

# Energy diffusion and dissipation processes in ocean surface

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Dedicated in Memmory to Prof  
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# Small Scale Turbulence (J.M. Redondo)

For the small scale, as basic indicator of the potential energy to kinetic energy ratio, we use the **Richardson's Number**, either its flux (Rf) or gradient form (Rg). This is a parameter able to distinguish between different stratification types.

Rf is defined as the ratio of buoyancy and the production term in the turbulent kinetic energy:

$$\frac{\partial K}{\partial t} = \overline{u' j} \frac{\partial}{\partial t} \left( \frac{p'}{\rho} + \frac{1}{2} u'_i u'_i \right) + \overline{u' u} \frac{\partial \bar{u}}{\partial Z} - \frac{g}{\rho} \overline{\rho' w'} - \nu \nabla \overline{u'_j} \nabla \overline{u'_j} \quad \rightarrow \quad Rf = \frac{g}{\rho} \frac{\overline{\rho' w'}}{\overline{u' w'} \frac{\partial \bar{u}}{\partial Z}} \quad (\text{Mixing Efficiency})$$

Considering the Boussinesq relationships between turbulent fluxes and average gradients as well as other expression:

$$K_m \overline{\rho' w'} = - \frac{\partial \rho}{\partial Z}$$

$$\frac{1}{\theta_0} \frac{\partial \theta}{\partial Z} = - \frac{1}{\rho} \frac{\partial \rho}{\partial Z}$$

$$Ri = \frac{K_m}{K_h} Rf$$

Obtaining the gradient Richardson number:

$$Ri = \frac{\frac{g}{\theta_0} \frac{\partial \theta}{\partial Z}}{\left[ \left( \frac{\partial U}{\partial Z} \right)^2 + U^2 \left( \frac{\partial \alpha}{\partial Z} \right)^2 \right]}$$



Where:  $u, \alpha$ -speed and direction wind;  $\theta$ -potential temperature;  $g$ -gravity acceleration and  $Z$ -height.

# Dynamic Model

- Filter again resolved scales
- Assume same form of model can be applied to both subgrid and resolved stresses

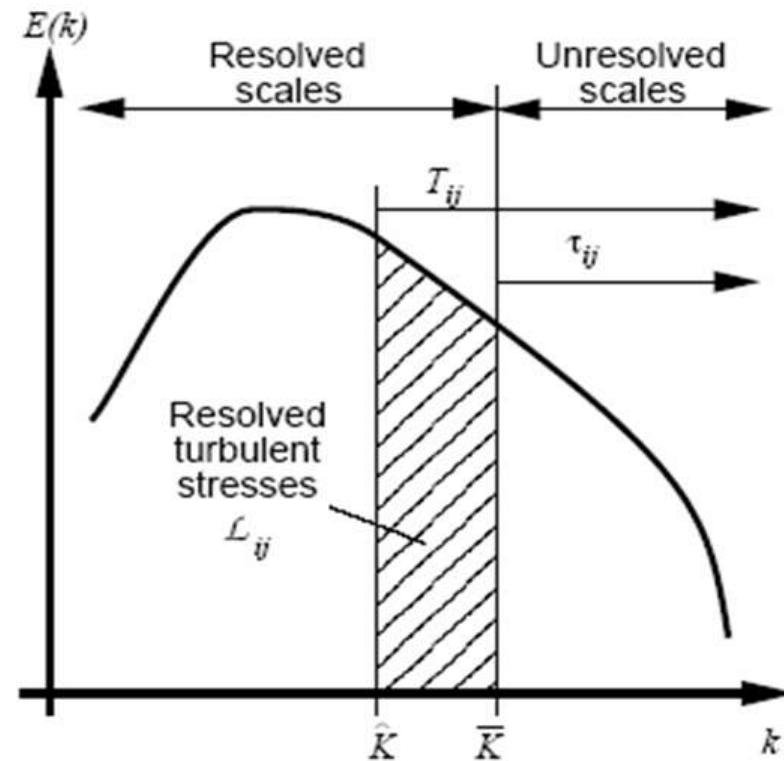
$$T_{ij} = \tilde{\bar{u}_i \bar{u}_j} - \tilde{\bar{u}_i} \tilde{\bar{u}_j}$$

$$L_{ij} = \tilde{\bar{u}_i \bar{u}_j} - \tilde{\bar{u}_i} \tilde{\bar{u}_j}$$

$$L_{ij} = T_{ij} - \tilde{\tau}_{ij}$$

- Compute Smagorinsky constant

Tamay Özgökmen 2011



GERMANO, M. 1992 Turbulence: the filtering approach. *J. Fluid Mech.* **238**, 325.

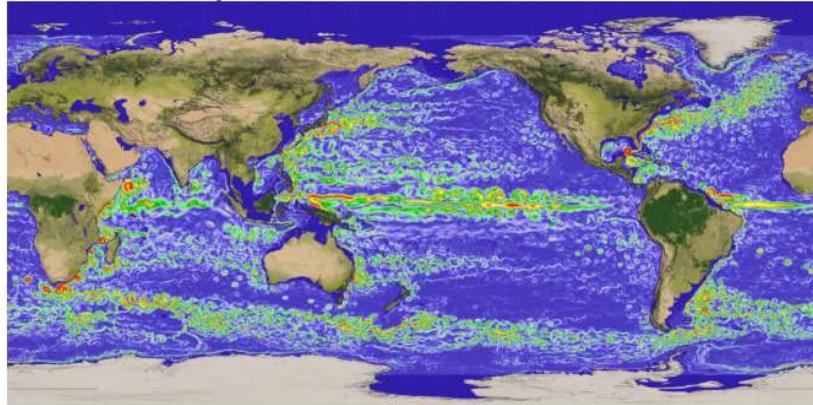
PIOMELLI, U. AND LIU, J. 1995 Large eddy simulation of rotating channel flows using a localized dynamic model. *Phys. Fluids A* **7**, 839.

# Large scales (climate)

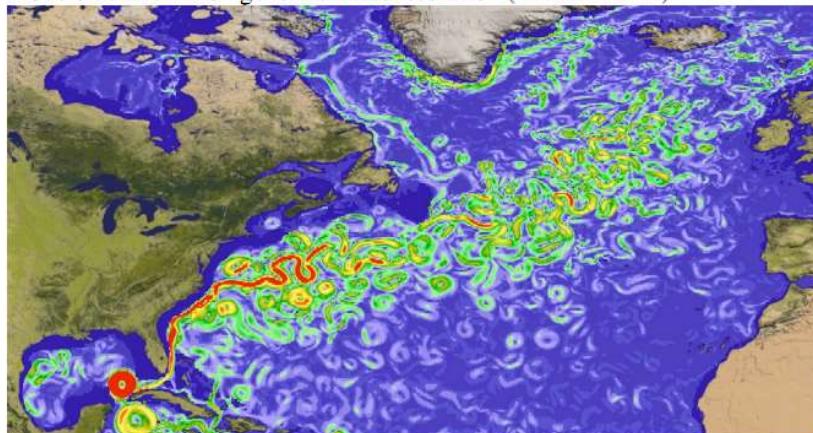
(El Niño – la Niña)

Time scale of horizontal diffusion

Simulations numériques sur le « Earth-Simulator »

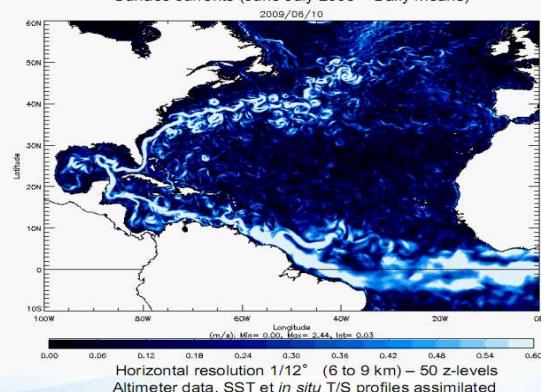


Modèle de circulation générale haute résolution (~10kmx10km)



Capacity illustration: example of products obtained with the 1/12° Atl. Nord + Mediterranean

Surface currents (June July 2009 – Daily means)



Gula et al, J.Phys.Oceanogr.,  
2015

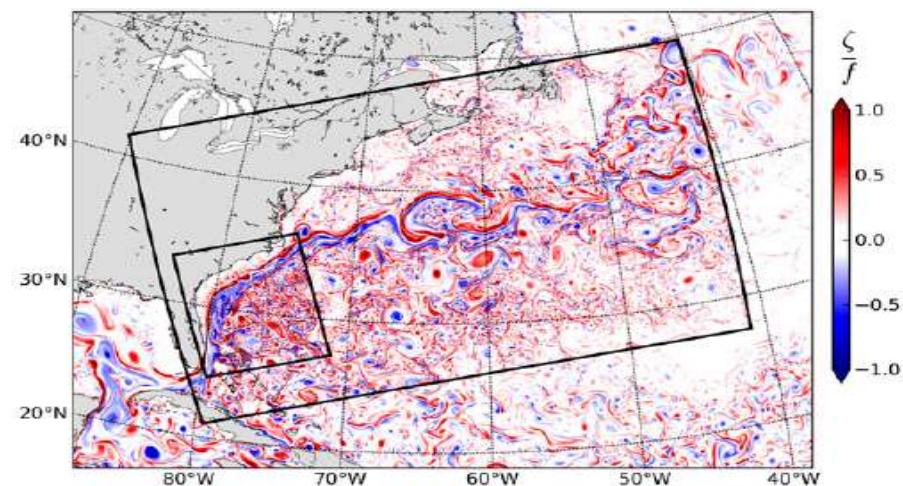
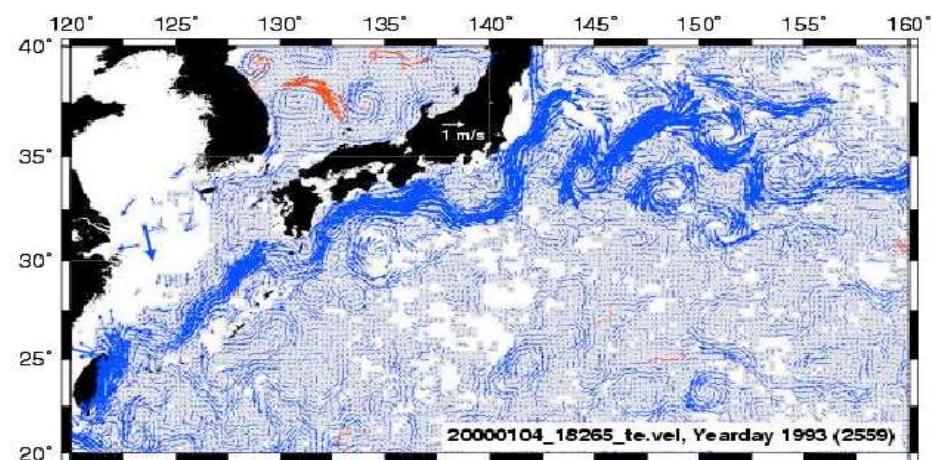
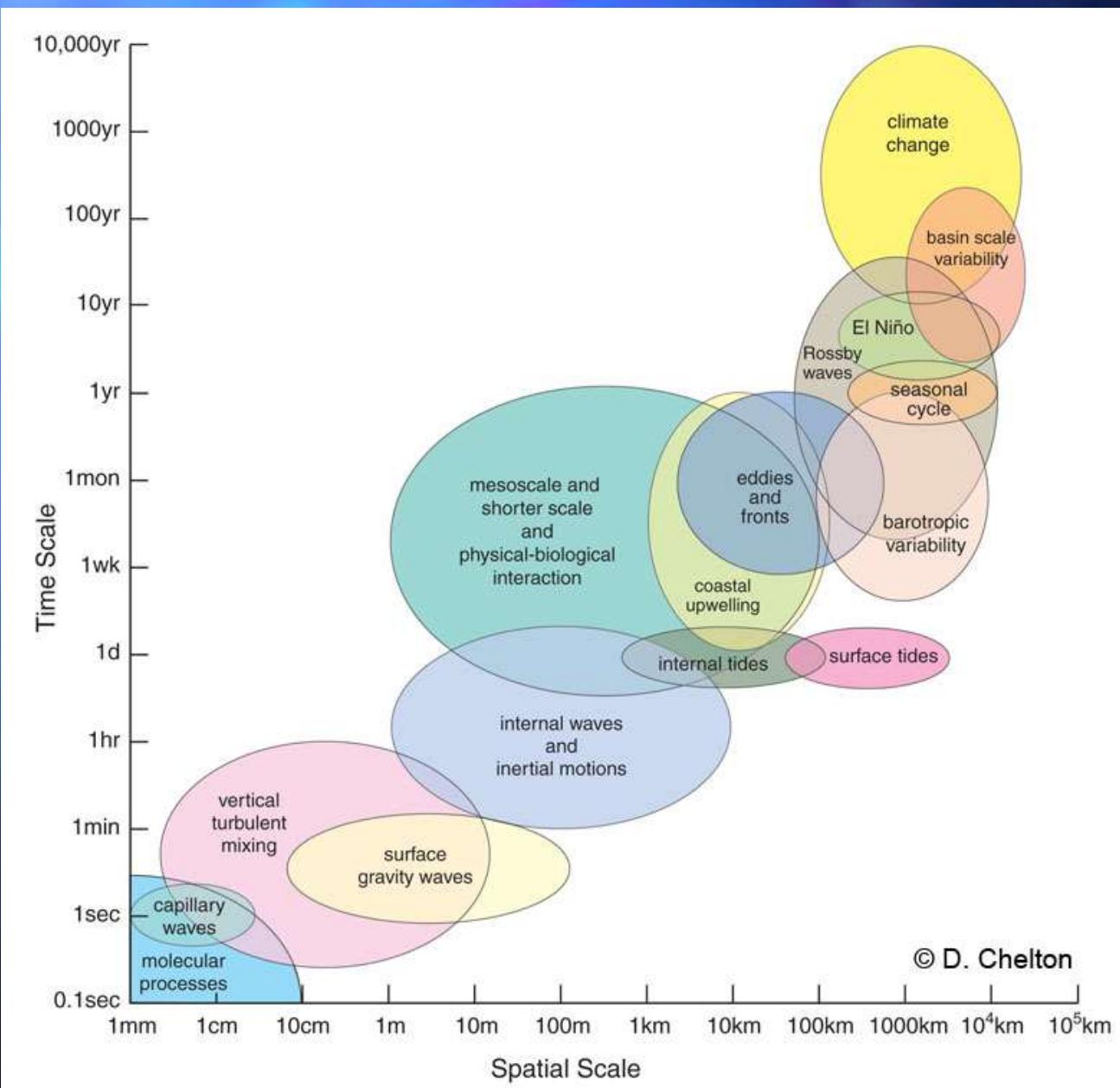


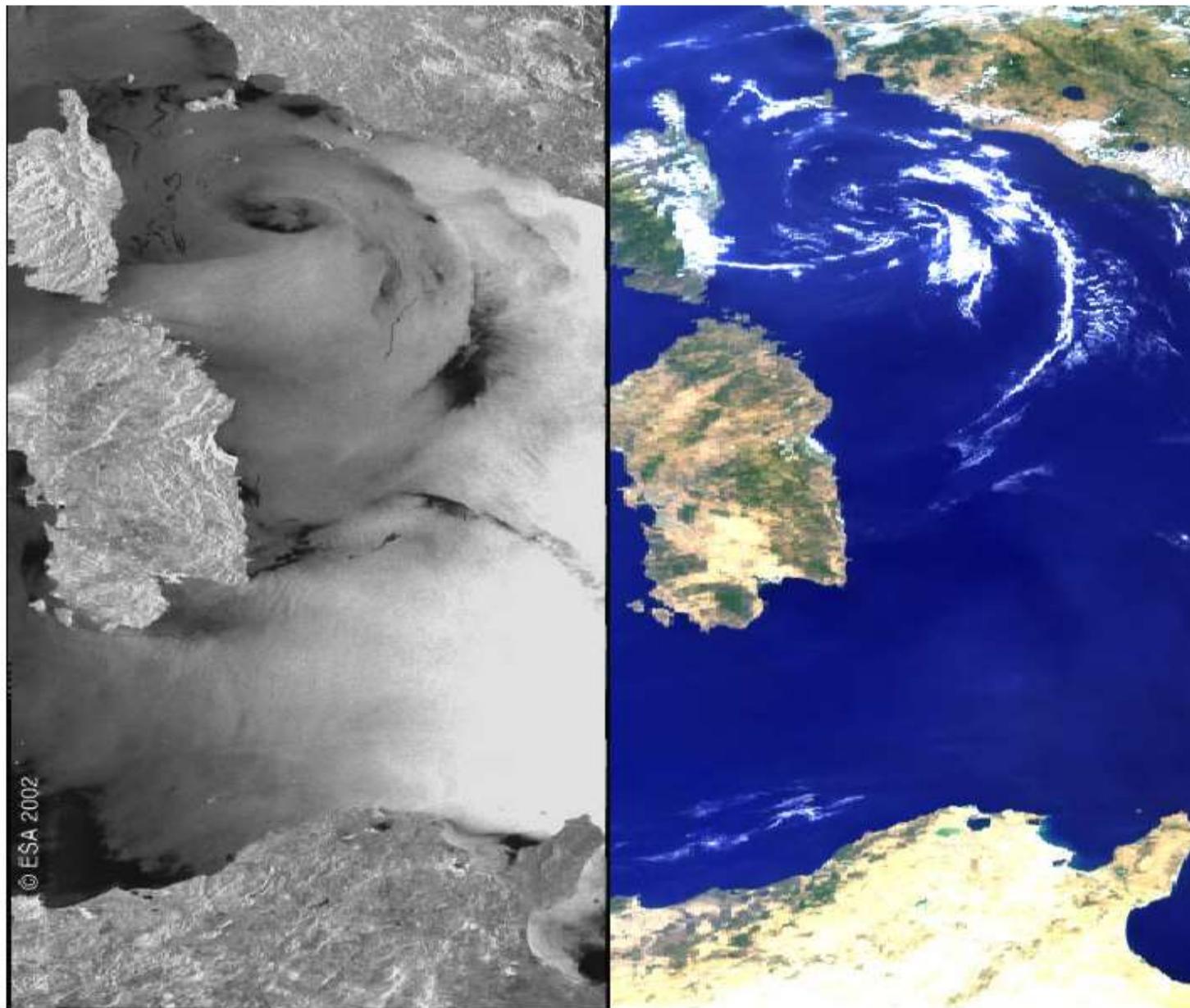
FIG. 3. Instantaneous surface relative vorticity  $\zeta = v_x - u_y$  in the region of the Gulf Stream at the end of winter as simulated by ROMS. The parent domain ROMS0 ( $\Delta x \approx 6$  km) covers most of the Atlantic Ocean (not shown). The boundaries of the successive nested domains ROMS1 ( $\Delta x = 2.5$  km) and ROMS2 ( $\Delta x = 750$  m) are delineated by thick black lines. The relative vorticity plotted



Champ de vitesse déduit de la hauteur de la mer (TOPEX-Poseidon)

# Time / Space scaling

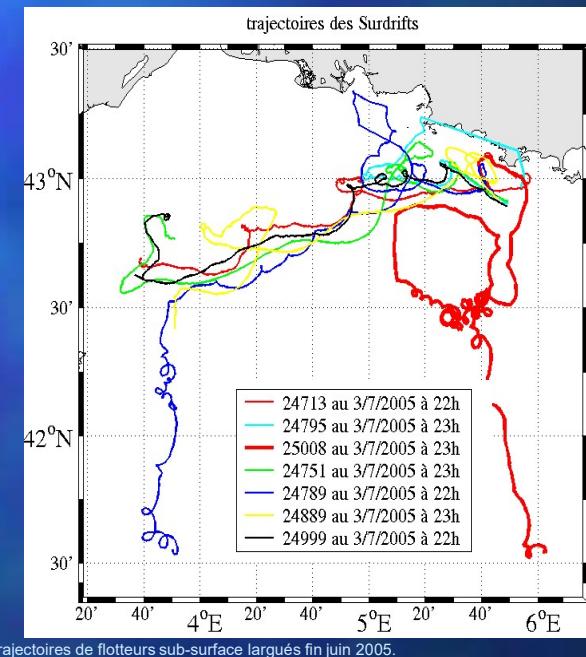
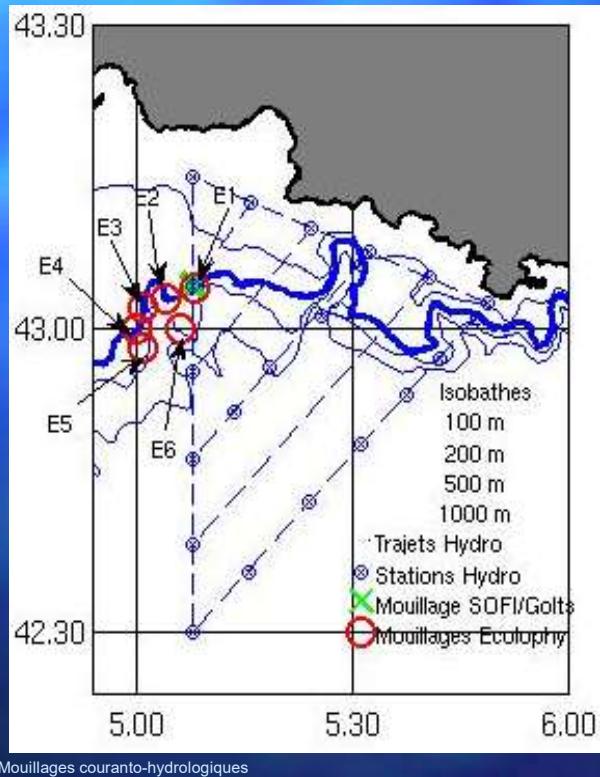


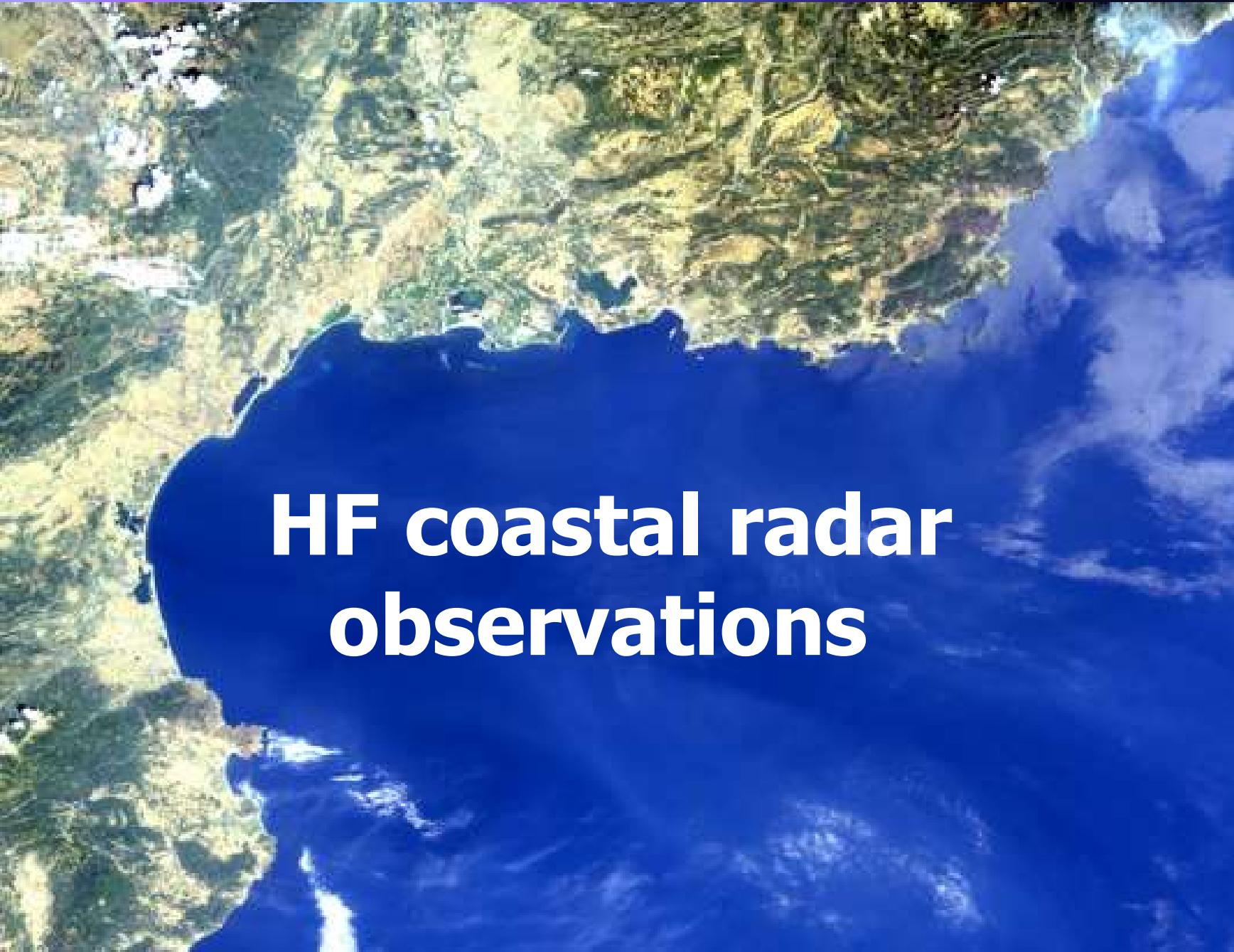


**Comparación de imagen satélite de SAR y de MERIS (visible) de ENVISAT (ESA).**

# Field experiments 2005

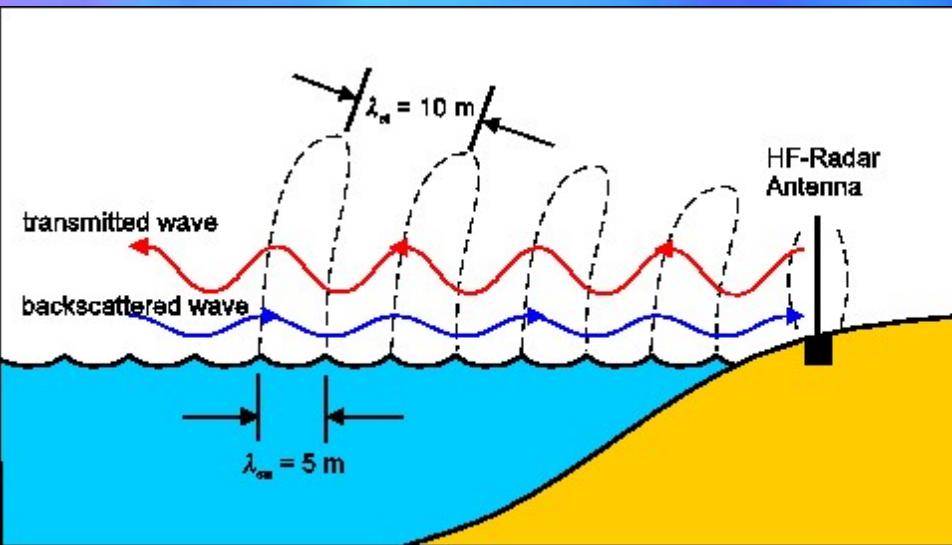
## ECOLOPHY cruise RV THETYS II INSU



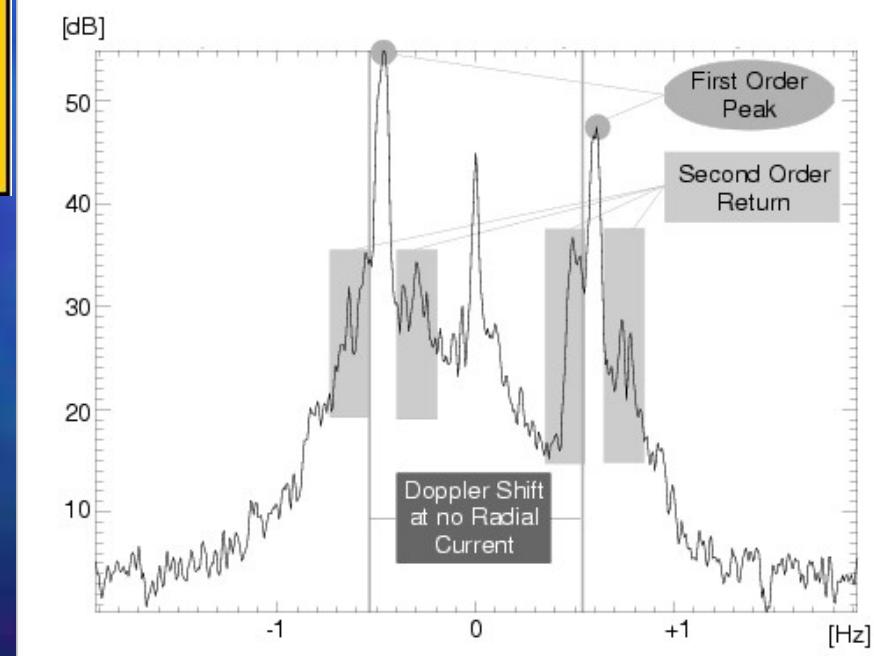


**HF coastal radar  
observations**

# |Coastal radar WERA : how does it work, what do we measure?

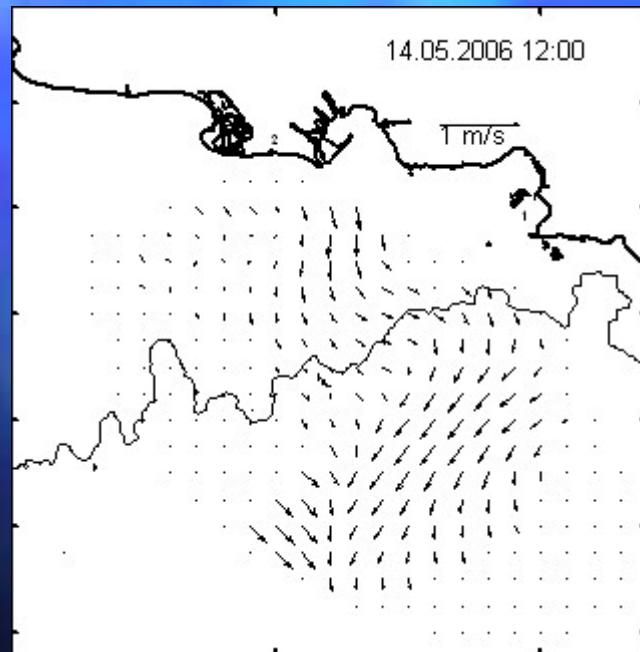


<http://ifmaxp1.ifm.uni-hamburg.de/WERA.shtml>

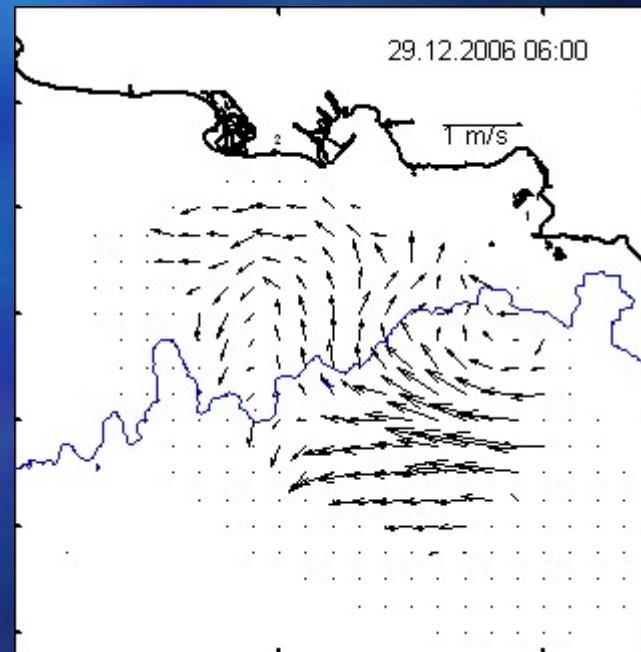


# Radar observations

Inertial circulation



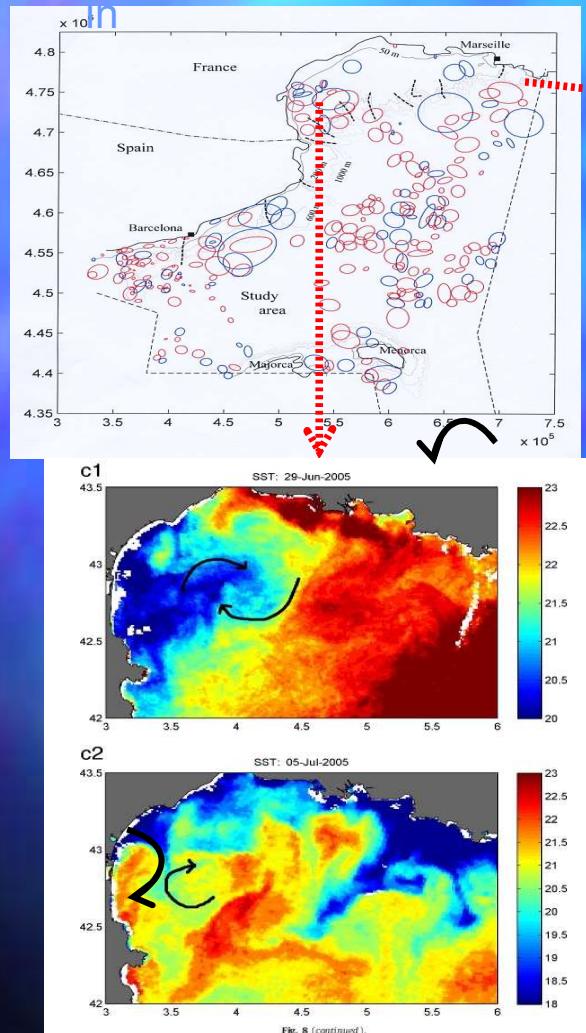
Anticyclonic circulation



By courtesy of P. Forget, Y. Barbin (LSEET)

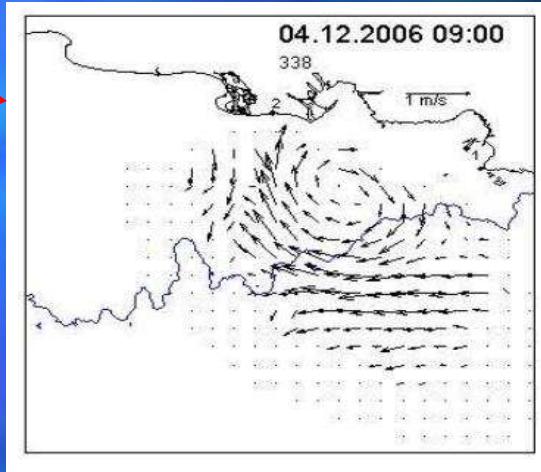
# Observational evidences...

From SAR images...  
(Fedorov et al., 2007,



From SST images...  
(Rubio et al., 2009)

From HF radar experiment,  
the ECOLO campaign...  
(Forget et al.,



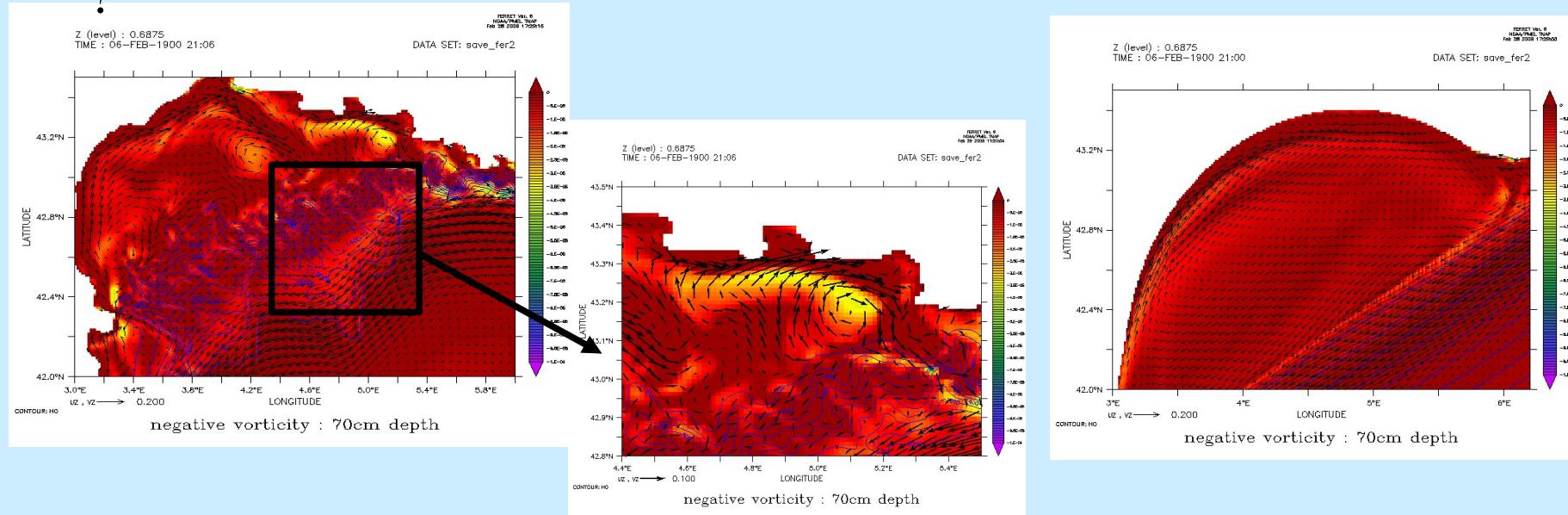
## ... main problematics

- ✓ eddies and vortex activity in the coastal zone
- ✓ interactions with the basin-scale circulation (the NC)
- ✓ toward better estimate of inshore/offshore exchanges

# Driving mechanisms of eddy formation

Respective effects of wind (inertia) / bathymetry

?



Similar length-scale and position

- > northern winds pulses (Mistral, Tramontane)
- > realistic bathymetry and coast line

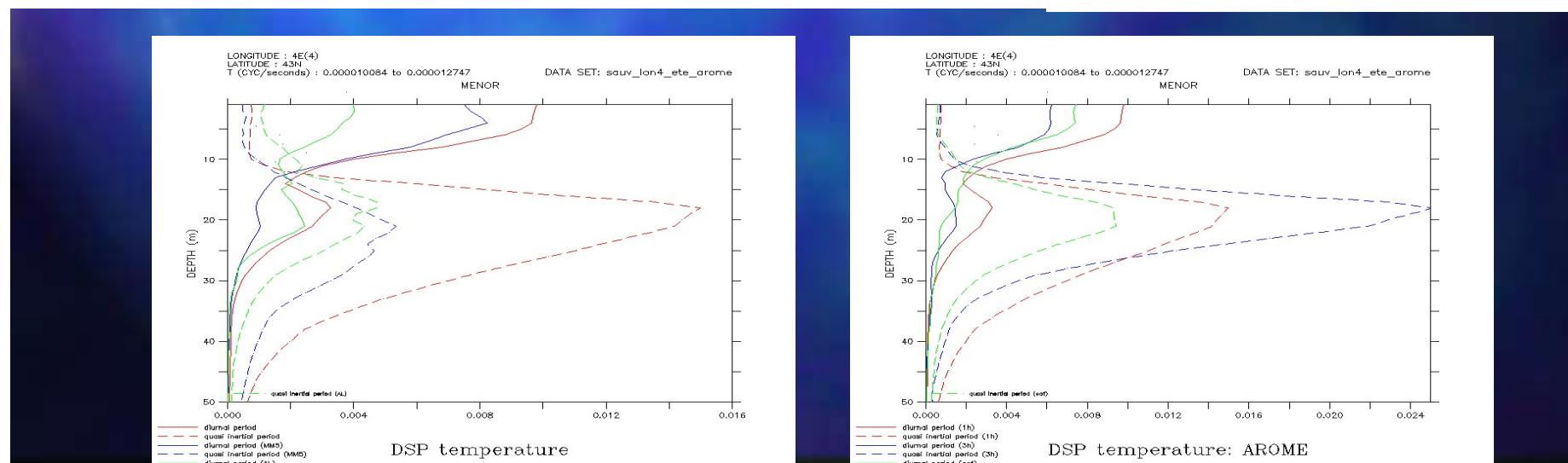
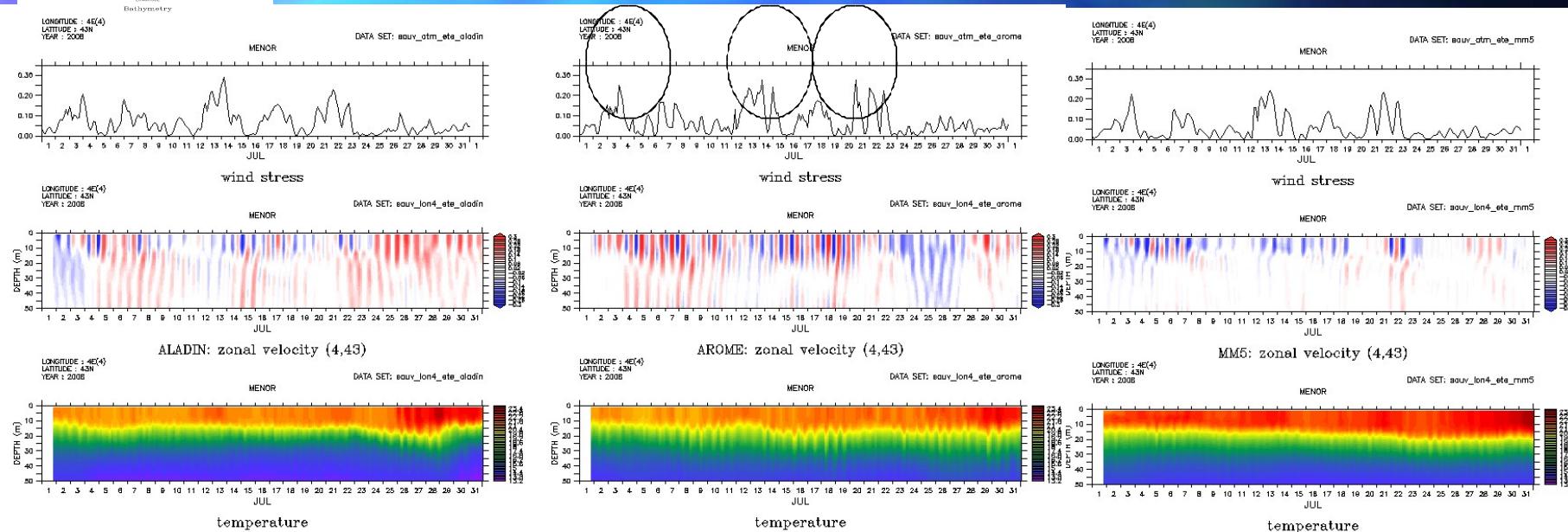
Other possible forcings :

- Liguro-Provencal current intrusions on the shelf
- Rhône river plume
- Thermohaline

-> interactions



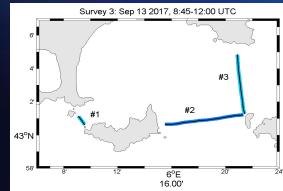
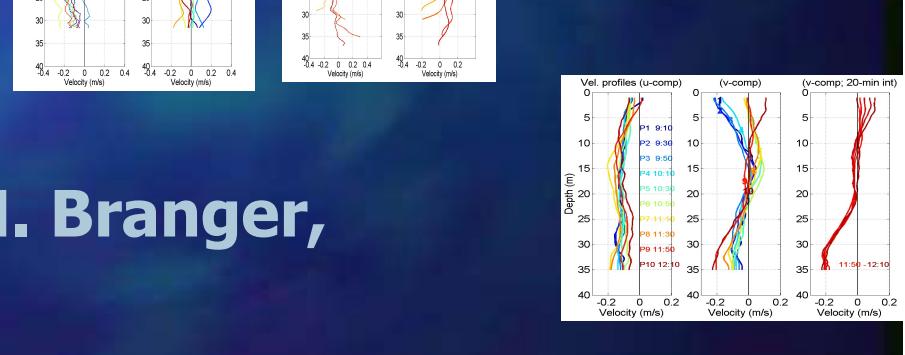
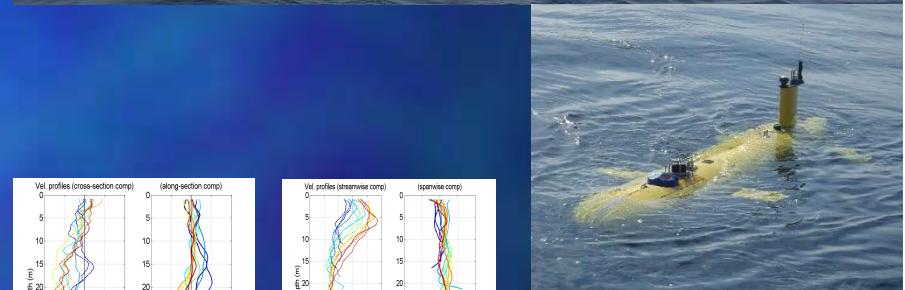
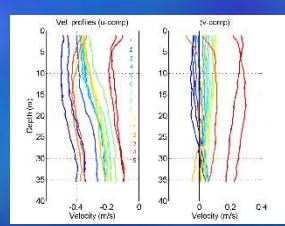
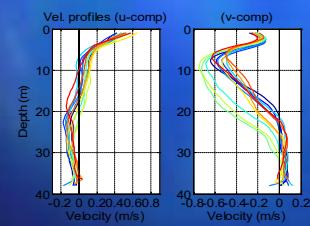
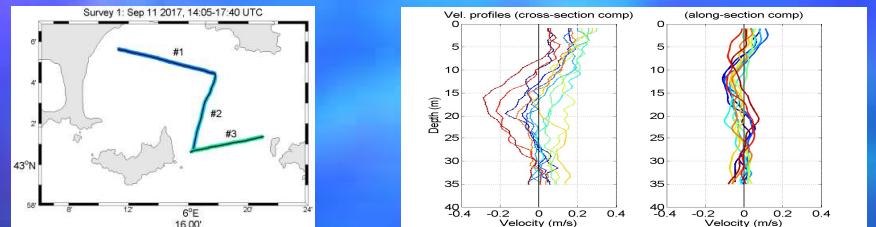
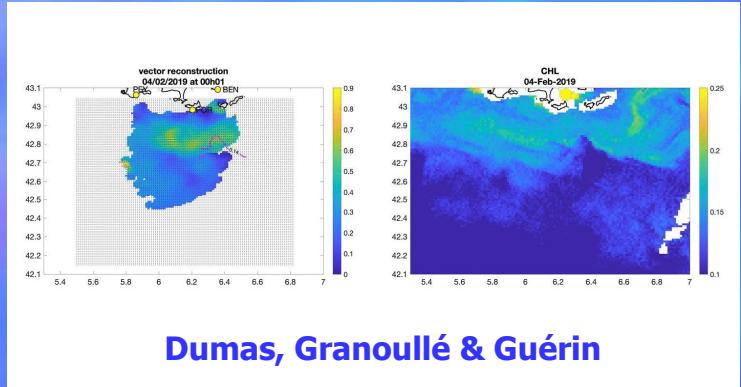
# Internal waves: summer 2008



# OCARINA (D. Bourras)

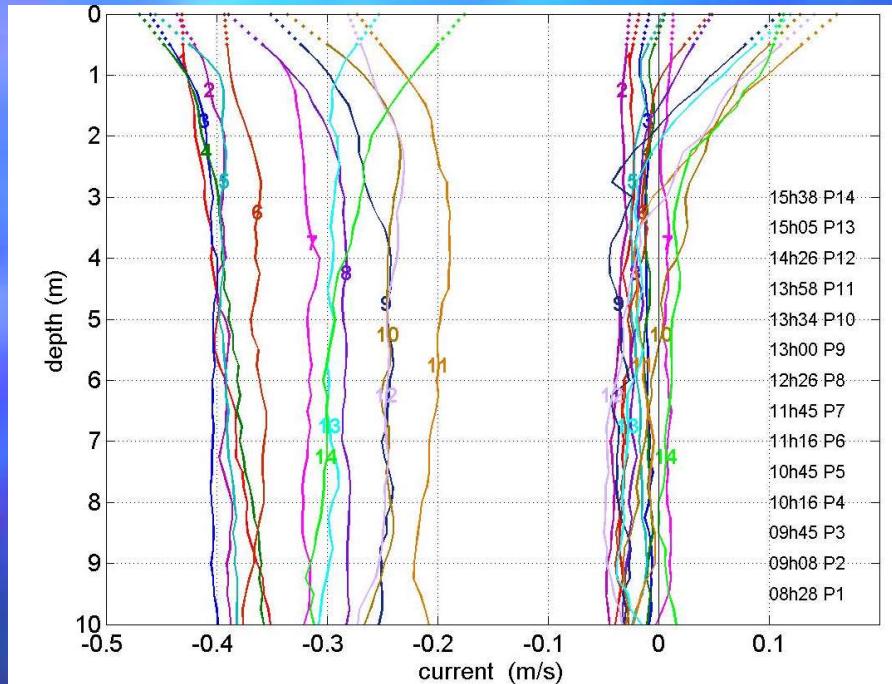


# TURBIDENT campaign 2018



D. Bourras, H. Branger,  
A. Sentchev

# SUBCORAD campaign 2013



Posters EGU2018-5188,  
EGU2018-14417

# Turbulent diffusion (J.M. Redondo)



Figura 3-5. Imágenes de la rotura del oleoducto de BP en el Golfo de México



**Figura 3-3. Accidente del Prestige**

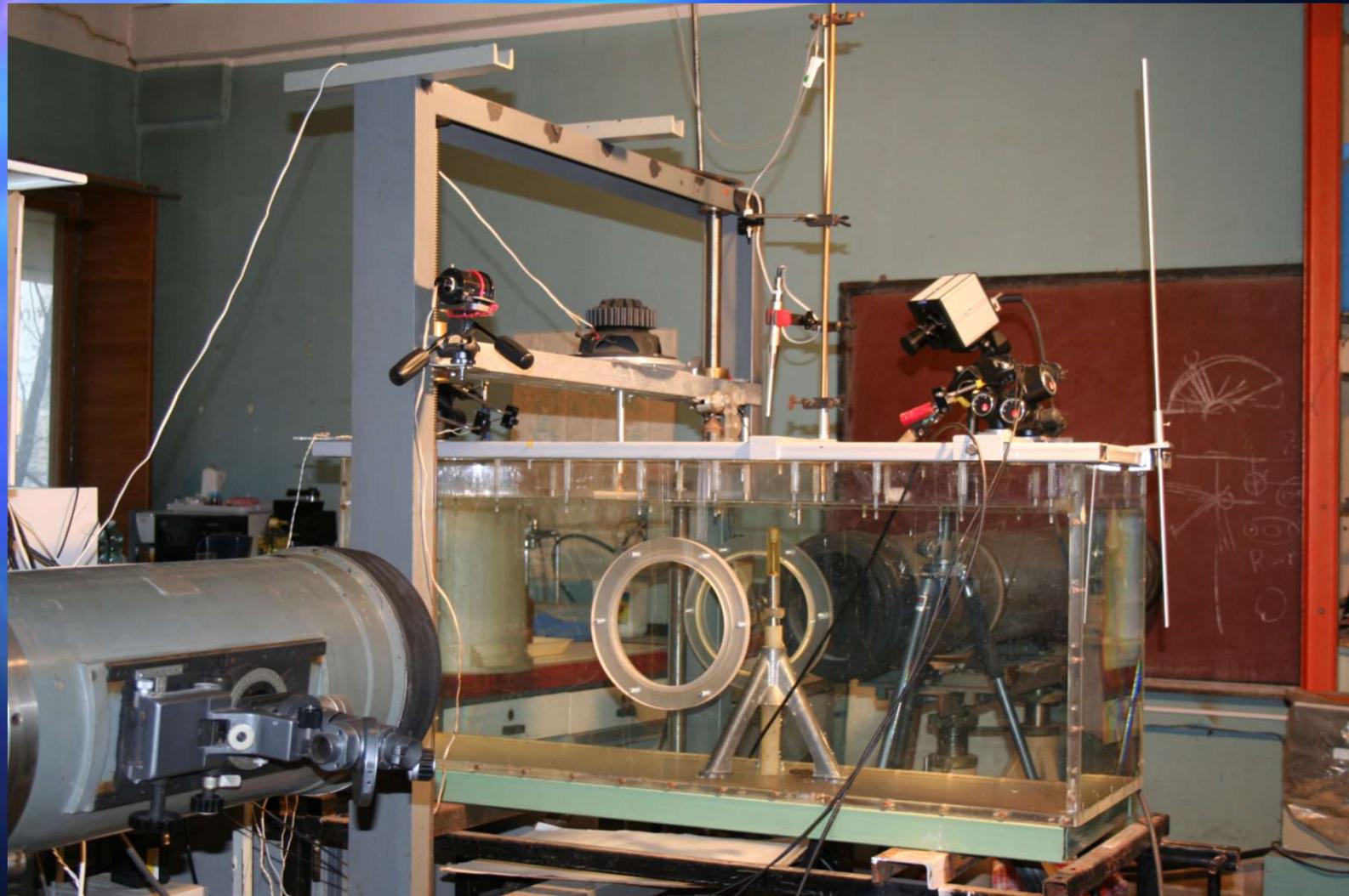
El buque frigorífico 'Sierra Nava' quedó varado en la bahía de Cádiz el 28 de enero de 2007 al intentar entrar en el puerto de Algeciras en medio de una tormenta, con 350 tn, afectando 10 km de costa del Estrecho.



**Figura 3-4 El Sierra Nava varado en la playa.**

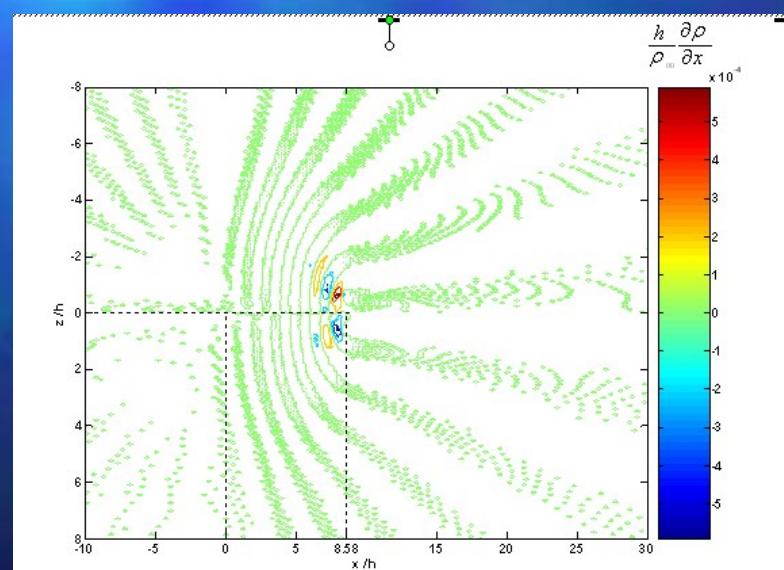
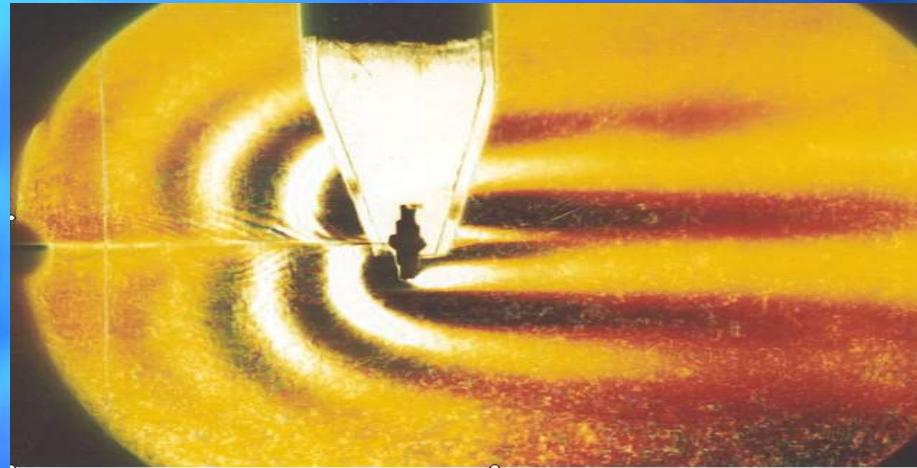
# Turbulence scales

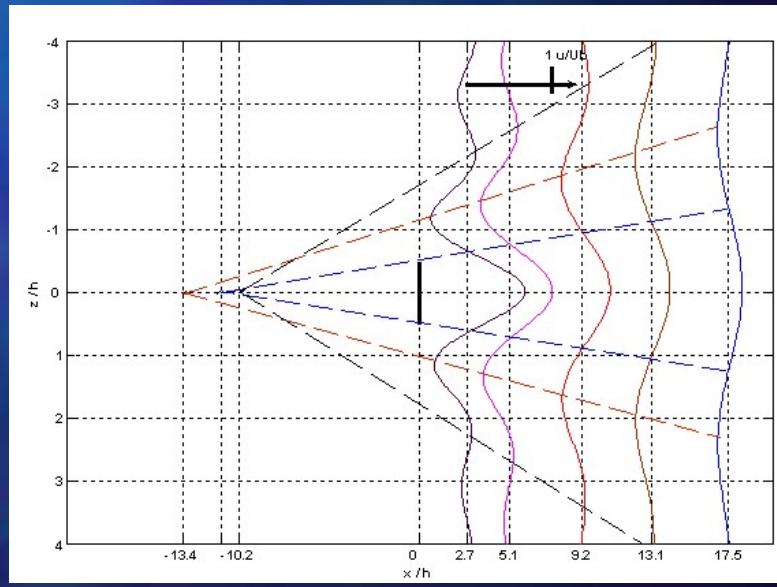
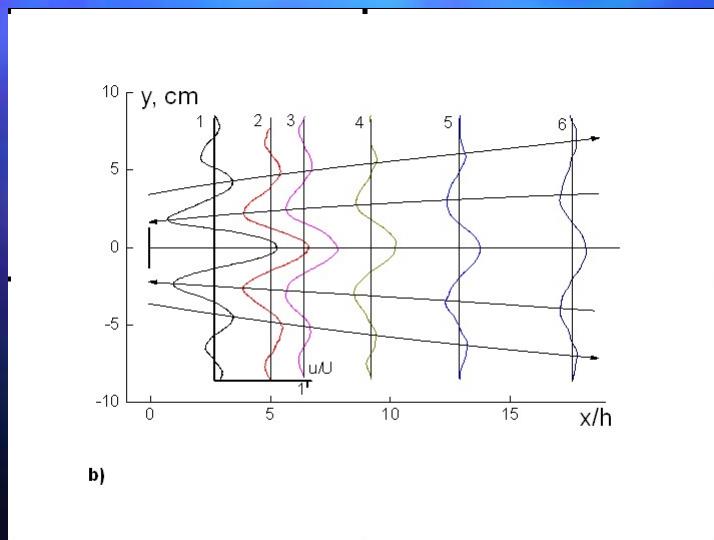
	Atmosphère	Laboratoire Soufflerie	Canal	Océan
$\nu = \mu/\rho \text{ (m}^2\text{s}^{-1}\text{)}$	$10^{-5}$ en tropo. /1 en méso.	$10^{-5}$	$10^{-6}$	$10^{-6}$
$Re = U \cdot L/\nu$	$10^8 - 10^9$ Tropo. Strato. / $10^4$ méso.	$(410^3 - 2510^3)^2$	$(10^3)^1$	$10^4 - 10^7$
$N_{BV} = \frac{g}{\rho} \frac{\partial \rho}{\partial z} \text{ (s}^{-1}\text{)}$	$10^{-2} - 210^{-2}$	$(2.5)^2$	$(0.24 - 0.98)^1$	$(0.1 - 1.74)^4$
$\epsilon \text{ (m}^2 \cdot \text{s}^{-3}\text{)}$	$(10^{-7} - 10^{-4})^7$ $(10^{-5} - 10^{-2})^5$ $(3 - 710^{-2})^8$	$(10^{-4} - 1)^1$	$(10^{-6} - 10^{-3})^2$	$(10^{-10} - 10^{-8})^3$ thermo. $(10^{-7} - 10^{-4})^3$ fjord $(10^{-8} - 10^{-6})^4$
$\eta = \left( \frac{\nu^3}{\epsilon} \right)^{1/4} \text{ (m)}$	$(10^{-2} - 10^{-3})^6$	$(310^{-4} - 1710^{-4})^2$	$(210^{-4} - 710^{-4})^1$	$(310^{-3} - 10^{-2})^3$ thermo. $(310^{-4} - 210^{-3})^3$ fjord $(10^{-3})^4$
$l_o = \sqrt{\left( \frac{\epsilon}{N_{BV}^3} \right)} \text{ (m)}$	10-100	$(210^{-3} - 0.2)^2$	$(510^{-3} - 0.6)^1$	$(1 - 4)^4$
$l_T = \sqrt{\langle d^2(z) \rangle} \text{ (m)}$	$(25 - 120)^7$	$(10^{-2})^2$	$(10^{-2})^1$	$1 - 4)^4$



Mo

# Lab Exp / Numerics





# NUMERICAL SIMULATION OF NEAR WAKE STRATIFIED FLOW AT LOW REYNOLDS NUMBER



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(3) Institute for Problems in Mechanics of the Russian Academy of Sciences, Moscow, Russia.

(4) Department of Applied Physics - U.P.C. Campus Nord, Barcelona 08034 Spain.



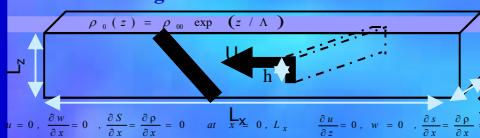
## Abstract

Numerical modeling of a flow past vertical and horizontal strips towed in a linearly stratified tank are performed by comparison to laboratory experiments using Schlieren visualization, density marker and probe measurements of internal wave fields. Both parts of the wave fields including upstream transient and downstream stationary waves were resolved. Analysis is here focusing on observed near wake singular components.

## Key Words

Stratified flows, wakes, internal waves, Large eddy simulation, thin strip

### Geometric configuration



### Dimensionless parameters

$$\begin{aligned} \text{Re} &= h / \delta_v = U_{b,0} h / v \\ \text{Fr} &= \lambda / 2\pi h = U_{b,0} / Nh \\ C_\lambda &= \Lambda / h \end{aligned}$$

### Mathematical formulations

#### Governing equations to the Boussinesq approximation

$$\begin{aligned} \frac{\partial \bar{u}_j}{\partial x_j} &= 0 \\ \frac{\partial \bar{u}_j}{\partial t} + \bar{u}_j \frac{\partial \bar{u}_i}{\partial x_i} - \frac{1}{\rho_0} \frac{\partial \bar{p}}{\partial x_i} + \frac{\rho}{\rho_0} g \delta_{ij} + 2 \frac{\partial}{\partial x_j} [(\nu + \nu_t) \bar{v}_y] &= -\bar{u}_j \frac{\partial \bar{S}}{\partial x_i} - \frac{\partial}{\partial x_i} \left[ (\kappa + \kappa_t) \frac{\partial \bar{S}}{\partial x_j} \right] \\ \frac{\partial \bar{S}}{\partial t} &= -\bar{u}_j \frac{\partial \bar{S}}{\partial x_i} + \frac{\partial}{\partial x_i} \left[ (\kappa + \kappa_t) \frac{\partial \bar{S}}{\partial x_j} \right] \\ \text{Smagorinsky model: } \nu_t &= (C_s \delta)^2 |\tau| \end{aligned}$$

$$\tau_{ij} = \frac{1}{2} \left( \frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} \right)$$

$$\rho = \rho_0 [1 + \beta_s (\bar{S} - \bar{S}_0)]$$

$$|\tau| = |2 (\epsilon_y \tau_y)|^{1/2}$$

### Thin vertical strip

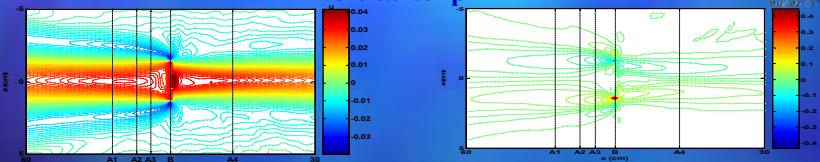


Fig.1. Isotachs of horizontal velocity; Tb = 17.4 s, N = 0.36 s<sup>-1</sup>, U = 0.033 cm/s, Fr = 0.036, Re = 8.25.

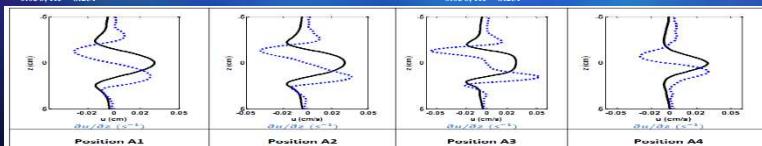


Fig.2. Iso-gradient lines of horizontal velocity; Tb = 17.4 s, N = 0.36 s<sup>-1</sup>, U = 0.033 cm/s, Fr = 0.036, Re = 8.25.

Fig.3. Profiles of the velocity u and its shear du/dz; Tb = 17.4 s, N = 0.36 s<sup>-1</sup>, U = 0.033 cm/s, Fr = 0.036, Re = 8.25.

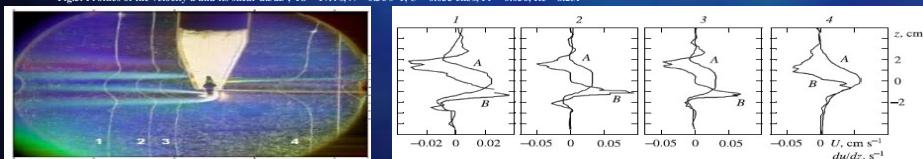


Fig.4. Shadowgraph of the leading disturbance and the plate wake with density markers (1-4) and the corresponding profiles of the velocity u (A) and its shear du/dz (B); Tb = 17.4 s, N = 0.36 s<sup>-1</sup>, U = 0.033 cm/s, Fr = 0.036, Re = 8.25.

### Thin horizontal strip

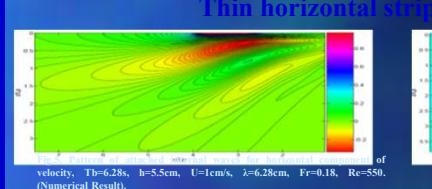


Fig.5. Numerical result for the wake of a thin horizontal strip. Tb = 6.28 s, h = 5.5 cm, U = 1 cm/s, λ = 6.28 cm, Fr = 0.18, Re = 550. (Numerical Result).

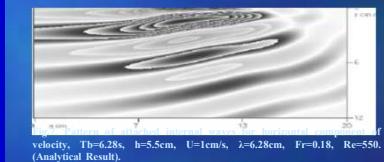


Fig.6. Numerical result for the wake of a thin horizontal strip. Tb = 6.28 s, h = 5.5 cm, U = 1 cm/s, λ = 6.28 cm, Fr = 0.18, Re = 550. (Analytical Result).

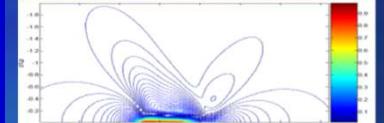


Fig.7. Numerical result for the wake of a thin horizontal strip moving from right to left, Tb = 14 s, h = 2 cm, U = 1 cm/s, λ = 14 cm, Fr = 1.12, Re = 200. (Numerical Result).

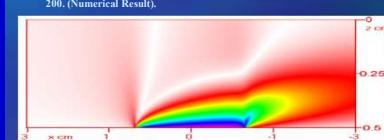


Fig.8. Numerical result for the wake of a thin horizontal strip moving from right to left, Tb = 14 s, h = 2 cm, U = 1 cm/s, λ = 14 cm, Fr = 1.12, Re = 200. (Numerical Result).

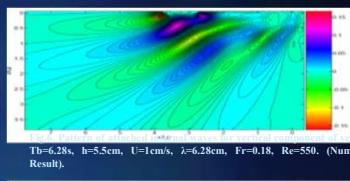


Fig.9. Numerical result for the wake of a thin horizontal strip. Tb = 6.28 s, h = 5.5 cm, U = 1 cm/s, λ = 6.28 cm, Fr = 0.18, Re = 550. (Numerical Result).

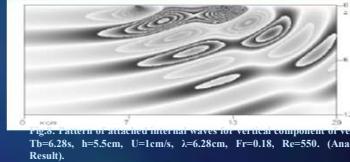


Fig.10. Numerical result for the wake of a thin horizontal strip moving from right to left, Tb = 6.28 s, h = 5.5 cm, U = 1 cm/s, λ = 6.28 cm, Fr = 0.18, Re = 550. (Analytical Result).

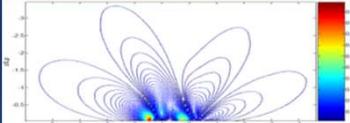


Fig.11. Numerical result for the wake of a thin horizontal strip moving from right to left, Tb = 14 s, h = 2 cm, U = 1 cm/s, λ = 14 cm, Fr = 1.12, Re = 200. (Numerical Result).

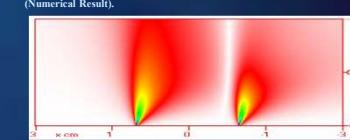


Fig.12. Numerical result for the wake of a thin horizontal strip moving from right to left, Tb = 14 s, h = 2 cm, U = 1 cm/s, λ = 14 cm, Fr = 1.12, Re = 200. (Analytical Result).

### Conclusions

- The present work provided useful physical insights on internal wave generation in a towed body system.
- The motion of flow near moving horizontal strips agrees with analytical results, attached internal waves near the body and details of the fine structure of the boundary layer are distinguished, one peak band for the horizontal component and two bands for the vertical component are clearly visualized in the flow around the horizontal strip.

### References

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# Conclusions

- New interest in sub-mesoscale and stratified flow for climate and safety purposes
- Need of lab experiments and HR observations
- Large (even limited) computing facilities allow to value previous theoretical and experimental works