

Studies of ERM Models with Correlated Disorder

by Tom Folgmann

Bachelor Thesis Presentation, 2024

Upfront: Goals.

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Foundations

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- What is ERM?

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 - Is there a visual approach to calculations?

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- What did we **conclude**?

What is ERM?

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Imagine a system of $N \in \mathbb{N}$ particles.

What is ERM?

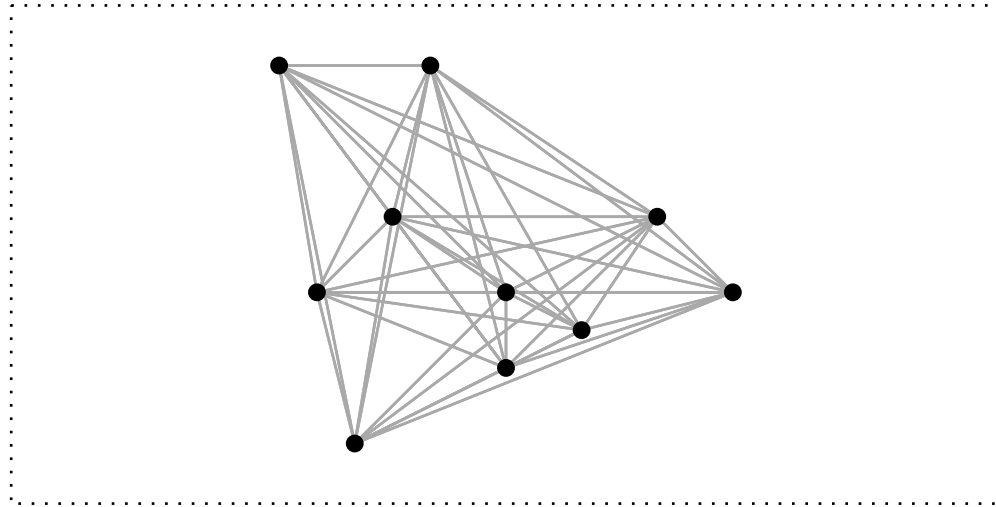
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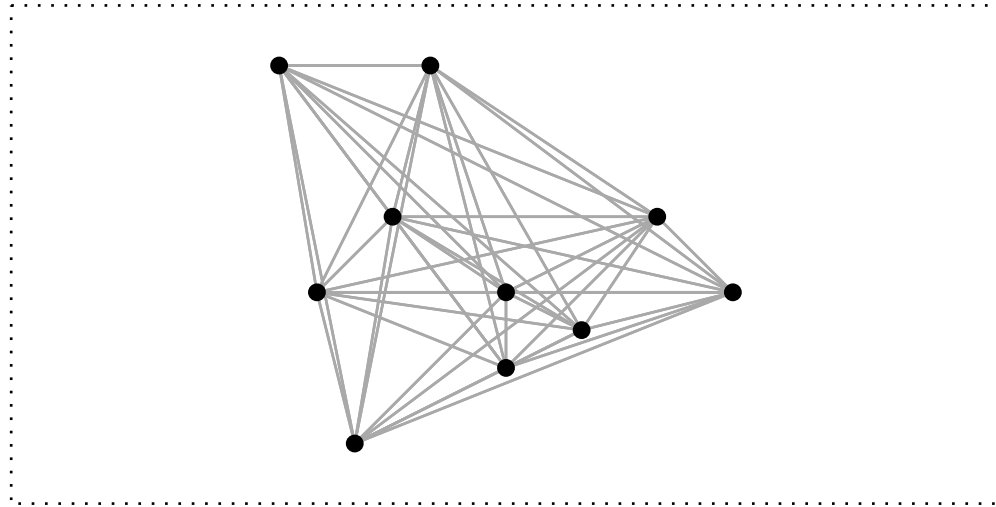
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→ How would you describe their relations?



A system with $N \in \mathbb{N}$ (related) particles can be described by a mathematical *Graph*.

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$$D(G) := \text{diag}(d), \quad d_i := \#\{e \in E : v_i \in e\},$$

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.. special case is the *Adjacency Matrix* A , where $w_{i,j} \in \{0, -1, 1\}$.

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$$\tilde{U}(f, r) := \begin{pmatrix} \Sigma(f, 1) & -f_{12} & \dots & -f_{1N} \\ -f_{21} & \Sigma(f, 2) & \dots & -f_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ -f_{N1} & -f_{N2} & \dots & \Sigma(f, N) \end{pmatrix} = D(G) - W(G).$$

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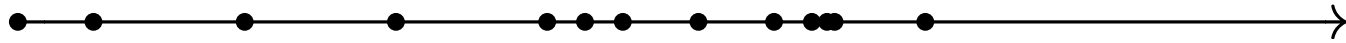
- *Interaction strength* given by $f_{ij} \stackrel{\text{m}}{=} f(r_i - r_j)$
- *Self-interaction* given by $\Sigma(f, i) \stackrel{\text{m}}{=} \sum_{j \in [N] \setminus \{i\}} f_{ij}$

How to measure Eigenvalues?

Let $\Lambda : [k] \rightarrow \sigma_P(\tilde{U}(f, r))$ map bijectively into the *point spectrum* of the ERM Laplacian.

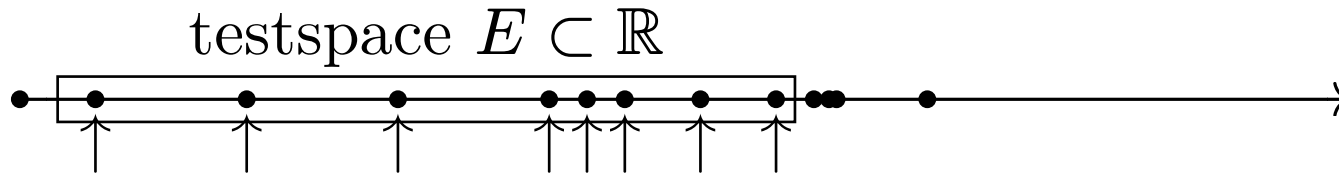
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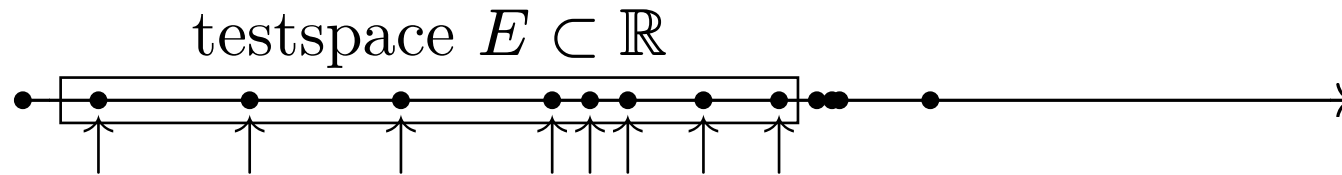
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... results in an (unnormalized) density function

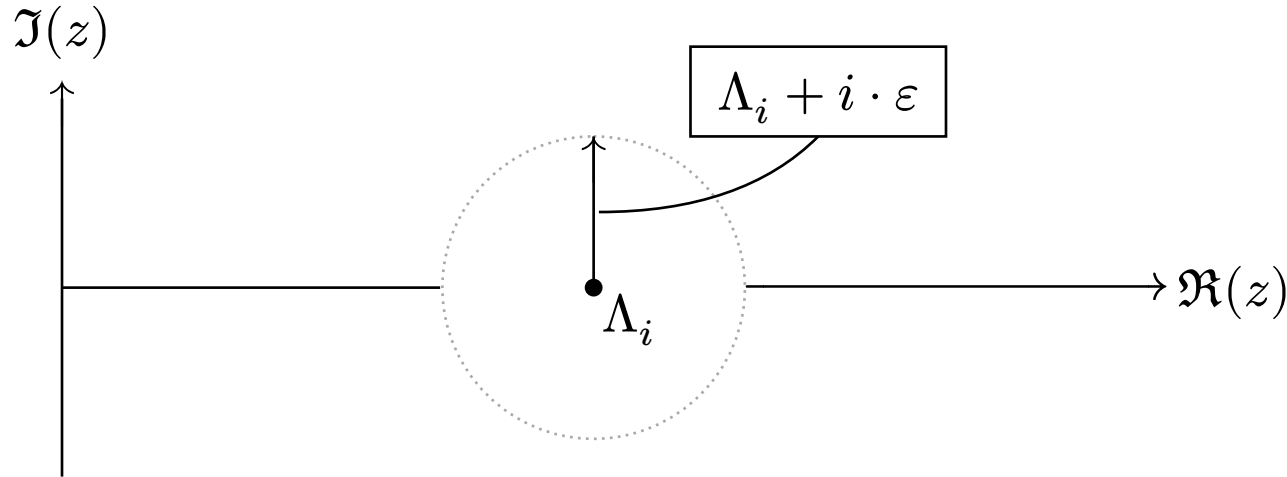
$$E \mapsto \sum_{i \in [k]} \delta_{\Lambda_i}(E) \quad \in \{0, \dots, k\}$$

The Resolvent Eigenvalue Approximation

.. by an example point Λ_i at $i \in [k]$.

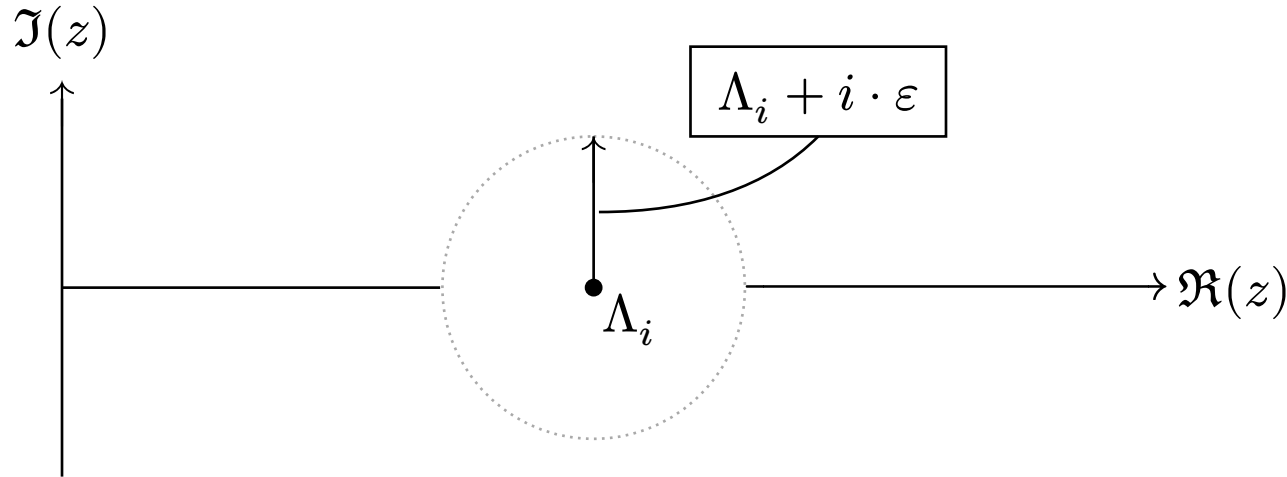
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↪ Usecase is the resolvent with a singularity at Λ_i .

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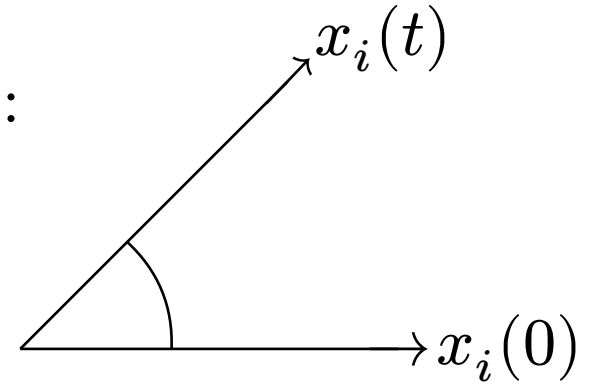
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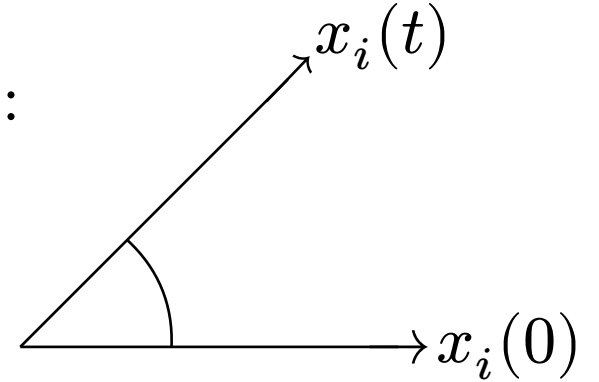
This is just $f_{i,j}$!
.. for $i \neq j$.

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In a visual approach $\langle x_i(t), x_i(0) \rangle$ represents:

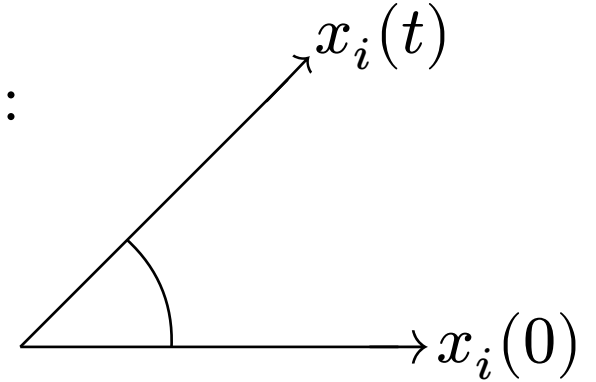


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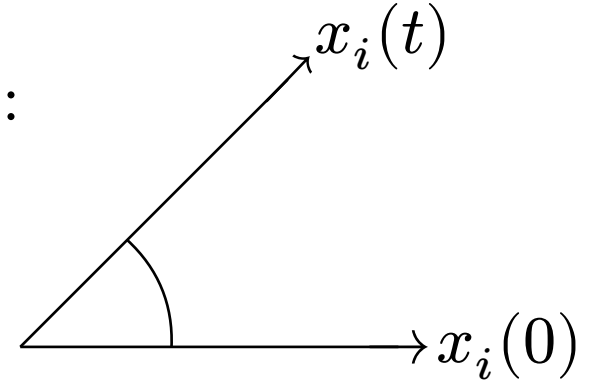
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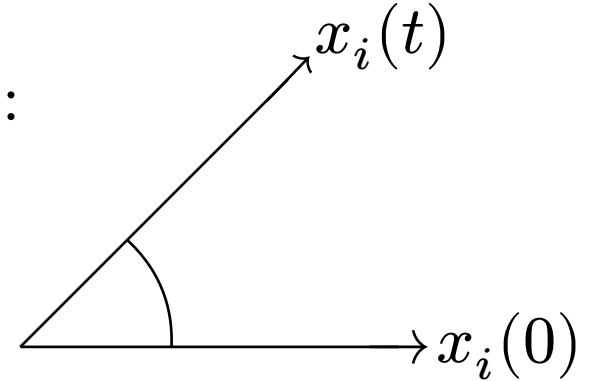
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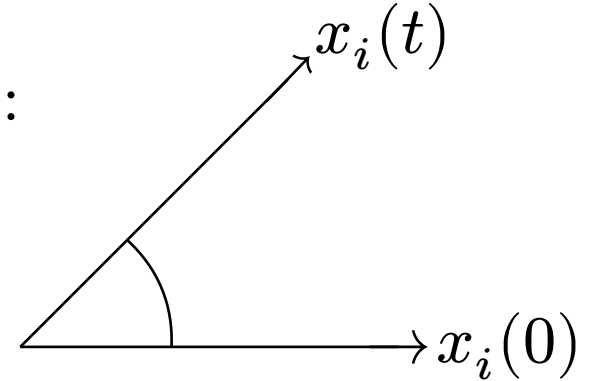
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$$(\mathcal{L}F_{j,i})(s) = \pm \frac{1}{\tilde{U}(f, x^*(t))_{i,j} - \delta_{ij} \cdot \lambda_i^2}.$$

⁶With $x^*(t) = (i \mapsto x_{i(t)})$ and $F_{j,i}(t) := \langle x_j(t), x_i(0) \rangle$.

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$$G_N(\mathbf{p}, z) = \sum_{(i,j)} \int \pm \frac{1}{\tilde{U}(f, r)_{i,j} - \delta_{ij} \cdot z} \cdot e^{i \cdot \mathbf{p} \cdot (x_i - x_j)} dx.$$

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Q: What are we integrating over?

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Ev. step	Meaning
R	Random variable, abstract
$R(\omega)$	Vector of time dep. pos.
$R(\omega)_i$	i -th particle position, time dep. path
$R(\omega)_i(t)$	Position of i -th particle at time t (fixed for us.)

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$$\begin{array}{l} \text{Value in } \mathbb{R} \dots \\ \text{(simplification!)} \end{array} \quad \xrightarrow{\quad} \quad \int_{\mathbb{R}^d} \varphi_i \cdot \varphi_j \left(\overbrace{e^{-\frac{\beta}{2} \cdot \langle (\tilde{U}(f,r) - z) \cdot \varphi, \varphi \rangle} \cdot \lambda}^{\text{Gaussian measure}} \right) (d\varphi). \quad \underbrace{\hspace{10em}}_{\text{Boltzmann density}}$$

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Boltzmann density

This is already a good starting point to understand our *Correlated Disorder* modification!

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.. missing key elements:

¹¹Expansion to a functional can be argued, see thesis p. 19.

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- The *moment generating function* $Z_{z,R_\omega}[J]$. It requires the *force field* J .

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.. coming from ..

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This needs explanation.

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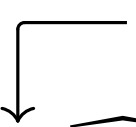
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$$\int_{(\mathbb{R}^d)^2} \mu_z(\mathbf{p}_1, \mathbf{p}_2) \cdot \left(\frac{\delta}{\delta \hat{J}(-\mathbf{p}_1)} \circ \frac{\delta}{\delta \hat{J}(\mathbf{p}_1 + \mathbf{p}_2)} \right) (\lambda \otimes \widehat{\delta \rho_{R_\omega}})(d\mathbf{p})$$

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→ Looking at different Taylor expansion terms of $\underbrace{\exp\left(\int_{(\mathbb{R}^d)^2} \dots d\mathbf{p}\right)}_{\text{Ex}_{\mathcal{L}_f}}$ yields different powers of integral operators.

Feynman Diagrammatics - Edges

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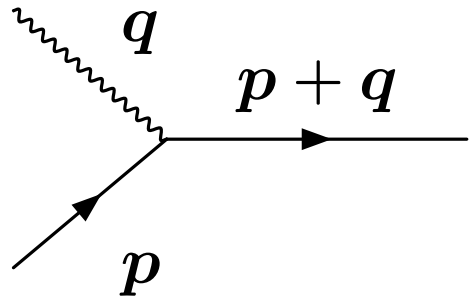
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.. possible connections of these edges are given by *vertices*:

What is ERM?

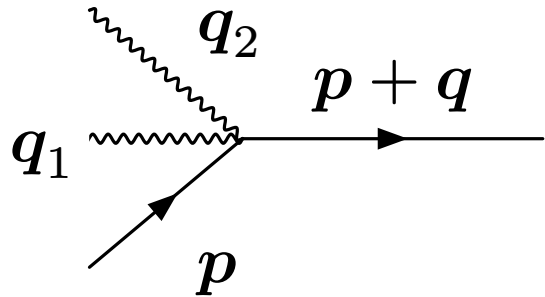
Feynman Diagrammatics - Vertices

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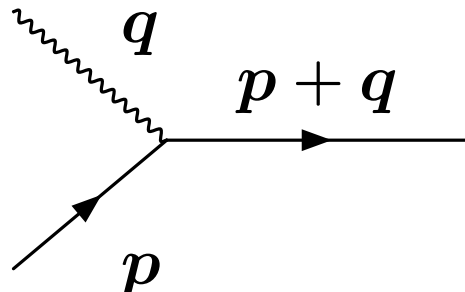
\Rightarrow Three-point Vertex

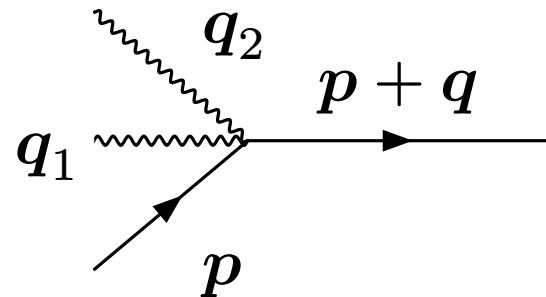


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Feynman Diagrammatics - Vertices


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.. which completes the set of Feynman rules.

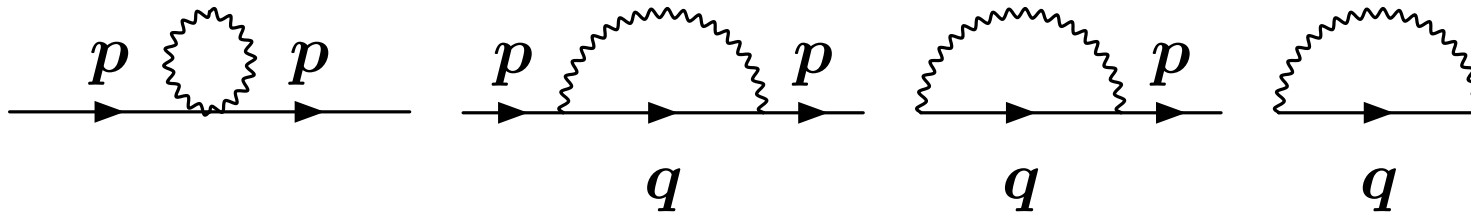
How can we use diagrammatics?

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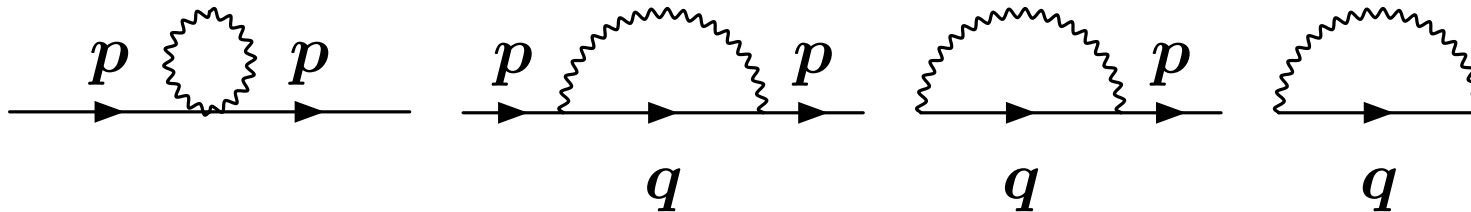
Observe **one** loop diagrams:



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Observe **one** loop diagrams:



.. represented diagrams are *irreducible*: $Z_{z,R_\omega}[J] \propto \exp\left(\sum_{C \in \mathcal{C}} C\right)$.

Integral representations¹⁶

¹⁶Attention! The terms have been simplified. For more details see Thesis sec. 2.4.2.

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$$\begin{aligned}
\text{Diagram 1} &= \frac{G_0(\mathbf{p}, z)^2}{\rho_*} \cdot \int_{\mathbb{R}^d} G_0(\mathbf{q} - \mathbf{p}, z) \cdot \mu_z(\mathbf{p}, -\mathbf{q})^2 d\mathbf{q}, \\
\text{Diagram 2} &= -\frac{2 \cdot G_0(\mathbf{p}, z)}{\rho_*} \cdot \int_{\mathbb{R}^d} G_0(\mathbf{p} - \mathbf{q}, z) \cdot \mu_{z(\mathbf{p}, -\mathbf{q})} d\mathbf{q}, \\
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- Utilization of vertex' and propagator symmetries.

What is Correlated Disorder?

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Main question to solve:

How can we include *structure* in our probability density?

The (radial) Particle Distribution Density

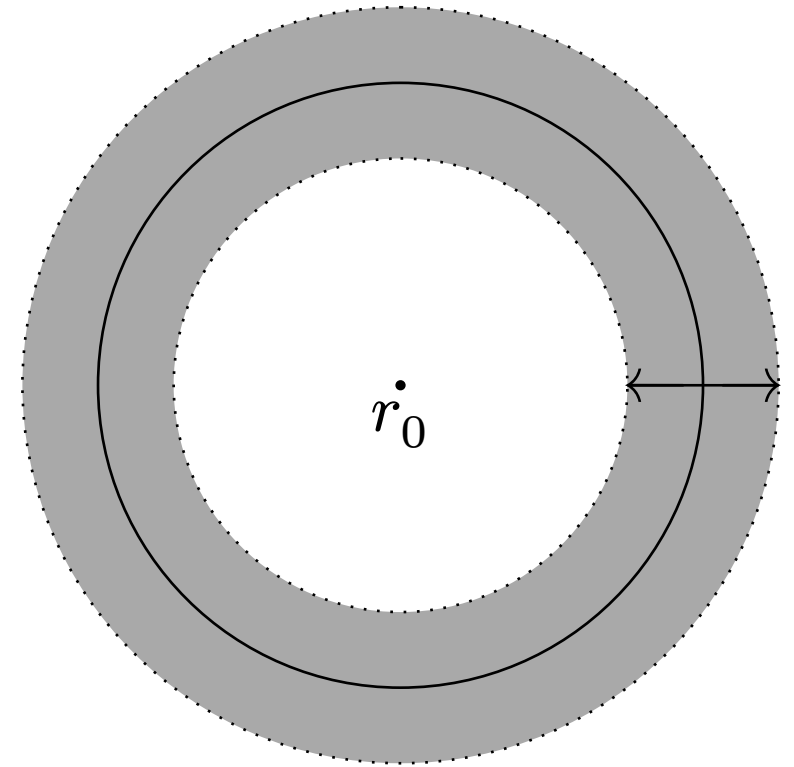
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$$g_\varepsilon(r_0, x_*) = \int_{B_{x_*, \varepsilon}(r_0)} \rho_N^{(2)}(r_0, r) dr,$$

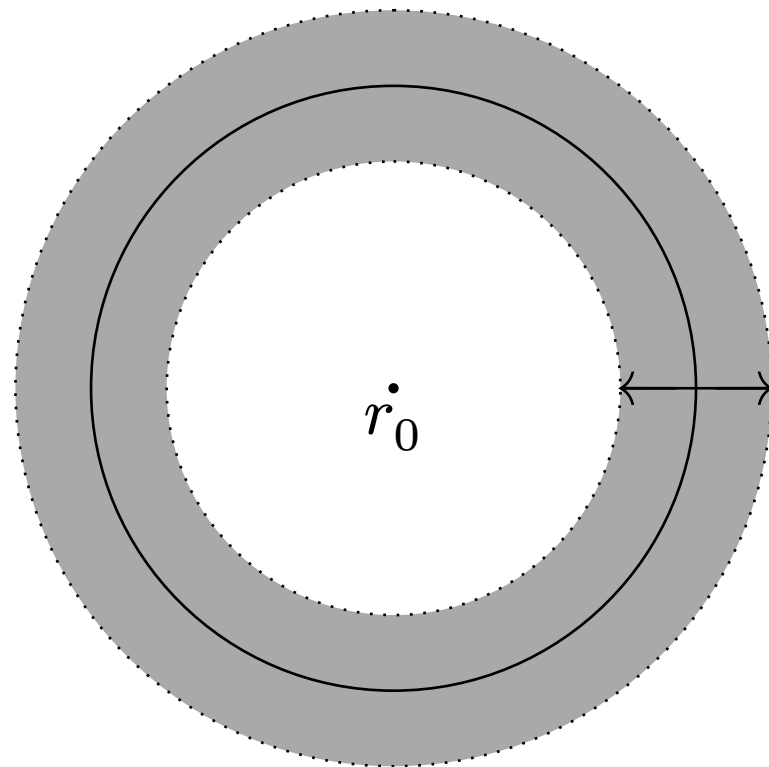


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while $\rho_N^{(2)}$ reflects integration of $\exp(-\beta \cdot H(r, \cdot))$ for remaining particles.



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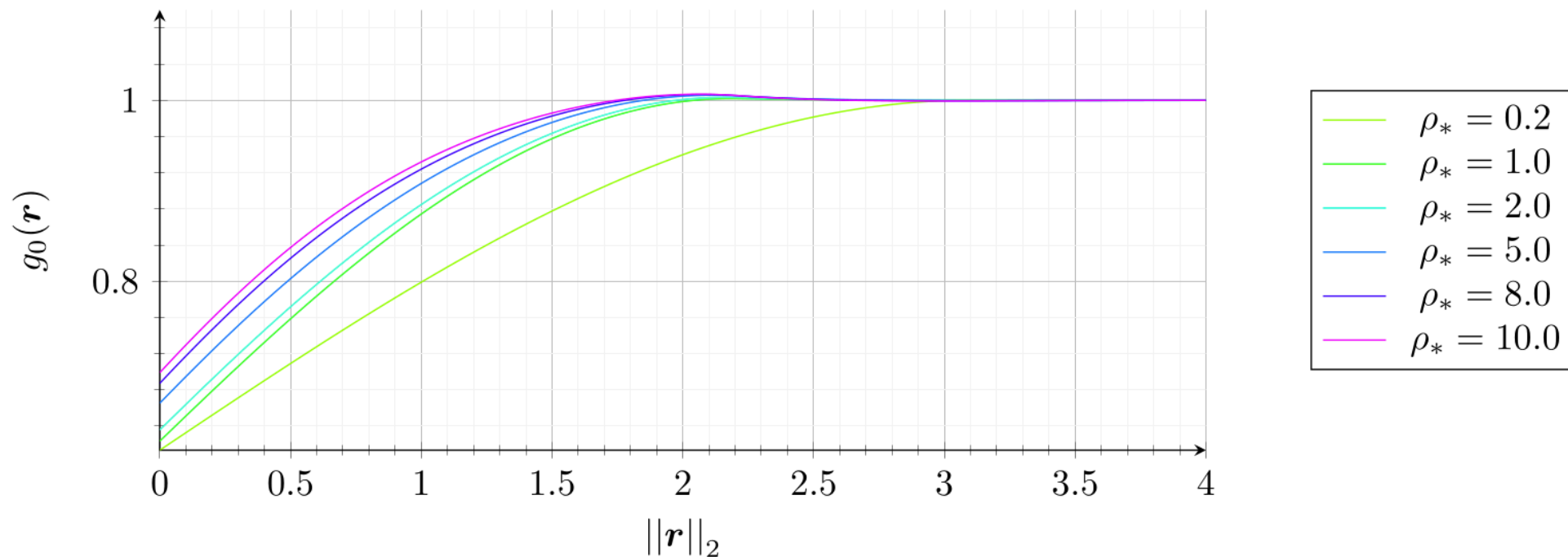
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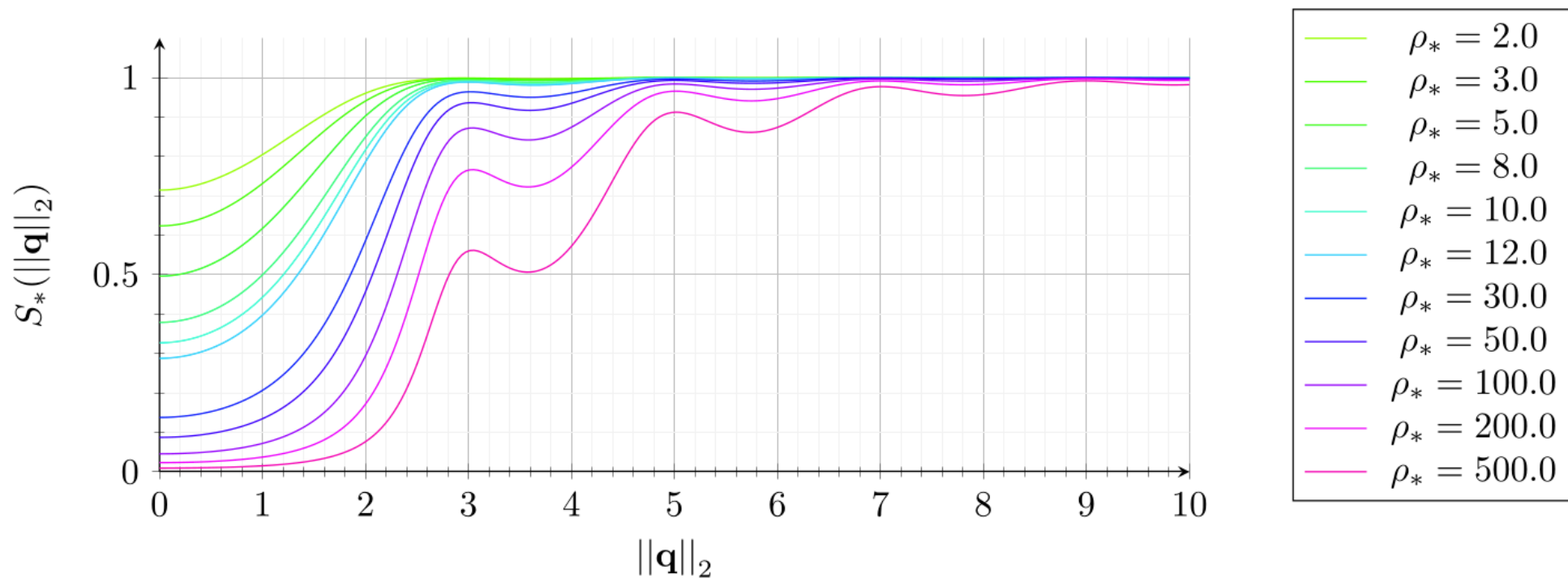
↑
normalized to $r_0 = 0$

What does g_0 look like?

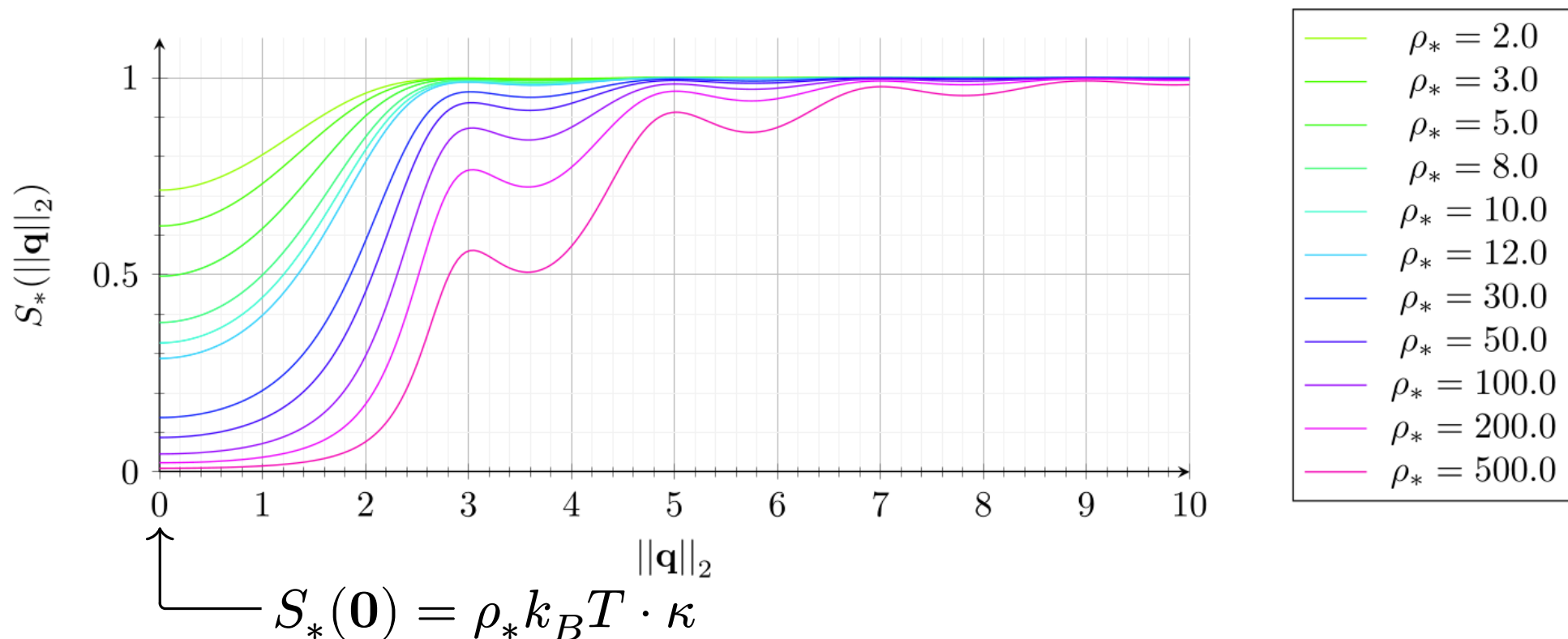
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Resulting in the Static Structure Factor



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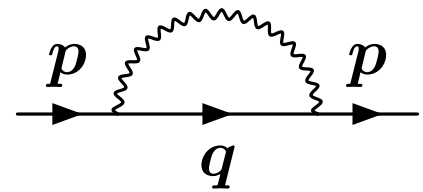
$$\begin{array}{c} \text{Diagram: A horizontal line with an arrow pointing right, labeled } \mathbf{p} \text{ at the left end and } \mathbf{p} \text{ at the right end. A wavy line (representing a disorder) connects the line to itself, forming a loop. Below the loop, the momentum } \mathbf{q} \text{ is indicated.} \\ \hline \end{array} = \frac{G_0(\mathbf{p}, z)^2}{\rho_*} \cdot \int_{\mathbb{R}^d} G_0(\mathbf{q} - \mathbf{p}, z) \cdot \mu_z(\mathbf{p}, -\mathbf{q})^2 \cdot S_*(\mathbf{q}) d\mathbf{q},$$

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This has an explicit approximation built into the spring function!

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- We did not change the zeroth order term in the propagator.

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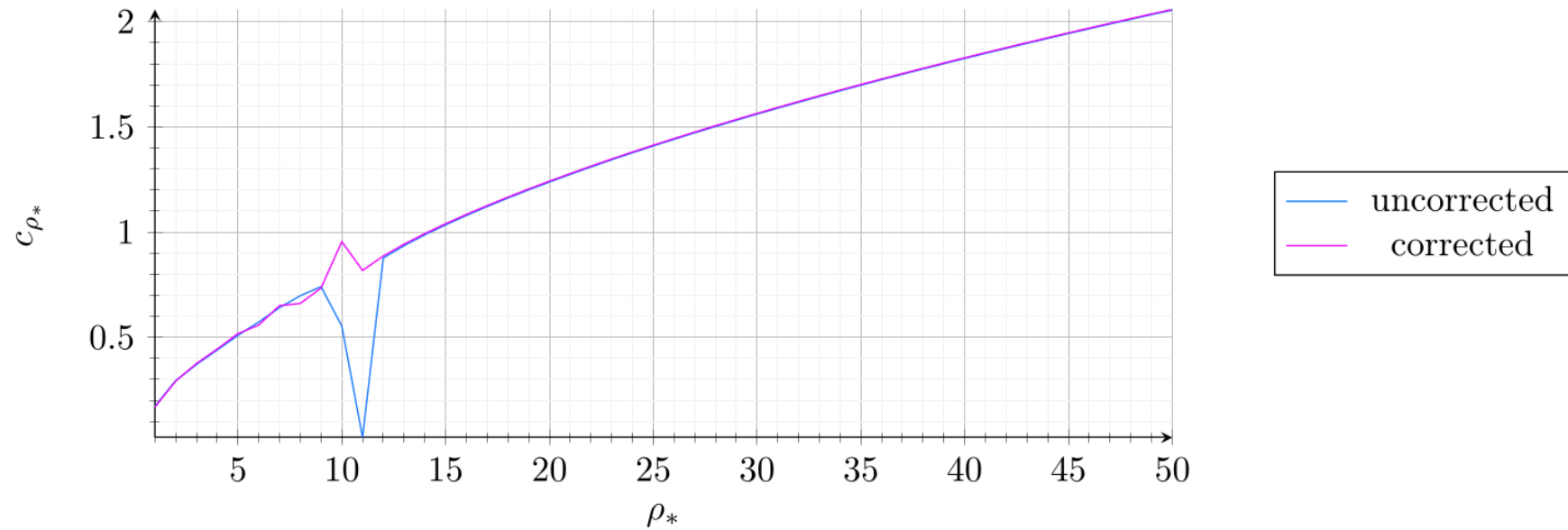
$$V_{d,N} \ni R \mapsto U_a^{(num)}(R) = \sum_{(i,j) \in [N]^2} \begin{cases} \frac{1}{2} \cdot (\|R_i - R_j\| - a)^2 & \text{if } \|R_i - R_j\| < a, \\ 0 & \text{else.} \end{cases}$$

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Results using the Hypernetted Chain

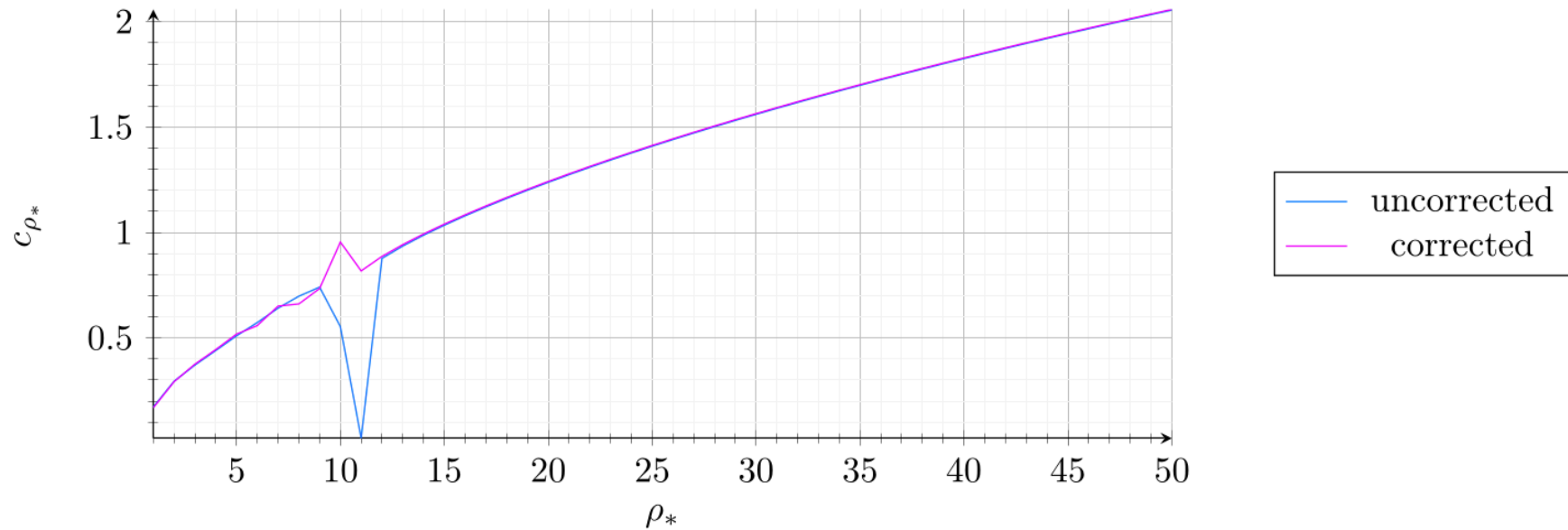
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Results using the Hypernetted Chain



→ No major differences in the velocity of sound noticeable.

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$$\exp(-\beta \cdot U(r)) \approx \exp\left(-\beta \cdot (r - \nu)^\perp \cdot A \cdot (r - \nu)\right), (\rho \text{ mediocre})$$

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Thank you for your attention!

This presentation was only possible thanks to **Typst**.

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Sources used in this presentation do not differ from the ones in the thesis.

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`unb3rechenbar/BA24-CorDis.git`

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