1 Rotating Frames

• Apply twice to Newton II: e.o.m in rot. frame:

$$m \frac{\mathrm{d}^2 \mathbf{r}}{\mathrm{d}r^2} \bigg|_R = \mathbf{F} \underbrace{-2m\boldsymbol{\omega} \times \mathbf{v}_R}_{\text{Coriolis Force}} \underbrace{-m\boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})}_{\text{Centrifugal Force}}$$

2 Rigid Bodies

1. to obtain centre of mass at origin, solve:

$$\mathbf{R} = \frac{1}{M} \sum_{a} m_a \mathbf{r}_a = 0, \qquad \mathbf{R} = \frac{1}{M} \int_{V} \rho(\mathbf{r}) \mathbf{r} \, \mathrm{d}^3 r = 0$$

2. tensor of inertia:

discrete
$$I_{ij} = \sum_{a} m_a (|\mathbf{r}_a|^2 \delta_{ij} - r_{ai} r_{aj})$$

continuous $I_{ij} = \int_{V} \rho(\mathbf{r}) (|\mathbf{r}|^2 \delta_{ij} - r_{i} r_{j}) d^3 r$

hint: use symmetry \rightarrow diagonal in principal axes basis. I_1 along rot. axis; $I_2 = I_3$ along any perpendiculars

strategy: calculate
$$X = \sum_a m_a x_a^2$$
, $Y = \dots$, $Z = \dots$ then $I_x = Y + Z$, $I_y = X + Y$, $I_z = \dots$

hint: if the body is totally symmetric: $I_1=I_2=I_3,$ calculate the sum $3I=2\rho\int \mathrm{d}V\,r^2$

3. angular momentum: $m{L} = \overline{m{ar{I}}} m{\omega}$ and torque $m{ au} = \dot{m{L}}|_I = m{r} imes m{F}$

state: in object's frame $\dot{\overline{I}}|_R = 0 \implies \dot{L}|_R = \overline{\overline{I}}\dot{\omega}$ use: relation between inertial and rotating frame:

$$\frac{\mathrm{d}\boldsymbol{L}}{\mathrm{d}t}\bigg|_{I} = \frac{\mathrm{d}\boldsymbol{L}}{\mathrm{d}t}\bigg|_{R} + \boldsymbol{\omega} \times \boldsymbol{L} \tag{1}$$

4. kinetic energy: $T = \frac{1}{2}m\dot{R}^2 + \frac{1}{2}\sum_i I_i\omega_i^2$

3 Lagrangian Mechanics

1. Conservative system: L = T - V, H = T + V

T: total kinetic energy: $\frac{1}{2}mv^2 + \frac{1}{2}I\omega^2$

V: total potential energy

2. conservation laws

$$\begin{split} \frac{\partial L}{\partial t} &= 0 &\iff H = \sum_i p_i \dot{q}_i - L \quad \text{is conserved} \\ \frac{\partial L}{\partial q_i} &= 0 &\iff p_i = \frac{\partial L}{\partial \dot{q}_i} \quad \text{is conserved} \end{split}$$

3. Euler-Lagrange equations (no need to derive)

$$\frac{\delta L}{\delta q_i} = \frac{\partial L}{\partial q_i} - \frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{\partial L}{\partial \dot{q}_i} \right) = 0$$

3.1 Small Oscillations

- equilibrium points ${m q}(t)=\overline{m q}$ satisfy $\left.\frac{\partial V}{\partial q_i}\right|_{\overline{m q}}=0$
- perturbations about eq. points: $q_i(t) = \overline{q}_i + \varepsilon \delta q_i(t)$
 - 1. expand L to second order in ε , $\mathbf{q} = (\delta q_1, \dots, \delta q_N)$ hint: ignore \overline{q}_i and ε cancels, so only $O(\varepsilon^2)$ matters
 - 2. diagonalise kinetic term $T = \frac{1}{2}\dot{q} \cdot \dot{q}$ hint: build perfect squares, define normalised q
 - 3. Lagrangian allows identification of k

$$L = \frac{1}{2}\dot{\boldsymbol{q}}\cdot\dot{\boldsymbol{q}} - \frac{1}{2}\boldsymbol{q}\cdot\underline{\underline{\boldsymbol{k}}}\boldsymbol{q}, \qquad k_{ij} = \left.\frac{\partial^2 V}{\partial q_i\partial q_j}\right|_{\boldsymbol{q}=\bar{\boldsymbol{q}}}$$

- 4. EL eqns: $\ddot{q} = -\underline{k}q$ have soln: $q = Ae^{\pm i\omega t}$
- 5. look for eigenvalues: $\underline{\underline{k}} A_n = \omega_n^2 A_n$ normal modes: eigenvectors A_n normal frequencies: root of eigenvalues ω_n^2 stability: \overline{q} only stable if all eigenvalues $\omega_n^2 > 0$ hint: 2×2 matrix, stable if

$$\operatorname{Tr} \underline{\mathbf{k}} = \sum \omega_n^2 > 0$$
 AND $\det \underline{\mathbf{k}} = \prod \omega_n^2 > 0$

4 Hamiltonian Mechanics

Hamilton's equations:

$$\dot{q}_i = \{q_i, H\} = \frac{\partial H}{\partial p_i} \qquad \dot{p}_i = \{p_i, H\} = -\frac{\partial H}{\partial q_i}$$
 (2)

Poisson Brackets:

$$\{F, H\} = \sum_{i=1}^{N} \left(\frac{\partial F}{\partial q_i} \frac{\partial H}{\partial p_i} - (p_i \leftrightarrow q_i) \right)$$
 (3)

• time-evolution of F(q, p, t):

$$\frac{\mathrm{d}F}{\mathrm{d}t} = \frac{\partial F}{\partial t} + \{F, H\}$$

• canonical coordinates: $\{q_i, p_i\} = \delta_{ij}$

4.1 Noether's theorem

Any phase-space function G(q, p) generates transformations:

$$\delta q_i = \frac{\partial G}{\partial p_i} \delta \lambda, \qquad \delta p_i = -\frac{\partial G}{\partial q_i} \delta \lambda,$$

If G is a generator of symmetry, then

$$\frac{\mathrm{d}H}{\mathrm{d}\lambda} = \{H, G\} = 0 \iff \frac{\mathrm{d}G}{\mathrm{d}t} = \{G, H\} = 0 \tag{4}$$

then G(q, p) is called a conserved charge.

5 Relativistic EM [(-+++), c=1]

1. field strength tensor: $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$

$$A_{\mu} = (-\phi, \mathbf{A})$$
: photon vector field

2. Lorentz law (tensor eq: valid in any frame)

$$\frac{\mathrm{d}p^{\mu}}{\mathrm{d}\tau} = qF^{\mu}_{\ \nu}u^{\nu} \quad \text{(provided in exam)}$$

3. Gauss' & Ampèere's laws:

$$\partial_{\mu}F^{\mu}_{\ \nu} = -\mu_0 J_{\nu}$$
 (provided in exam)

$$J_{\mu} = (-\rho, \boldsymbol{J})$$
: charge density

4. Lorentz & Gauge transformation

- physics invariant under these
- can choose frame / gauge that simplifies calculations

5.1 Lorentz Transforms

$$x'^{\mu'} = \Lambda^{\mu'}_{\ \mu} x^{\mu}$$

Need to be able to write down LT matrix for boosts in any direction (x, y, z)

$$\Lambda^{\mu}_{\ \nu} = \begin{pmatrix} \gamma & -v\gamma & 0 & 0 \\ -v\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \qquad \gamma = \frac{1}{\sqrt{1-v^2}}$$

check: if v=0, same frame, so $\Lambda^{\mu'}_{\ \mu}=\delta^{\mu'}_{\mu}$

 $\bullet \,$ can just state that $(\Lambda^{\mu}{}_{\nu})^{-1} = \Lambda_{\mu}{}^{\nu}$ has $v \to -v$

hint: raising both indices changes sign in first row & col:

$$F_{\mu\nu} = \left(\begin{array}{c|c} 0 & x \\ \hline -x & y \end{array}\right) \longleftrightarrow F^{\mu\nu} = \left(\begin{array}{c|c} 0 & -x \\ \hline x & y \end{array}\right)$$

- ullet need to be able to derive how $E\ \&\ B$ transform
- matrix multiplication: contract second with first index:

$$C = A \cdot B \iff C_{\mu\nu} = A_{\mu}{}^{\alpha} B_{\alpha\nu}$$

5.2 Gauge Transforms

$$A_{\mu} \xrightarrow{\chi} A_{\mu} + \partial_{\mu}\chi, \qquad F_{\mu\nu} \xrightarrow{\chi} F_{\mu\nu}$$

Lorentz Gauge: $\partial_{\mu}A^{\mu} = 0$ (by using $\Box \chi = -\partial_{\mu}A^{\mu}$)

 $\partial_{\mu}F^{\mu}_{\ \nu}=-\mu_{0}J_{\nu}$ becomes wave equation: $\Box A_{\nu}=0$

Coulomb Gauge: $\nabla \cdot \mathbf{A} = 0$

Weyl / Temporal Gauge: $A^0 = \phi = 0$