



Direct CP violation in the decay $B^+ \rightarrow K^+ \pi^0$ at LHCb

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on behalf of LHCb Collaboration

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[Phys. Rev. Lett. 126, 091802](#)

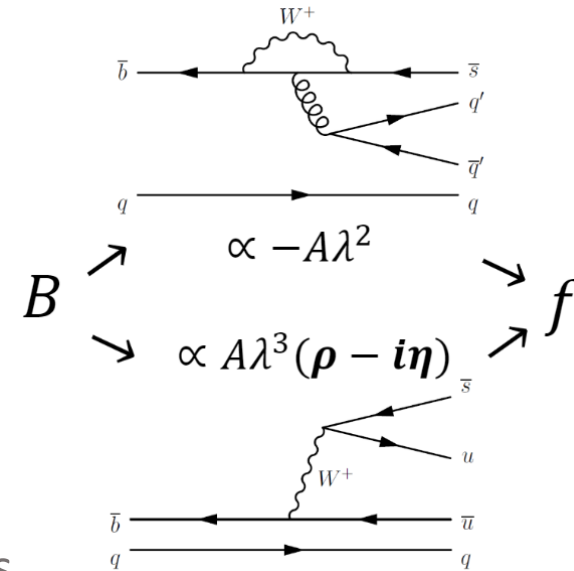
CP Violation in B System

- The breaking of charge-parity symmetry arises from the interference between two contributing amplitudes with *different strong and weak phases*

- Direct CP violation: interference in decay amplitudes

$$A_{CP} = \frac{\Gamma(\bar{B}^0 \rightarrow \bar{f}) - \Gamma(B^0 \rightarrow f)}{\Gamma(\bar{B}^0 \rightarrow \bar{f}) + \Gamma(B^0 \rightarrow f)}.$$

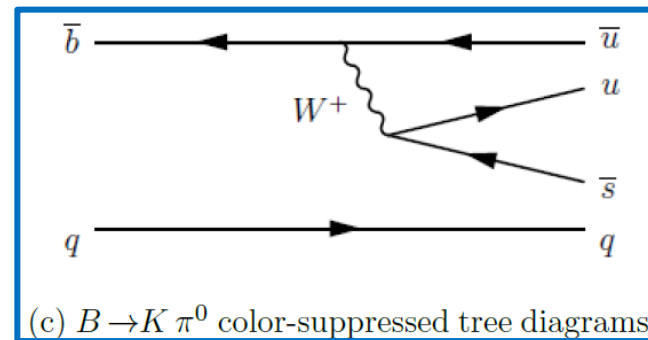
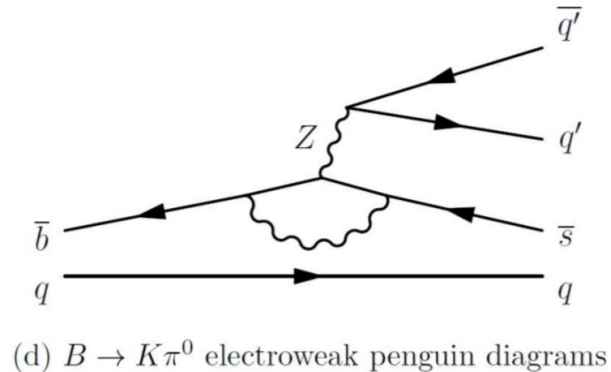
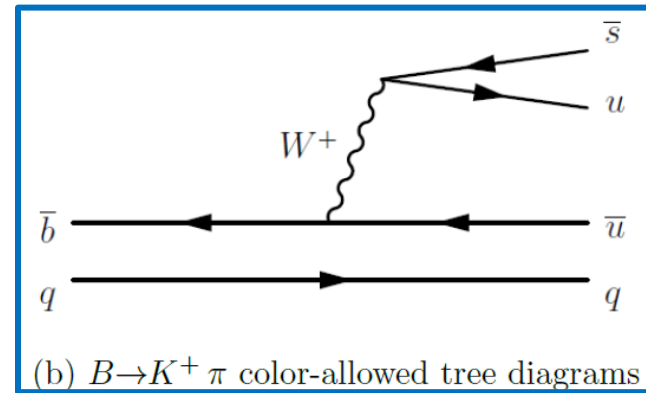
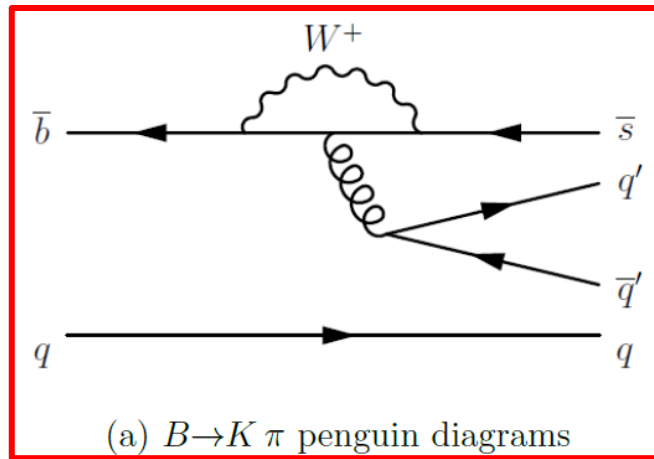
- CP violation in mixing: interference in mixing amplitudes
- Indirect CP violation: interference between mixing and decay amplitudes



- Because CKM matrix is the only CPV source in SM, any deviation could indicate new physics

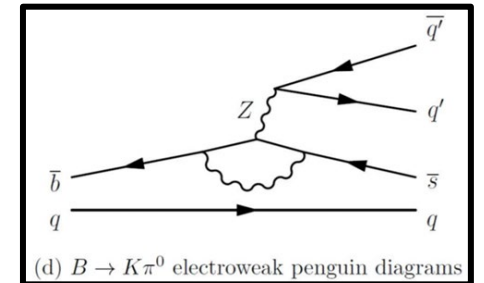
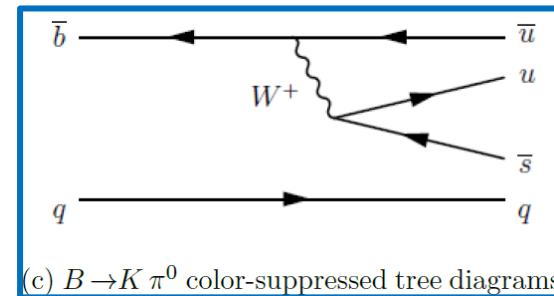
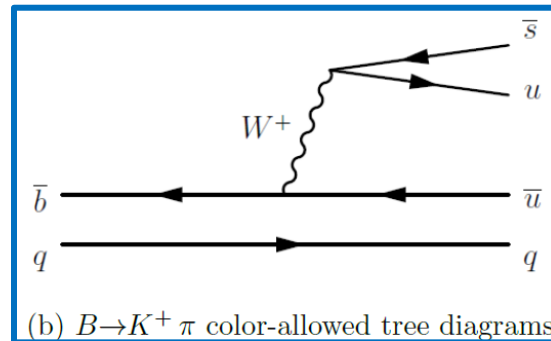
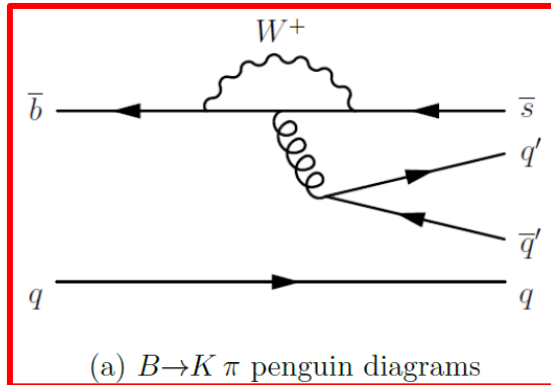
The $B \rightarrow K \pi$ System

- There are four $B \rightarrow K \pi$ decays: $B^0 \rightarrow K^+ \pi^-$, $B^+ \rightarrow K^+ \pi^0$, $B^0 \rightarrow K^0 \pi^0$, $B^+ \rightarrow K^0 \pi^+$
- Particularly interesting because:
 - Tree amplitudes are suppressed by the CKM matrix element $V_{ub} = A\lambda^3(\rho - i\eta)$.
 - Leading contribution is from QCD penguin amplitudes



The $B \rightarrow K \pi$ System

- $B^0 \rightarrow K^0 \pi^0$ and $B^+ \rightarrow K^0 \pi^+$
 - dominated by QCD penguin amplitude
- $B^+ \rightarrow K^+ \pi^0$ and $B^0 \rightarrow K^+ \pi^-$
 - leading contribution is QCD penguin, non-negligible tree amplitudes



- interesting to see if the pattern of decays follow SM predictions -- extensive studies done at B factories and now LHCb

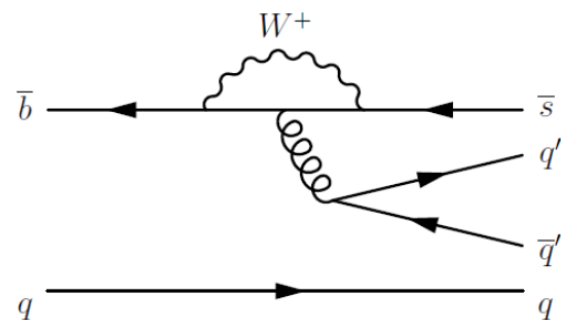
Origin of the $B \rightarrow K \pi$ Puzzle

- The current status of A_{CP} measurements:

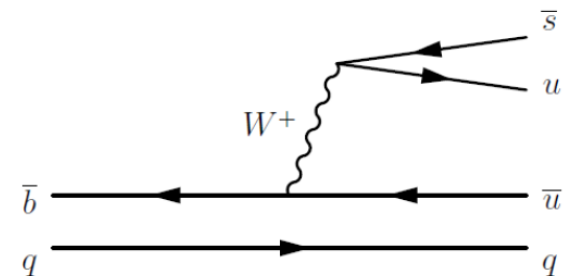
decay	<i>BaBar</i>	<i>Belle</i>	<i>LHCb</i>
$B^0 \rightarrow K^+ \pi^-$	$-0.107 \pm 0.016^{+0.006}_{-0.004}$	$-0.069 \pm 0.014 \pm 0.007$	$-0.080 \pm 0.007 \pm 0.003$
$B^+ \rightarrow K^+ \pi^0$	$+0.030 \pm 0.039 \pm 0.010$	$+0.043 \pm 0.024 \pm 0.002$	n/a
$B^0 \rightarrow K^0 \pi^0$	$-0.13 \pm 0.13 \pm 0.03$	$+0.14 \pm 0.13 \pm 0.06$	n/a
$B^+ \rightarrow K^0 \pi^+$	$-0.029 \pm 0.039 \pm 0.010$	$+0.043 \pm 0.024 \pm 0.002$	$-0.022 \pm 0.025 \pm 0.010$

- $B^+ \rightarrow K^+ \pi^0$ and $B^0 \rightarrow K^+ \pi^-$
 - Both penguin and tree diagrams differ only in the spectator quark (u or d)

$$\Rightarrow A_{CP}^{K^+ \pi^-} \sim A_{CP}^{K^+ \pi^0}$$



(a) $B \rightarrow K \pi$ penguin diagrams



(b) $B \rightarrow K^+ \pi$ color-allowed tree diagrams

Origin of the $B \rightarrow K \pi$ Puzzle

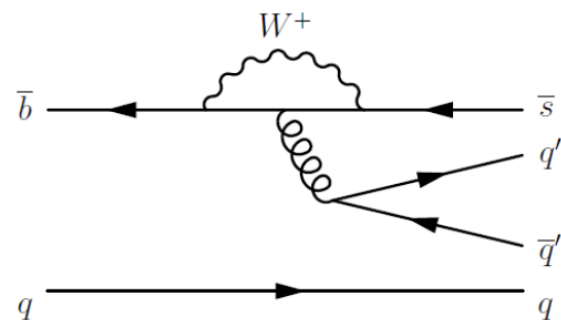
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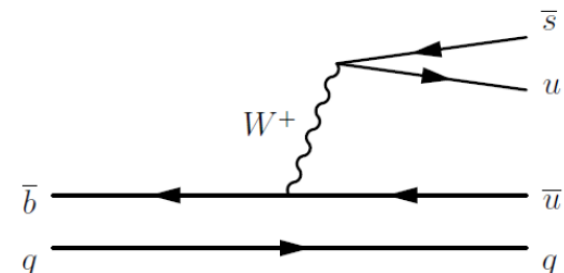
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Not as expected
(disagreement above 5σ)



(a) $B \rightarrow K \pi$ penguin diagrams



(b) $B \rightarrow K^+ \pi$ color-allowed tree diagrams

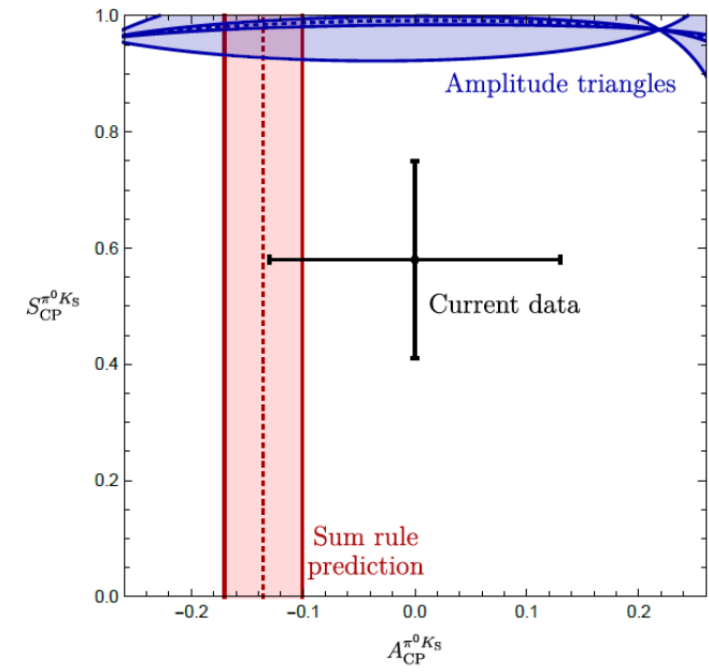
Broader Picture of the $B \rightarrow K \pi$ Puzzle

- Guided by isospin symmetry, a [sum rule](#) was derived:

$$\left[A_{CP}^{K^0 \pi^+} \frac{\mathcal{B}^{K^0 \pi^+}}{\mathcal{B}^{K^+ \pi^-}} - A_{CP}^{K^+ \pi^0} \frac{2\mathcal{B}^{K^+ \pi^0}}{\mathcal{B}^{K^+ \pi^-}} \right] \frac{\tau_{B^0}}{\tau_{B^\pm}} + \left[A_{CP}^{K^+ \pi^-} - A_{CP}^{K^0 \pi^0} \frac{2\mathcal{B}^{K^0 \pi^0}}{\mathcal{B}^{K^+ \pi^-}} \right] = 0$$

[arXiv:0508.047](#)
[arXiv:0608.040](#)
[arXiv:1712.02323](#)

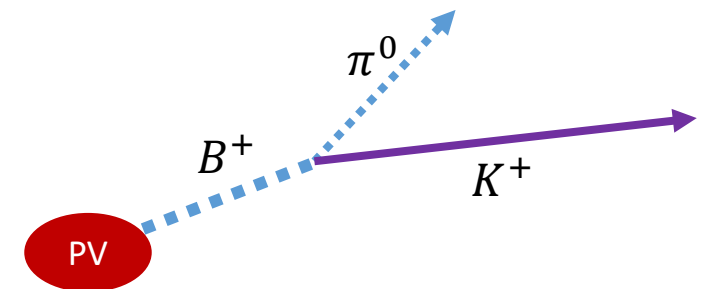
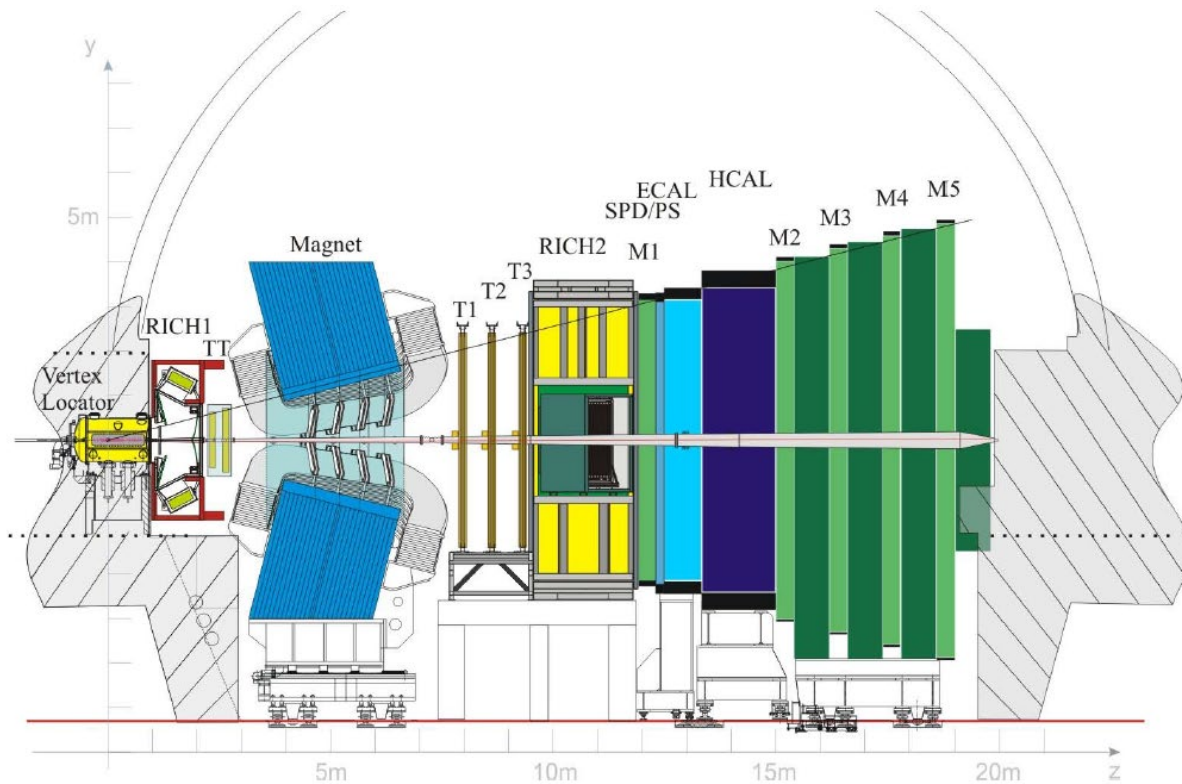
- $B^0 \rightarrow K^0 \pi^0$ is least precisely measured among all decay modes:
 - Sum rule predicts $A_{CP}(K_S^0 \pi^0) = -0.150 \pm 0.032$
 - Measured $A_{CP}(K_S^0 \pi^0) = 0.01 \pm 0.10$
- Global study of $B \rightarrow K \pi$ system indicates need for significant enhancement of color-suppressed tree or EW penguin [arXiv:1812.02672](#)
- $B \rightarrow K \pi$ system need to be measured more precisely (room to probe for new physics)



[arXiv:1812.02672](#)

$B^+ \rightarrow K^+ \pi^0$ at LHCb

- LHCb is designed to study CP violation and search for new physics in the heavy flavor sector
- Challenging mode with no secondary vertex (1 track + 1 calorimeter cluster)
 - Run1 software triggers relied on secondary vertex
 - Dedicated trigger developed for Run2

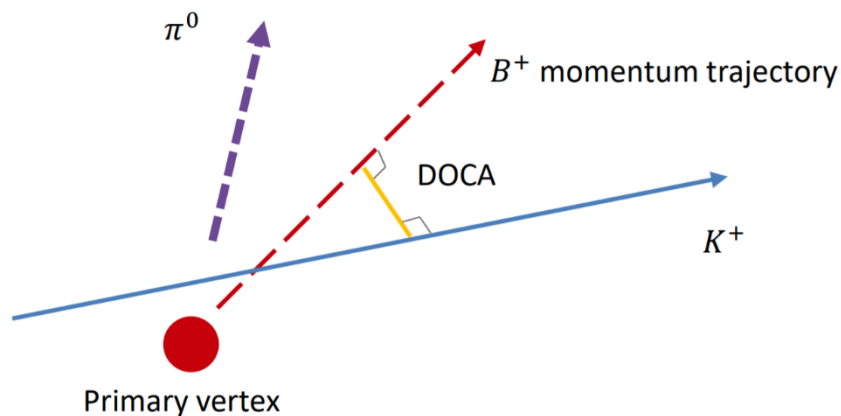


Key Variables

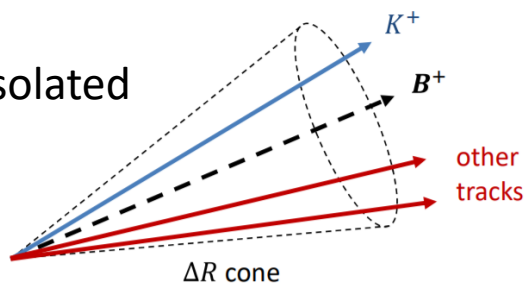
- No secondary vertex, but still take advantage of LHCb's precision tracking
 - initial $S/B \sim 10^{-7}$

- K^+ should be originating from the trajectory of B
 - Cut on distance of closest approach

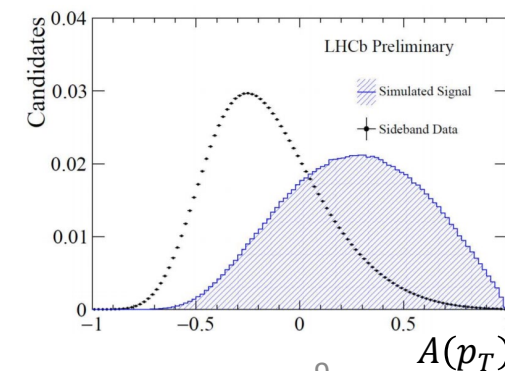
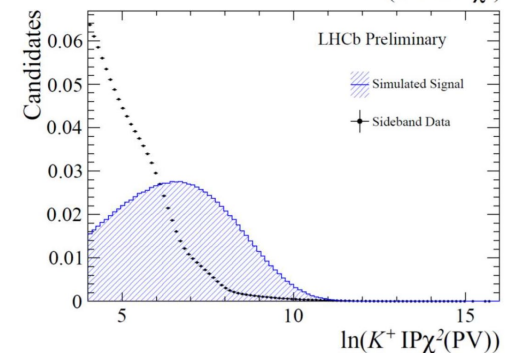
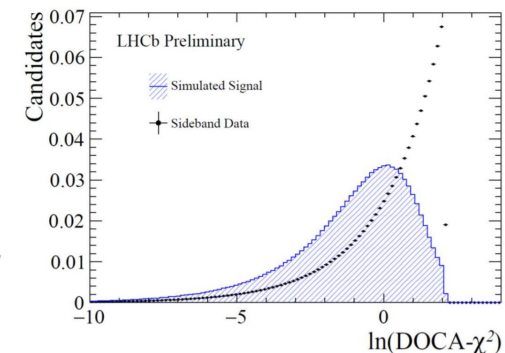
- K^+ should be inconsistent with originating from any primary vertex
 - Cut on impact parameter w.r.t. PV, when K^+ is included in primary vertex fit



- B candidate should be more isolated from other high- p_T tracks:



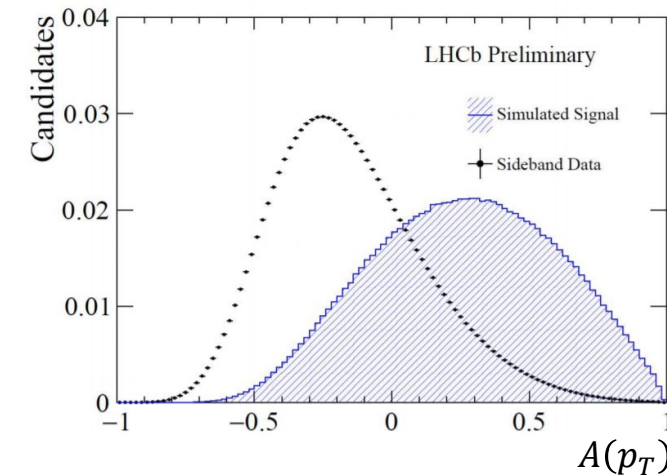
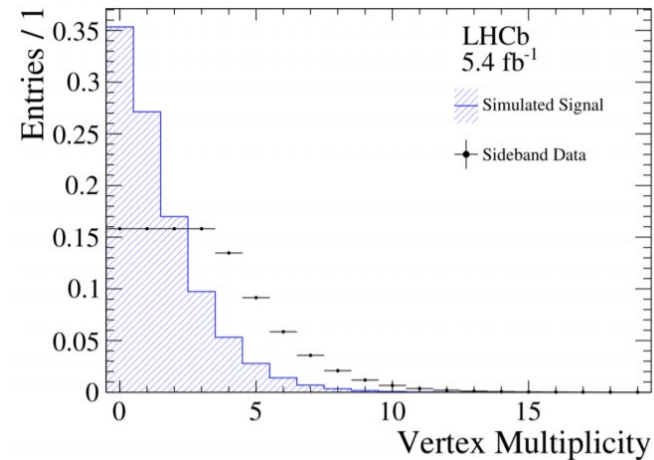
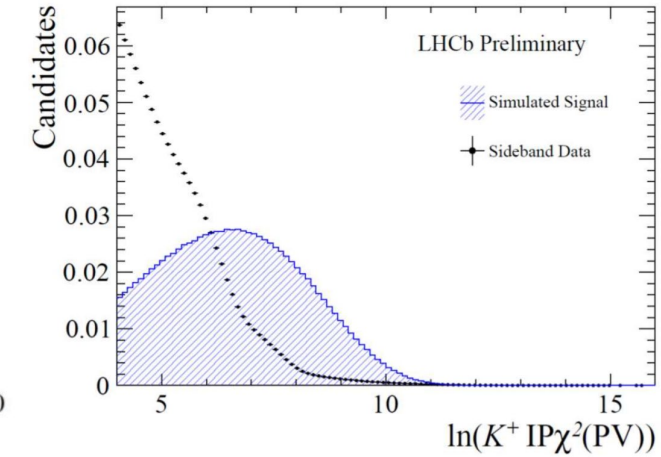
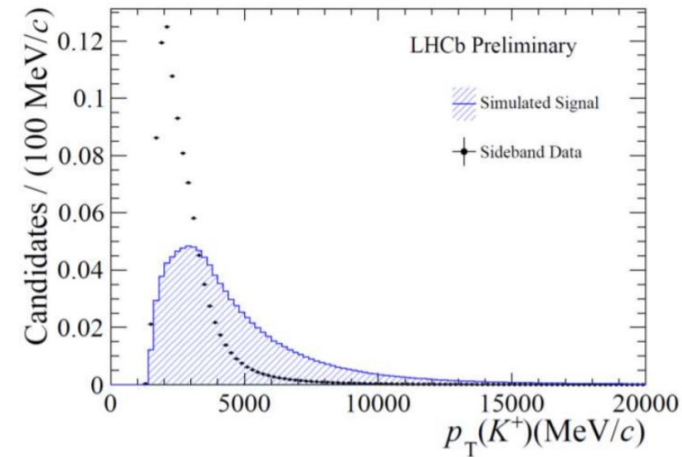
$$A(p_T) \equiv \frac{p_T(B) - p_T(\text{cone})}{p_T(B) + p_T(\text{cone})}$$



Multivariate Analysis

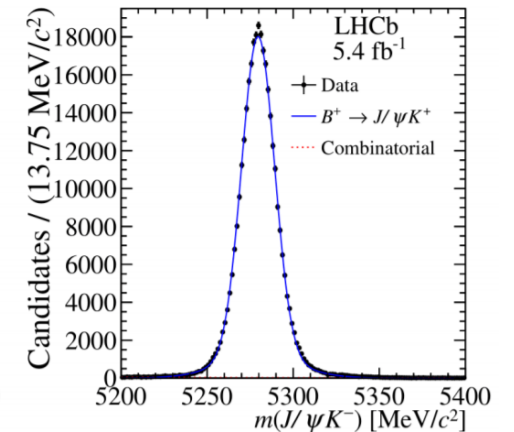
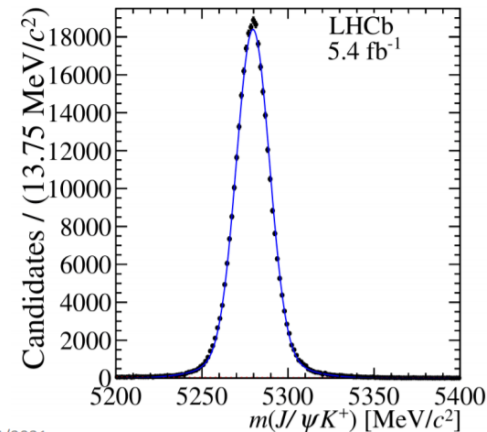
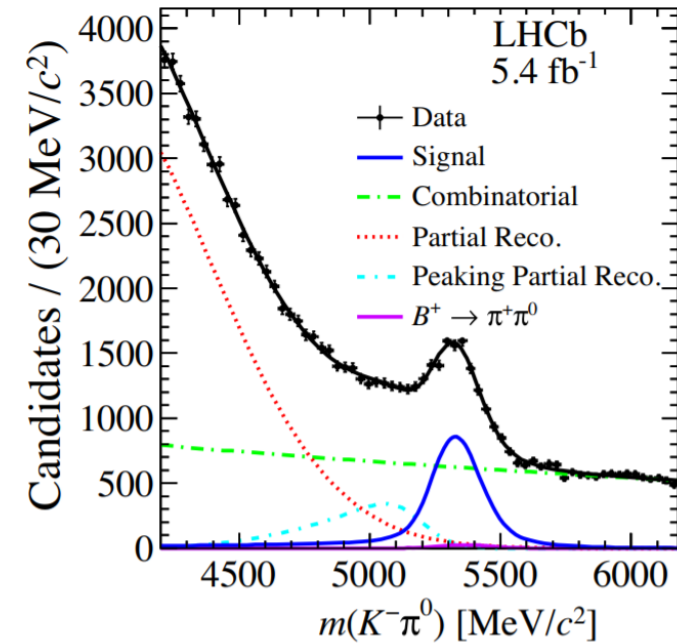
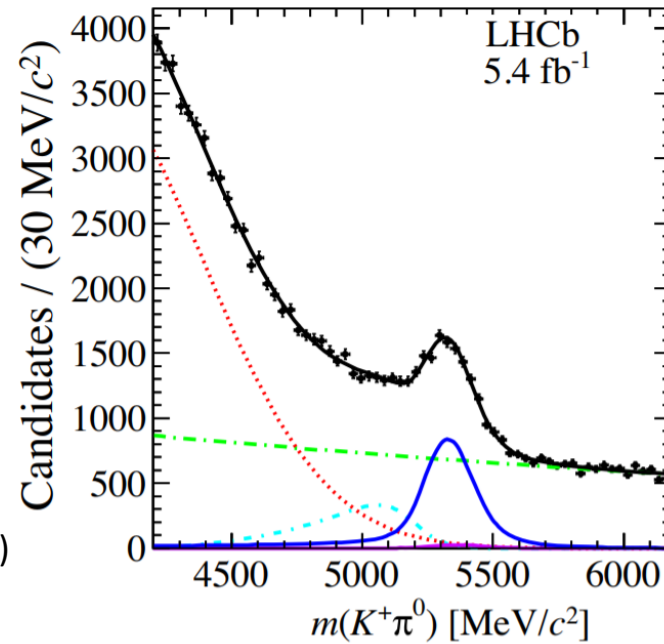
- Using discriminant variables including
 - impact parameters
 - kinematics
 - isolation variables
- Train Boosted Decision Trees to optimize

$$\frac{\epsilon_{MC}}{\sqrt{S+B}}$$



Invariant Mass Fit

- Two more background categories:
 - $B^+ \rightarrow \pi^+ \pi^0$ where π^+ is misidentified as K
 - Peaking partial reconstructed, e.g.
 - $B^{+/-} \rightarrow (K^{*+/-} \rightarrow K^+ \pi^{0/-}) \pi^0$
- Total $N(B^\pm \rightarrow K^\pm \pi^0) \approx 16500$
- Raw $A_{cp}(B^+ \rightarrow K^+ \pi^0) = 0.005 \pm 0.022$ (Magnet Up),
 0.019 ± 0.021 (Magnet Down)
- Apply corrections due to B production asymmetry and K detection asymmetry
 - directly measured using $B^\pm \rightarrow J/\psi K^\pm$ data
- Raw $A_{cp}(B^+ \rightarrow J/\psi K^+) = -0.009 \pm 0.002$ (Magnet Up),
 -0.012 ± 0.002 (Magnet Down)



Invariant Mass Fit

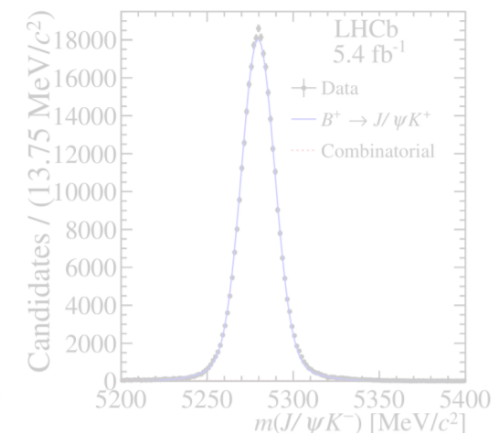
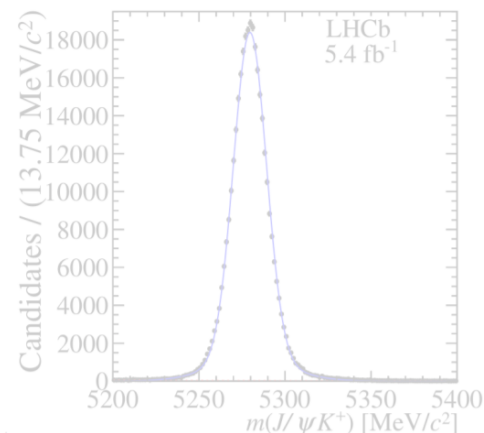
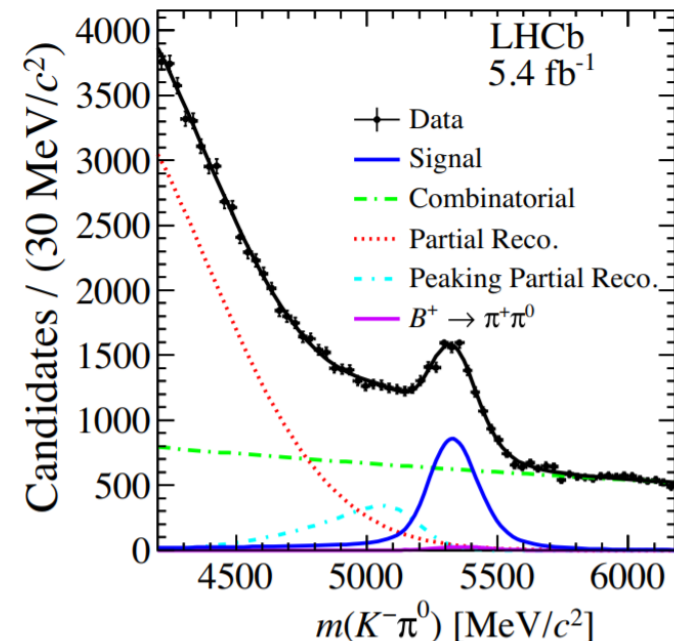
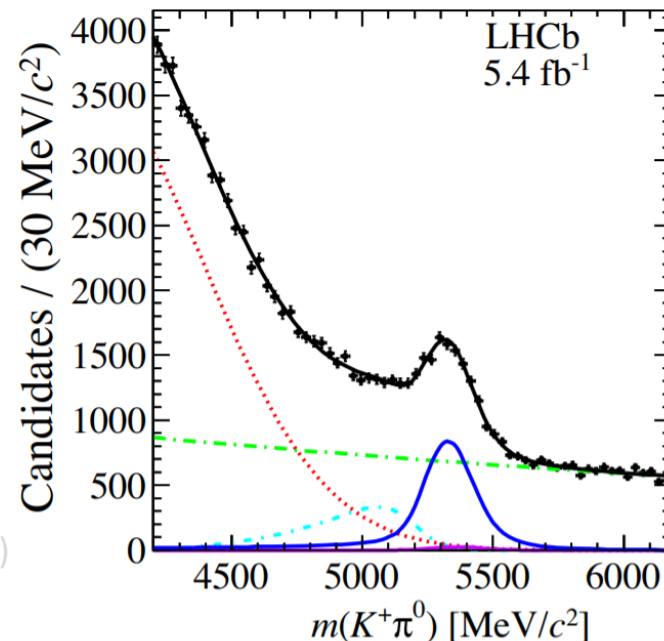
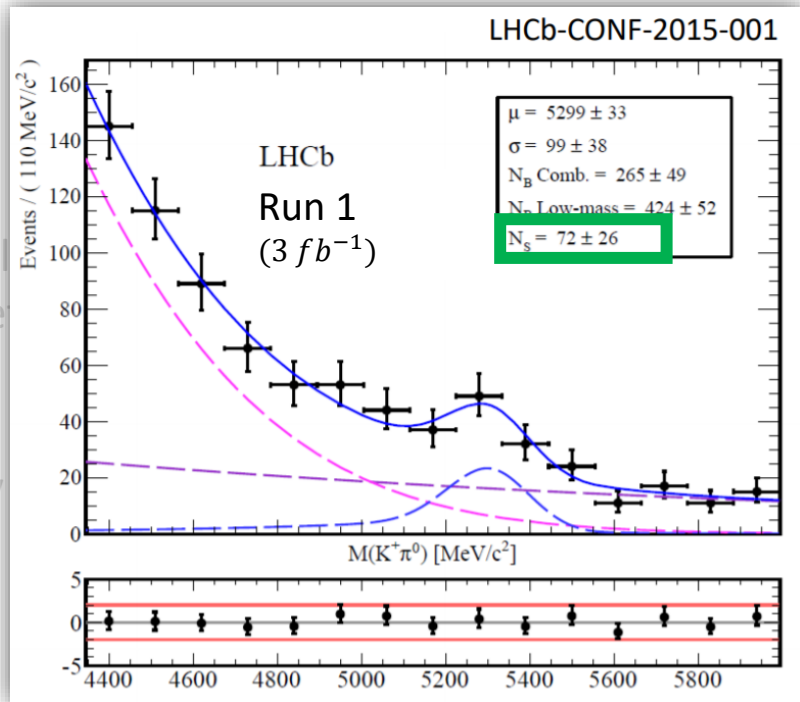
- Two more background categories:
 - $B^+ \rightarrow \pi^+ \pi^0$ where π^+ is misidentified as K
 - Peaking partial reconstructed, e.g.
 - $B^{+/-} \rightarrow (K^{*+/-} \rightarrow K^+ \pi^{0/-}) \pi^0$

• Total $N(B^\pm \rightarrow K^\pm \pi^0) \approx \boxed{16500}$

• Raw $A_{cp}(B^+ \rightarrow K^+ \pi^0) = 0.005 \pm 0.022$ (Magnet Up),

• Appl
K de

• Raw



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Results

- After corrections, the measurement result is

$$A_{CP}(B^+ \rightarrow K^+ \pi^0) = 0.025 \pm 0.015_{stat} \pm 0.006_{syst} \pm 0.003_{ext}$$

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$B^+ \rightarrow K^0 \pi^+$	$-0.029 \pm 0.039 \pm 0.010$	$+0.043 \pm 0.024 \pm 0.002$	$-0.022 \pm 0.025 \pm 0.010$

- Consistent with current world average, *the most precise to date*
- $A_{CP}(B^+ \rightarrow K^+ \pi^0) - A_{CP}(B^0 \rightarrow K^+ \pi^-) = 0.115 \pm 0.014$, non-zero at 8.2σ
 - previously 0.124 ± 0.021 , 5.9σ
- Updated sum rule prediction for $A_{CP}(K^0 \pi^0) = -0.138 \pm 0.025$, non-zero at 5.5σ

Conclusion

- We have measured the direct CP violation in $B^+ \rightarrow K^+ \pi^0$,

$$A_{CP}(B^+ \rightarrow K^+ \pi^0) = 0.025 \pm 0.015_{stat} \pm 0.006_{syst} \pm 0.003_{ext}$$

- Most precise measurement to date

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- Confirms and strengthens the $B \rightarrow K \pi$ puzzle
- PRL paper recently published: [Phys. Rev. Lett. 126, 091802](#)
- After Run3, LHCb would finish collecting $50 fb^{-1}$: projected statistical uncertainty $\sim \pm 0.005$
- Highlights LHCb's potential in modes with neutral particles ($B^0 \rightarrow K^0 \pi^0$ and beyond)

Backup slides

The CKM Mechanism

- In the Standard Model (SM), the Cabibbo-Kobayashi-Maskawa (CKM) mechanism describes the weak interactions between the three generations of quarks:
 - one universal weak coupling strength for quarks and leptons, but
 - the quarks' **weak eigenstates** differ from their **mass eigenstates**,
 - related by the 3x3 unitary CKM matrix V_{CKM} ,

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

The CKM Mechanism

- In the Wolfenstein parametrization, the CKM matrix is written in terms of four real parameters, λ , A , ρ , and η :

$$V_{\text{CKM}} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

- $\lambda = \sin \theta_C \approx 0.225$, where θ_C is the Cabibbo angle
- η is associated with the *single irreducible phase* in the CKM matrix

The *only* known source
of CP violation in the SM

CP Violation

- Three possible CP-violating effects:
 - CP violation in decay (direct CPV)
 - CP violation in mixing
 - CP violation in the interference between mixing and decay
- can be measured in the decay rate difference between a process and its charge conjugate

e.g. $|A(\bar{B}^0 \rightarrow \bar{f})| \neq |A(B^0 \rightarrow f)|.$

$$\begin{aligned} M &= |M_1| e^{i\theta_1} e^{i\phi_1} + |M_2| e^{i\theta_2} e^{i\phi_2} \\ \bar{M} &= |M_1| e^{i\theta_1} e^{-i\phi_1} + |M_2| e^{i\theta_2} e^{-i\phi_2} \end{aligned}$$
$$|M|^2 - |\bar{M}|^2 = -4|M_1||M_2|\sin(\theta_1 - \theta_2)\sin(\phi_1 - \phi_2)$$

CP Violation

- Three possible CP-violating effect:
 - CP violation in decay
 - CP violation in mixing
 - CP violation in the interference between mixing and decay
- seen in the transition rate difference between a neutral meson **to** and **from** its antiparticle,

$$\text{e.g. } |A(B^0 \rightarrow \bar{B}^0)| \neq |A(\bar{B}^0 \rightarrow B^0)|.$$

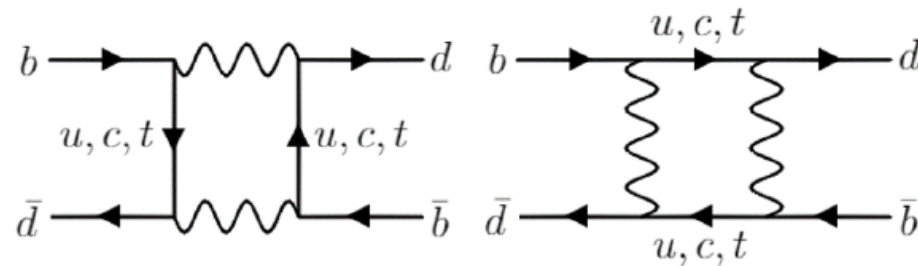


Figure 1: One-loop box diagrams for B^0 - \bar{B}^0 mixing

CP Violation

- Three possible CP-violating effect:
 - CP violation in decay
 - CP violation in mixing
 - CP violation in the interference between mixing and decay
- arises from the interference between decays **with** and **without** mixing.

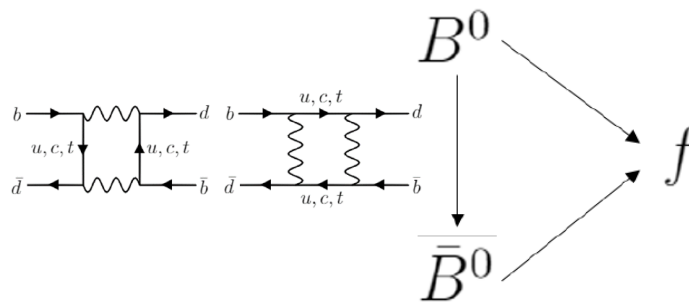


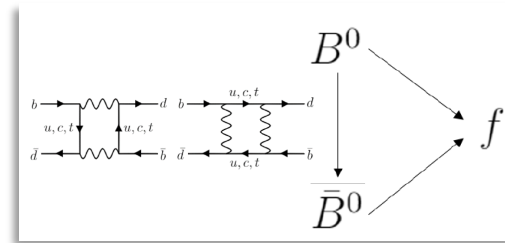
Figure 2: Interference between $B^0 \rightarrow \bar{B}^0 \rightarrow f$ and $B^0 \rightarrow f$

The Broader Picture of the $B \rightarrow K \pi$ Puzzle

- Uniquely, $B^0 \rightarrow K^0 \pi^0$ allows CP violation in the interference between mixing and decay.
- The observable $S_{CP}^{K_S^0 \pi^0}$ can be measured in the time-dependent decay rate asymmetry,

$$\frac{\Gamma(\bar{B}^0(t) \rightarrow K_S^0 \pi^0) - \Gamma(B^0(t) \rightarrow K_S^0 \pi^0)}{\Gamma(\bar{B}^0(t) \rightarrow K_S^0 \pi^0) + \Gamma(B^0(t) \rightarrow K_S^0 \pi^0)} = A_{CP}^{K_S^0 \pi^0} \cos(\Delta m t) + S_{CP}^{K_S^0 \pi^0} \sin(\Delta m t)$$

- Because a single penguin amplitude is expected to dominate,
 - $S_{CP}^{K_S^0 \pi^0}$ should arise mostly from the phase 2β of B^0 - \bar{B}^0 mixing ($\beta \equiv \arg(V_{td}^*)$)
 - $\sin(2\beta)$ has been precisely measured
 - new physics at loop level will deviate $S_{CP}^{K_S^0 \pi^0}$ from $\sin(2\beta)$
- Precise measurement of $S_{CP}^{K_S^0 \pi^0}$ can probe for new physics at loop level



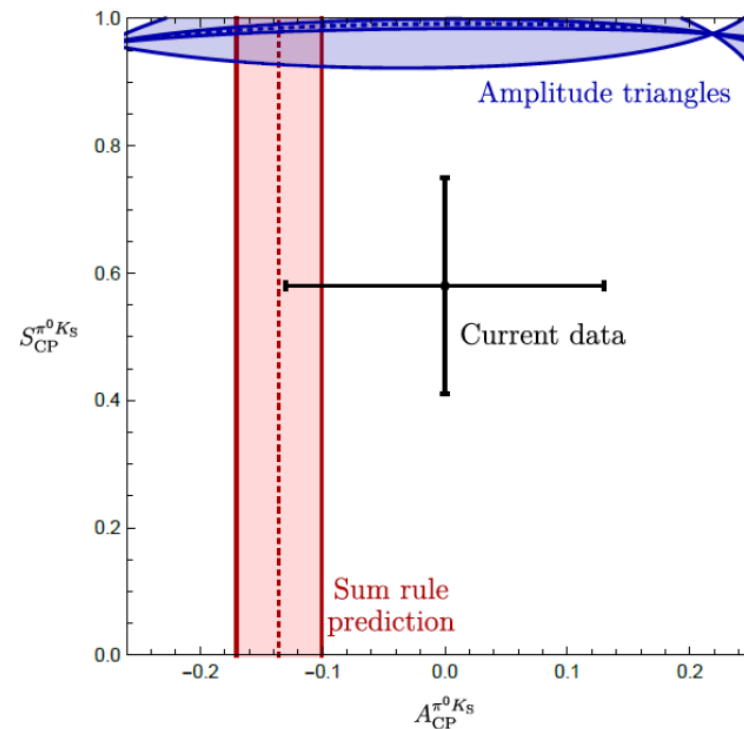
The Broader Picture of the $B \rightarrow K \pi$ Puzzle

- More importantly, the correlation between $S_{CP}^{K_S^0 \pi^0}$ and $A_{CP}^{K_S^0 \pi^0}$ can be derived as:

$$S_{CP}^{K_S^0 \pi^0} = \sin(\phi_d - \phi_{00}) \sqrt{1 - (A_{CP}^{K_S^0 \pi^0})^2}$$

- ϕ_{00} is the angle between amplitude $A(B^0 \rightarrow K^0 \pi^0)$ its charge conjugate
- $\phi_d = (43.2 \pm 1.8)^\circ$ is the B^0 - \bar{B}^0 mixing phase

- The correlation, the sum rule prediction, and the current experimental values are shown on the $S_{CP}^{K_S^0 \pi^0}$ - $A_{CP}^{K_S^0 \pi^0}$ plane:

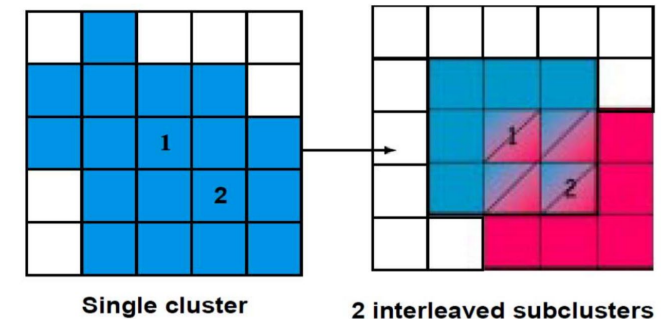
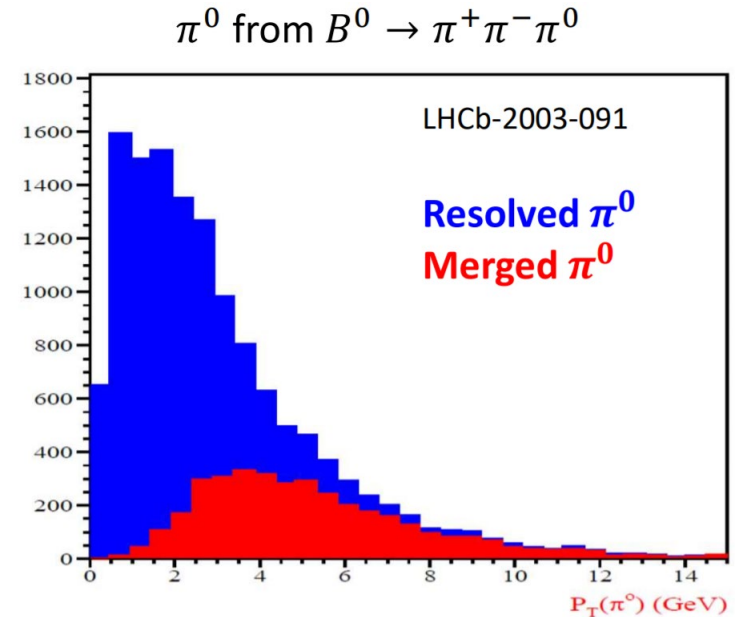


(2.5 σ discrepancy)

R. Fleischer, R. Jaarsma, E. Malami, K. K. Vos, Probing New Physics in $B \rightarrow \pi K$ Decays, PoS LHCP2018 (2018) 176.

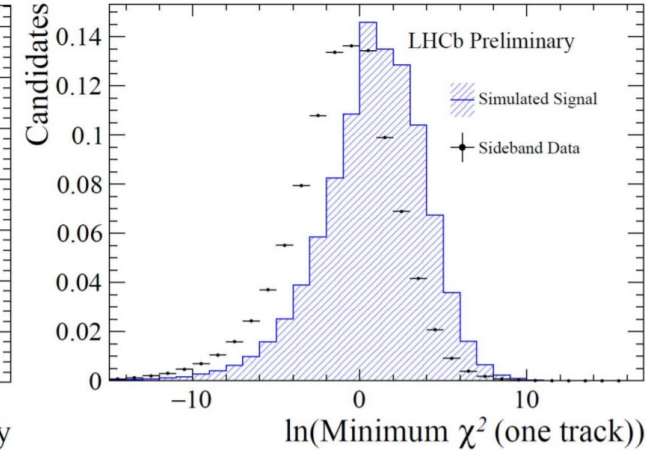
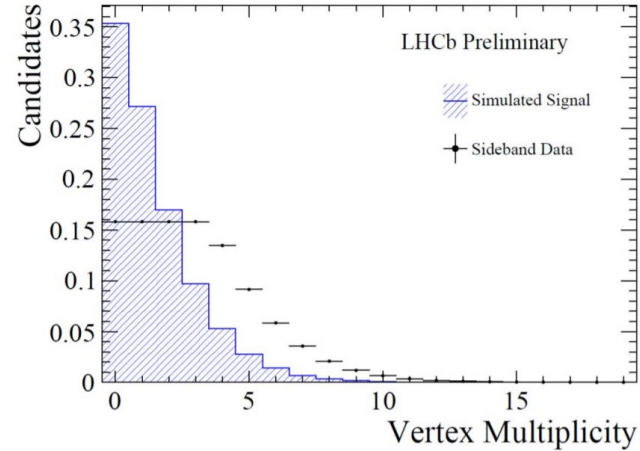
π^0 reconstruction

- π^0 identified by decay to two photons
- At higher p_T , two photon showers merge into a single cluster
- Reconstructed as two subclusters centered on highest energy deposits according to expected transverse profile
- Despite wider mass resolution, use only merged π^0 to reduce combinatorial background and preserve bandwidth



Isolation Variables

- Events with other tracks pointing back to B candidate are unlikely to be signal
- Combine each track in the event individually with K to form vertices



- Consider tracks within cone $\Delta R = 1.7$ around B momentum
- Define cone pT asymmetry $A(p_T) \equiv \frac{p_T(B) - p_T(\text{cone})}{p_T(B) + p_T(\text{cone})}$

