Complex Numbers in the UNNS Substrate: Derived Convenience or Foundational Necessity?

Abstract

UNNS (Unbounded Nested Number Sequences) models mathematics and physics via recursive operators acting on nested numeric states. Do such systems require complex numbers, or can complex behavior be reconstructed from integer—algebraic lattices and cyclotomic embeddings? We argue that complex numbers are not foundational for UNNS: oscillations, phases, and spectra can be realized within algebraic integer rings and their cyclotomic fields. Nevertheless, $\mathbb C$ offers an efficient shorthand for resonance and spectral geometry. We formalize embeddings, projection operators (inlaying), bounds for Gaussian and Eisenstein lattices, and show when $\mathbb C$ is convenient, when it is optional, and how it interfaces with UNNS thermodynamics, gauge phases, and discrete field theories.

1 UNNS Structures and Operators (Brief)

A UNNS structure is a quadruple $\mathcal{S} = (A, \mathcal{O}, \mathcal{N}, \mathcal{R})$ with seeds A, operators \mathcal{O} (Collapse, Inlaying, Inletting, Normalize, Evaluate, Adopt), nesting depth \mathcal{N} (recursion index), and a resonance/stability map \mathcal{R} . Time corresponds to recursion depth; space to lattice embedding.

2 Complex Numbers as Lattice Geometry

2.1 Integer lattices and algebraic lattices

Definition 1 (Algebraic lattices). The Gaussian integers $\mathbb{Z}[i]$ form a square lattice in \mathbb{R}^2 and the Eisenstein integers $\mathbb{Z}[\omega]$, $\omega = e^{2\pi i/3}$ form a hexagonal lattice. More generally, rings of integers \mathcal{O}_K in cyclotomic fields $K = \mathbb{Q}(\zeta_n)$ embed as full-rank lattices in $\mathbb{R}^{\phi(n)}$.

Definition 2 (Inlaying (projection)). Given $x \in \mathbb{R}^2$, the Gaussian inlaying $G : \mathbb{R}^2 \to \mathbb{Z}[i]$ maps x to its nearest lattice point in $\mathbb{Z}[i]$ (componentwise rounding). The Eisenstein inlaying $E : \mathbb{R}^2 \to \mathbb{Z}[\omega]$ maps x to the nearest hexagonal lattice point.

Proposition 1 (Projection error bounds). For Gaussian inlaying G, $||x - G(x)||_{\infty} \leq \frac{1}{2}$ for all $x \in \mathbb{R}^2$. For Eisenstein inlaying E, $||x - E(x)||_2 \leq r$, where r is the circumradius of the hexagonal Voronoi cell (a fixed numerical constant). Thus inlaying perturbs values by a uniformly bounded amount.

2.2 Embedding chain and operator view

Remark 1. Complex numbers appear here as coordinates on algebraic lattices. They serve as a convenient parameterization, not as an axiom of the substrate.

3 Cyclotomic Phases and UNNS Oscillations

3.1 Roots of unity as discrete rotations

Definition 3 (Cyclotomic inlaying). For $n \geq 1$, let $\mu_n = \{\zeta_n^k : 0 \leq k < n\}$ be the n-th roots of unity. Define the cyclotomic snap $\Pi_n : \mathbb{C} \to \mu_n$ by projecting any $z \neq 0$ to the nearest ζ_n^k by angle.

Lemma 1 (Finite-phase capture). Any finite-phase recursion (periodic phase increments) can be represented within a cyclotomic field $\mathbb{Q}(\zeta_n)$ for suitable n; its values lie in the ring of integers $\mathbb{Z}[\zeta_n]$ up to bounded projection error.

Proof (sketch). Phase increments rational in 2π land in a cyclotomic subgroup; values are polynomials in ζ_n with integer coefficients (after clearing denominators), hence lie in $\mathbb{Z}[\zeta_n]$ up to scaling. Projection errors are uniformly bounded.

3.2 Oscillations without invoking \mathbb{C} as primitive

Proposition 2 (2D real recursion for complex eigenpairs). Let $a_{k+1} = \alpha a_k - \beta a_{k-1}$ with $\alpha, \beta \in \mathbb{Q}$ and $\alpha^2 - 4\beta < 0$. Then the complex eigenvalues $\lambda_{\pm} = \rho e^{\pm i\theta}$ give rise to a real 2×2 rotation-dilation:

$$\begin{bmatrix} a_{k+1} \\ a_k \end{bmatrix} = \underbrace{\begin{bmatrix} \alpha & -\beta \\ 1 & 0 \end{bmatrix}}_{R(\rho,\theta)} \begin{bmatrix} a_k \\ a_{k-1} \end{bmatrix}, \quad R = \rho \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix} \quad in \ a \ suitable \ basis.$$

Thus oscillations are realizable by real UNNS dynamics; \mathbb{C} is a compact encoding.

4 UNNS Projection Pipelines and Spectral Stability

Let T be a linear recursion on \mathbb{R}^d with spectral radius $\rho(T) > 1$. Consider a UNNS pipeline with inlaying and collapse:

$$x_{k+1} = \mathcal{T}(x_k) := \underbrace{\prod (C_{\varepsilon}(Tx_k))}_{\text{project \& repair}},$$

where Π is a fixed lattice projection (e.g. G or E), and C_{ε} zeros strict subthreshold values.

Lemma 2 (Uniform perturbation). There exists a norm and constant B > 0 such that $\|\mathcal{T}(x) - Tx\| \leq B$ for all x. In particular, for $k \geq 0$, $x_k = T^k x_0 + \sum_{j=0}^{k-1} T^{k-1-j} \Delta_j$ with $\|\Delta_j\| \leq B$.

Theorem 1 (Asymptotic spectral stability). If T is diagonalizable with simple dominant eigenvalue λ_{\star} and $|\lambda_{\star}| > \max_{i \neq \star} |\lambda_{i}|$, then the UNNS pipeline satisfies

$$\lim_{k \to \infty} \frac{\|x_{k+1}\|}{\|x_k\|} = |\lambda_{\star}|, \qquad \frac{x_k}{\|x_k\|} \to v_{\star}/\|v_{\star}\|,$$

i.e. growth factor and projective direction of the dominant eigenpair are preserved.

Proof (outline). Standard perturbation: the additive O(1) error is negligible compared to $||T^kx_0|| \approx |\lambda_{\star}|^k$; projective convergence follows from spectral dominance.

5 Gauge Phases and Discrete Connections in UNNS

5.1 Phase inletting as U(1) connection

Definition 4 (Phase inletting). Let $U_{\theta} = \operatorname{diag}(e^{i\theta_1}, \dots, e^{i\theta_m})$ act on a block decomposition of $x \in \mathbb{C}^m$ (or its real lift). The phase inletting step is $x \mapsto U_{\theta}x$ before projection Π .

Proposition 3 (Spectral invariance under constant phase). If U_{θ} is constant across steps, then the spectral radius of the effective transfer is unchanged: $\rho(U_{\theta}T) = \rho(T)$; projective limits are rotated but stable.

Remark 2. Slowly varying phases (adiabatic inletting) modulate transients but preserve the dominant growth rate; fast variations imprint quasi-random residues that are absorbed by collapse/repair unless they resonate with subdominant modes.

6 When \mathbb{C} is Useful (but Optional)

6.1 Spectral geometry and concise notation

- Complex exponentials compactly represent rotations: $e^{i\theta}$.
- Eigenvalues naturally live in \mathbb{C} ; describing spectra and resonance is simplest in \mathbb{C} .
- Fourier/character sums (Gauss, Jacobi, Eisenstein) enter as arithmetic constants on UNNS lattices; C packages their phases.

6.2 Hilbert embeddings and measurement

Definition 5 (UNNS Hilbert space). Associate to a UNNS structure S the Hilbert space $\mathcal{H}_{S} = \operatorname{span}_{\mathbb{C}}\{|a,n\rangle : a \in A, n \in \mathbb{N}\}, \text{ with } \langle a,n|a',n'\rangle = \delta_{a,a'}\delta_{n,n'}. \text{ Operators in } \mathcal{O} \text{ lift to linear maps } \widehat{O} \text{ on } \mathcal{H}_{S}.$

Proposition 4. If O is norm-preserving under \mathcal{R} , then \widehat{O} is (approximately) unitary; collapse/repair maps are contractions. Spectra of \widehat{O} encode UNNS resonance constants.

7 Worked Examples

7.1 Fibonacci under Gaussian inlaying

Let $A = \begin{bmatrix} 1 & 1 \\ 1 & 0 \end{bmatrix}$, $v_{k+1} = G(C_{\varepsilon}(Av_k))$. Then $|\lambda_{\star}| = \varphi = (1 + \sqrt{5})/2$ governs growth, and the projective direction approaches the φ -eigenline. Rounding error per step is bounded by $\frac{1}{2}$ in $\|\cdot\|_{\infty}$.

7.2 Complex eigenpair without \mathbb{C}

Let $R = \rho \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$, $x_{k+1} = \Pi(C_{\varepsilon}(Rx_k))$. Then $\lim ||x_{k+1}|| / ||x_k|| = \rho$, and the rotation angle θ is visible in the phase of the 2D recursion even if we never mention \mathbb{C} .

8 Do We *Need* Complex Numbers?

Theorem 2 (Complex as a derived envelope). Let $C = \{\mathbb{Q}(\zeta_n) : n \geq 1\}$ be the filtered system of cyclotomic fields, with embeddings $\mathbb{Q}(\zeta_m) \hookrightarrow \mathbb{Q}(\zeta_{mn})$. For any finite-phase UNNS process with rational coefficients and bounded inlaying error, there exists n such that the process is representable inside $\mathbb{Q}(\zeta_n)$ (up to bounded projection). Moreover, the collection of all such representations is dense on the unit circle in the angular variable.

Proof (sketch). Finite-phase increments land in a finite subgroup of \mathbb{T} , hence in μ_n for some n. Values are algebraic over \mathbb{Q} ; after clearing denominators we obtain elements of $\mathbb{Z}[\zeta_n]$. Density follows from $\bigcup_n \mu_n$ being dense in \mathbb{T} in the angular sense for rational approximations. \square

Corollary 1. For UNNS pipelines built from algebraic operators and cyclotomic inlayings, \mathbb{C} is not foundational but serves as a convenient completion capturing limits and spectra succinctly.

9 Diagrams

Embedding chain and spiral phases



cyclotomic phases on unit circle

10 Interfaces with UNNS Physics

10.1 Time-harmonic fields and DEC/FEEC

Time-harmonic fields $e^{i\omega t}$ are conveniently expressed in \mathbb{C} . In UNNS, ω -phases can be approximated cyclotomically; FEEC/DEC estimates carry over since inlaying errors are bounded and do not affect stability constants at leading order, while complex notation remains a concise calculus for waves.

10.2 Thermodynamics and entropy of phases

Phase dispersion under inlaying contributes to effective *entropy* in UNNS thermodynamics; collapse acts as dissipation, while constant phase inletting is entropy-neutral (spectrally inert).

11 Conclusions

- Complex numbers are *not* foundational for UNNS. Oscillations, phases, and spectra arise from algebraic integer lattices and cyclotomic inlayings.
- C is an *efficient shorthand*: it packages rotations, spectra, and wave phenomena succinctly, and is invaluable for analysis and exposition.
- In practice, one may implement UNNS pipelines purely over $\mathbb{Z}[i]$, $\mathbb{Z}[\omega]$, or $\mathbb{Z}[\zeta_n]$, and only lift to \mathbb{C} for spectral summaries, Fourier views, or Hilbert embeddings.