \, dx$$

The mean can be interpreted as the long-run average value of the random variable *X*. It can also be interpreted as a measure of centrality of the probability density function.

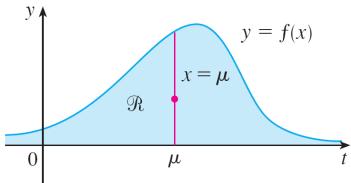
The expression for the mean resembles an integral we have seen before.

If  $\Re$  is the region that lies under the graph of f, we know that the x-coordinate of the centroid of  $\Re$  is

$$\overline{x} = \frac{\int_{-\infty}^{\infty} x f(x) dx}{\int_{-\infty}^{\infty} f(x) dx} = \int_{-\infty}^{\infty} x f(x) dx = \mu$$

because of Equation 2.

So a thin plate in the shape of  $\Re$  balances at a point on the vertical line  $x = \mu$ . (See Figure 4.)



 $\mathfrak R$  balances at a point on the line  $x = \mu$ 

# Example 3

Find the mean of the exponential distribution of Example 2:

$$f(t) = \begin{cases} 0 & \text{if } t < 0 \\ ce^{-ct} & \text{if } t \ge 0 \end{cases}$$

#### Solution:

According to the definition of a mean, we have

$$\mu = \int_{-\infty}^{\infty} t f(t) \, dt$$

$$= \int_0^\infty tce^{-ct} dt$$

### Example 3 – Solution

To evaluate this integral we use integration by parts, with u = t and  $dv = ce^{-ct} dt$ :

$$\int_0^\infty tce^{-ct} dt = \lim_{x \to \infty} \int_0^x tce^{-ct} dt$$

$$= \lim_{x \to \infty} \left( -te^{-ct} \right]_0^x + \int_0^x e^{-ct} dt$$

$$= \lim_{x \to \infty} \left( -xe^{-cx} + \frac{1}{c} - \frac{e^{-cx}}{c} \right)$$

$$= \frac{1}{c}$$

## Example 3 – Solution

The mean is  $\mu = 1/c$ , so we can rewrite the probability density function as

$$f(t) = \begin{cases} 0 & \text{if } t < 0 \\ \mu^{-1} e^{-t/\mu} & \text{if } t \ge 0 \end{cases}$$

Another measure of centrality of a probability density function is the *median*.

That is a number m such that half the callers have a waiting time less than m and the other callers have a waiting time longer than m. In general, the **median** of a probability density function is the number m such that

$$\int_{m}^{\infty} f(x) \ dx = \frac{1}{2}$$

This means that half the area under the graph of f lies to the right of m.

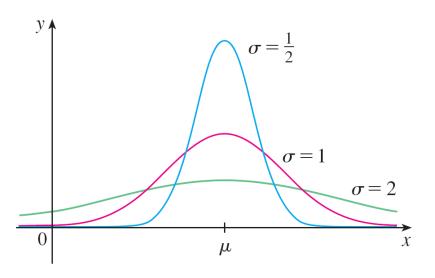
Many important random phenomena—such as test scores on aptitude tests, heights and weights of individuals from a homogeneous population, annual rainfall in a given location—are modeled by a **normal distribution**.

This means that the probability density function of the random variable *X* is a member of the family of functions

3 
$$f(x) = \frac{1}{\sigma\sqrt{2\pi}} e^{-(x-\mu)^2/(2\sigma^2)}$$

You can verify that the mean for this function is  $\mu$ . The positive constant  $\sigma$  is called the **standard deviation**; it measures how spread out the values of X are.

From the bell-shaped graphs of members of the family in Figure 5, we see that for small values of  $\sigma$  the values of X are clustered about the mean, whereas for larger values of  $\sigma$  the values of X are more spread out.



Normal distributions

Figure 5

Statisticians have methods for using sets of data to estimate  $\mu$  and  $\sigma$ .

The factor  $1/(\sigma\sqrt{2\pi})$  is needed to make f a probability density function. In fact, it can be verified using the methods of multivariable calculus that

$$\int_{-\infty}^{\infty} \frac{1}{\sigma \sqrt{2\pi}} e^{-(x-\mu)^2/(2\sigma^2)} dx = 1$$

#### Example 5

Intelligence Quotient (IQ) scores are distributed normally with mean 100 and standard deviation 15. (Figure 6 shows the corresponding probability density function.)

(a) What percentage of the population has an IQ score between 85 and 115?

(b) What percentage of the population has an IQ above 140?

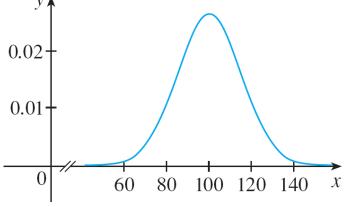


Figure 6

#### Example 5(a) - Solution

Since IQ scores are normally distributed, we use the probability density function given by Equation 3 with  $\mu$  = 100 and  $\sigma$  = 15:

$$P(85 \le X \le 115) = \int_{85}^{115} \frac{1}{15\sqrt{2\pi}} e^{-(x-100)^2/(2\cdot15^2)} dx$$

The function  $y = e^{-x^2}$  doesn't have an elementary antiderivative, so we can't evaluate the integral exactly.

#### Example 5(a) – Solution

But we can use the numerical integration capability of a calculator or computer (or the Midpoint Rule or Simpson's Rule) to estimate the integral.

Doing so, we find that

$$P(85 \le X \le 115) \approx 0.68$$

So about 68% of the population has an IQ between 85 and 115, that is, within one standard deviation of the mean.

#### Example 5(b) – Solution

The probability that the IQ score of a person chosen at random is more than 140 is

$$P(X > 140) = \int_{140}^{\infty} \frac{1}{15\sqrt{2\pi}} e^{-(x-100)^2/450} dx$$

To avoid the improper integral we could approximate it by the integral from 140 to 200. (It's quite safe to say that people with an IQ over 200 are extremely rare.)

Then

$$P(X > 140) \approx \int_{140}^{200} \frac{1}{15\sqrt{2\pi}} e^{-(x-100)^2/450} dx$$
$$\approx 0.0038$$

Therefore about 0.4% of the population has an IQ over 140.