

# 1 Abstract Group Theory

Lagrange's Theorem:

- For a subgroup  $H$  of  $G$ :  $\frac{|G|}{|H|} = i \in \mathbb{N}^+$

**index  $i$ :** number of distinct cosets

$\implies$  Groups of prime order are cyclic

**pitfall:** you can have none or more than one subgroup of the same order

Abelian groups:

- Group containing only second order elements are Abelian.
- Cyclic groups are Abelian.
- $\implies$  A group of prime order is Abelian.
- all subgroups of an Abelian group are invariant

First group isomorphism theorem:

- $f : G \rightarrow L$  is injective (isomorphism) if  $\ker(f) = E$
- a matrix representation is faithful / isomorphic to  $G$  if  $D(g) \neq I$  for  $g \neq E$

$\implies$  no other element has the same character as  $E$

Schur's Lemmas:  $D^{(i)}(G)M = MD^{(j)}(G) \implies M = \delta_{ij}\alpha I$

## 1.1 Conjugate Classes

- identity  $E$  forms a class of its own
- conj. elements have same order:  $g^k = E \implies (g')^k = E$
- $G$  is abelian iff every conjugacy class has size one

Point Groups:

- conjugation changes axis of rotation
- $\implies$  classes consist of (im)proper rotations of the same angle, such that their axes can be mapped onto each other
- conjugating a rotation by a reflection yields the inverse rotation
- $\implies C_n$  and  $C_n^{-1}$  belong to the same class if the group contains  $\sigma_v$  or  $S_n$

## 1.2 Generators

Finite groups: Elements whose products generate the whole group. Often not unique.

**example:** generate  $C_{3v}$  from products of  $C_3$  and  $\sigma_v$

# 2 Character Table Completion

1. # IRREPs = # Conjugate Classes
2. First column  $\chi(E)$  using  $\sum_j n_j^2 = |G| = \sum_C n_C$
3. homomorphism:  $g^n = E \iff D(g)^n = I$ 
  - Abelian  $\implies$  1D IRREP  $\implies \chi(g) = \exp(i\frac{2\pi}{n}k)$
  - $g, g^{-1}$  in same class  $\implies \chi(g) = \chi(g)^*$  is real
  - then, if  $n$  odd (e.g.  $C_5$ ) then it is  $+1$
4. Row and column orthogonality

Row orthogonality

$$(\chi^{(i)}, \chi^{(j)}) = \sum_g \chi^{(i)}(g)^* \chi^{(j)}(g) = |G| \delta_{ij}$$

Column orthogonality

$$\sum_i n_C (\chi^{(i)}(C))^* \chi^{(i)}(C') = |G| \delta_{CC'}$$

The basis functions  $|k_i\rangle$  of an IRREP  $k$  are defined as

$$\hat{T}(g) |k_i\rangle = \sum_j D(g)_{ji}^{(k)} |k_j\rangle$$

# 3 Degeneracy and Lifting

A  $n \times n$  IRREP has  $n$  degenerate eigenfunctions  $\psi$  as a basis.

- Perturbation can only decrease symmetry  $G' \subset G$
- new rep: select from  $D(G)$  the matrices  $D(g' \in G')$ 
  - If new rep. is reducible: lift degeneracy; split level
  - guess or use character orthogonality to decompose

**trick:** In questions about lifting degeneracy, use only rotational subgroup  $SO(3)$ . For selection rules, use full  $O(3)$ .

**pitfall:** Example  $\Gamma^2 = A_1 \oplus 2E$ , then the 2  $E$  have different energies in general (each with 2-fold degeneracy).

# 4 Normal Modes in Molecules

Normal modes transform as (form basis of) IRREPs of the symmetry group of the potential

- basis: 3N-component N-atom displacement vectors
- representation matrices  $D^{(3N)} = R \otimes A$ 
  - $R$  exchanges equilibrium positions of N atoms
  - $A(g)$  is action on each atom (e.g. rotation)
- $n_g$ : number of atoms that stay in place after  $g$

In practice, find normal modes IRREPs as follows:

1. Write down table (group elements,  $n_g$ ,  $\varphi$ , (im)proper)

2. Decompose into IRREPs:  $\chi^{(\text{vib})} = \sum_i a_i \chi^{(i)}$

$$a_i = \frac{1}{m} \sum_g \chi^{(i)}(g)^* \chi^{(\text{vib})}(g)$$

3. if IRREP turns up twice ( $a_i \neq 1$ ), cannot deduce normal modes / eigenstates

- normal mode is lin. comb. of projection op. results

Projection operators:

$$\hat{P}_{jk}^{(i)} \propto \sum_g D_{jk}^{(i)}(g)^* \hat{T}(g)$$

## 5 Selection Rules

- $\langle \psi_i | \hat{T}'_\alpha | \psi_j \rangle \neq 0$  only if  $D^{(i)} \otimes D' \otimes D^{(j)}$  contains  $A_{1g}$ 
  - Calculate  $\chi^{(D^{(i)} \otimes D' \otimes D^{(j)})}(g) = \chi^{(i)}(g) \chi'(g) \chi^{(j)}(g)$
  - Scalar product with  $\chi^{(A_1)}$  is 0  $\Rightarrow$  forbidden
- fundamental transition  $|0\rangle \rightarrow |1\rangle$ : ground state transforms as totally symmetric IRREP  $A_{1g}$

**trick (PS10)** use direct product parity rules

$$\text{gerade: } \chi_g \otimes \chi_g = \chi_u \otimes \chi_u = \bigoplus \chi_g$$

$$\text{ungerade: } \chi_g \otimes \chi_u = \bigoplus \chi_u \neq A_{1g}$$

### 5.1 Perturbation of spherically symmetric $\hat{H}$

For degeneracy, only use  $\text{SO}(3)$ , but for transitions, need  $\text{O}(3)$

- Since  $\hat{i}Y_{lm} = (-1)^l Y_{lm}$ , have extra factor  $(-1)^l$

$$\chi^{(l)}(i \otimes R(\varphi)) = (-1)^l \chi^{(l)}(R(\varphi))$$

- Can then calculate characters with

$$S_n(\varphi) = iC_n(\varphi + \pi)$$

e.g.:

$$\chi(\sigma_h) = \chi(i \otimes C_2) = \chi(i \otimes R(\pi))$$

$$\chi(S_4) = \chi(i \otimes C_4^{-1}) = \chi(i \otimes R(\frac{\pi}{2}))$$

$$\chi(S_n) = \chi(i \otimes C_2 \otimes C_n) = \chi(i \otimes R(\pi + \frac{2\pi}{n}))$$

$\text{O}(3)$	$R(\varphi)$	$i = i \otimes E$	$i \otimes R(\varphi)$	basis
$\chi^{(l)}$	$\frac{\sin[(2l+1)\varphi/2]}{\sin(\varphi/2)}$	$(-1)^l(2l+1)$	$(-1)^l \chi^{(l)}(R(\varphi))$	$Y_{lm}$

- perturbation example: Octahedral arrangement
  - new  $\hat{H}$  has symmetry of rotational subgroup of  $O_h$
  - calculate characters  $\chi^{(l)}(O_h)$  via table
  - decompose into IRREPs of  $O_h$ 
    - \*  $l = 2$  level is split into  $\chi^{(2)} = E \oplus T_2$
  - selection rules between split states
    - \* use symmetry (g / u) arguments

## 6 Selection Rules and Wigner Eckart

Perturbation  $V = V_0 \underbrace{xy \dots}_{\omega \text{ factors}}$

Selection rules:

Have matrix elements  $\langle j'm' | V | jm \rangle$

1. write  $V \propto \sum_\mu Y_{\omega,\mu}$  using spherical coordinates
2. Clebsch-Gordan series

$$D^{(\omega)} D^{(j)} = \bigoplus_{L=|\omega-j|}^{\omega+j} D^{(L)} = D^{|\omega-j|} \oplus \dots \oplus D^{\omega+j} \Rightarrow |j' - j| \leq \omega$$

3. selection rules for  $\langle j'm' | Y_\mu^{(\omega)} | jm \rangle$  are

$$j' = j + \omega, \dots, |j - \omega|$$

$$\text{C-G coeff.: } m' = \mu + m \\ (\text{left}) = (\text{center}) + (\text{right})$$

4. finally, using parity  $(-1)^j$ ,

$$\mathbf{V \text{ even: }} 1 = (-1)^{j+j'}, \text{ so } \Delta j = 0, 2, \dots$$

$$\mathbf{V \text{ odd: }} 1 = (-1)^{j+j'+1}, \text{ so } \Delta j = 1, 3, \dots$$

Wigner-Eckert theorem:

$$\langle N'j'm' | \hat{T}_\mu^{(\omega)} | Njm \rangle = C_{m\mu m'}^{j\omega j'} \underbrace{\langle N'j' | \hat{T}^{(\omega)} | Nj \rangle}_{\text{reduced matrix element}}$$

where

$$\hat{T}_\mu^{(\omega)} = r^\omega Y_{\omega,\mu}$$

1. Find  $\langle 0, 0 | T^\omega | N, \omega \rangle \propto C^{-1} r_{0N}^\omega$
2. then other matrix elements are written down as (e.g.)

$$\langle 0, 0, 0 | \underbrace{r^2 Y_{2,\mp 2}}_{T^{(2)\mp 2}} | N, 2, \pm 2 \rangle = C_{0\pm 2\mp 2}^{022} \langle 0, 0 | T^{(2)} | N, 2 \rangle$$

### 6.1 Spherical Harmonics

- First spherical harmonic:  $Y_{0,0} = \frac{1}{\sqrt{4\pi}}$
- symmetry under inversion:  $(-1)^l$

	$C_n$	$\sigma_v$	$\sigma_h$
$\varphi$	$\varphi + \frac{2\pi}{n}$	$-\varphi$	$\pi - \theta$
$\theta$			

Table 1: Action of symmetry operations