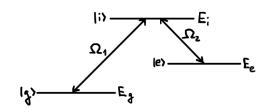
Sheet 1 - The Rule of Two, Episode I

(Due: 01.11.2021, 18:00)

Exercise 1 - Tricks of the Trade I (12 points)

- (a) Show that on a Hilbert space \mathcal{H} , any unitary operator $\hat{U}: \mathcal{H} \mapsto \mathcal{H}$ preserves norms, i.e. $\forall |\psi\rangle \in \mathcal{H}$ the relation $\|\hat{U}|\psi\rangle\| = \||\psi\rangle\|$ holds. Moreover, show that all eigenvalues of \hat{U} lie on the complex unit circle. [2P]
- (b) Consider the so-called Λ -system illustrated in the figure on the right (for simplicity all displayed quantities can be assumed to be real-valued). Write the Hamiltonian \hat{H} of this system both as an abstract Hilbert-space operator involving the sum of dyadic products and as a marix with respect to the basis $\{|g\rangle, |i\rangle, |e\rangle\}$. [4P] Hint: Horizontal lines indicate (field-free) energy levels, e.g., $\langle g|\hat{H}|g\rangle = E_g$. Connecting arrows indicate couplings between these levels, e.g., $\langle g|\hat{H}|i\rangle = \Omega_1$.



- (c) Consider the Hamiltonian $\hat{H} = \alpha |g\rangle \langle g| + \beta |e\rangle \langle e| \ (\alpha, \beta \in \mathbb{R})$ on a Hilbert space \mathcal{H} with orthonormal basis states $\{|g\rangle, |e\rangle\}$. Write an explicit expression for the time-evolution operator $\hat{U} = e^{-\frac{i}{\hbar}\hat{H}t}$ as a sum of dyadic products on \mathcal{H} . [2P]
- (d) The Heisenberg equation of motion allows to find the time evolution of expectation values via an ordinary first-order differential equation. Specifically, given a Hamiltonian \hat{H} and an operator \hat{A} without explicit time dependence, the expectation value $\langle \hat{A} \rangle_{|\psi(t)\rangle} \equiv \langle \psi(t) | \hat{A} | \psi(t) \rangle$ behaves as

$$\frac{d}{dt}\langle \hat{A}\rangle_{|\psi(t)\rangle}(t) = \frac{i}{\hbar}\langle [\hat{H}, \hat{A}]\rangle_{|\psi(t)\rangle}.$$

Which statements can you make if the commutator $[\hat{H}, \hat{A}]$ vanishes? How does the solution for $\langle \hat{A} \rangle_{|\psi(t)\rangle}(t)$ look if $[\hat{H}, \hat{A}] = \hbar \kappa \hat{A}$ for some parameter $\kappa \in \mathbb{C}$? [4P]

Exercise 2 - Like a Hamster in the Wheel (22 points)

From a quantum information point of view, the most fundamental quantum object is the two-level system, or *qubit*. Any such qubit can be mathematically described by a Hilbert space \mathcal{H}_2 , formed by the so-called *canonical basis* states $|0\rangle$ and $|1\rangle$.

- (a) Qubits can be physically realised through a wide variety of implementations. Provide two examples for such a realisation and explain which physical states of the system are encoded in the canonical basis. [2P]

 Hint: You are welcome to research the internet for examples! That being said, providing references for your sources is a critical part of good scientific practice.
- (b) The Hilbert space \mathcal{H}_2 contains fewer states than the "naive" vector space $\mathbb{C}^2 \equiv \operatorname{span}_{\mathbb{C}} \{|0\rangle, |1\rangle\}$. Show that without any loss of physical generality, the state of a qubit can be written as $(\theta, \varphi \in \mathbb{R})$

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\varphi}\sin\frac{\theta}{2}|1\rangle$$
 (1)

Compare the number of real-valued parameters in this expression with the number of real-valued parameters required to describe a general vector in \mathbb{C}^2 . Explain the physical origin for the discrepancy. [3P] Hint: How would a state's global phase $\chi \in [0, 2\pi)$, i.e., $|\psi\rangle \longrightarrow e^{i\chi} |\psi\rangle$, affect expectation values?

(c) The three Pauli matrices (here and in the following interpreted as matrices written with respect to the canonical basis),

$$\hat{\sigma}_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \hat{\sigma}_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \hat{\sigma}_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

allow to completely characterise the state of a qubit. Show that for the state in Eq. (1) one obtains [4P]

$$\langle \hat{\sigma}_x \rangle_{|\psi\rangle} = \sin \theta \cos \varphi, \quad \langle \hat{\sigma}_y \rangle_{|\psi\rangle} = \sin \theta \sin \varphi, \quad \langle \hat{\sigma}_z \rangle_{|\psi\rangle} = \cos \theta.$$
 (2)

(d) Due to Eq. (2), the state of a qubit can be bijectively mapped onto a point on the surface of the three-dimensional unit sphere, where - as it is usually done in physics - θ is the polar angle and φ is the azimuthal angle. This is known as the *Bloch sphere representation*. Which (sets of) states correspond to (i) the north pole, (ii) the south pole, and (iii) the equator of the Bloch sphere? [2P]

(e) Consider the following Hamiltonian,

$$\hat{H} = \frac{\hbar\Omega_0}{2}\hat{\sigma}_x + \frac{\hbar\Delta}{2}\hat{\sigma}_z. \tag{3}$$

Write this Hamiltonian as a matrix with respect to the canonical basis and determine its eigenstates and eigenvalues. [3P]

(f) Show that, given the Hamiltonian \hat{H} in Eq. (3) for $\Delta = 0$, a state with $|\psi(t=0)\rangle = |0\rangle$ transforms under time evolution as follows,

$$|\psi(t)\rangle = \cos\left(\frac{\Omega_0 t}{2}\right)|0\rangle - i\sin\left(\frac{\Omega_0 t}{2}\right)|1\rangle$$
.

Sketch and describe the state's evolution on the Bloch sphere. For which times t is the initial state $|0\rangle$ transformed to the state $|1\rangle$. At which times t does the initial state $|0\rangle$ return to itself? [8P]

Exercise 3 - Like a Cat on a Swing (26 points)

In this exercise we consider a one-dimensional harmonic oscillator with frequency ω , corresponding to the Hamiltonian

$$\hat{H} = \hbar\omega \left(\hat{a}^{\dagger} \hat{a} + \frac{1}{2} \right),\tag{4}$$

where \hat{a} is the so-called annihilation operator, and \hat{a}^{\dagger} is its Hermitian adjoint, the so-called creation operator.

- (a) Show that the so-called number operator, $\hat{n} \equiv \hat{a}^{\dagger} \hat{a}$, is Hermitian and all its eigenvalues are non-negative. [2P]
- (b) The Hamiltonian \hat{H} acts on the infinite-dimensional Hilbert space $\mathcal{H} = L^2(\mathbb{R})$. An orthonormal basis of \mathcal{H} is given by the infinite set $\{|0\rangle, |1\rangle, |2\rangle, \dots\}$, the so-called *Fock states*, with $\hat{a}|0\rangle = 0$ and

$$\forall n \neq 0: \ \hat{a} |n\rangle = \sqrt{n} |n-1\rangle \ . \tag{5}$$

Show that $\hat{a}^{\dagger} | n \rangle = \sqrt{n+1} | n+1 \rangle$. [2P]

Hint: Multiply Eq. (5) from the left with $|l\rangle$ for an arbitrary $l \in \mathbb{N}$.

(c) Show that the expectation values of position and momentum in a Fock state vanish, i.e., $\langle \hat{x} \rangle_{|n\rangle} = \langle \hat{p} \rangle_{|n\rangle} = 0$. [2P]

Hint: In suitable units (which we will be using in this exercise) the relations $\hat{x} = \frac{1}{\sqrt{2}} (\hat{a} + \hat{a}^{\dagger})$ and $\hat{p} = \frac{1}{\sqrt{2}i} (\hat{a} - \hat{a}^{\dagger})$ hold.

- (d) Show that for an arbitrary $\alpha \in \mathbb{C}$ the state $|\alpha\rangle = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} C_{\alpha} |n\rangle$ is an eigenstate of \hat{a} these states are called coherent states. Determine the corresponding eigenvalue λ_{α} and the normalization constant C_{α} . Why are the eigenvalues of the annihilation operator not necessarily real-valued? [5P]
- (e) Show that the expectation values of position and momentum operator in a coherent state with $\alpha = |\alpha| e^{i\varphi}$ take on the values [3P]

$$\langle \hat{x} \rangle_{|\alpha\rangle} = \sqrt{2} \, |\alpha| \cos{(\varphi)} \,\,, \quad \langle \hat{p} \rangle_{|\alpha\rangle} = \sqrt{2} \, |\alpha| \sin{(\varphi)} \,\,.$$

(f) Consider the time evolution of the harmonic oscillator according to the Hamiltonian in Eq. (4) via the Heisenberg equation of motion for an initial coherent state, $|\alpha(t=0)\rangle = |\alpha\rangle$, with $\alpha = |\alpha|e^{i\varphi}$. Show that the time-dependent expectation values of position and momentum in this state obey the following equations, [5P]

$$\langle \hat{x} \rangle_{|\alpha(t)\rangle} = \sqrt{2} \, |\alpha| \cos \left(\varphi - \omega t\right) \,, \quad \langle \hat{p} \rangle_{|\alpha(t)\rangle} = \sqrt{2} \, |\alpha| \sin \left(\varphi - \omega t\right) \,.$$

(g) A superposition of two coherent states with opposite sign is called a coherent cat state,

$$|\langle \alpha \rangle \rangle = N_{\alpha} (|\alpha\rangle + |-\alpha\rangle)$$
.

Usually in quantum mechanics, a cat state is a superposition of two orthogonal states. Is this also true for coherent cat states? If not, is it approximately true for some values of α ? Finally, determine the normalisation constant N_{α} . [4P]

(h) Consider a cat state in the harmonic oscillator according to Eq. (4) at time zero, i.e., $|\psi(t=0)\rangle = |\{\xi\}_{\alpha}\rangle$. Calculate the time-dependent expectation values of position and momentum operator in this state. [3P] Hint: The calculation and result for this exercise is comparatively boring. Looking only at the two expectation values $\langle \hat{x} \rangle$ and $\langle \hat{p} \rangle$ evidently does not reveal what is actually going on in this continuous-variable system with its infinite-dimensional Hilbert space. On the next sheet, we will learn about an interesting way to visualise the state of a harmonic oscillator!