

# Studies of ERM Models with Correlated Disorder

by Tom Folgmann

Bachelor Thesis Presentation, 2024

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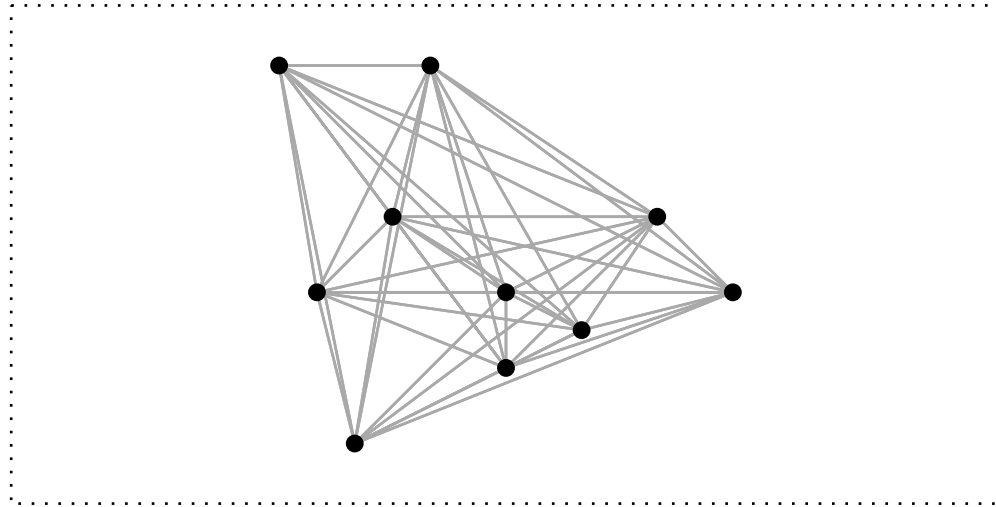
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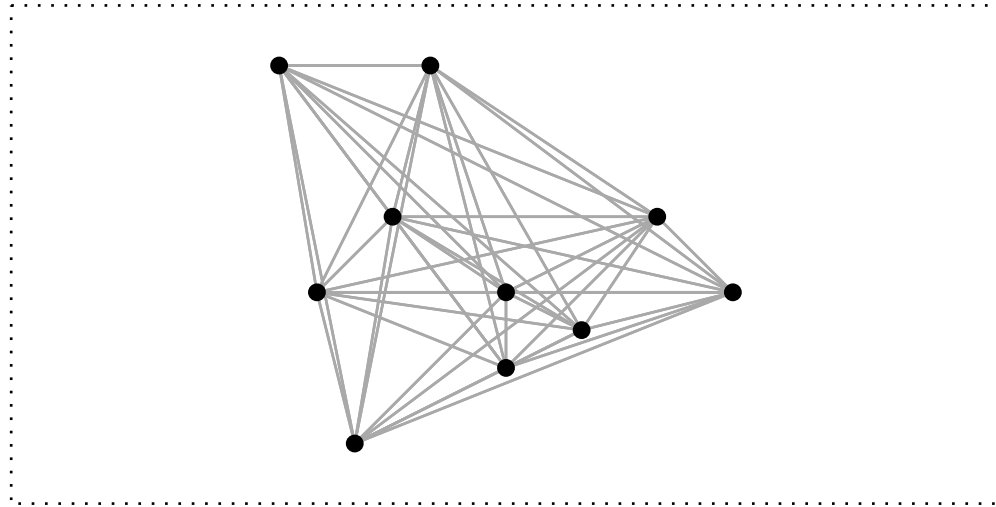
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→ How would you describe their relations?



A system with  $N \in \mathbb{N}$  (related) particles can be described by a mathematical *Graph*.

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.. special case is the *Adjacency Matrix*  $A$ , where  $w_{i,j} \in \{0, -1, 1\}$ .

## Definition of the ERM Laplacian Matrix

In the ERM model the Laplacian matrix is defined as

$$\tilde{U}(f, r) := \begin{pmatrix} \Sigma(f, 1) & -f_{12} & \dots & -f_{1N} \\ -f_{21} & \Sigma(f, 2) & \dots & -f_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ -f_{N1} & -f_{N2} & \dots & \Sigma(f, N) \end{pmatrix},$$

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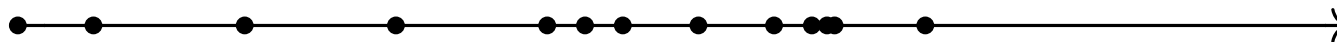
- *Interaction strength* given by  $f_{ij} \stackrel{\text{m}}{=} f(r_i - r_j)$
- *Self-interaction* given by  $\Sigma(f, i) \stackrel{\text{m}}{=} \sum_{j \in [N] \setminus \{i\}} f_{ij}$

## How to measure Eigenvalues?

Let  $\Lambda : [p] \rightarrow \sigma_P(\tilde{U}(f, r))$  map bijectively into the *point spectrum* of the ERM Laplacian.

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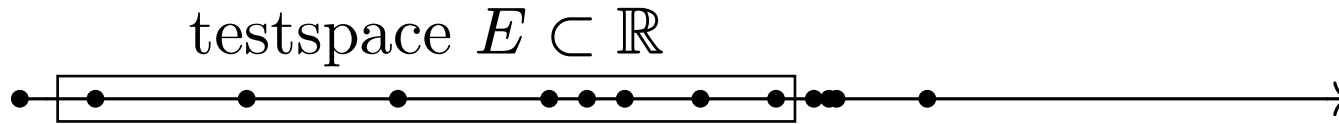
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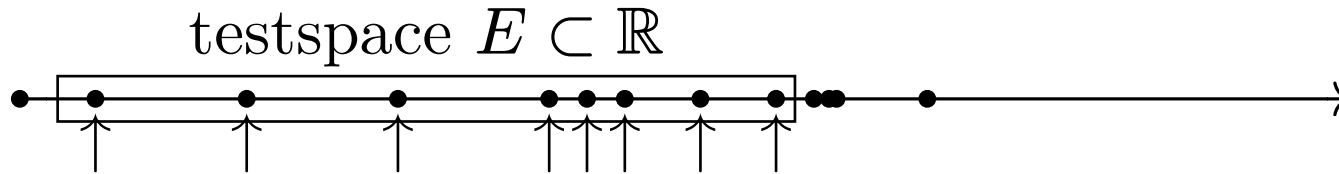
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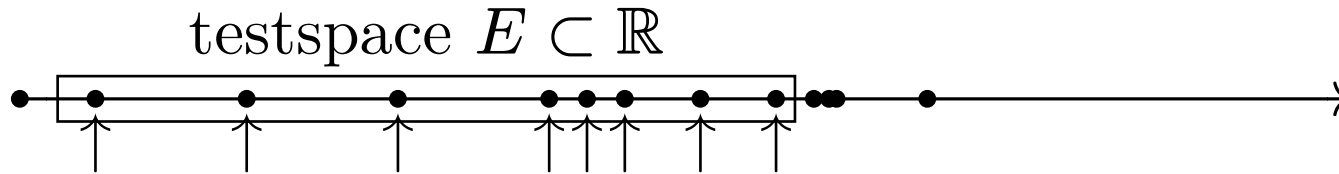
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... results in an (unnormalized) density function

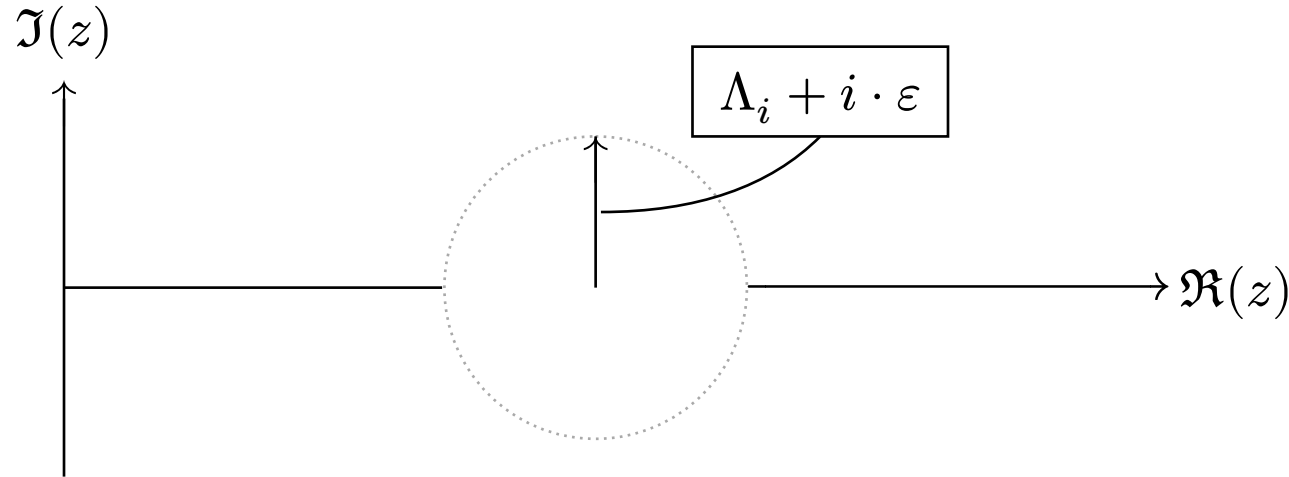
$$E \mapsto \sum_{i \in [p]} \delta_{\Lambda_i}(E) \quad \in \{0, p\}$$

# The Resolvent Eigenvalue Approximation

.. by an example point  $\Lambda_i$  at  $i \in [p]$ .

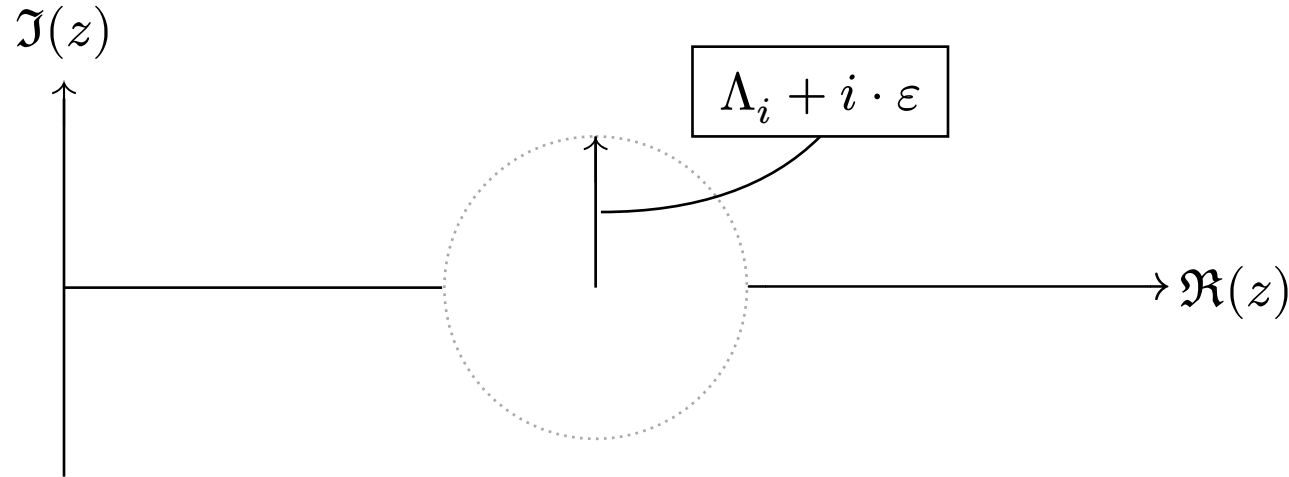
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↪ Usecase is the resolvent with a singularity at  $\Lambda_i$ .

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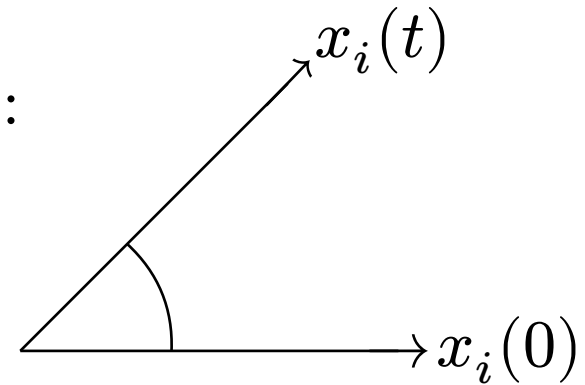
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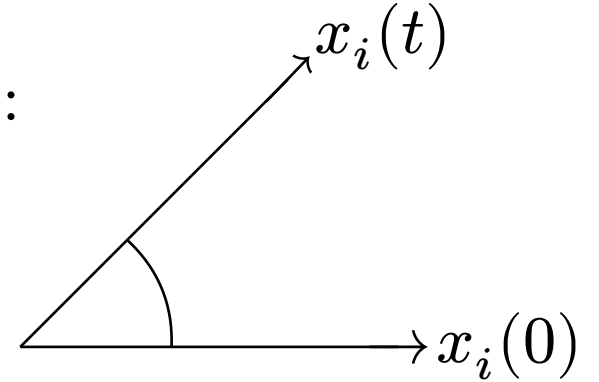
.. looking at the behaviour with regard to the initial conditions:

$$\left(\frac{d}{dt}\right)^2 \langle x_i(t), x_i(0) \rangle = -\tilde{U}(f, i \mapsto x_i(t))_{i,j} \cdot \langle x_j(t), x_i(0) \rangle.$$

In a visual approach  $\langle x_i(t), x_i(0) \rangle$  represents:

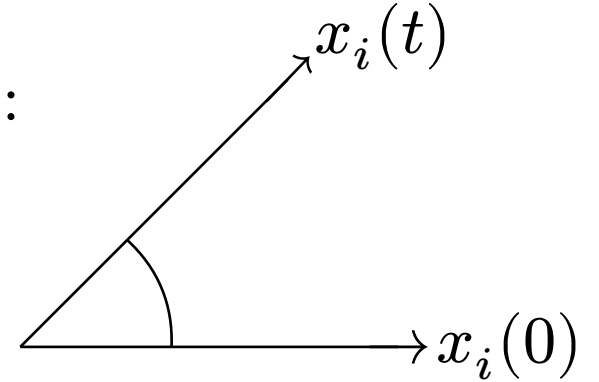


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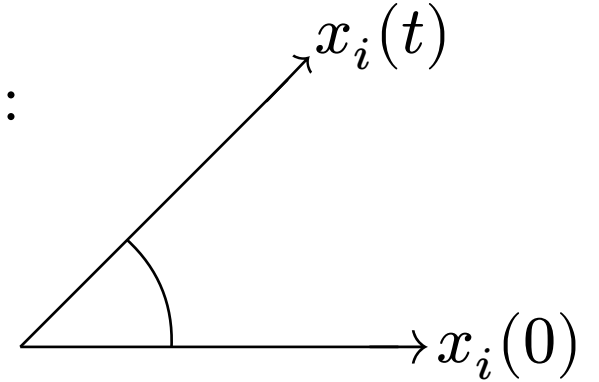
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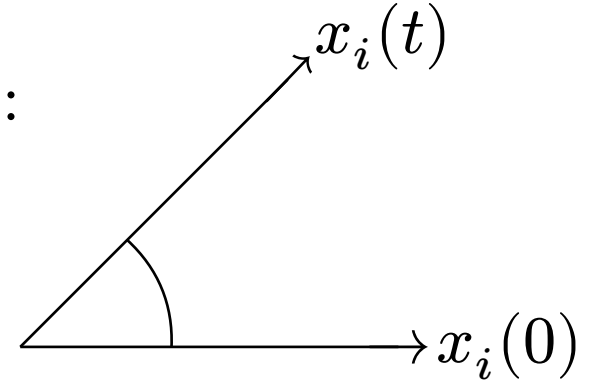
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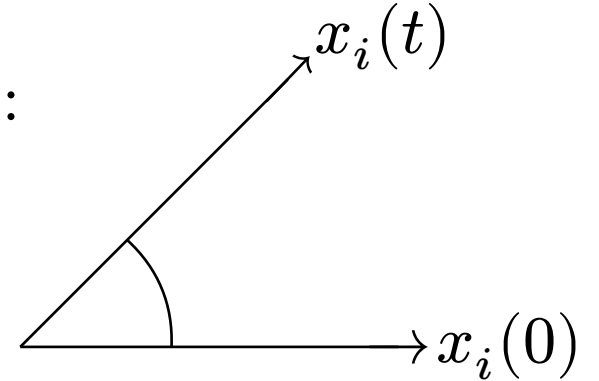
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$$(\mathcal{L}F_{j,i})(s) = \pm \frac{1}{\tilde{U}(f, x^*(t))_{i,j} - \delta_{ij} \cdot \lambda_i^2}.$$

---

<sup>6</sup>With  $x^*(t) = (i \mapsto x_{i(t)})$  and  $F_{j,i}(t) := \langle x_j(t), x_i(0) \rangle$ .

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Ev. step	Meaning
$R$	Random variable, abstract
$R(\omega)$	Vector of time dep. pos.
$R(\omega)_i$	$i$ -th particle position, time dep. path
$R(\omega)_i(t)$	Position of $i$ -th particle at time $t$

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This is already a good starting point to understand our *Correlated Disorder* modification!



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.. missing key elements:

---

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$$-\frac{\beta}{2} \cdot S_{z,R_\omega}(\varphi) := -\frac{\beta}{2} \cdot \left\langle \left( \tilde{U}(f, r) - z \right) \cdot \varphi, \varphi \right\rangle$$

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- The *moment generating function*  $Z_{z,R_\omega}[J]$ . It requires the *force field*  $J$ .

---

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**Definition 2.22.** External Field Shift.

For  $R : \Omega \rightarrow V_{d,N}$  and  $\Phi \in \mathbb{F}_{d,N}$  we define

$$J \mapsto -\frac{1}{2} \cdot S_{z,R_\omega}^{(0)}(\Phi) + \int_{\mathbb{R}^d} J(x) \cdot \Phi(-x) + J(-x) \cdot \Phi(x) \lambda(dx)$$

the *field shifted action*  $S_{z,R_\omega}^{(0)}$  by an external field  $J \in \mathcal{S}(\mathbb{R}^d)$ .

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$$\hookrightarrow \frac{\delta}{\delta J(x)} S_{z,R_\omega}^{(0)}[\Phi] = i \cdot \Phi(-x).$$

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$$Z_{z,R_\omega}[J] = \int_{\mathbb{F}_{d,N}} e^{\left(S_{z,R_\omega}^{(0)}\Phi + S_{z,R_\omega}^{(int)}\Phi\right)[J]} d\Phi = \underbrace{\left[ \mathbb{E}_{\mathcal{L}_f} \left[ \int_{\mathbb{F}_{d,N}} e^{\left(S_{z,R_\omega}^{(0)}\Phi\right)[\cdot]} d\Phi \right] \right]}_{\text{Generative Part}} [J].$$

→ Looking at different Taylor expansion terms yields different integrals.

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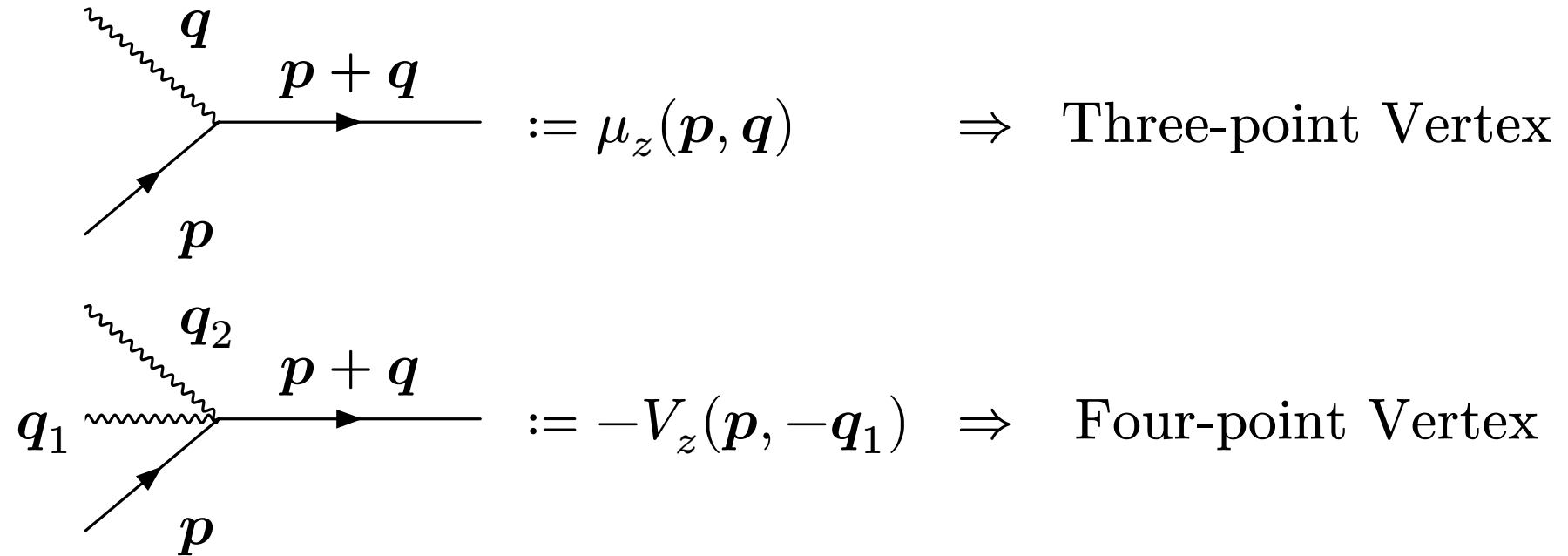
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.. possible connections of these edges are given by *vertices*:

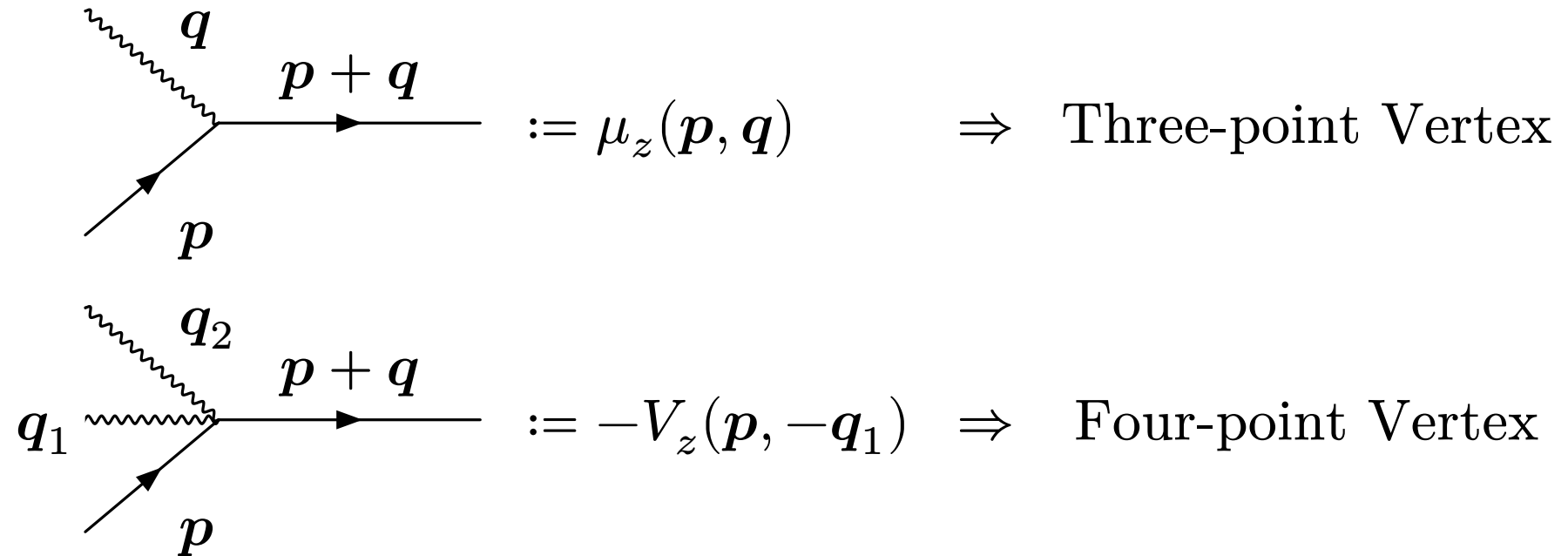
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.. which completes the set of Feynman rules.

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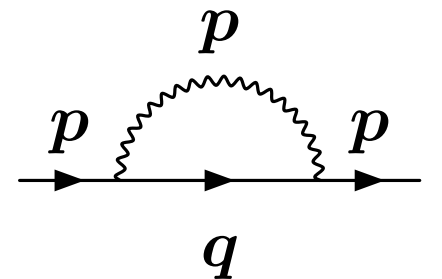
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We begin with **one** loop diagrams:


$$= \frac{G_0(\mathbf{p}, z)^2}{\rho_*} \cdot \int_{\mathbb{R}^d} G_0(\mathbf{q} - \mathbf{p}, z) \cdot \mu_z(\mathbf{p}, -\mathbf{q})^2 d\mathbf{q}$$

# What is Correlated Disorder?

What did we implement?